

Observation of metallic surface states in the strongly correlated Kitaev-Heisenberg candidate Na_2IrO_3

Nasser Alidoust,^{1,*} Chang Liu,^{1,2} Su-Yang Xu,¹ Ilya Belopolski,¹ Tongfei Qi,³ Minggang Zeng,^{4,5}
Daniel S. Sanchez,¹ Hao Zheng,¹ Guang Bian,¹ Madhab Neupane,^{1,6} Yu-Tzu Liu,^{4,5} Stephen D. Wilson,⁷ Hsin Lin,^{4,5}
Arun Bansil,⁸ Gang Cao,³ and M. Zahid Hasan^{1,†}

¹*Laboratory for Topological Quantum Matter and Spectroscopy (B7), Department of Physics, Princeton University,
Princeton, New Jersey 08544, USA*

²*Department of Physics, South University of Science and Technology of China,
No 1088, Xueyuan Rd., Xili, Nanshan District, Shenzhen, Guangdong, China 518055*

³*Center for Advanced Materials, Department of Physics and Astronomy, University of Kentucky, Lexington, Kentucky 40506, USA*

⁴*Centre for Advanced 2D Materials and Graphene Research Centre, National University of Singapore, 6 Science Drive 2, Singapore 117546*

⁵*Department of Physics, National University of Singapore, 2 Science Drive 3, Singapore 117542*

⁶*Department of Physics, University of Central Florida, Orlando, Florida 32816, USA*

⁷*Materials Department, University of California, Santa Barbara, California 93106, USA*

⁸*Department of Physics, Northeastern University, Boston, Massachusetts 02115, USA*

(Received 23 October 2014; revised manuscript received 20 May 2016; published 14 June 2016)

We report high-resolution angle-resolved photoemission spectroscopy measurements on the honeycomb iridate Na_2IrO_3 . Our measurements reveal the existence of a metallic surface band feature crossing the Fermi level with nearly linear dispersion and an estimated surface carrier density of $3.2 \times 10^{13} \text{ cm}^{-2}$, which has not been theoretically predicted or experimentally observed, and provides the first evidence for metallic behavior on the boundary of this material, whereas the bulk bands exhibit a robust insulating gap. We further show the lack of theoretically predicted Dirac cones at the $\bar{M}$ points of the surface Brillouin zone, which confirms the absence of a stacked quantum spin Hall phase in this material. Our data indicates that the surface ground state of this material is exotic and metallic, unlike as predicted in theory, and establishes Na_2IrO_3 as a rare example of a strongly correlated spin-orbit insulator with surface metallicity.

DOI: [10.1103/PhysRevB.93.245132](https://doi.org/10.1103/PhysRevB.93.245132)

Over the last few years, iridium oxides (iridates) have received considerable attention in condensed matter physics [1–11]. This has been mainly due to their rich ground state phase diagrams, which involve spin-orbit coupling, on-site Coulomb interactions, and direct and indirect hopping between different elements. All of these various types of interactions have comparable energy scales, creating a unique platform to study their interplay in real materials. For instance, in many of the iridates, electron correlations lead to a Mott or many-body insulating ground state. Furthermore, spin-orbit coupling, an important interaction in iridates, has come to play an increasingly important role in condensed matter physics since the discovery of topological insulators [12–19], because it leads to the nontrivial topology of the electronic band structure in this class of materials. However, almost all topological insulators discovered to date are weakly correlated materials, and only recently some experimental evidence for the existence of topological phases in the strongly correlated Kondo insulator SmB_6 has been found [20–22]. The A_2IrO_3 honeycomb iridate is one of the recently debated series among the iridate compounds, and in this series Na_2IrO_3 is especially interesting. There have been proposals of a possible Kitaev-Heisenberg magnetic ground state in this compound and the existence of a quantum spin liquid phase [10,23–26]. Moreover, and more relevant to our present work, this material has been proposed to possess topologically nontrivial band

structure with protected metallic surface states [1,3,4,27]. Specifically, a quantum spin Hall (QSH) state has been predicted to exist in this compound, whose hallmark signature is the existence of Dirac cones at the $\bar{M}$ points of the Brillouin zone (edge Kramers' points, unlike Bi_2Se_3) in a stacked 3D configuration of this material [3,4]. It has experimentally been suggested that Na_2IrO_3 is a relativistic Mott insulator [28–30]. There have also been indications that this compound is a Slater insulator [31]. However, the issue is not fully resolved yet and moreover, there has not been any experimental evidence supporting the emergence of any surface states in this intriguing Mott-like insulating phase.

