

Transition from stress-driven to thermally activated stress relaxation in metallic glassesJ. C. Qiao,¹ Yun-Jiang Wang,^{2,3,*} L. Z. Zhao,⁴ L. H. Dai,^{2,3} D. Crespo,⁵ J. M. Pelletier,⁶ L. M. Keer,⁷ and Y. Yao^{1,†}¹*School of Mechanics, Civil Engineering and Architecture, Northwestern Polytechnical University, Xi'an 710072, China*²*State Key Laboratory of Nonlinear Mechanics, Institute of Mechanics, Chinese Academy of Sciences, Beijing 100190, China*³*School of Engineering Science, University of Chinese Academy of Sciences, Beijing 101408, China*⁴*Institute of Physics, Chinese Academy of Science, Beijing 100190, China*⁵*Departament de Física, EETAC, Universitat Politècnica de Catalunya, 08860-Castelldefels, Barcelona, Spain*⁶*Université de Lyon, MATEIS, UMR CNRS5510, Bat. B. Pascal, INSA-Lyon, F-69621 Villeurbanne cedex, France*⁷*Department of Mechanical Engineering, Northwestern University, 2145 Sheridan Rd., Evanston, Illinois 60208, USA*

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The short-range ordered but long-range disordered structure of metallic glasses yields strong structural and dynamic heterogeneities. Stress relaxation is a technique to trace the evolution of stress in response to a fixed strain, which reflects the dynamic features phenomenologically described by the Kohlrausch-Williams-Watts (KWW) equation. The KWW equation describes a broad distribution of relaxation times with a small number of empirical parameters, but it does not arise from a particular physically motivated mechanistic picture. Here we report an anomalous two-stage stress relaxation behavior in a $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$ metallic glass over a wide temperature range and generalize the findings in other compositions. Thermodynamic analysis identifies two categories of processes: a fast stress-driven event with large activation volume and a slow thermally activated event with small activation volume, which synthetically dominates the stress relaxation dynamics. Discrete analyses rationalize the transition mechanism induced by stress and explain the anomalous variation of the KWW characteristic time with temperature. Atomistic simulations reveal that the stress-driven event involves virtually instantaneous short-range atomic rearrangement, while the thermally activated event is the percolation of the fast event accommodated by the long-range atomic diffusion. The insights may clarify the underlying physical mechanisms behind the phenomenological description and shed light on correlating the hierarchical dynamics and structural heterogeneity of amorphous solids.

DOI: [10.1103/PhysRevB.94.104203](https://doi.org/10.1103/PhysRevB.94.104203)**I. INTRODUCTION**

As a relatively new family of amorphous materials, metallic glasses (MGs) exhibit excellent physical and chemical properties. MGs are usually obtained by rapid quenching from their melts, and show distinctive mechanical behaviors beyond the conventional oxide and polymer glasses due to the unique dense packing of disordered atomic structure [1–4]. Structural inhomogeneity [5–10] usually leads to dynamic heterogeneity [11–15] which is closely correlated with the viscoelasticity and plasticity of MGs [16], both facts of engineering importance and scientific curiosity. Unraveling the atomic-scale structure-dynamic relationship is a paramount challenge [17–22] as well as a pressing necessity towards eventual applications of MGs [23].

Stress relaxation is a technique to trace the temporal variation of stress in response to a fixed magnitude of strain generated in the structure, thus reflecting the characteristics of viscoelastic and plastic deformation of the material. It is an important clue to link the dynamic responses to the intrinsic atomic-level structure, which in turn determines the thermodynamic status (aged or rejuvenated) [24,25] and consequently the macroscopic mechanical performance of MGs [26–29]. Considerable research has been done on the stress relaxation nature of various MGs [30–34], which demonstrated that localized plastic flow could be activated

during the viscoelastic and plastic deformation induced by stress relaxation. Stress relaxation induces a simultaneous excitation at all timescales, thus allowing one to understand the fundamental science behind the potential applications as an engineering material [30–34]. However, the understanding of the microscopic mechanisms accommodating the stress relaxation in amorphous alloys is limited. In particular, the physical connection between the macroscopic stress relaxation behavior and the structural relaxation (primary α , and secondary β) usually explored by the internal friction experiment remains quite mysterious and elusive [22,35,36].

The whole relaxation spectrum of viscoelastic materials is commonly fitted to the empirical stretched exponential Kohlrausch-Williams-Watts (KWW) function [30–34]. However, the stretched exponential is only a phenomenological description of the behavior and does not arise from a particular physically motivated mechanistic picture. It describes a broad distribution of relaxation times, but with a small number of empirical fitting parameters to describe the distribution. Here, experiments uncover an anomalous but universal two-stage stress relaxation in a prototypical $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$ glassy system. Thermodynamic and discrete analysis on the stress relaxation spectra rationalize and explain the transition from fast stress-driven to slow thermally activated mechanisms with stress decreasing in stress relaxation. Atomic simulations clarify the atomic-scale mechanisms accommodating stress relaxation and demonstrate their competition. The comprehensive insights from combined experiment, theory, and simulations drive one to reexamine the complex correlation between stress, energetic, and structural relaxations of amorphous alloys,

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and may rationalize the physical mechanisms underlying the empirical KWW function.

