

Magnetotransport in double quantum well with inverted energy spectrum: HgTe/CdHgTeM. V. Yakunin,^{1,2,*} A. V. Suslov,³ M. R. Popov,¹ E. G. Novik,⁴ S. A. Dvoretzky,^{5,6} and N. N. Mikhailov^{5,7}¹*Institute of Metal Physics, 18 S. Kovalevskaya Str., Ekaterinburg 620990, Russia*²*Ural Federal University, 19 Mira Str., Ekaterinburg 620002, Russia*³*NHMFL, FSU, 1800 East Paul Dirac Drive, Tallahassee, Florida 32310, USA*⁴*Physical Institute, University of Würzburg, D-97074 Würzburg, Germany*⁵*Institute of Semiconductor Physics, 13 Lavrentyev Ave., Novosibirsk 630090, Russia*⁶*Tomsk State University, 36 Lenin Ave., Tomsk 634050, Russia*⁷*Novosibirsk State University, 2 Pirogov Str., Novosibirsk 630090, Russia*

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We present an experimental study of the double-quantum-well (DQW) system made of two-dimensional layers with inverted energy band spectrum: HgTe. The magnetotransport reveals a considerably larger overlap of the conduction and valence subbands than in known HgTe single quantum wells (QW), which may be regulated here by an applied gate voltage V_g . This large overlap manifests itself in a much higher critical field B_c separating the range above it with a plain behavior of the Hall magnetoresistance $\rho_{xy}(B)$, where the quantum peculiarities shift linearly with V_g , and the range below with a complicated behavior. In the latter case, specific structures in $\rho_{xy}(B)$ are formed like a double- N -shaped $\rho_{xy}(B)$, reentrant sign-alternating quantum Hall effect with transitions into a zero-filling-factor state, etc., which are clearly manifested here due to better magnetic quantization at high fields, as compared to the features seen earlier in a single HgTe QW. The coexisting electrons and holes were found in the whole investigated range of positive and negative V_g as revealed (i) from fits to the low-field N -shaped $\rho_{xy}(B)$, (ii) from the Fourier analysis of oscillations in $\rho_{xx}(B)$, and (iii) from a specific behavior of $\rho_{xy}(B)$ at high positive V_g . A peculiar feature here is that the found electron density n remains almost constant in the whole range of investigated V_g while the hole density p drops down from the value a factor of 6 larger than n at extreme negative V_g to almost zero at extreme positive V_g passing through the charge-neutrality point. We show that this difference between n and p stems from an order of magnitude larger density of states for holes in the lateral valence subband maxima than for electrons in the conduction subband minimum. We analyze our observations on the basis of a calculated picture of magnetic levels in a DQW and suggest that their specificity is due to (i) a nonmonotonic course of the valence subband magnetic levels and an oscillating behavior of the valence subband top versus field due to lateral maxima in the energy spectrum, (ii) a reduced gap between the lowest-electron and the highest-hole magnetic levels where the electron- and hole-type localized states are superposed, and (iii) a possible formation of the interlayer electron-hole excitons.

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Spatially quantized energy spectrum of the HgTe layer wider than $d_w = d_c \approx 6.3$ nm is markedly specific due to its inversion character [1]. In this case, the conduction and valence subbands are close in energy due to their quasidegeneracy, and for $d_w \gtrsim 8.3$ nm both are built of the Γ_8 -symmetry wave functions contrary to the traditional semiconductor structures with their finite intrinsic gap and a conduction subband of the Γ_6 character. Moreover, the conduction subband may overlap here with the lateral maxima of the valence subband [2] leading to formation of a semimetal with coexisting electrons and holes. The overlap occurs in a relatively wide HgTe quantum well (QW) depending on the heterostructure characteristics, of which probably the main is the built-in deformation [3,4]. Diverse new physics arises in such a semimetal. For example, correlations between the electron and hole subsystems may cause a transition into a state of exciton insulator [5] while the enhanced electron-hole friction in this situation results in modified temperature dependencies of transport coefficients [6]. The picture of magnetic levels in

this semimetal is very intricate. It consists of two overlapped fan charts oppositely directed in energy. The lowest Landau level (so-called zero-mode level [7,8]) of the conduction band contains pure heavy-hole states which do not mix with the light-particle states. Accordingly, this level lowers its energy linearly with increasing magnetic field [8] like the hole levels do. Another indication of its holelike nature comes from that it lowers its energy as it approaches the sample edge [9,10]. Similar anomalous level exists in the valence subband, which contains, although not a pure state, but a truncated set of states, resulting in that it monotonously moves up with field, thus indicating its electronlike nature. Its energy increases on approaching the sample edge as it should be for electrons. The infinite increase of its energy with field distinguishes this level from all the other valence subband levels, which, although move up starting from zero field, then turn down after reaching the energy of the valence subband lateral maxima. A gap is opened between the lowest electron-magnetic level and the highest hole level at the fields higher than some critical magnetic field B_c , thus creating a possibility to realize a unique zero-filling-factor state when the Fermi level gets into this gap [11,12]. This state has also been revealed in graphene and is actively debated [13]. The specific shape of the valence subband with a central minimum and lateral maxima

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causes an oscillatinglike behavior of the subband top profile versus quantizing magnetic field. These oscillations may be felt in the quantum magnetotransport once their amplitudes are comparable with the typical distances between the valence subband levels for the case when the quantum features connected with these levels are resolved experimentally. Physics gets especially unusual here because of a possible competition between the states of a heavy-electron character in the central valence subband minimum and the hole states around its lateral maxima and at larger wave vectors.

We present our investigations of a technique to change and regulate *in situ* the subband overlap using a system of two relatively wide HgTe quasi-two-dimensional (quasi-2D) layers separated by a thin CdHgTe barrier, i.e., a double quantum well (DQW) with inverted energy spectrum. Application of a gate voltage V_g causes a shift of the whole energy spectrum picture in the upper HgTe layer (which is closer to the gate) with respect to a similar energy spectrum of the lower HgTe layer since the free charges in the upper layer screen the lower layer so that its energy is almost insensitive to V_g [14]. It results in an increase of the overlap between the conduction (c) subband in one layer and the valence (v) subband in the other. The HgTe/CdHgTe system is specific in this sense differing from a traditional DQW, e.g., in the GaAs/AlGaAs system, since for the normal gapped semiconductors the interlayer overlap occurs only for the same kind of bands, thus causing an interlayer interaction between the same kind of carriers: e - e or h - h . In a DQW of traditional semiconductors the exhaustion of the upper layer may be achieved, after what the charge density in the lower layer starts to vary with V_g [14]. In contrast, the exhaustion of the upper layer would never come for the HgTe layers with inverted spectrum, only the kind of free charges will be changed here when the Fermi level E_F moves between the touched or overlapped c and v subbands.

An exclusion is the InAs/GaSb system where two normal spectra are intrinsically so much shifted in energy that the InAs c band falls below the GaSb v -band top [15,16]. However, a substantial interlayer hybridization gap appears in this system, which considerably modifies the collective energy spectrum. Also, a balanced state with equal electron and hole densities (a charge-neutrality point, CNP), which is one of the most remarkable moments in a system with overlapped c and v bands, has not been reached in InAs/GaSb system so far [17].

Although the additional overlap of c and v subbands occurs in our case *between* the layers, it would affect interlayer electron-hole (e - h) interactions in the DQW with a thin enough barrier, thus influencing the magnetotransport in the whole system. Variation of V_g in a DQW provides some new opportunities to scan the complicated structure of the valence subband top by the Fermi level. Information about the valence subband top in the case of overlapped bands is especially actual since some controversy about its structure has been revealed recently [18]. A possibility to considerably increase the total overlap of c and v subbands in the DQW will allow to shift B_c to much higher fields from its position in a single QW, thus moving all the effects connected with the transition of E_F through CNP to a robust quantum Hall (QH) regime with well-resolved magnetic levels.

The paper is organized as follows: In Sec. II, description of samples and the measurement techniques along with empirical

description of results are presented. Detailed analysis of the results is presented in Sec. III starting from a simplified model of the Fermi-level dynamics within some abstract intersecting electron and hole levels. A technique to extract parameters of electron and hole gases from the low-field magnetoresistances for two and three kinds of carriers along with the obtained results are given in Sec. III A. These results are supplemented by the Fourier analysis of oscillations, and summarized in Sec. III B. Section III C is based on the calculated energy spectrum, first for a single QW. Then, in Sec. III C 1, the energy structure for the DQW is built based on obtained free-carrier parameters. In Sec. III C 2, the magnetic level pattern calculated for the single QW and the one built for DQW are presented. The observed experimental peculiarities in the QH range of fields are analyzed on this basis. The outcomes are summarized in Conclusions.