In this paper we report systematic angle-resolved photoemission spectroscopy (ARPES) studies on Na_2IrO_3 single crystals mapping the first Brillouin zone (BZ) entirely. We report the observation of a metallic linearlike surface band feature crossing the Fermi level at the $\bar{\Gamma}$ point (center of the surface BZ), providing the first evidence for metallic behavior on the surface boundary of this material, despite the fact that our experimentally observed bulk bands exhibit a robust gap. This is in contrast to the theoretical prediction of Dirac cones at the $\bar{M}$ points of the BZ. Therefore, our results suggest that the surface ground state electronic structure of this system is not in a stacked QSH phase as predicted theoretically, but rather reflects an exotic many-body metallic ground state on the surface. It is worth mentioning that measuring the electronic band structure of Na_2IrO_3 single crystals using ARPES is very challenging due to the small size of the samples and the low success rate in cleaving flat surfaces from them. However, we have managed to overcome these obstacles and obtain reproducible

*Corresponding author: nassera@princeton.edu

†Corresponding author: mzhasan@princeton.edu

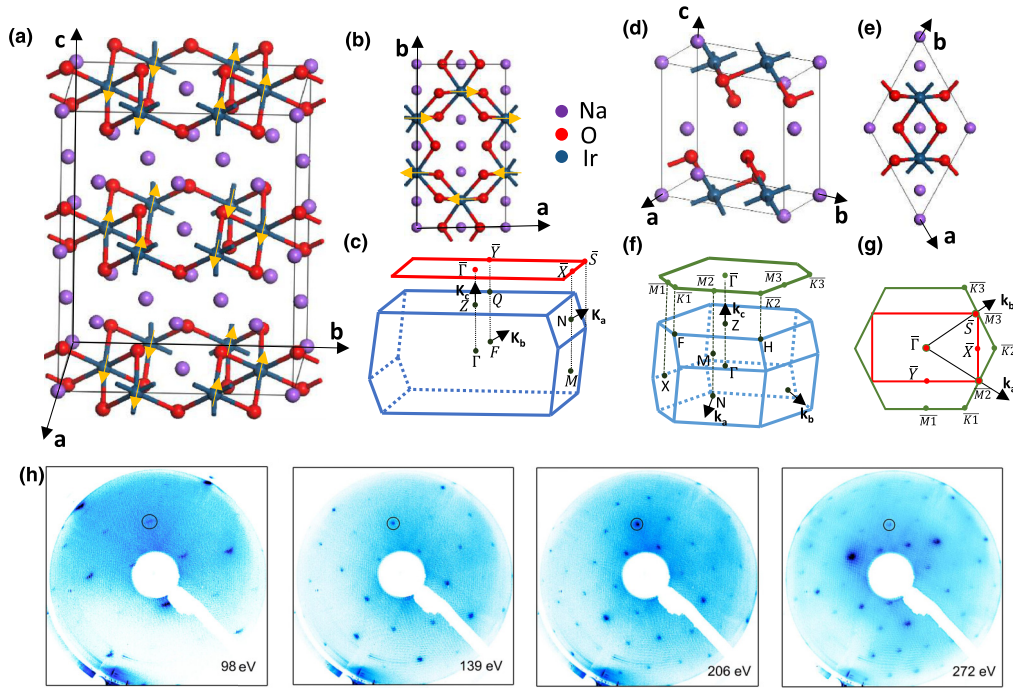


FIG. 1. (a) Side and (b) top views of the crystal structure and (c) the corresponding Brillouin zone (BZ) of magnetic Na_2IrO_3 with the zigzag-AFM spin configuration. The arrows indicate magnetic moments of Ir atoms. (d)–(f) Same as (a)–(c) for nonmagnetic Na_2IrO_3 . The projection of the bulk BZ on the (001) surface is indicated by the red and green rectangles in (c) and (f), respectively. (g) Relationship between the (001) surface BZ of the nonmagnetic (green) and zigzag-AFM magnetic (red) Na_2IrO_3 . (h) Measured low-energy electron diffraction (LEED) patterns of the crystals used in our studies at various electron energies as marked on each panel.

data. Our data indicates that the surface ground state of this material is metallic and more exotic than that predicted in theory and establishes Na_2IrO_3 as a rare example of a strongly correlated spin-orbit insulator with surface metallicity.

A self-flux method was used to grow single crystals of Na_2IrO_3 from off-stoichiometric quantities of IrO_2 and Na_2CO_3 . Technical details can be found on the preparation of related iridate compounds in Refs. [25,32–34]. The studied crystals have typical sizes of $2 \text{ mm} \times 1 \text{ mm} \times 0.5 \text{ mm}$. ARPES measurements were performed with incident photon energies of 60–100 eV, and low-energy electron diffraction (LEED) measurements were performed with electron energies of 90–300 eV, both at beamline 4.0.3 of the Advanced Light Source (ALS) in the Lawrence Berkeley National Laboratory (LBNL). Samples were cleaved *in situ* at 300 K in chamber pressure better than 5×10^{-11} torr resulting in shiny surfaces. Energy resolution was better than 15 meV and momentum resolution was better than 1% of the surface BZ. The measurements were also conducted at 300 K, which excludes the possibility of an unconventional low-temperature surface reconstruction reported on this material using scanning tunneling microscopy [35]. We note that we also found the ARPES signal to disappear entirely at those low temperatures where the STM measurements were performed. But at room temperature the metallic linearlike surface states that cross the Fermi level are clearly seen in our data (see figures in this paper).