II. EXPERIMENTS AND THEORY

A. Stress relaxation experiment

Benefitting from the excellent thermal stability of the CuZr-based MG, in our work, MG ribbons (thickness $\sim 30 \mu\text{m}$ and width $\sim 1 \text{ mm}$) with atomic composition $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$ were processed by the melt-spinning technique in an inert argon atmosphere. The thermal properties of the samples were determined by differential scanning calorimetry (DSC), where the glass transition temperature is 697 K at a heating rate of 20 K/min.

The tensile stress relaxation experiments were performed using a TA Q800 dynamic mechanical analyzer (DMA). In order to rule out any influences from initial states, the MG ribbons were heated above the glass transition temperature prior to the measurements. The tensile stress relaxation, carried out at a constant strain of 0.4%, was loaded on the model alloy for 30 minutes after an initial 3 min equilibration.

B. Discrete analysis

In order to determine the presence of discrete relaxation times, the relaxation response was fitted to a distribution of discrete relaxation times given by

$$\begin{aligned} \sigma(t)/\sigma_0 &= \sum_{i=0}^n a_i \exp\left(-\frac{t}{\tau_i}\right) \\ &= a_0\delta(t) + \sum_{i=1}^{n-1} a_i \exp\left(-\frac{t}{\tau_i}\right) + a_\infty. \end{aligned} \quad (1)$$

The relaxation times, namely 0 (instantaneous relaxation), τ_i , and ∞ (purely elastic response) are fixed. We fit a single functional form depending on a discrete intensity coefficient $a(\tau)$. This is a variational fitting method using nonorthogonal fitting functions, which is only convergent with very small relative tolerances. In order to span it on a proper timescale, the successive relaxation times were chosen in logarithmic ratio, that is $\frac{\tau_i}{\tau_{i-1}} = ct$. Thus the fitting parameters are only the $n + 1$ non-negative intensity coefficients, $a_0 \dots a_n$. The minimization was performed by using the sequential quadratic programming method implemented in Maple [37], with $\varphi = 1.1$ and $n = 151$, and a relative tolerance lower than 10^{-11} . The smooth form of the fitting functions and the non-negativity of the a_i coefficients allowed us to obtain a low number, physically significant coefficients corresponding to the discrete relaxation times.

C. Atomistic simulations

The MD stress relaxation tests were conducted by LAMMPS [38]. A model glass of $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$ with 13 500 atoms was prepared by a heating-quenching technique with a cooling rate of 10^{10} K/S . The force field was described by an embedded-atom method (EAM) potential [39]. To simulate the relaxation tests, we first pulled the simulation box to the tensile strain of 5% (the yield strain is 6.4%) with a constant

strain rate 10^8 s^{-1} , and then fixed the uniaxial strain to trace the time evolution of stress tensor for 50 ns. The simulations were carried out with an isothermal-isobaric ensemble. The visualization is performed by OVITO [40].

III. RESULTS

A. Anomalous stress relaxation at low temperature

The isothermal stress relaxation spectra of $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$ MG, obtained at increasing temperatures from 393 K to 693 K, are displayed in Fig. 1(a). The relaxation curves were fitted to a phenomenological KWW equation for the entire temperature range. The KWW equation can characterize the evolution of stress as a function of time during the stress relaxation test,