II. EXPERIMENTAL DETAILS

The samples are grown on a (013)-oriented GaAs substrate with their active part consisting of two HgTe layers, 20 nm wide, separated by a 6–10 nm Cd_{0.65}Hg_{0.35}Te barrier. The DQW is symmetrically doped with In in the adjacent outer barriers at the distances of 10 nm from the outer interfaces. Here, we outline the results for two samples, which are representatives of two investigated groups differed mainly by the hole densities p . The first one (sample 1) having a moderate $p = 6.5 \times 10^{15} \text{ m}^{-2}$ at $V_g = 0$ and the other (sample 2) with a larger p due to a larger amount of Hg vacancies. The samples are shaped in the double Hall bridge and covered with a gate above the silicon insulating layer, as in Ref. [12]. The results on the DQW are compared with those obtained on a single QW with identical parameters grown in the same technological cycle (20 nm wide, the same symmetrically doped). Diagonal and Hall magnetoresistances (MR), $\rho_{xx}(B)$ and $\rho_{xy}(B)$, were measured simultaneously by the direct current reversal technique using an improved version of original experimental setup [19] on a dc current of 0.25–2 μA under magnetic fields up to 31 T at a temperature of 0.31 K at a fixed or sweeping V_g (Figs. 1 and 2). In the latter case, a continuous sequence of bipolar triangular or sawtooth pulses was applied to the gate with amplitudes corresponding to the predetermined range of V_g at a frequency an order of magnitude smaller than the measuring rate at a sufficiently low-field sweep. Based on these data, $\rho_{xy}(B, V_g)$ and $\rho_{xx}(B, V_g)$ were built as continuous functions of two variables [see Figs. 2(c) and 2(d)].

The curves for $\rho_{xy}(B)$ and $\rho_{xx}(B)$ at weak fields for negative and small positive voltages $-3 \text{ V} \leq V_g \lesssim +1 \text{ V}$ have a structure relevant to the classical model of two kinds of free charges, a small quantity of high-mobility electrons and a large quantity of low-mobility holes, i.e., the N -shaped $\rho_{xy}(B)$ and a concomitant parabolic $\rho_{xx}(B)$ (Figs. 1 and 2), as it was first found in Ref. [2] for this heterosystem. MR oscillations are superimposed on this classical background. Within this range of V_g , the positive $\rho_{xy}(B)$ reaches the QH plateau at h/e^2 for the hole component of conductivity in the highest fields simultaneously with $\rho_{xx}(B)$ dropping to zero. Also, some feature exists in $\rho_{xy}(B)$ at $h/2e^2$ representing a distorted $i = +2$ QH state of holes.

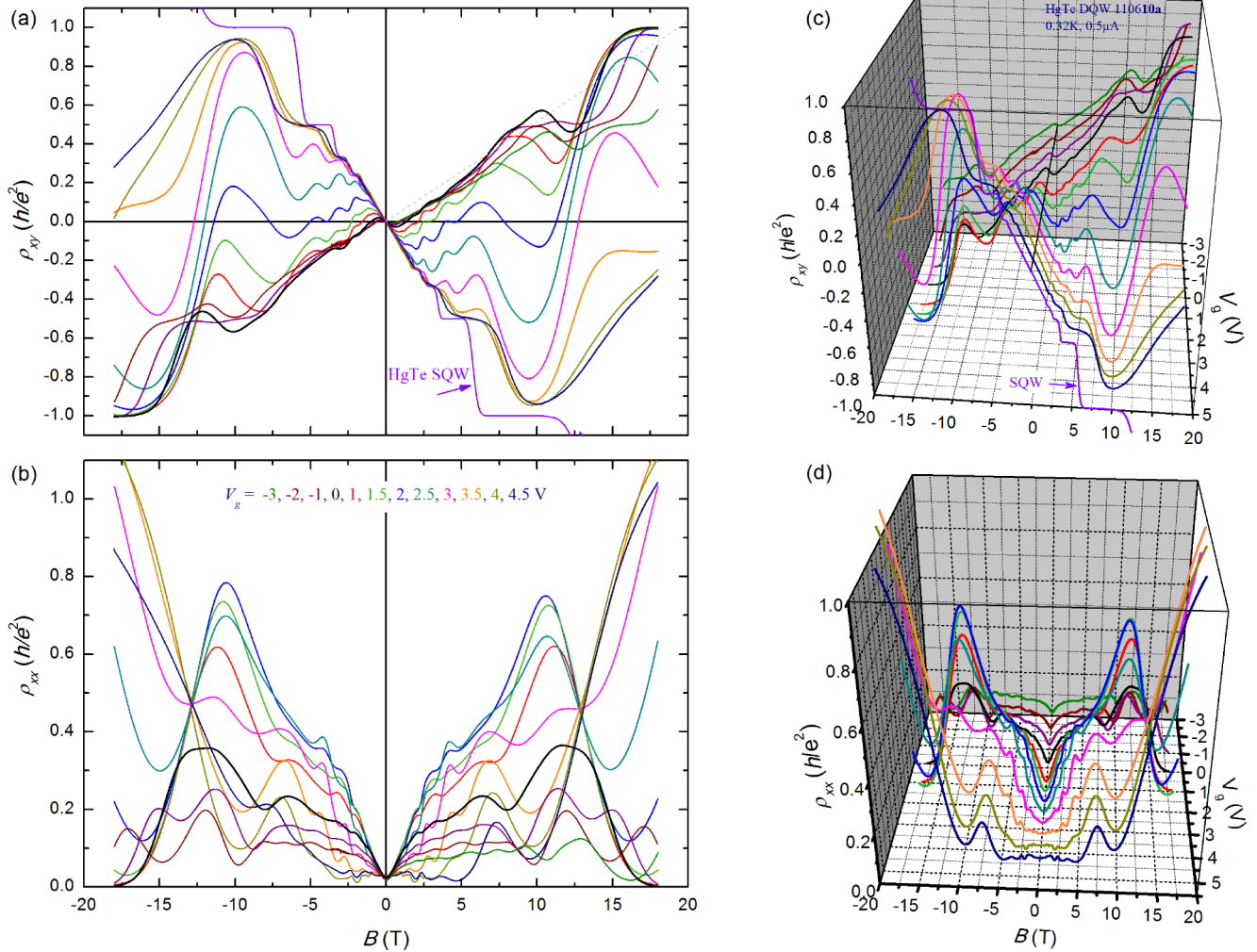


FIG. 1. (a) The Hall $\rho_{xy}(B)$ and (b) diagonal $\rho_{xx}(B)$ magnetoresistances of the DQW sample 1 in a series of gate voltages $V_g = (-3 - +4.5)$ V. The colors of digits for the V_g values in the figure correspond to the colors of relevant curves, and this color system is retained in other corresponding figures. (c), (d) The same in the bulk graph presentations. $\rho_{xy}(B)$ at large positive V_g are compared with the QH curve for an identical single QW.

There is a striking difference in the evolution of $\rho_{xy}(B)$ with V_g within the QH range of fields for negative and small positive V_g on one side and for the large V_g on the other. The QH features in the hole conductivity component move in field linearly with V_g in the former case as it should be for a regular picture of magnetic levels and a constant capacitance between the gate and the two-dimensional conducting system. To illustrate this linear dependence, the dashed inclined line is drawn through the middle of the $i = +1$ QH plateau area for sample 2 in Fig. 2(c). Also, such a linear shift is seen for the $i = +1$ to $+2$ plateau-plateau transition (PPT) in $\rho_{xy}(B)$ for both samples in Figs. 1(a) and 2(a). But, more complicated behavior shows up for the further increase of V_g . Instead of further shrinkage of the field scale as is expected in a simple case for the decreasing p , we observe that the mentioned PPT stops moving to weaker fields at ~ 15 T, then $\rho_{xy}(B)$ gradually sinks to negative values, passing through the states where several sign reversal points of $\rho_{xy}(B)$ exist (Fig. 1). Simultaneously, the $\rho_{xx}(B)$ minimum at the highest fields moves up from zero. This behavior is quite different from that observed in a single QW where the transition from a positive to negative $\rho_{xy}(B)$

of an almost undistorted QH shape occurs rather sharply (see, e.g., Ref. [10] for a narrow QW and Ref. [18] for a 20-nm-wide well).