The recently revised monoclinic structure of Na_2IrO_3 (space group: $C2/m$, No. 12) consists of edge-sharing and c-stacking IrO_6 octahedra [24,25]. Ir atoms form a nearly ideal honeycomb lattice with the centers occupied by Na.

Different from the rhombohedral (ABC) stacking, an in-plane offset of successive Ir honeycombs gives rise to the monoclinic distortion, and one Na layer is sandwiched by the neighboring IrO_6 octahedra layers. The existence of trigonal distortions of oxygen ions in IrO_6 octahedra can split the t_{2g} manifold. Na_2IrO_3 undergoes a magnetic phase transition into zigzag-antiferromagnetic (AFM) order at $T_N \sim 15 \text{ K}$. Therefore, crystal cells with eight [Fig. 1(a) (side view) and 1(b) (top view)] and two [Fig. 1(d) (side view) and 1(e) (top view)] formula units are used to describe the zigzag-AFM and nonmagnetic phases, respectively. In Figs. 1(a) and 1(b) the long-range magnetic order is indicated by the local magnetic moments of Ir atoms, which adopt the zigzag spin configuration with Ir zigzag chains. The corresponding Brillouin zones for the magnetic and nonmagnetic phases are illustrated in Figs. 1(c) and 1(f), respectively. A doubled surface unit cell, and thus a halved surface BZ, is adopted to consider the zigzag-AFM spin configuration. For comparison, as shown in Fig. 1(g), we plot the hexagonal surface BZ of the nonmagnetic Na_2IrO_3 in green, along with that of the zigzag-AFM Na_2IrO_3 in red.

We conduct LEED measurements at various electron energies on the single crystals of Na_2IrO_3 used in our ARPES experiments. The results are shown in Fig. 1(h), and their analysis in Fig. S1 [36]. As revealed from our LEED measurements no detectable surface reconstruction is present in these crystals. We also measure the core level spectra of these crystals [Fig. 2(a)] and confirm the chemical composition of this compound by observing sharp Na and Ir peaks in the