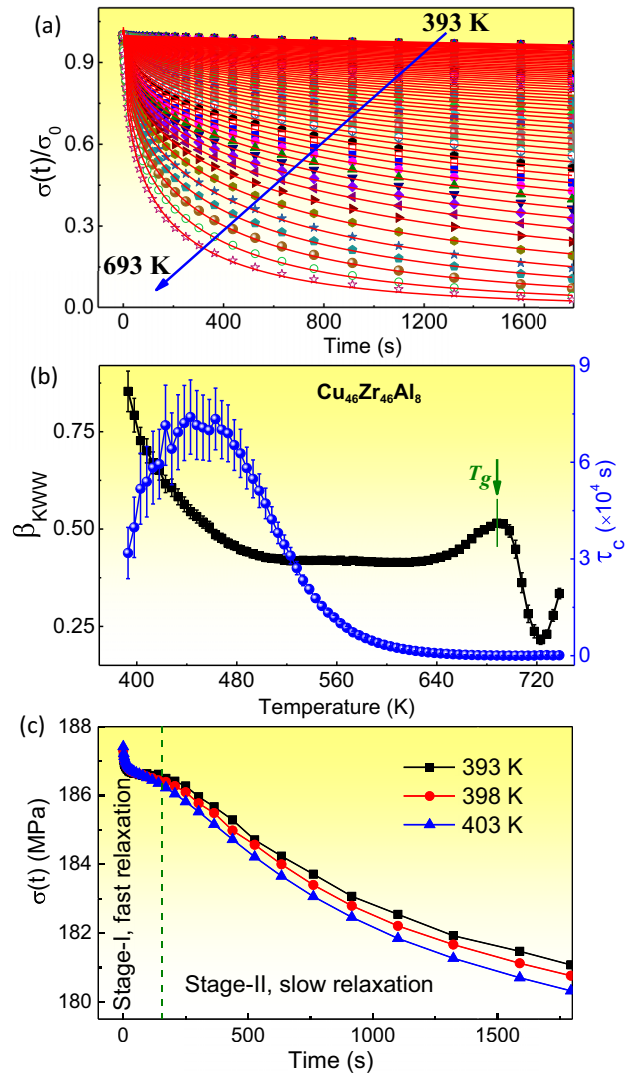


FIG. 1. Stress relaxation of a $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$ MG. (a) Stress-relaxation spectra from 393 to 693 K ($T_g = 697 \text{ K}$). The stress levels have been normalized by the initial stresses. The red lines denote the phenomenological KWW fittings. (b) The fitted KWW parameters β_{KWW} and τ_c as a function of temperature. Variation of τ_c with T indicates an anomalous stress-relaxation regime at low temperatures. (c) Original stress-relaxation curves demonstrate clearly two-stage relaxation, i.e., fast and slow relaxation, at low temperatures.

i.e.,

$$\sigma(t) = \sigma_\infty + \sigma_0 \exp \left[- \left(\frac{t}{\tau_c} \right)^{\beta_{\text{KWW}}} \right], \quad (2)$$

where σ_0 is the initial stress and σ_∞ is a relaxed reference stress at infinite time. Closely related to the structural heterogeneity of glass [30], β_{KWW} is a stretched exponent parameter, and τ_c is a characteristic time that usually separates different relaxation mechanisms [34]. By using the empirical KWW model to fit the stress relaxation data we can compare with previous results and show a qualitative trend of the mechanism transition by tracing the variation of the effective relaxation time τ_c with change in temperature.

The fitted parameters as a function of temperature are shown in Fig. 1(b). The parameter β_{KWW} lies within [0.4, 1] as usually found in the stress relaxation of amorphous alloys [30,33,34]. However, the characteristic time τ_c shows an unexpected behavior, first increasing and then decreasing with increasing temperature, with a critical temperature of ~ 450 K. This is contradictory to the common sensing that relaxation time decreases with increasing temperature for a specific relaxation mechanism. Similar experimental results have been observed in other typical metallic glasses (see Fig. S1 in the Supplemental Material [41]). To examine the anomalous relaxation behaviors, we provide the original stress relaxation spectra of three low temperature cases in Fig. 1(c)

where it is obvious that the stress relaxation at low temperature can be divided into two stages: a fast relaxation and slow relaxation process, with a boundary at ~ 200 seconds. The fast and slow relaxation may indicate different relaxation mechanisms during stress relaxation.

B. Thermodynamic analyses

We use a strategy of thermodynamic analysis to characterize the fast and slow events, which are commonly believed to be thermally activated, and stress-driven rate-controlling processes during stress relaxation (detailed theory in Fig. S2 in the Supplemental Material [41]) [42,43]. The results are summarized in Fig. 2. In Fig. 2(a), the logarithmic minus stress rate is plotted against stress at different temperatures. We plot it in this way since the slope of $\ln(-\dot{\sigma})$ with respect to σ is related to the activation volume of an event as

$$\Omega = \sqrt{3} k_B T \frac{\partial \ln(-\dot{\sigma})}{\partial \sigma}. \quad (3)$$

Here Ω is the activation volume which is proportional to the size of the event, T is temperature, and k_B is the Boltzmann constant. As shown in Fig. 2(a), the slope changes sharply from high to low stress with a critical magnitude at around 186.5 MPa (the critical stress and time of mechanism transition are shown in Fig. S3 in the Supplemental Material [41]). The