III. DISCUSSION OF THE RESULTS

The described behavior may be explained on a simplified scheme of some abstract overlapped pure electron and hole magnetic levels (Fig. 3). Here, when at large enough p (corresponding to our case of negative V_g) E_F is sufficiently deep in the valence subband, the superposed electron levels are empty, as if they are absent at all, and we have only a regular fan chart of the populated hole magnetic levels. In this case, QH features would move linearly with decreasing p and, consequently, with increasing V_g . We modeled a finite width Γ of the levels by introducing a tentative Lorentzian density of states. Using Gaussian or semielliptic shape would not radically change the results. A QH feature in this picture corresponds to the range of fields where E_F is sufficiently distant from both the neighboring levels when it passes through the gap between them. Introducing the finite-level width makes

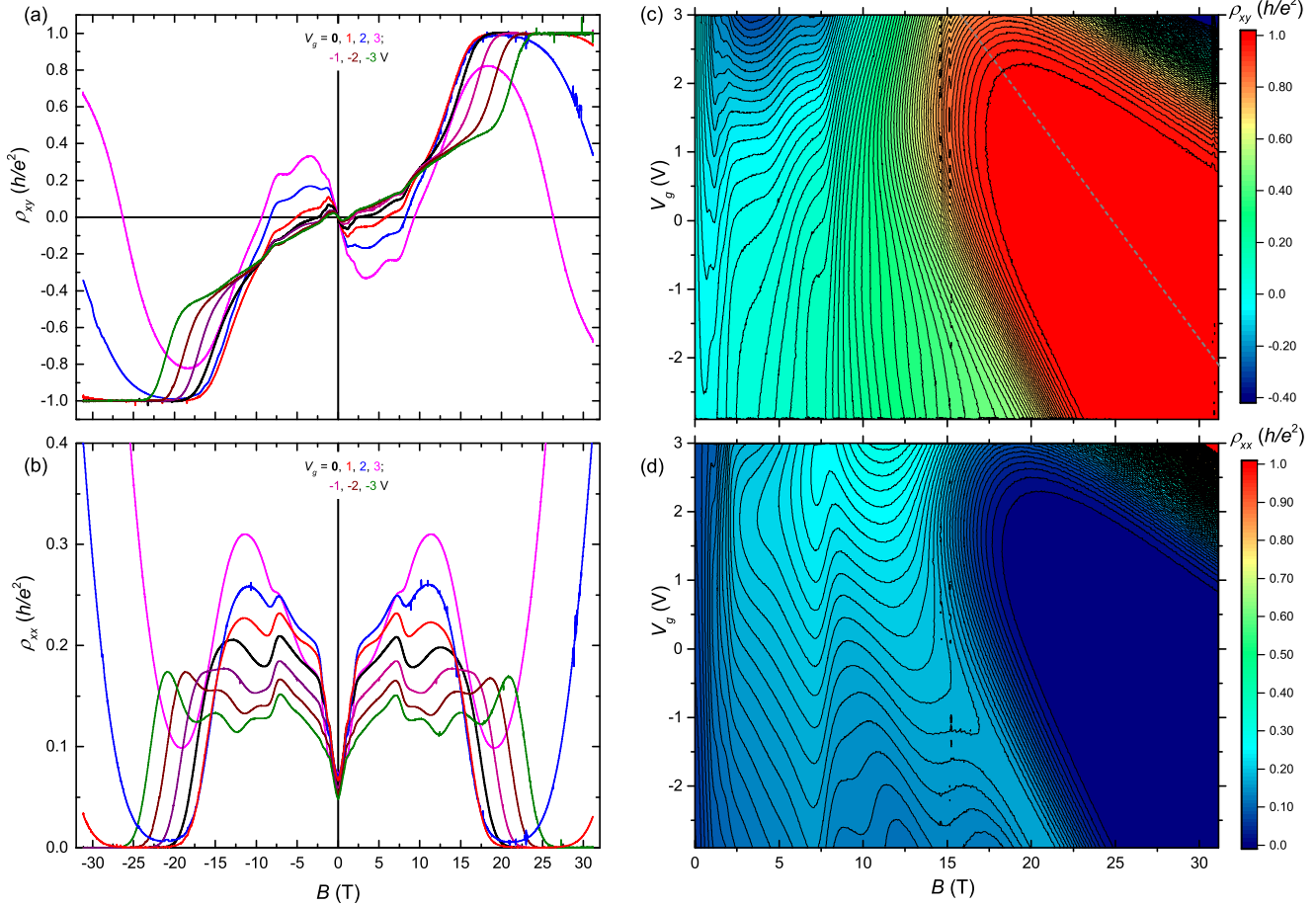


FIG. 2. (a), (c) The Hall $\rho_{xy}(B)$ and (b), (d) diagonal $\rho_{xx}(B)$ magnetoresistances of the DQW sample 2. Figures (a) and (b) are for the series of fixed values of $V_g = -3$ to $+3$ V while (c) and (d) are obtained by the continuous coordinated sweeps of V_g and B and presented in the color scale map view. The inclined dashed line in (c) marks the middle of the $i = +1$ QH plateau as it linearly shifts in B with V_g .

it possible to display the finite sizes of the QH features, but the main thing is that these calculations allow to represent automatically the behavior of E_F when the electron levels come into play with the hole level. In the picture of overlapping hole and electron levels, a hole level is populated by free holes above E_F simultaneously with free electrons occupying the electron level below E_F (Fig. 3, inset).

When, on decreasing p , the E_F comes into touch with an electron level and this level starts to populate with free electrons, the corresponding QH feature stops moving towards lower fields. To elucidate this behavior, the E_F curves calculated for the full set of levels (solid lines in Fig. 3) are compared with those calculated for the same values of $p - n$ but for the case as if there are no electron levels: $n = 0$, dotted lines. It is seen that population of electron levels considerably modifies the positions of QH features stopping them in a vicinity of the point where the corresponding hole level crosses an electron level. When E_F , on moving along a hole level with decreased field, reaches the crossing electron level, then a part of electrons occupying the hole level below E_F starts to move into the electron level (Fig. 3, inset) so that n appears and increases. That is why, at fixed $p - n$, p increases and the corresponding QH feature within a set of hole levels appears at the fields higher than the position where it would be without the electron level. This process hampers

the further shift of QH features with increasing V_g to lower fields within a set of hole levels. Thus, the experimental field where the QH features stop moving with V_g indicates the approximate position of some crossing point in the picture of levels. Another important consequence appearing when E_F approaches the electron level is an increased contribution of electron states into the magnetotransport that would manifest in a development of $\rho_{xy}(B)$ to negative values.

The reentrant sign-alternating behavior of $\rho_{xy}(B)$ has been revealed in a single HgTe QW in Ref. [12] at much lower fields of several Tesla, and a value of the band overlap in their sample was estimated as of the order of 1 meV. Much higher fields of the $\rho_{xy}(B)$ sign reversal in our DQW structures indicate much larger overlaps and much higher values of characteristic fields for the level crossings.

For the largest V_g , $\rho_{xy}(B)$ of sample 1 becomes fully negative and reveals a series of QH plateaus at intermediate fields due to electron conductivity component. However, $\rho_{xy}(B)$ turns back towards zero at the highest fields [Figs. 1(a) and 1(c)] reaching a plateaulike feature close to zero. To emphasize the difference of this anomalous behavior of $\rho_{xy}(B)$ in the DQW from the traditional one, the trace for a single n -type HgTe QW (measured without a gate) is added to Figs. 1(a) and 1(c). Concomitantly with the drop of $\rho_{xy}(B)$ in a DQW, a sharp growth appears in $\rho_{xx}(B)$ contrary to the

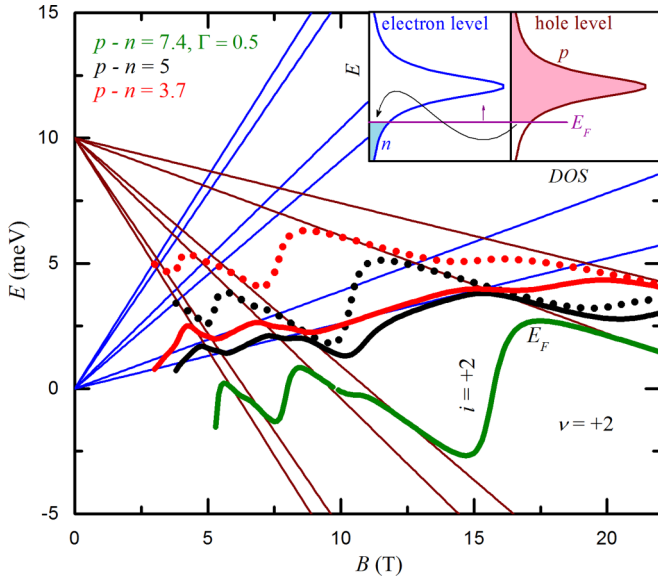


FIG. 3. Calculated evolution of the Fermi level within the modeled overlapped fan charts of the abstract pure hole and electron magnetic levels with decreasing hole density ($p - n$) (solid lines) compared to the corresponding cases without electron levels: $n = 0$ (dotted lines). The densities are indicated in the picture in 10^{15} m^{-2} . Lorentzian shape of the density of states in the level of the width $\Gamma = 0.5 \text{ meV}$ is supposed. Inset: populations of the electron and hole levels by the free electrons and holes, correspondingly. The arrow illustrates the transition of electrons from the lower part of the hole level into the electron level as the hole level is moving up with respect to the electron level with decreasing field.

behavior in the traditional QH effect (QHE). This quasizero plateau in $\rho_{xy}(B)$ is well reproduced in the conductivity (Fig. 4) and may be an indication of a zero-filling-factor state [11]. We do not see the electron QHE and its further development in sample 2 because of a smaller achievable $V_g = +3 \text{ V}$ due to a less reliable insulator under the gate in it. Nevertheless, the reentrant sign change of $\rho_{xy}(B)$ combined with the sharp growth of $\rho_{xx}(B)$ are well seen in Fig. 2 at maximal fields and the maximum positive V_g .