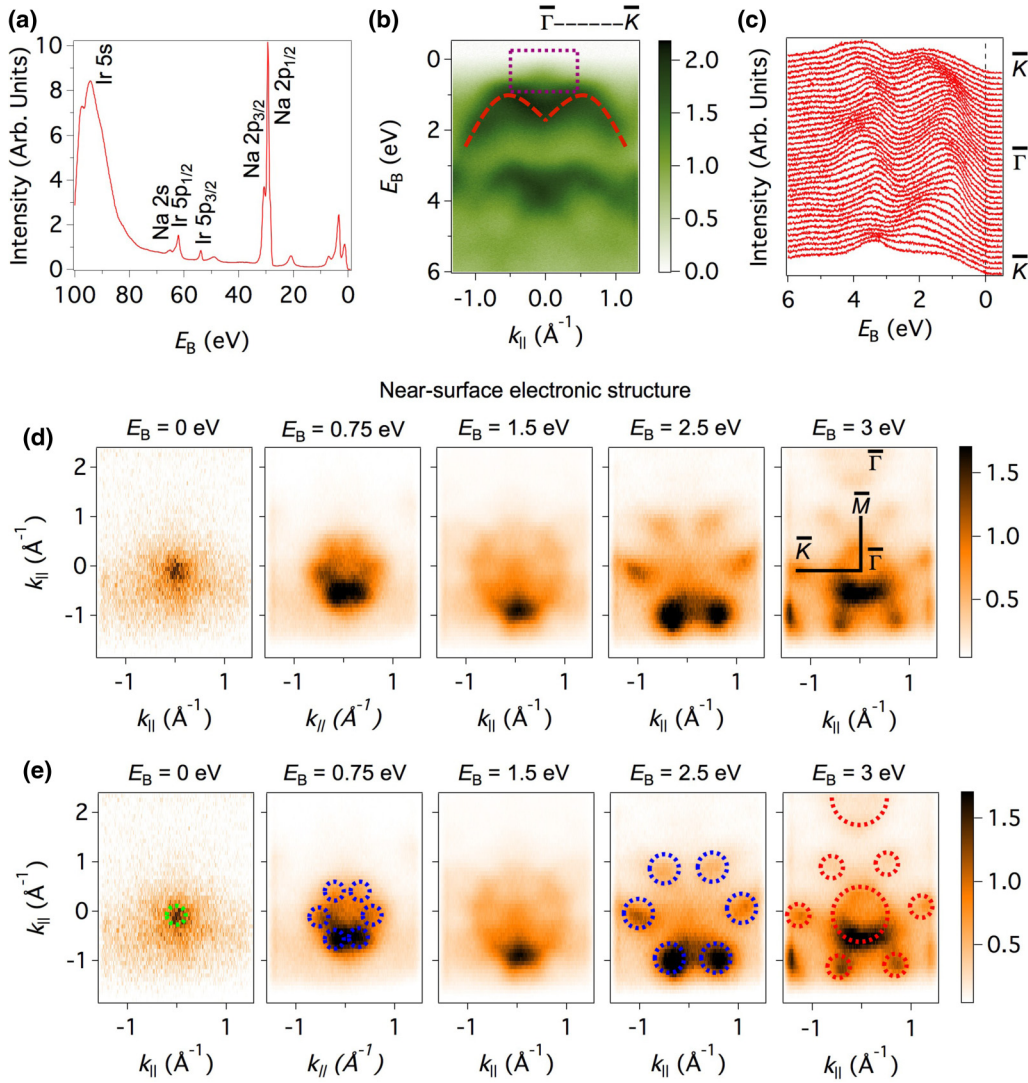


FIG. 2. (a) Core level spectra of Na₂IrO₃. (b) ARPES electronic band structure of Na₂IrO₃ along the high symmetry direction $\bar{\Gamma}$ - $\bar{K}$ of the BZ. (c) Energy distribution curves (EDCs) of the k - E intensity profile in (b). The dashed line indicates the position of E_F . (d) Constant binding energy cuts at various binding energies between E_F and $E_B = 3$ eV. The corresponding binding energies are indicated at the top of each panel. The high symmetry directions of the BZ, $\bar{\Gamma}$ - $\bar{K}$ and $\bar{\Gamma}$ - $\bar{M}$, are marked in the last panel. (e) Replica of (d) with dashed guide-to-the-eye lines drawn at the place of each Fermi surface pocket to emphasize the sixfold symmetry of the BZ. Green lines mark the pockets formed by the intensities crossing the E_F , blue lines the pockets formed by the lower binding energy bulk valence bands, and red lines the pockets formed by the higher binding energy bulk valence bands. These measurements were conducted at $T = 300$ K with $h\nu = 90$ eV.

core level intensity plot, which is suggestive of the high quality of the Na₂IrO₃ samples used in our measurements.

In Fig. 2(b) we present a k - E ARPES intensity profile of Na₂IrO₃ over a wide binding energy scale measured with a synchrotron-based ARPES system using a photon energy of $h\nu = 90$ eV at a temperature of $T \sim 300$ K. We observe two bulk bands in the binding energy range from the Fermi level to binding energy $E_B = 6$ eV: the first extending from a binding energy of ~ 0.5 eV to ~ 3 eV, and the second from a binding energy of ~ 3 eV to ~ 6 eV. Our measurements indicate that the top of the bulk valence band is located approximately 500 meV below the Fermi level. The red dashed lines are guides to the eye highlighting the M shape of the bulk valence band, indicative of its inverted nature. The stack plot of the corresponding energy distribution curves (EDCs)

is shown in Fig. 2(c), which is in qualitative agreement with the previously reported EDCs of this material [29] (see Fig. S2 [36]). However, we highlight that in the present work, bands with clear dispersion have been observed on Na₂IrO₃. These bands are relatively broader as compared to the ARPES data on many weakly interacting systems (such as graphene or Bi₂Se₃). We note that this is likely due to intrinsic correlation effects rather than lack of crystalline order of our samples, whose core level spectra demonstrate their high quality. The kinetic energy value of the Fermi level is determined by measuring that of the polycrystalline gold separately at the same photon energies and under the same experimental conditions [see Fig. 3(f)].

Figures 2(d) and 2(e) show the measured momentum intensity maps at selected constant binding energies from the Fermi level E_F to $E_B = 3$ eV, with the high symmetry