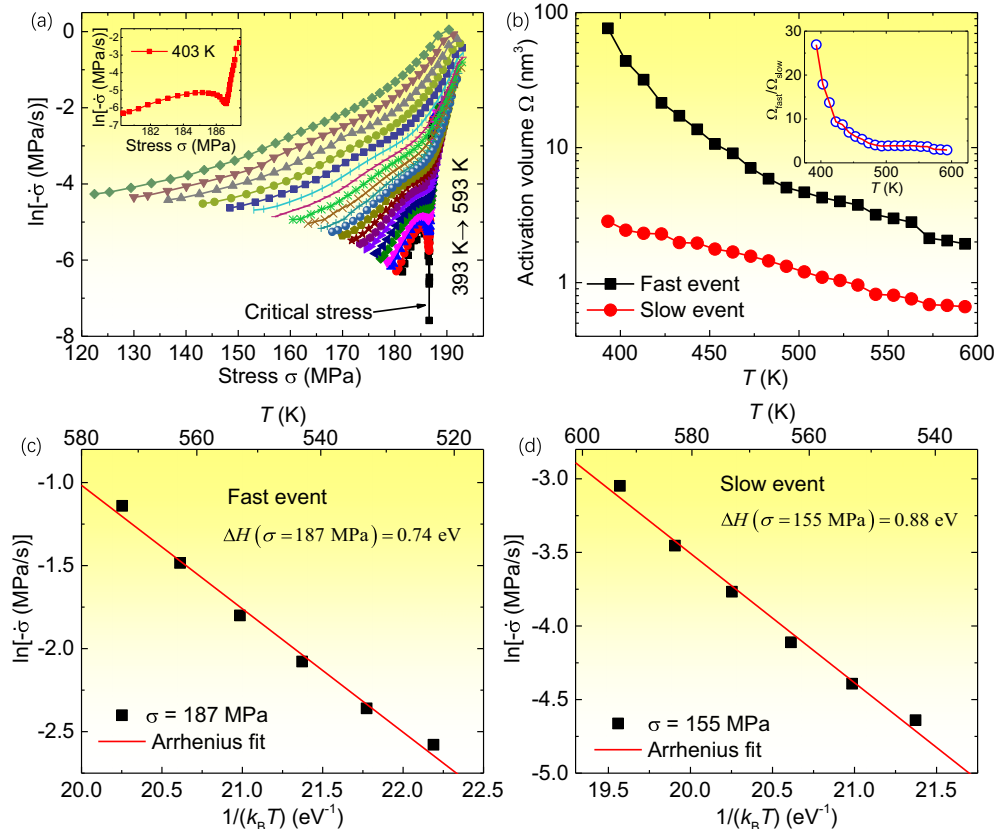


FIG. 2. Transition from stress-driven to thermally activated stress relaxation. (a) Dependence of $\ln(-\dot{\sigma})$ on stress. The inset shows an example at 403 K, which indicates a clear transition of relaxation mechanisms with variation in stress. (b) The derived activation volume of the fast and slow events versus temperature. The inset shows the ratio between them. (c) and (d) Arrhenius plots indicate the fast event presents an activation enthalpy of 0.74 eV at 187 MPa; while it is 0.88 eV at 155 MPa for the slow event.

transition is particularly pronounced at low temperature, as it is exemplified in the case of 403 K shown in the inset of Fig. 2(a). The fast event takes place at the early stage of the stress relaxation process with higher stress level, while the slow event occurs at a lower stress. The calculated activation volumes of the two events are shown in Fig. 2(b) as a function of temperature. A noted reasonable trend is that the activation volumes decrease with increasing temperature, which means that the relaxation event becomes less collective at higher temperature. An interesting fact is that the fast event has a very large activation volume, of the order of $\sim 100 \text{ nm}^3$ at low temperature and decreases rapidly with increasing temperature. But the slow event has a much smaller activation volume of the order of 1 nm^3 , and decreases slightly with increasing temperature compared with that of the fast event. The inset of Fig. 2(b) shows the ratio of activation volumes of the fast and slow event. We find that the ratio decreases with increasing temperature, and becomes nearly constant above a critical temperature of approximately 450 K, which corresponds to the critical temperature of the KWW parameter τ_c as shown in Fig. 1(b) and may indicate that the fast event dominates stress relaxation at the initial stage, while the slow event governs the later one. Finally, there seems to be left a puzzle that a negative activation volume might appear in the transition region of the two categories of events, in particular at low temperature, as shown by the inset of Fig. 2(a). It is not intuitive since a thermally-activated, and stress-driven rate-controlling physical process should have a positive activation volume [44,45]. Actually, this is an artifact since it is not reasonable anymore to deduce the activation volume using Eq. (3) if two or more processes are strongly coupled and hard to decompose.

Furthermore, the activation enthalpies $\Delta H(\sigma) = \Delta Q - (\sigma/\sqrt{3})\Omega$ of the two events at specific stresses are determined by the Arrhenius plot of $\ln(-\dot{\sigma})$ against inverse temperature, as shown in Figs. 2(c) and 2(d), respectively. The unbiased activation energy is ΔQ in the case of $\sigma = 0$. For the fast event, the activation enthalpy is 0.74 eV at 187 MPa, while it is 0.88 eV at 155 MPa for the slow event. Considering that the activation volume at 550 K is about 3 nm^3 , the unbiased activation energy of the big event is of the order of $\Delta Q = \Delta H(\sigma) + (\sigma/\sqrt{3})\Omega = 2.76 \text{ eV}$, which is about $46k_B T_g$ ($T_g = 697 \text{ K}$) for $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$. Thus the activation free energy is significantly reduced by the work done by the stress $(\sigma/\sqrt{3})\Omega$. The fast event is actually a stress-driven process occurring at the initial stage of the stress relaxation (see Fig. S4 in the Supplemental Material for the sketch of stress-driven and thermally activated processes). However, the slow event has an activation energy of 1.2 eV, considering the activation volume of 0.7 nm^3 at 570 K. The activation energy is about $22k_B T_g$, which is of the order of the activation energy for the usual β relaxation (Fig. S5 in the Supplemental Material [41]) in MGs [35,46]. Considering the relatively small activation volume, the free energy barrier is overcome mainly by the thermal energy instead of stress. The slow event is mainly a thermally activated process which takes place on a longer timescale. There exists a transition from the stress-driven to thermally activated mechanism during the stress relaxation. The fast event has higher activation energy but is significantly biased by the work done by stress, while the

slow event is less collective with a smaller activation volume, which is mainly activated by thermal energy.