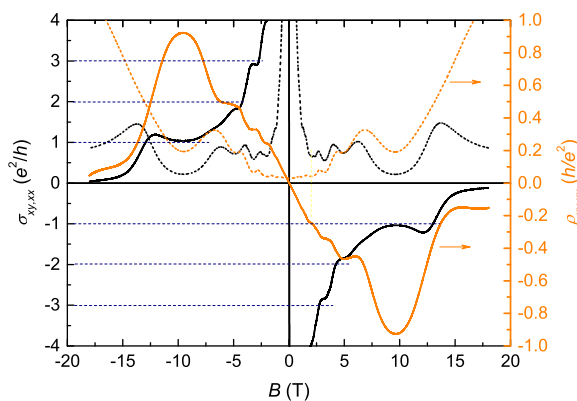


FIG. 4. MR and conductivities of sample 1 for $V_g = +3.5 \text{ V}$. Notice a drop of $\sigma_{xy}(B)$ to zero at the highest fields with concomitant minimum of $\sigma_{xx}(B)$.

In the range of the largest V_g , when a considerable electron part of $\rho_{xy}(B)$ is formed up to intermediate fields, remarkable is that the weak-field slope of this curve does not change with V_g , contrary to the data for single QWs [10,18] where the increase of free-charge density with V_g manifests in a monotonic decrease of its slope, i.e., in a monotonic decrease of the Hall voltage.

A. Classical treatment

To estimate the parameters of free charges in our structures, we first used a simple classical model for magnetotransport with two kinds of carriers (see, e.g., Ref. [20]), electrons with density n and mobility μ_e and holes with density p and mobility μ_h :

$$\rho_{xy} = \frac{B}{|e|} \frac{(p-n)\mu_h^2\mu_e^2B^2 + (p\mu_h^2 - n\mu_e^2)}{(p-n)^2\mu_h^2\mu_e^2B^2 + (p\mu_h + n\mu_e)^2}, \quad (1a)$$

$$\rho_{xx} = \frac{1}{|e|} \frac{(p\mu_e + n\mu_h)\mu_h\mu_e B^2 + (p\mu_h + n\mu_e)}{(p-n)^2\mu_h^2\mu_e^2B^2 + (p\mu_h + n\mu_e)^2}. \quad (1b)$$

To extract these four parameters, the following four quantitative experimental features in $\rho_{xy}(B)$ and $\rho_{xx}(B)$ were used: (i) a zero field resistivity ρ_0 ; (ii) a slope of the Hall resistivity at high fields R_H^∞ where it should approach the value $1/e(p-n)$ in the classical model; (iii) a slope of the Hall resistivity around zero R_H^0 ; and (iv) a field of the first inversion B_0 . Then, a system of four equations for four unknowns reads as

$$p - n = \frac{1}{eR_H^\infty}, \quad (2a)$$

$$n\mu_e + p\mu_h = \frac{1}{e\rho_0}, \quad (2b)$$

$$\mu_e\mu_h = \frac{\sqrt{-R_H^0 R_H^\infty}}{\rho_0 B_0}, \quad (2c)$$

$$n\mu_e^2 - p\mu_h^2 = -\frac{R_H^0}{e\rho_0^2}, \quad (2d)$$

which was solved numerically.

The classical description should be valid at weak fields while deviations from it may appear at higher fields caused by unconsidered complications due to quantization. This fit indeed works in our samples at the negative V_g . Thus, for sample 1, the weak-field fit is excellent for $V_g = -3 \text{ V}$ (Fig. 5) with $n = 1.89 \times 10^{15} \text{ m}^{-2}$, $\mu_e = 3.8 \text{ m}^2/\text{V} \cdot \text{s}$, $p = 8.87 \times 10^{15} \text{ m}^{-2}$, $\mu_h = 0.66 \text{ m}^2/\text{V} \cdot \text{s}$, while at higher fields the Shubnikov–de Haas oscillations (SHO) are superposed on the classical horizontal line for $\rho_{xx}(B)$ and a deviation from the inclined line for $\rho_{xy}(B)$ appears around the quantum Hall value of $\rho_{xy} = h/2e^2$. The fits look reliable until around $V_g = 0$ [Fig. 6(a)], where the obtained set is $n = 1.53 \times 10^{15} \text{ m}^{-2}$, $\mu_e = 4.42 \text{ m}^2/\text{V} \cdot \text{s}$, $p = 6.48 \times 10^{15} \text{ m}^{-2}$, $\mu_h = 0.47 \text{ m}^2/\text{V} \cdot \text{s}$. As seen from these and further estimations, the changes in MR with increasing V_g are mainly connected here with the drop in the hole density while the changes in n are not substantial and limited within the range of $n = (1.5-2) \times 10^{15} \text{ m}^{-2}$ for the whole set of negative and positive V_g . The weakness of the electron density dependence

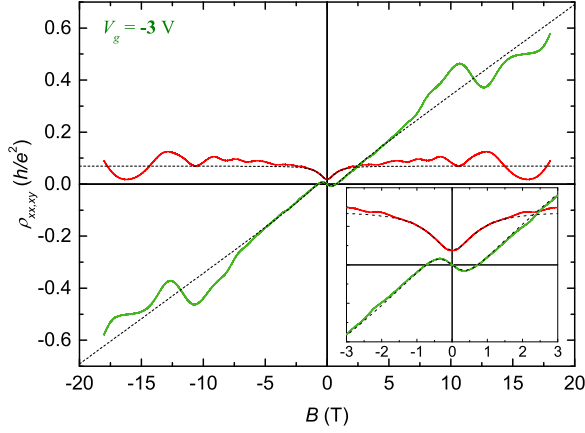


FIG. 5. Two-carrier fit to MR of sample 1 for $V_g = -3$ V with an enlargement in the inset.

on V_g is well seen for large $V_g > \sim 3.5$ V where the weak-field part of $\rho_{xy}(B)$ shows a monotonous linear behavior due to a single kind of carriers, electrons only: $\rho_{xy} = B/en$, with the slope almost independent on V_g (Fig. 1).

With a further increase of V_g , some factors of uncertainty appear in the overall MR. The first warning is that the QH features shift below the fundamental values of $\rho_{xy} = h/ie^2$, initially for the distorted feature at $i = 2$ for $V_g \geq +1.5$ V

$$\rho_{xy} = \frac{B}{|e|} \frac{(p - n_1 - n_2)B^4 + [p(b_1^2 + b_2^2) - n_1(b_2^2 + b_h^2) - n_2(b_1^2 + b_h^2)]B^2 + pb_1^2b_2^2 - n_1b_2^2b_h^2 - n_2b_1^2b_h^2}{(p - n_1 - n_2)^2B^4 + [(n_1b_2 + n_2b_1)^2 + c.p.1 + 2(n_1n_2b_h^2 - n_2pb_1^2 - n_1pb_2^2)]B^2 + (n_1b_2b_h + c.p.2)^2}, \quad (3a)$$

$$\rho_{xx} = \frac{1}{|e|} \frac{(n_1b_1 + n_2b_2 + pb_h)B^4 + [n_1b_1(b_2^2 + b_h^2) + n_2b_2(b_1^2 + b_h^2) + pb_h(b_1^2 + b_2^2)]B^2 + b_1b_2b_h(n_1b_2b_h + n_2b_1b_h + pb_1b_2)}{(p - n_1 - n_2)^2B^4 + [(n_1b_2 + n_2b_1)^2 + c.p.1 + 2(n_1n_2b_h^2 - n_2pb_1^2 - n_1pb_2^2)]B^2 + (n_1b_2b_h + c.p.2)^2}, \quad (3b)$$

with

$$c.p.1 \equiv (n_2b_h + pb_2)^2 + (n_1b_h + pb_1)^2,$$

$$c.p.2 \equiv n_2b_1b_h + pb_1b_2,$$

$$b_{1,2,h} \equiv 1/\mu_{1,2,h}.$$

Using Eqs. (3), it is indeed possible to obtain a second reversal of $\rho_{xy}(B)$ (Fig. 7) for a realistic set of parameters: $n_1 = 1 \times 10^{15} \text{ m}^{-2}$, $\mu_1 = 4.53 \text{ m}^2/\text{V} \cdot \text{s}$, $n_2 = 0.14 \times 10^{15} \text{ m}^{-2}$, $\mu_2 = 0.28 \text{ m}^2/\text{V} \cdot \text{s}$, $p = 1.13 \times 10^{15} \text{ m}^{-2}$, $\mu_h = 1.14 \text{ m}^2/\text{V} \cdot \text{s}$ (n_1 , n_2 , μ_1 , and μ_2 are the densities and the mobilities of two kinds of electrons). Of course, these values are less reliable than in the case of a single inversion as here are six fitting parameters and we used the overall fitting within a range of weak fields. Thus, this approach is rather an illustration of possible manifestation of heavy electrons. The next inversion of $\rho_{xy}(B)$ to positive values in Fig. 7 cannot be modeled in the three-carrier fit. A second kind of heavy holes with a still larger mass is needed for this. Thus, the task will be transformed into a four-carrier model. The necessity to add extra holes is seen from that $\rho_{xy}(B)$ is positive at the highest fields, thus indicating that the sum of the hole densities is larger

[Figs. 6(a) and 6(b)], then for the plateau with $i = 1$ for $V_g \geq +2.5$ V [Figs. 6(c) and 6(d)]. The development of this tendency leads to appearing additional inversion points in $\rho_{xy}(B)$: Fig. 6(c) for $V_g = +2$ V and, with the further lowering of $\rho_{xy}(B)$, for $V_g = +2.5$ V [Fig. 6(d)] and $V_g = +3$ V (Fig. 1). The nature of these inversion points is certainly different from the discussed above so that using them as a fitting point becomes problematic. Another characteristic feature used in fitting, i.e., a linear slope in $\rho_{xy}(B)$ at high classic fields, disappears at large V_g . Thus, description by the classical two-carrier model becomes doubtful in this range of V_g . The origin of these transformations may be in the admixture of additional electron states to the states of the hole magnetic levels as in speculations around Fig. 3.