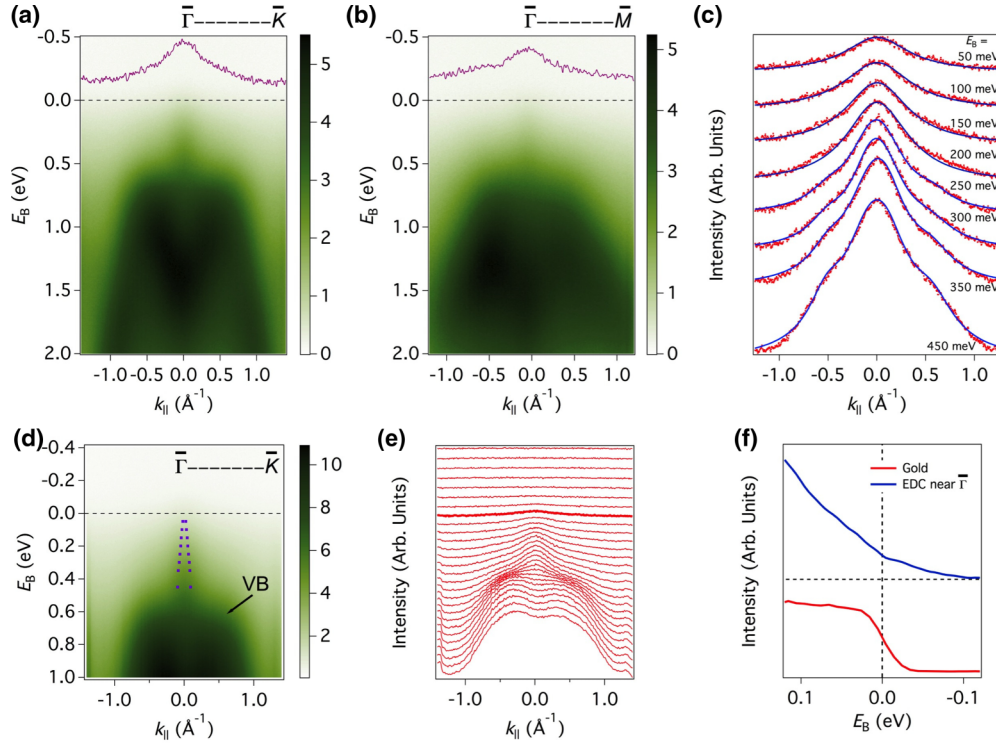


FIG. 3. ARPES electronic band structure of Na_2IrO_3 along the high symmetry directions of (a) $\bar{\Gamma}-\bar{K}$ and (b) $\bar{\Gamma}-\bar{M}$ of the BZ, respectively. The momentum distribution curves (MDCs) at the Fermi level are plotted on top of each panel in purple. (c) MDCs (red) and heuristic fits to the MDCs (blue) of the metallic in-gap states around the $\bar{\Gamma}$ point of the Brillouin zone using two Lorentzians. (d) Zoomed-in version of the metallic in-gap states and (e) their corresponding MDCs. The bold curve in (e) indicates the MDC at E_F , and the purple dots in (d) show the peak positions of the Lorentzians used to fit to the MDCs of these metallic in-gap states. The horizontal black dashed lines in (a), (b), and (d) denote the position of E_F . (f) Comparison between the EDCs of polycrystalline gold and that of the Na_2IrO_3 samples in our study around the $\bar{\Gamma}$ point of the BZ. Note that the numerical values on the color bars may not be compared across figures, since they are in arbitrary units specific to a certain color plot.

directions of the BZ, $\bar{\Gamma}-\bar{K}$ and $\bar{\Gamma}-\bar{M}$, marked on the last panel of Fig. 2(d). At E_F we observe a dotlike pocket at the center of the BZ ($\bar{\Gamma}$ point), and upon going to larger binding energies (e.g., $E_B = 2.5$ eV) the constant binding energy momentum cuts evolve to six circular pockets at the corners of the BZ ($\bar{K}$ point). As demonstrated by the dashed guide-to-the-eye lines drawn at the place of each Fermi surface pocket in Fig. 2(e), these constant binding energy maps clearly demonstrate a hexagonal Brillouin zone, consistent with the sixfold symmetry of the (001) surface of this compound.

An important feature of the electronic band structure of this compound revealed by our measurements is the observation of spectral intensities at the time-reversal invariant Kramers' $\bar{\Gamma}$ point of the BZ extending from the Fermi level to the top of the valence band ($E_B \sim 0.5$ eV) as can be seen in Fig. 2(b). The purple rectangular dashed box in this figure highlights the momentum region in which these intensities are observed. These states form a metallic Fermi surface pocket at E_F as shown in the first panel of Fig. 2(d). In an attempt to understand these states in more detail, we next focus on studying the low-energy states of this material lying close to E_F .

In Figs. 3(a) and 3(b) we present k - E ARPES intensity profiles of the low-energy electronic states observed close to

the Fermi level in this material, along the two high symmetry directions $\bar{\Gamma}-\bar{K}$ and $\bar{\Gamma}-\bar{M}$, respectively. These high symmetry directions are shown in the fifth panel of Fig. 2(d). We focus on a narrow binding energy window between the Fermi level and $E_B = 2$ eV. These k - E intensity plots clearly show the existence of in-gap electronic states between the top of the valence band at $E_B \sim 0.5$ eV and E_F . For clarity, the momentum distribution curves (MDCs) at E_F are plotted at the top of each panel in purple, which further confirm the existence of finite intensities at E_F around the time-reversal invariant Kramers' $\bar{\Gamma}$ point. To study the dispersion of these metallic states in more detail, we perform a quantitative analysis by heuristically fitting the MDCs of these metallic in-gap states in the binding energy range between 50 meV and 450 meV around the $\bar{\Gamma}$ point using two Lorentzians [Fig. 3(c)]. We then track the momentum positions of these peaks at the corresponding binding energies [purple dots in Fig. 3(d)], which reveals the nearly linear dispersion of these states (for details of this quantitative analysis see Figs. S3 and S4 [36]). Figure 3(e) displays the stack plot of the corresponding MDCs in Fig. 3(d). Our analysis suggests a Fermi velocity of $v_F \sim (8.10 \pm 0.75) \times 10^5 \text{ ms}^{-1}$ ($5.33 \pm 0.49 \text{ eV} \cdot \text{\AA}$) for these observed metallic linearlike states, larger than that of the topological surface states of Bi_2Se_3 [16], but smaller than that of the Dirac bands of graphene [37].