IV. DISCUSSION

A. Discrete analysis of stress relaxation

Though widely used, the stretched exponential can only describe a unimodal distribution of relaxation times with a main (or average) relaxation time. However, there is experimental evidence that several relaxation times may exist in metallic glasses [47]. The concurrent relaxation events give rise to a multimodal distribution, with well-defined relaxation times (See Fig. S6 in the Supplemental Material for the generalized Maxwell model and foundation of KWW equation). These relaxation times may be different by orders of magnitude, and in this case the fit with a single stretched exponential fitting describes uniquely the dominant relaxation event.

Figure 3 shows the results of the fit of a multimodal distribution to the relaxation data, shown in Fig. 1(a). Two well-defined relaxation times are determined below 475 K, and three relaxation times are obtained above that temperature. In order to show the intensity at each relaxation time τ_i , the radius of the symbol is proportional to the corresponding intensity coefficient a_i . The faster relaxation time is always lower than the threshold between the two relaxation events determined above. The characteristic time of this fast relaxation event is $\sim 10 \text{ s}$, and shows a very weak temperature dependence, confirming it to be an almost instantaneous, stress-driven process. More details of discrete analysis are provided in Figs. S6 and S7 in the Supplemental Material [41].

The slow relaxation times, larger than the threshold time, show a markedly different behavior above and below 670 K. Above 670 K, close to the glass transition temperature, both relaxation times show almost the same temperature dependence, which is fitted to an Arrhenius dependence with an activation energy of 1.47 eV. It was unexpected to find two discrete relaxation times well below the glass transition. To check the validity of this fit, the ‘‘apparent’’ relaxation time τ_{app}

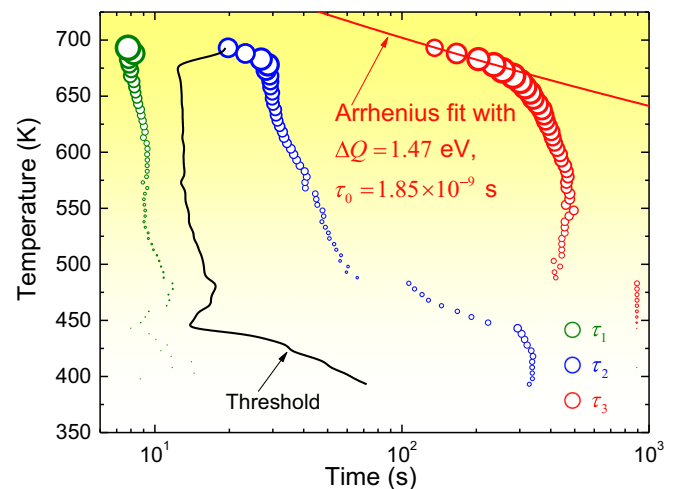


FIG. 3. The threshold time and relaxation times τ_i fitted with Eq. (1). The size of the symbol is proportional to the corresponding intensity a_i .

was computed (Table S1 in the Supplemental Material [41]). It can be seen that the “apparent” relaxation time falls in the scale of the glass transition time window. This result may not be specific to this particular metallic glass, and more experimental work is needed to check whether it appears in other metallic or nonmetallic glasses.

Below 670 K both relaxation times show much weaker temperature dependence. A detailed analysis shows that this behavior is due to the length of the experimental data series. The stress relaxation data was obtained for 600 s, and it can be seen that both relaxation times remain restricted to the experimental window. Consequently, it is deduced that the multimodal fit gives only adequate results while the relaxation times fall within the experimental window. This fact is also responsible for the vanishing of the largest relaxation time below 475 K; below that temperature only an average relaxation time is obtained. However, a noticeable change in the slope of this average relaxation time can be seen at 450 K, which correlates to the critical temperature of the KWW relaxation time.

B. Atomistic understanding of stress relaxation

In order to understand the microscopic mechanisms and their transition during stress relaxation, we carried out molecular dynamics (MD) simulations. The results are summarized in Figs. 4 and 5. The MD stress relaxation curves are displayed in Fig. 4(a). KWW fits are performed on those curves and the

parameters are shown in Fig. 4(b). We find β_{KWW} fluctuates around 0.4, reflecting a considerable distribution of local activation barriers during stress relaxation and resembling quantitative agreement with experiments. It is also interesting to note a variation of τ_c with temperature similar to that found in experiment.

Note that the timescale of MD is orders of magnitude shorter than that of experiment. It could be argued that the relaxation is too fast to be representative of the macroscopic measurement; but the application of stress is also several orders of magnitude higher than in the experiment, and thus the relaxation response of the simulation is accordingly fast. Thus, the qualitative behavior of the simulations is close to that of the stress relaxation, e.g., a reasonable β_{KWW} value indicating the discrete distribution of relaxation times, and the variation in τ_c which also resembles the experimental data.