A guess to explain this complicated behavior in classical terms is that it is a manifestation of an additional kind of particles with lower mobility. Their nature may be connected with the complicated structure of the valence subband top where an electronlike curvature of $E(k_{\parallel})$ exists in some range of wave-vector values around zero between the lateral maxima. These particles should behave like heavy electrons as the $E(k_{\parallel})$ curvature is lower here than in the conduction subband. The addition of low-mobility electron component to conductivity will cause a modification of Eqs. (1) into Eqs. (3), obtained by the reversal of conductivity tensor into the resistivity tensor:

than that for the electrons, contrary to the presented values obtained in the three-carrier fit. A nature of the second kind of holes may be suggested as being due to an asymmetry of the lateral maxima in the valence subband, thus causing different values of dE/dk on its opposite slopes. Some asymmetry in the DQW potential profile will lead to that different parts of the valence subband top would participate in the magnetotransport in the two layers, thus providing heavy electrons from one layer and the second type of holes from the other one. We did not perform a four-carrier fit because of complexities and unreliable values expected for the eight variable fits, but qualitatively the situation seems feasible once the classical approach works.

B. Oscillations in $\rho_{xx}(B)$

The uncertainty in obtaining parameters of free-charge carriers from the classical magnetotransport at $V_g \gtrsim 0$ is a motivation to look for other ways to find them. Experimental curves for $\rho_{xx}(B)$ contain rich pictures of oscillations, although masked by the large and nonmonotonic backgrounds [Figs. 1, 2 and 8(a)]. We extract the oscillations by taking their second derivative $-d^2\rho_{xx}/dB^2(B)$ [Fig. 8(b)] that considerably

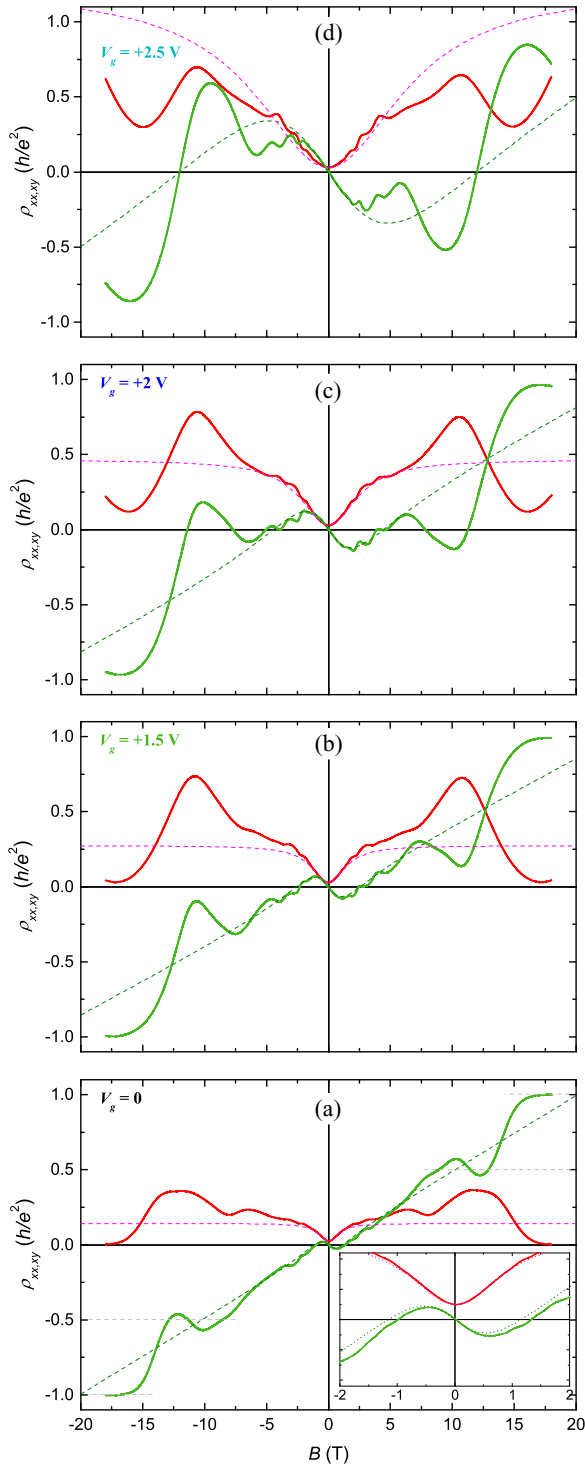


FIG. 6. Two-carrier fits to MR of sample 1 for (a) $V_g = 0$ with an enlargement in the inset; (b) $V_g = +1.5$ V; (c) $V_g = +2$ V; and (d) $V_g = +2.5$ V. Notice the ambiguities in two-carrier fits for the large positive V_g values.

enhances fast oscillations in the weak-field region. Oscillations demonstrate several periods in reciprocal magnetic fields $-d^2\rho_{xx}/dB^2(1/B)$ [Fig. 8(c)], that is expected for several populated subbands. To do the analysis feasible, considering weak initial amplitudes and a limited number of oscillations, we performed separate fast Fourier transformations (FFT) in

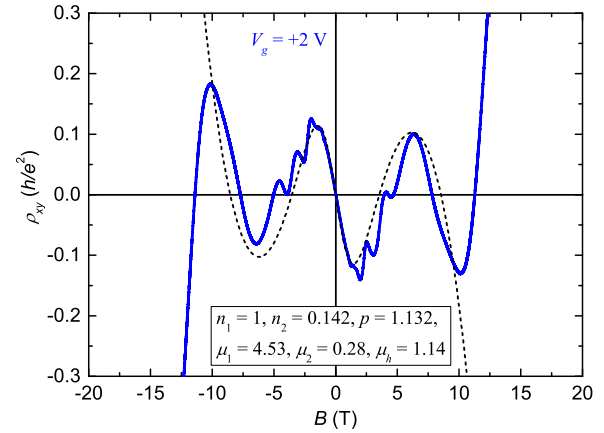


FIG. 7. Three-carrier fit to $\rho_{xy}(B)$ for $V_g = +2$ V.

the sectors of high fields, where the high-density holes manifest themselves in the high-field oscillations for the negative V_g , and within a low-field window, where the oscillations should be due to the low-density electrons [Figs. 8(d) and 8(e)]. The upper axis is added for the electron and hole densities related to the FFT frequency f as $n, p = 2fe/h$ for the unresolved spin splittings.

In the low-frequency range, the FFT curves reveal two main peaks: one around $n_1 = 1.5 \times 10^{15} \text{ m}^{-2}$ and the other at $n_2 \approx 0.5 \times 10^{15} \text{ m}^{-2}$ as depicted by two dashed verticals in Fig. 8(d). No directed movement of these peaks with V_g is seen. In the high-frequency window [Fig. 8(e)], FFT peaks are revealed for the set of V_g between -1 and -3 V. One peak is close in its positions to the hole densities found from the two-carrier fits (indicated by arrows), if evaluations are performed for the spin-unresolved case as $p = 2fe/h$. The second peak seen in Fig. 8(e) at approximately double frequencies should be connected with the spin-resolved component of the same oscillations and yields the same p values if estimated for the spin-resolved case as $p = fe/h$. The stable directed shift of the hole FFT amplitudes is seen with the changes of V_g to negative values.

The results of all estimations for all kinds of carriers are collected in Fig. 9: for two-carrier fits that yield summary densities for electrons and holes and their mobilities for $V_g \leq 0$; for simple weak-field classical linear $\rho_{xy}(B)$ yielding n at maximum positive V_g ; and for Fourier treatment of oscillations. Comparing FFT results for electrons and the estimates for n from the two-carrier fit, it should be considered that in a latter case the fit would not feel the possible distribution of electrons between two subbands if they have similar mobilities in the subbands, while FFT feels the difference in their densities. Then, the sum $n_1 + n_2$ obtained from FFT should be compared with n obtained from two-carrier fit. As seen from Fig. 9, the dots for $n_1 + n_2$ are between the data for two-carrier fits and simple classical linear weak field $\rho_{xy}(B)$ in a reasonable agreement. The general conclusion about electrons is that n is almost independent on V_g .