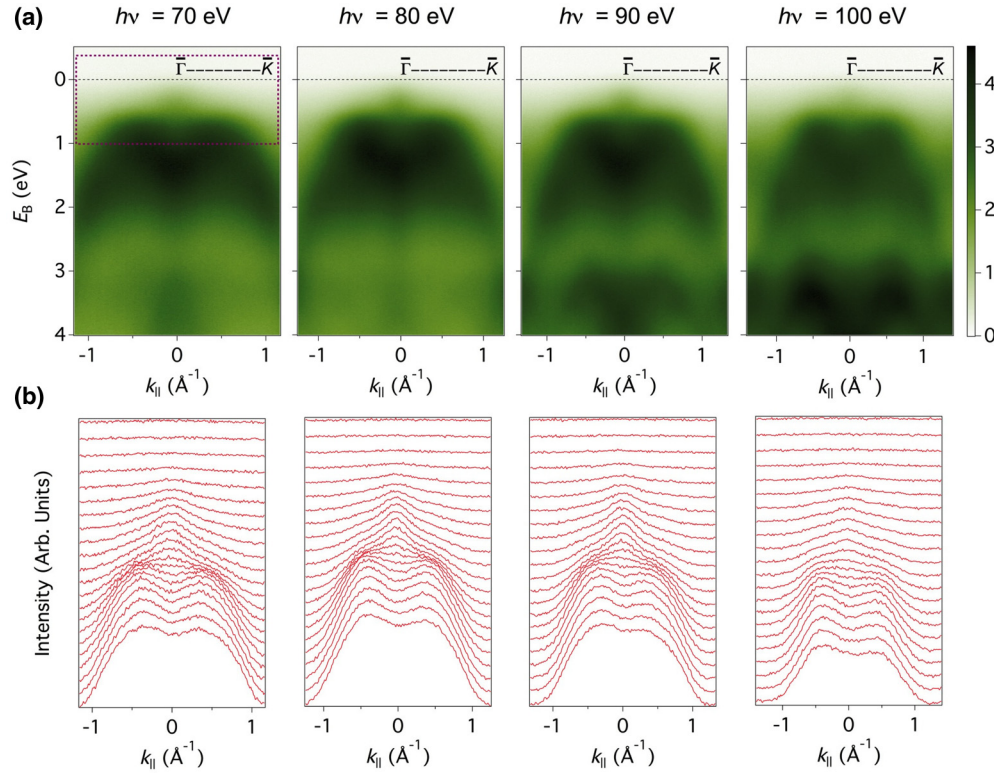


FIG. 4. (a) Photon energy $h\nu$ ($\propto k_z$) dependence measurements of the observed metallic linearlike in-gap electronic states and (b) the corresponding MDCs in the region indicated by the purple rectangular dashed box in the first panel of (a). The horizontal black dashed lines in (a) show the position of E_F .

In order to investigate whether these states exhibit two- or three-dimensional dispersion, we perform k_z -dependent measurements of the observed electronic bands by measuring their dispersion upon changing the photon energy $h\nu$ in our ARPES experiments. As can be observed in Fig. 4(a), the metallic states crossing the Fermi level survive at the $\bar{\Gamma}$ point of the BZ at all measured photon energies. Plotting the MDCs at each $h\nu$ in Fig. 4(b) [in the region denoted by the purple rectangular dashed box in the first panel of Fig. 4(a)], we can show that these states lack any dispersion along k_z . Furthermore, based on the dotlike pocket that these metallic states form on the Fermi surface [see Fig. 2(c)] we estimate a surface carrier density of $\sim 3.2 \times 10^{13} \text{ cm}^{-2}$ for our samples. Our LEED measurements indicate that no detectable surface reconstruction is present in these crystals within our measurements precision, and thus the possibility of a change in the surface stoichiometry, which may induce new states at the surface due to polarity effects, is ruled out. These evidences, taken collectively, convincingly imply the observation of metallic in-gap states on the surface boundary of Na_2IrO_3 at room temperature.

Now we discuss the implications of our experimental observation of Na_2IrO_3 electronic bands. We have resolved the electronic band structure of this honeycomb iridate compound at room temperature and have shown its sixfold symmetry structure along its natural cleavage plane consistent with its previously reported crystal structure in the high temperature nonmagnetic phase. Our data indicates that the low-energy physics of this material is governed by metallic in-gap states