For a deeper understanding of the thermodynamic status during stress relaxation, we plot the variation of potential energy during tension and stress relaxation in Fig. 4(c). The reference potential energy is taken before loading. Tension introduces extra potential energy to the samples. But the subsequent stress relaxation process would definitely relax the glass structures energetically, and the glass continues to be aged. In particular, at higher temperature and a long enough time, the sample after stress relaxation can reach an even more relaxed state than that of the as-quenched sample. The change in the energetic state would in turn affect the stress relaxation behavior as characterized by the phenomenological KWW

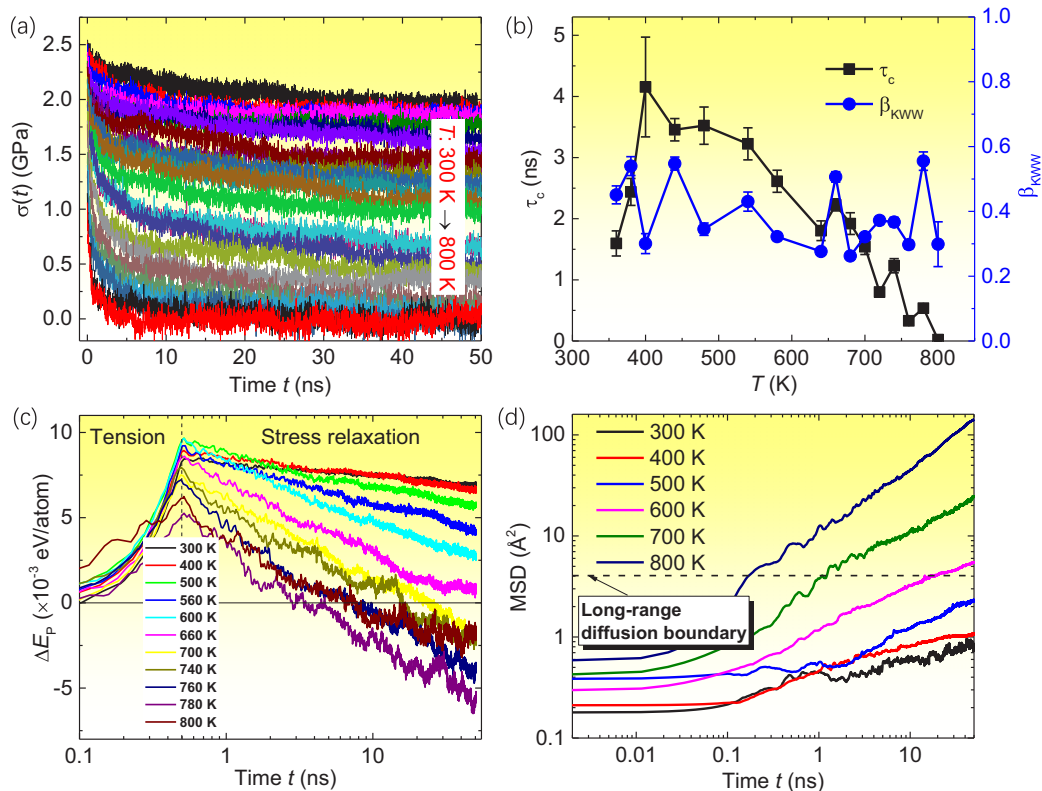


FIG. 4. Atomistic simulations of stress relaxation in $\text{Cu}_{46}\text{Zr}_{46}\text{Al}_8$. (a) The MD stress relaxation curves at various temperatures. (b) The KWW parameters as a function of temperature based MD. (c) Variation in potential energies accompanying the stress relaxation processes. The reference potential energies are those of the samples before loading. (d) MSD as a function of time during stress relaxation.