We have reliable data on p only for $V_g \leq 0$ both for two-carrier fits and FFT analysis. However, we can add one supplementary estimation for the positive V_g . The Hall MR $\rho_{xy}(B)$ in its evolution with increasing V_g becomes fully

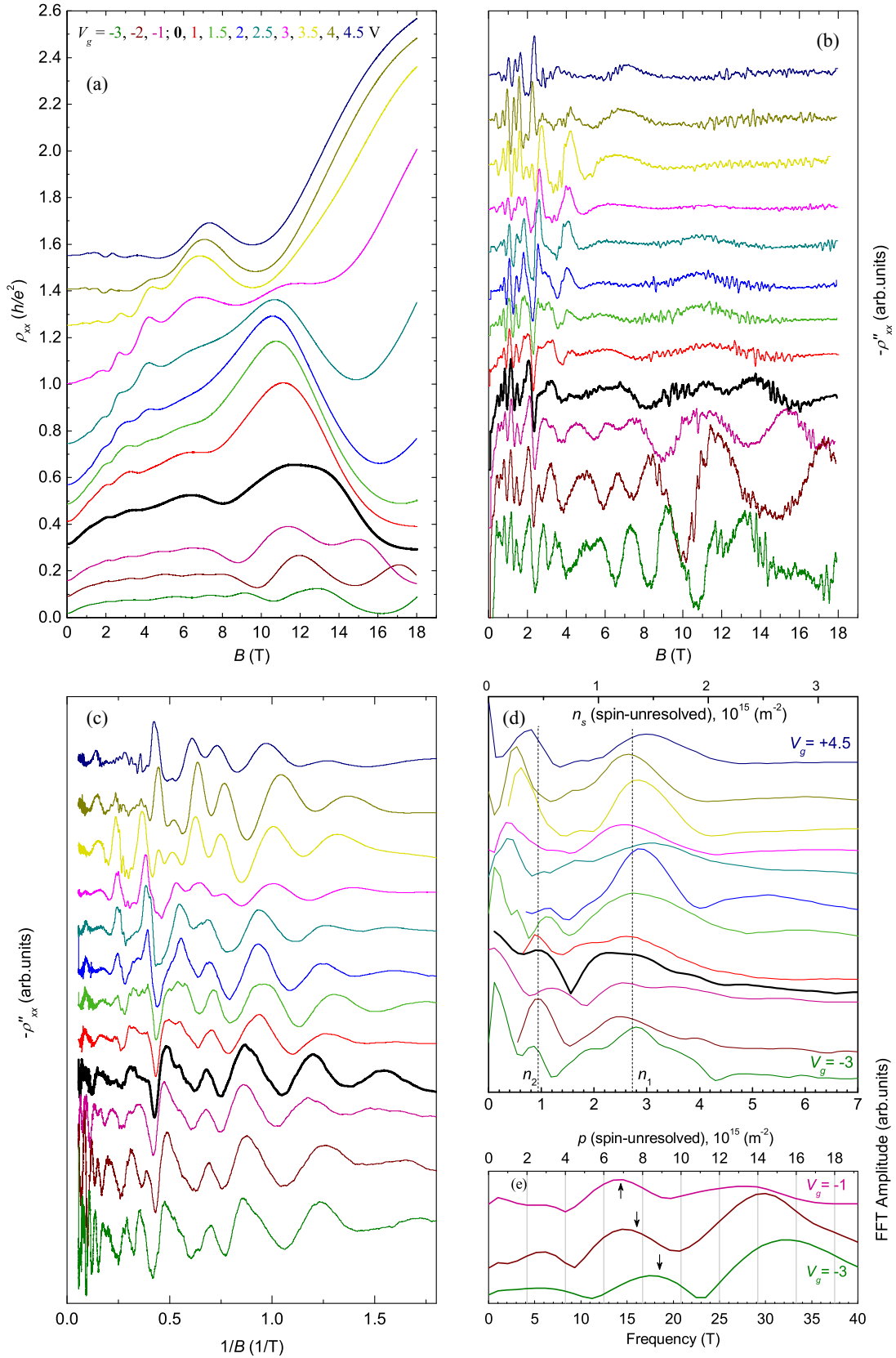


FIG. 8. (a) $\rho_{xx}(B)$ of sample 1 for different V_g , vertically shifted for clarity from the lowest $V_g = -3$ V at the bottom to the highest $V_g = +4.5$ V at the top; (b) their second derivatives $-d^2\rho_{xx}/dB^2$; (c) $-d^2\rho_{xx}/dB^2$ versus reciprocal field and (d), (e) their FFTs. (d) The FFT results for the low-field windows corresponding to the low densities of electrons. Two dashed verticals correspond to some averaged positions of FFT peaks for n_1 and n_2 . (e) The FFT results for the high-field windows corresponding to the high densities of holes. The arrows indicate p values obtained from the classical two-carrier treatment.

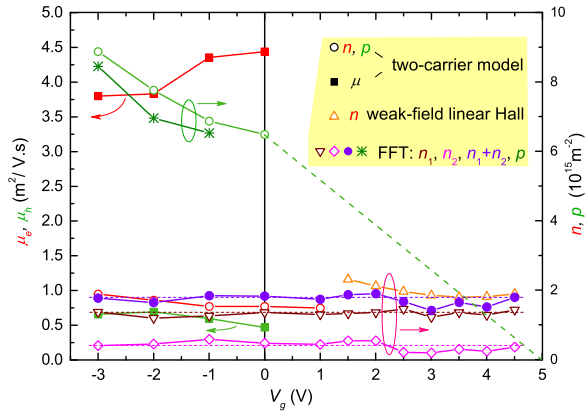


FIG. 9. The results on the electron (lines and symbols of the red tint) and hole (green tints) densities (circles, triangles, diamonds, asterisks) as well as mobilities (squares) obtained for sample 1 in the two-carrier fits (empty circles for densities), from weak-field linear Hall (up triangles) and FFT (down triangles and diamonds for two electron-components and solid circles for their sum, asterisks for holes).

negative at $V_g = 3.5$ V (Fig. 1). It should happen when n exceeds p . Thus, we can add this CNP to the function $p(V_g)$: $p = n$ at $V_g = 3.5$ V. The dashed line in Fig. 9 depicts this rough prolongation of $p(V_g)$ to the positive V_g . The general conclusion on p : it pronouncedly decreases with V_g from $p \approx 9 \times 10^{15} \text{ m}^{-2}$ at $V_g = -3$ V to zero at about $V_g = +5$ V.

C. Analysis based on the energy band calculations

The energy band structure of a single fully strained QW, of which our DQW is built, was calculated in an envelope-function approach within the framework of eight-band $\mathbf{k} \cdot \mathbf{p}$ theory self-consistently with the Poisson equation for the charge distribution [21]. We used the simple [001] orientation for calculations in spite of [013] growth direction in our structures hoping that the 18° deviation between these directions would not influence considerably the results of our analysis, which is mainly of qualitative character. The resultant band structure is presented in Fig. 10.

In the valence subband, $E_v(k_{\parallel})$ differs for the (1,0) and (1,1) directions in the (k_x, k_y) plane (bold solid and dashed curves in Fig. 10, respectively), so that the constant energy contours are warped. We ignored this warping in our rough estimations of the carrier densities and used the averaged curves for the valence subband dispersion (thin green curves in Fig. 10) in the further numerical analysis.

A substantial overlap about 6.5 meV of the valence and conduction subbands is obtained when the strain is considered in calculations (while the overlap is absent at all without strain). But, this overlap would not be felt experimentally in a single QW at electron densities $n > \sim 1 \times 10^{15} \text{ m}^{-2}$ since the Fermi level is above the overlap region. Thus, attempts to use this picture without any modifications for interpretation of our experiments on DQWs would not be successful as we have $n > 1.5 \times 10^{15} \text{ m}^{-2}$.

Calculations of the energy spectra for HgTe/CdHgTe DQW [22] show that the tunneling gap between the Γ_8 states is

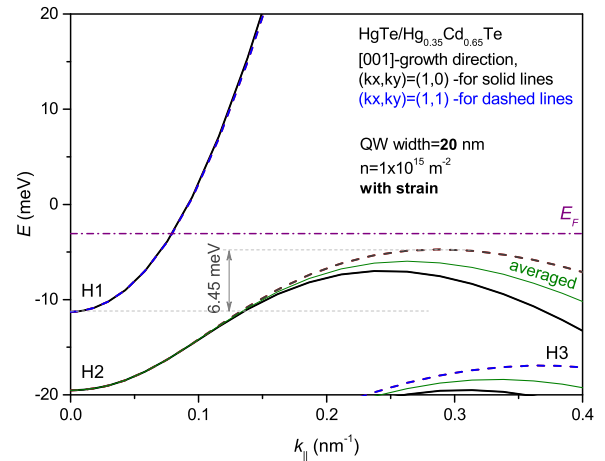


FIG. 10. The calculated band structure for a single quantum well with parameters for the layers of which our DQW is built. Bold solid lines are for $(k_x, k_y) = (1, 0)$, dashed lines for $(k_x, k_y) = (1, 1)$. Thin solid curves are for the averaged structures of the bands. Dotted-dashed horizontal line is an example of the Fermi-level position for $n = 1 \times 10^{15} \text{ m}^{-2}$.

vanishingly small. In this case, the DQW energy spectrum may be obtained by a simple superposition of pictures calculated for single QWs composing the DQW, with a relative shift in energy if some electric field presents across the structure. Here, remarkable is that we see the QH plateau with $i = 1$ in the DQW even for zero gate voltage. For zero tunneling gap, the $i = 1$ QH state is impossible in a DQW with symmetric potential profile. In this case, the $i = 1$ QH plateaus in each of the composing equal single QWs would yield the resulting Hall resistivity of the whole structure twice as small as for each layer switched in parallel, thus corresponding to $i = 2$ in the measurements on the DQW sample. In other words, the equal QWs would have equal sets of magnetic levels and combining these layers into a DQW structure of symmetric potential profile would simply lead to the doubling of all levels. Thus, the only possibility for appearing the gaps with odd filling factors in this case is an energy shift between the two layers. The existence of such a shift at $V_g = 0$ may be due to asymmetric distribution of fixed charges: ionized dopants and uncontrolled impurities, surface charges, etc., or due to some differences in the shape of QWs.