crossing the Fermi level that have linearlike dispersion and form a dotlike Fermi surface pocket at the center of the BZ (the time-reversal invariant Kramers' $\bar{\Gamma}$ point). Photon energy dependence measurements further confirm the 2D nature of these states and provide strong evidence that they are indeed surface states. Our measurements are performed at room temperature, much higher than the temperature at which the antiferromagnetic order sets in. Thus, time-reversal symmetry is preserved in the temperature regime at which our studies are carried out, consistent with the Z_2 topologically nontrivial phase. In spite of the fact that our experimentally observed bulk bands exhibit a robust gap, our observation of a metallic linearlike surface band feature crossing the Fermi level at the $\bar{\Gamma}$ point provides the first evidence for metallic behavior on the surface boundary of this material, which is in contrast to the theoretical prediction of Dirac cones at the $\bar{M}$ points of the BZ. Therefore, our results suggest that the surface ground state electronic structure of this system is not in a stacked QSH phase as predicted theoretically, but rather reflects an exotic many-body metallic ground state on the surface. The disagreement between our experimental data and theoretical results may reflect the difficulty of modeling strongly correlated systems in theory. These models are typically more accurate for weakly interacting systems rather than strongly correlated ones.

In transport measurements Na_2IrO_3 behaves as an insulator. Since transport measurements are mainly governed by bulk states, this further suggests that the observed metallic states in our measurements are surface states. In order to further investigate these observed surface states, more detailed transport

measurements are required in the future. The results presented here establish Na_2IrO_3 as a rare example of a strongly correlated spin-orbit Mott insulator with surface metallicity. Our study of near-surface electronic structure and the observation of anomalous metallicity reflecting a nontrivial ground state opens up many opportunities for observing exotic transport and Hall effect on the surface of this highly correlated spin-orbit material and paves the way to future studies exploring the possibility that topological insulators, Weyl semimetals, and other topological phases may arise in this class of materials.

We thank Jonathan D. Denlinger for technical assistance. Work at Princeton University is supported by the US National Science Foundation Grant No. NSF-DMR-1006492. M.Z.H. acknowledges visiting scientist support from Lawrence Berkeley National Laboratory and additional partial support from the A. P. Sloan Foundation and NSF-DMR-0819860. G.C. acknowledges NSF support via Grant No. DMR 1265162. The photoemission measurements using synchrotron x-ray facilities are supported by the Basic Energy Sciences of the US Department of Energy under DE-FG-02-05ER46200.