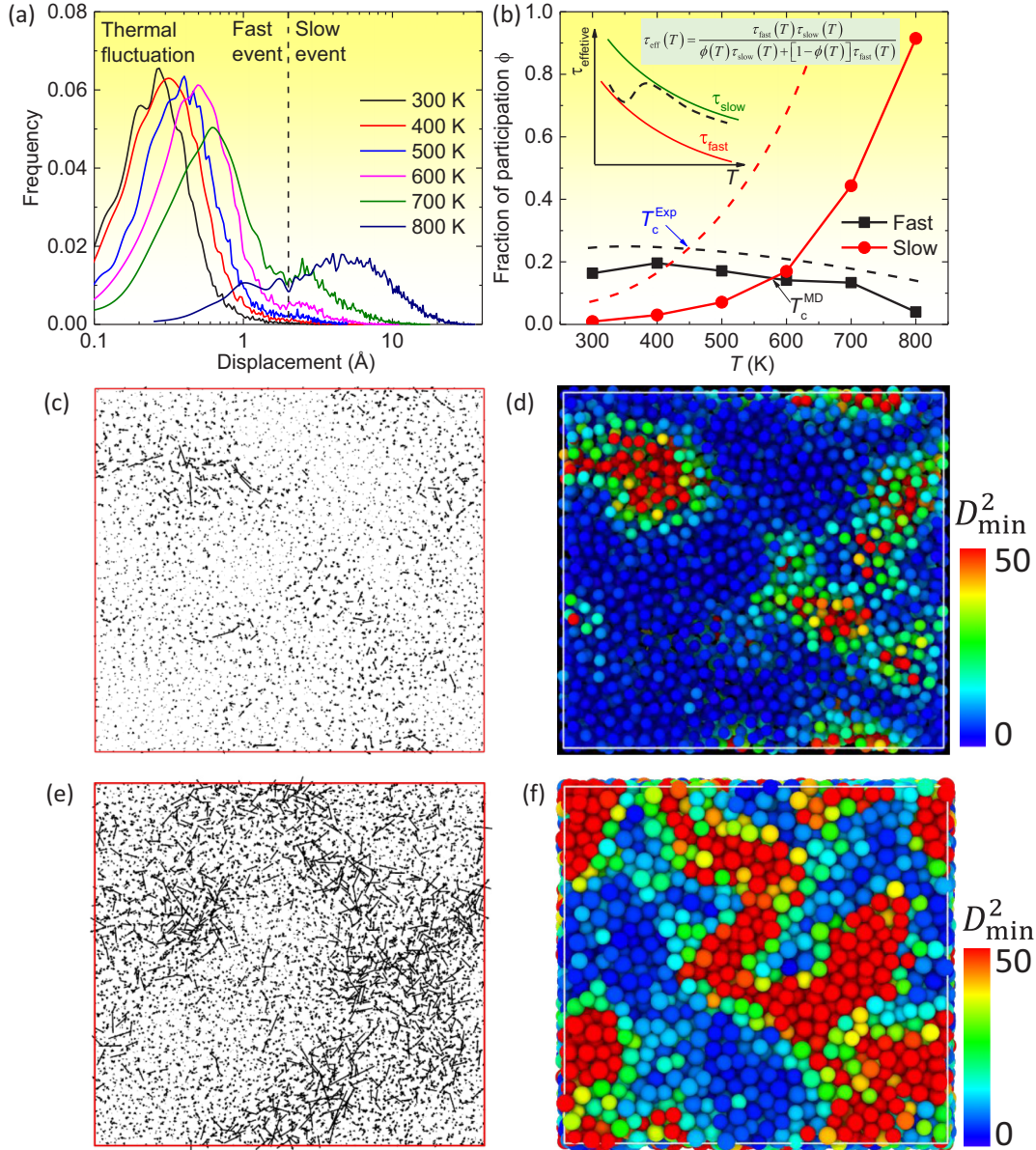


FIG. 5. Competition between the local atomic rearrangement and the long-ranged diffusive mechanism during stress relaxation. (a) Distribution of displacement at different temperatures after stress relaxation for 50 ns. The fast and slow event are defined by this plot. (b) Fraction of the atoms participating in fast and slow event as a function of temperature. Solid lines with symbols are the MD data, while the dashed lines are the schematic speculations at the experimental timescale. The inset is a sketch illustrating the variation of effective characteristic time with temperature. (c) and (e) Displacement vector, (d) and (f) nonaffine squared displacement of the stress-relaxed samples after 50 ns at 300 K, and 500 K, respectively. Only the atoms in a slice of 3 nm are shown for clarity. Fast event involves local collective atomic rearrangement, while the slow event is the percolation of the former accommodated by the long-range atomic diffusion.

parameters in Figs. 1(b) and 4(b). Note that the aging effect is not quite pronounced any more at 800 K, which is above the glass transition temperature ($T_g = 765$ K estimated from the present EAM potential). This is because long-range diffusion happens during stress relaxation which may interfere with the energetic relaxation process.

Figure 4(d) shows the variation of mean-squared displacement (MSD) during stress relaxation, which helps to distinguish the atomic-scale event. It has been shown that the effective jump distance of a long-range diffusion is of the order of 2 \AA [48], which corresponds to the MSD of 4 \AA^2 , as the

boundary shown by a dashed line in Fig. 4(d). Figure 4(d) indicates that most of the atoms experience a short-range displacement at lower temperature and shorter time, while long-range diffusion will occur if the time is long enough or the temperature is high enough. This atomic information indicates that the fast event is likely to be a local atomic rearrangement accommodated by hopping of energetic states between local basins, which did not contribute significantly to MSD; while the slow event has been verified by the MSD to be a long-range atomic diffusion assisted mechanism. The scenario of transition is consistent with the mode-coupling

theory interpretation of diffusion as “rattling in the cage” to “cage breaking” [49–52].

C. Competition between fast and slow event

Figure 5 provides the atomic-scale information about the competition between the fast and slow events, and provides a clue to recognize the distinct character of fast and slow events. Figure 5(a) displays the distribution of the magnitude of displacements after stress relaxation for 50 ns. The peak position (d_{peak}) accounts for the most probable thermal fluctuation of all the atoms at a specific temperature. In a first approximation, we categorize the atoms with displacement smaller than $2d_{\text{peak}}$ as thermal fluctuation. Atoms exhibiting displacements between $2d_{\text{peak}}$ and 2 \AA are responsible for local rearrangement and contribute to the fast event. The remaining atoms with displacement over 2 \AA are those that have experienced long-range diffusion. Except for the thermal fluctuation, we increasingly notice that the atoms experience long-range diffusion as denoted by the higher intensity of the curves after 2 \AA . Particularly above 500 K, one notices a peak for the long-range diffusive atoms. The thermal fluctuation peak vanishes at 800 K, as expected above the glass transition where the glass transforms into an undercooled liquid.