1. Free-charge distributions between subbands

We can reconstruct the energy spectrum of our DQW and its evolution with V_g using the data on the free charges in the subbands. First, we take an energy spectrum calculated for a single QW and determine where the Fermi level should be drawn in this picture to get the experimentally obtained p . For example, for $p = 8.87 \times 10^{15} \text{ m}^{-2}$ measured at $V_g = -3$ V, we calculate the area S_k occupied by holes in the two-dimensional k space using the relation $p = S_k/2\pi^2$ for the spin-degenerated states and adjust the E_F position with respect to the lateral maxima of the valence subband to get this area of the ring cut from it by E_F (see the solid red upper valence subband in Fig. 11 with the obtained E_F position and the corresponding ring in the inset). The derived E_F cuts a part

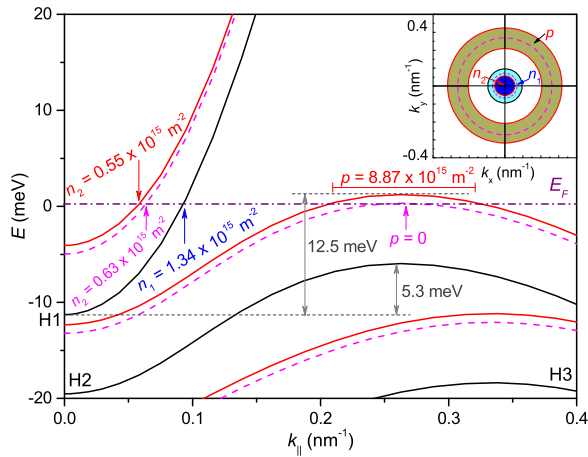


FIG. 11. The band structure of a double quantum well obtained by superposition of two similar pictures for the single QW (see Fig. 10, averaged) with an energy shift for the upper layer (upper curves of the red tint) with respect to the lower one (black lower curves). Energy of the upper layer is sensitive to V_g while E_F is fixed in the buffer structure below DQW. Solid red lines are for $V_g = -3$ V and the pink dashed lines are for the moment when becomes $p = 0$ (expected at $V_g = +5$ V). Notice that the simultaneous change in n is tiny. Inset: the corresponding Fermi contours in the (k_x, k_y) plane (for the averaged energy in the valence subband so that the warping is ignored). Notice a huge difference in the filled areas of the valence and conduction subbands.

of the conduction subband of this layer (see another solid red upper curve for the conduction subband), and a similar calculation yields $n_2 = 0.55 \times 10^{15} \text{ m}^{-2}$. This value is less than the measured $n = 1.89 \times 10^{15} \text{ m}^{-2}$. The missing part of electrons is in the conduction subband of the other (lower) layer of the DQW: $n_1 = 1.34 \times 10^{15} \text{ m}^{-2}$. To fulfill this condition, the energy spectrum of the lower layer is superposed onto that of the upper layer (black lower curves in Fig. 11) with a shift down on the energy of 7.2 meV that yields the area cut by E_F from the conduction subband of the lower layer for the needed value of n_1 . This way, the value of the energy shift between the layers is determined. This superposition results in an increased overlap between the conduction and valence subbands from 6.5 meV within a single layer to 13.7 meV between the layers at the expense of the inclusion of the interlayer overlap. The overlapped states of the two layers contribute to the common conductivity of the DQW with influence of all the possible interlayer interactions between them.

Variation of V_g causes a shift of the upper-layer energy with respect to the lower-layer energy while the latter is insensitive to V_g as the lower layer is screened from the upper gate by the free charges in the upper layer. The upper-layer energy also shifts with respect to E_F , which is fixed in the lower layer and the part of the structure below. As we obtained experimentally, the increase of V_g to positive values causes a drop of p from a substantial value to zero while n remains almost unchanged. The found DQW spectrum (Fig. 11) allows to explain this behavior. (i) The main part of the total electron density n is in the lower-layer conduction subband (n_1), which is insensitive to V_g . (ii) A very large density of states corresponds to the holes around the lateral maxima. These holes occupy a ring

(a warped ring, to be more exact) in the 2D k space with an area an order of magnitude larger than for circle areas corresponding to electrons [Fig. 11 (inset)]. A large effective mass corresponds to these holes. That is why a small decrease in the valence subband position with respect to E_F causes a large drop in p . So, for the drop from $p = 8.87 \times 10^{15} \text{ m}^{-2}$ (at $V_g = -3$ V) to zero (expected at $V_g = +5$ V), only a lowering of the upper-layer energy on 0.9 meV is needed (see the dashed pink curves in Fig. 11). This shift causes only a small increase of the electron density in the conduction subband of the upper layer, from $n_2 = 0.55 \times 10^{15} \text{ m}^{-2}$ to $0.63 \times 10^{15} \text{ m}^{-2}$, that amounts to only about 4% of the total increase in n , on the background of large n_1 . This is also illustrated in the inset to Fig. 11 for the same shift: while the occupied valence subband ring shrinks into a pink dashed circle line with the lowering of the upper-layer energy, the conduction circle for n_2 bordered with the red solid circle line expands only a little to the pink dashed circle contour whereas the black circle contour for n_1 remains unchanged.

This interpretation confirms the existence of two components in the electron conductivity seen in the Fourier spectra of oscillations where two FFT peaks for n_1 and n_2 were obtained [Fig. 8(d)]. The DQW spectrum presented in Fig. 11 explains why the $i = +1$ QH state exists in DQW: the valence subband of the lower layer is shifted much below E_F in the whole range of V_g . Thus, it is always empty of holes and do not contribute to the hole component of the QHE.

Why did we choose a scheme with the energy of the upper layer shifted above the energy of the lower layer? It is due to the observed dynamics of magnetotransport pictures with V_g . In the opposite case (swap the colors in Fig. 11 between the red tint and the black), p will be independent of V_g as it will be in the valence subband of the lower layer and all the changes with increasing V_g should be at the expense of large shifts down of the upper-layer conduction subband energy, i.e., for the large increase in n_2 and of the total electron density n , that is in drastic contrast with all our experimental data.

2. Peculiarities in the QH range of fields

The scheme for computing $E(k_{\parallel})$ was developed further for calculations of magnetic levels [21]. The magnetic level picture for the spectrum of Fig. 10 is presented in Fig. 12. Here, an example of the Fermi-level dynamics for $n = 1 \times 10^{15} \text{ m}^{-2}$ is shown for the simplest case of δ -shaped density of states in the levels. It is noteworthy that while E_F does not touch the valence subband in $E(k_{\parallel})$ for this n (Fig. 10), it nevertheless enters the set of valence magnetic levels between 4 and 5 T in Fig. 12. This may be connected with the anomalous nature of the lowest electron-magnetic level. It is a so-called zero-mode level [7,8]: It starts from the valence subband at $B = 0$ but monotonously moves up with field, as it should be for an electron level, while all the other valence subband levels, also starting from the valence subband at $B = 0$ and moving initially like electron levels, turn down later at a certain field as they reach the energy of the lateral maximum, thus behaving further like the hole levels. A nonmonotonic behavior of the main set of the valence subband levels is well understood from a quasiclassical view for the valence subband with the lateral maxima, but the anomalous behavior of the zero-mode level

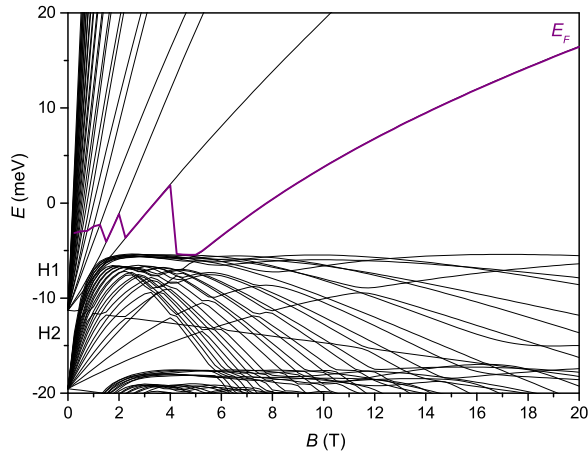


FIG. 12. Picture of magnetic levels calculated for the energy spectrum of Fig. 10 with an example of the schematic Fermi-level dynamics for $n = 1 \times 10^{15} \text{ m}^{-2}$.

does not have a quasiclassical interpretation. We believe that this zero-mode level yields a negative contribution to the Hall voltage, thus proving its electron character. It is confirmed by the calculated behavior of this level as it approaches the sample edge: it moves up [9,10] like all the other electron levels do, contrary to the hole levels, which move down. Another zero-mode level, but of the hole nature, is seen in Fig. 12: it starts at $B = 0$ from the conduction subband but monotonously moves down as it should for the hole level. Being the hole level it should not contribute to the conductivity as it should be empty of holes when E_F is above it. In fact, the movement of $E_F(B)$ is more complicated within the levels of finite width especially within the field ranges where the electron and hole levels overlap, according to the modeled picture of levels in Fig. 3.