-
- [1] D. Pesin and L. Balents, *Nat. Phys.* **6**, 376 (2010).
 - [2] X. Wan, A. M. Turner, A. Vishwanath, and S. Y. Savrasov, *Phys. Rev. B* **83**, 205101 (2011).
 - [3] A. Shitade, H. Katsura, J. Kuneš, X.-L. Qi, S.-C. Zhang, and N. Nagaosa, *Phys. Rev. Lett.* **102**, 256403 (2009).
 - [4] C. H. Kim, H. S. Kim, H. Jeong, H. Jin, and J. Yu, *Phys. Rev. Lett.* **108**, 106401 (2012).
 - [5] G. Cao, Y. Xin, C. S. Alexander, J. E. Crow, P. Schlottmann, M. K. Crawford, R. L. Harlow, and W. Marshall, *Phys. Rev. B* **66**, 214412 (2002).
 - [6] S. J. Moon, H. Jin, K. W. Kim, W. S. Choi, Y. S. Lee, J. Yu, G. Cao, A. Sumi, H. Funakubo, C. Bernhard, and T. W. Noh, *Phys. Rev. Lett.* **101**, 226402 (2008).
 - [7] B. J. Kim, H. Jin, S. J. Moon, J.-Y. Kim, B.-G. Park, C. S. Leem, J. Yu, T. W. Noh, C. Kim, S.-J. Oh *et al.*, *Phys. Rev. Lett.* **101**, 076402 (2008).
 - [8] G. Jackeli and G. Khaliullin, *Phys. Rev. Lett.* **102**, 017205 (2009).
 - [9] Y. Machida, S. Nakatsuji, S. Onoda, T. Tayama, and T. Sakakibara, *Nature (London)* **463**, 210 (2010).
 - [10] J. Chaloupka, G. Jackeli, and G. Khaliullin, *Phys. Rev. Lett.* **105**, 027204 (2010).
 - [11] I. Kimchi and A. Vishwanath, *Phys. Rev. B* **89**, 014414 (2014).
 - [12] M. Z. Hasan and C. L. Kane, *Rev. Mod. Phys.* **82**, 3045 (2010).
 - [13] X.-L. Qi and S.-C. Zhang, *Rev. Mod. Phys.* **83**, 1057 (2011).
 - [14] D. Hsieh, D. Qian, L. Wray, Y. Xia, Y. S. Hor, R. J. Cava, and M. Z. Hasan, *Nature (London)* **452**, 970 (2008).
 - [15] D. Hsieh, Y. Xia, L. Wray, D. Qian, A. Pal, J. H. Dil, F. Meier, J. Osterwalder, G. Bihlmayer, C. L. Kane *et al.*, *Science* **323**, 919 (2009).
 - [16] Y. Xia, D. Qian, D. Hsieh, L. Wray, A. Pal, H. Lin, A. Bansil, D. Grauer, Y. S. Hor, R. J. Cava, and M. Z. Hasan, *Nat. Phys.* **5**, 398 (2009).
 - [17] Y. L. Chen, J.-H. Chu, J. G. Analytis, Z. K. Liu, K. Igarashi, H.-H. Kuo, X. L. Qi, S. K. Mo, R. G. Moore, D. H. Lu *et al.*, *Science* **329**, 659 (2010).
 - [18] D. Hsieh, Y. Xia, D. Qian, L. Wray, J. H. Dil, F. Meier, J. Osterwalder, L. Patthey, J. G. Checkelsky, N. P. Ong *et al.*, *Nature (London)* **460**, 1101 (2009).
 - [19] S.-Y. Xu, M. Neupane, C. Liu, D. Zhang, A. Richardella, L. A. Wray, N. Alidoust, M. Leandersson, T. Balasubramanian, J. Sánchez-Barriga *et al.*, *Nat. Phys.* **8**, 616 (2012).
 - [20] M. Dzero, K. Sun, V. Galitski, and P. Coleman, *Phys. Rev. Lett.* **104**, 106408 (2010).
 - [21] M. Neupane, N. Alidoust, S.-Y. Xu, T. Kondo, Y. Ishida, D. J. Kim, C. Liu, I. Belopolski, Y. J. Jo, T.-R. Chang *et al.*, *Nat. Commun.* **4**, 2991 (2013).
 - [22] J. Jiang, S. Li, T. Zhang, Z. Sun, F. Chen, Z. R. Ye, M. Xu, Q. Q. Ge, S. Y. Tan, X. H. Niu *et al.*, *Nat. Commun.* **4**, 3010 (2013).
 - [23] X. Liu, T. Berlijn, W.-G. Yin, W. Ku, A. Tsvelik, Y.-J. Kim, H. Gretarsson, Y. Singh, P. Gegenwart, and J. P. Hill, *Phys. Rev. B* **83**, 220403 (2011).
 - [24] S. K. Choi, R. Coldea, A. N. Kolmogorov, T. Lancaster, I. I. Mazin, S. J. Blundell, P. G. Radaelli, Y. Singh, P. Gegenwart, K. R. Choi *et al.*, *Phys. Rev. Lett.* **108**, 127204 (2012).
 - [25] F. Ye, S. Chi, H. Cao, B. C. Chakoumakos, J. A. Fernandez-Baca, R. Custelcean, T. F. Qi, O. B. Korneta, and G. Cao, *Phys. Rev. B* **85**, 180403 (2012).
 - [26] J. Chaloupka, G. Jackeli, and G. Khaliullin, *Phys. Rev. Lett.* **110**, 097204 (2013).
 - [27] C. H. Sohn, H.-S. Kim, T. F. Qi, D. W. Jeong, H. J. Park, H. K. Yoo, H. H. Kim, J.-Y. Kim, T. D. Kang, D.-Y. Cho *et al.*, *Phys. Rev. B* **88**, 085125 (2013).
 - [28] Y. Singh and P. Gegenwart, *Phys. Rev. B* **82**, 064412 (2010).
 - [29] R. Comin, G. Levy, B. Ludbrook, Z.-H. Zhu, C. N. Veenstra, J. A. Rosen, Y. Singh, P. Gegenwart, D. Stricker, J. N. Hancock *et al.*, *Phys. Rev. Lett.* **109**, 266406 (2012).
 - [30] H. Gretarsson, J. P. Clancy, X. Liu, J. P. Hill, E. Bozin, Y. Singh, S. Manni, P. Gegenwart, J. Kim, A. H. Said *et al.*, *Phys. Rev. Lett.* **110**, 076402 (2013).
 - [31] H.-J. Kim, J.-H. Lee, and J.-H. Cho, *Sci. Rep.* **4**, 5253 (2014).
 - [32] S. Chikara, O. Korneta, W. P. Crummett, L. E. DeLong, P. Schlottmann, and G. Cao, *Phys. Rev. B* **80**, 140407 (2009).
 - [33] M. A. Laguna-Marco, D. Haskel, N. Souza-Neto, J. C. Lang, V. V. Krishnamurthy, S. Chikara, G. Cao, and M. van Veenendaal, *Phys. Rev. Lett.* **105**, 216407 (2010).
 - [34] M. Ge, T. F. Qi, O. B. Korneta, D. E. De Long, P. Schlottmann, W. P. Crummett, and G. Cao, *Phys. Rev. B* **84**, 100402 (2011).
 - [35] F. Lupke, S. Manni, S. C. Erwin, I. I. Mazin, P. Gegenwart, and M. Wenderoth, *Phys. Rev. B* **91**, 041405(R) (2015).
 - [36] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.93.245132> for analysis of the low-energy electron diffraction data, comparison to the bulk energy distribution curves previously reported, Lorentzians used to fit to the momentum distribution curves of the metallic in-gap states, and the linear fits to the peak positions of these Lorentzians.
 - [37] A. Bostwick, T. Ohta, T. Seyller, K. Horn, and E. Rotenberg, *Nat. Phys.* **3**, 36 (2007).