To quantitatively rationalize the interplay between the fast and slow events, we provide the fraction of atoms participating in the local (fast) and long-range (slow) movement as a function of temperature in Fig. 5(b). The solid lines with symbols represent MD simulations. At low temperature or short time scales, the slow event dominates the stress relaxation. The fraction of the fast event shows minor temperature dependence since it is a stress-driven process, which happens in a very short time scale, while the slow event is a thermally activated event and is very sensitive to temperature. With increasing temperature, the fraction of atoms participating in the slow event increases substantially. Note that the height of the first peak in Fig. 5(a) is yet higher than that of the slow event, while the former is only the fraction of thermal fluctuation, the atoms with displacement between $2d_{\text{peak}}$ and 2 \AA are the atoms participating in slow event. Up to a critical temperature of about 580 K in MD, the fast event dominates the stress relaxation. Note that there are several orders of magnitude between the experimental and MD timescales; and we anticipate that the longer timescale will boost the diffusive event as the red dashed line shown in Fig. 5(b). It is reasonable to assume that the critical temperature between the fast and slow events would be much earlier in the experimental scale than in the MD simulation, as we schematically illustrate by the dashed lines in Fig. 5(b). The critical temperature of the dominating mechanism transition could happen at around 450 K and could rationalize the transition of τ_c as shown in Fig. 1(b).

Last but not least, there seems to be a puzzle regarding why τ_c first increases, and then decreases, while both timescales of fast and slow events decrease monotonically with temperature. The inset of Fig. 5(b) rationalizes the transition. The effective relaxation time τ_{eff} is closely related to the participation fraction of both fast and slow events. In a first-order approximation, the effectively observable frequency of two events is expressed as $\Gamma_{\text{eff}}(T) = \phi(T)\Gamma_{\text{fast}}(T) + [1 - \phi(T)]\Gamma_{\text{slow}}(T)$, where $\phi(T)$

is the participation fraction of the fast event, and $[1 - \phi(T)]$ is therefore that of the slow event. The frequency is related to the relaxation time as $\Gamma = 1/\tau$, consequently, the effective relaxation time writes $\tau_{\text{eff}} = \frac{1}{\Gamma_{\text{eff}}} = \frac{\tau_{\text{fast}}(T)\tau_{\text{slow}}(T)}{\phi(T)\tau_{\text{slow}}(T) + [1 - \phi(T)]\tau_{\text{fast}}(T)}$. Although both $\tau_{\text{fast}}(T)$ and $\tau_{\text{slow}}(T)$ decrease with increasing temperature, $\phi(T)$ decreases simultaneously [as shown in Fig. 5(b)], which means that the effective relaxation time $\tau_{\text{eff}}(T)$ transitions from $\tau_{\text{fast}}(T)$ to $\tau_{\text{slow}}(T)$ in the two limits of $\phi(T) = 1$, or $\phi(T) = 0$. Therefore, the decrease of $\phi(T)$ with T induces the increase of τ_{eff} at intermediate temperatures as the slow events eventually dominate the relaxation dynamics at high temperature, as schematically illustrated in the inset of Fig. 5(b).

Finally, the atomistic simulations provide scenarios of both the fast and slow events, as demonstrated in Figs. 5(c)–5(f). Figures 5(c) and 5(d) display the displacement vector and the nonaffine squared displacement D_{min}^2 (42) of the snapshot after stress relaxation for 50 ns at 300 K. They correspond to the fast event. The displacement vector indicates local motion of atoms in a confined volume of the order of nanometers and involves local rearrangement of a large number of atoms as indicated by D_{min}^2 and the large activation volume shown in Fig. 2(b). On the other hand, Figures 5(e) and 5(f) show the counterparts for the slow event. In contrast to the fast event, the slow event requires a long-range diffusion. It is representative of the behavior of very few atoms characterized by the smaller activation volume. The atomic scenario of the slow event is actually the percolation of the local atomic rearrangement with a critical role of atomic diffusion [35,53,54], which is a thermally activated process at a low stress level that needs a longer time to be excited. The above scenario of the local atomic rearrangement and the long-range diffusion in stress relaxation is in accordance with the hypothesized atomic pictures of a recently discovered unusual ultrafast and usual fast β relaxation in MGs [55]. Here we further provide an atomic-scale evidence for the fast stress-driven and slow thermally activated processes during stress relaxation of MGs.

V. CONCLUDING REMARKS

To summarize, a mechanism transition from stress-driven to thermally activated stress relaxation in metallic glass is unveiled by a combined study of experiment, theory, and computer simulations. The stress-driven event is fast and characteristic of large activation volume which dominates the early stage of stress relaxation. The relaxation mechanism gradually transitions to a slow mechanism which is governed by the thermally-activated mechanism that is typical of small activation volume. Discrete analysis distinguishes the heterogenous dynamics and rationalizes the transition of stress relaxation mechanisms. Atomistic simulations demonstrate the interplay between fast and slow event and display similar results with experimental observations. Simulations reveal the atomic-scale mechanism of fast event as the local atomic rearrangement, while the slow event is the percolation of the former involving a significant contribution from the atomic diffusion, which needs longer timescale to occur. The combined insights provide comprehensive physical insights into the underlying mechanisms of the phenomenological KWW equation and may shed light on elucidating the mechanism of the percolation

of the secondary β relaxation to primary α relaxation in glassy materials.

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