The magnetic level pattern for the DQW is obtained similarly to that was done for the DQW energy spectrum in Fig. 11: the chart of levels calculated for the single QW in Fig. 12 is superimposed onto the same chart with the energy shift on the already known value of 7.2 meV. We used the so-obtained diagram of DQW levels for interpretation of the observed features in the QH range of fields. Of course, this interpretation is far from exact, but we hope at least to understand some of the main trends. The Fermi-level dynamics within the obtained plot of the DQW magnetic levels was calculated for the finite-width Lorentzian-shaped density of states in the levels similarly to that in Fig. 3. In these schematic calculations, the levels in the valence subband were assumed to be of the hole nature within the whole range of fields (in spite of their electronlike behavior at low fields), i.e., their contribution to $\rho_{xy}(B)$ was programed as positive, with exception of the zero-mode level, which was programed as electron level as well as regular levels in the conduction subband. More detailed consideration of the nature of states in the levels is left for the future. The results of our calculations are presented in Fig. 13 for different free-charge densities ($p - n$) corresponding to the values of V_g in experiment and compared to the measured $\rho_{xy}(B)$.

The $i = 1$ QH plateau observed at $V_g = 0$ corresponds to the position of E_F inside the triangle-shaped gap formed

by the arc-shaped course of a level near the lateral maximum at the top and the positively and negatively sloped levels on its flanks [Fig. 13(a)]. The field B_1 when E_F passes through the middle of the gap in this triangle correlates with QH plateau on $\rho_{xy}(B)$. With decrease of $V_g < 0$, the tendency is seen as the calculated B_1 values move to the higher fields in agreement with our experiment. The transition region between QH states for $i = 1$ and 2 corresponds to the left flank of the triangle. The movement of the ($i = 1$) – ($i = 2$) PPT to higher fields with decreased V_g correlates with the shifts of the point where E_F intersects the left flank.

The gap for filling factor $\nu = 2$ to the left from the $\nu = 1$ triangular gap is formed by two negatively sloped and one positively sloped levels. It differs from the $\nu = 1$ gap in shape and is much narrower. E_F passes through this gap in the vicinity of three points where the levels cross. Existence of groups of closely spaced levels was observed to cause a similar distortion of the QH plateau in the shape of a dip (e.g., in a traditional DQW [23,24]). A possible cause for appearance of such a dip on the QH plateau may be connected with an interplay between current carrying strips at the sample edges [25], and close positions of the neighboring levels may promote the effect.

With increase of $V_g > 0$, the hole density decreases and the calculated Fermi level shifts upwards, finally penetrating into a gap above the lateral valence subband maxima and below the lowest electron level [Figs. 13(b) and 13(c)]. In the joint picture of the hole and electron levels, population of the hole levels is described by the positive filling factors $\nu = +1, +2, \dots$ and population of the electron levels by the negative ones $\nu = -1, -2, \dots$. Then, the gap above the valence subband lateral maxima is formally designated by $\nu = 0$. It is characterized by an interference of the hole and electron states resulting in a drop of $\rho_{xy}(B)$ to zero in the vicinity of CNP. Simultaneously, $\rho_{xx}(B)$ diverges: Figs. 1 and 2, as it was observed in a single HgTe QW at much smaller fields in the local [12] and nonlocal [26] transport. Possible physics causing a sharp drop of $\rho_{xy}(B)$ to zero with simultaneous growth of $\rho_{xx}(B)$ may be connected with a formation of the interlayer electron-hole excitons [5] causing specific ratios for scattering probabilities between different channels [12].

E_F on passing with increased V_g through the valence subband edge at intermediate fields goes repeatedly across the field regions where it is either below or above the hole magnetic levels [Figs. 13(b) and 13(c)]. This is due to the oscillating profile of the valence subband top formed in quantized magnetic fields by nonmonotonic course of the hole levels with different numbers around the energy of lateral maximum. Because of that and the superimposed electron levels, the balance of the states on E_F having the hole or electron character is alternatively changed. The contribution of electron states is enhanced when E_F goes along the electron level between 8 and 9 T. All this results in several inversions in the course of $\rho_{xy}(B)$.

Speculations developed on the modeled level picture presented in Fig. 3 that the movement of the hole QH features with V_g is stopped when E_F reaches a crossing point between the hole and electron levels may be justified in the pictures of the calculated levels shown in Fig. 13. Here, the first stopping

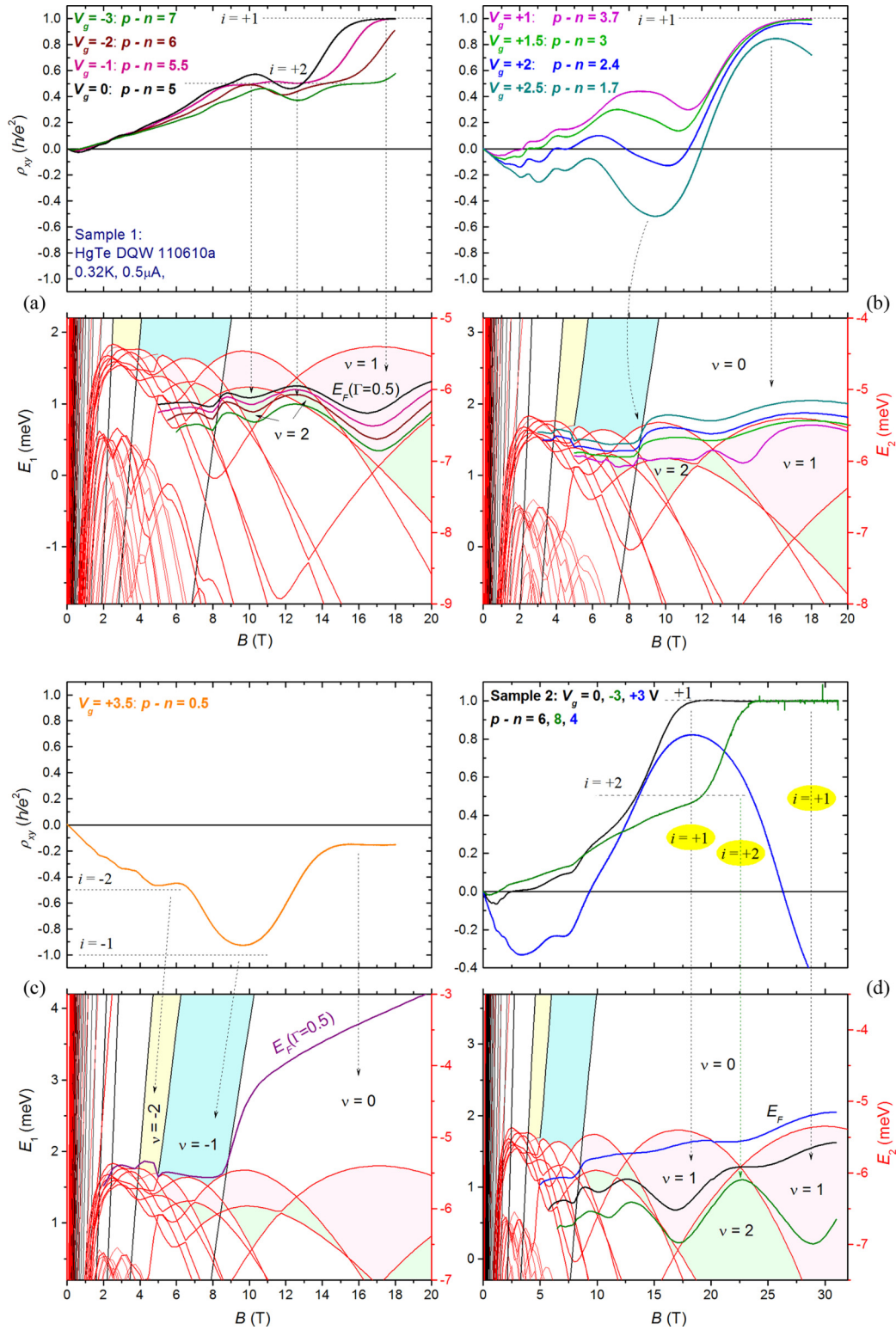


FIG. 13. Quantum Hall features in $\rho_{xy}(B)$ of sample 1 (a)–(c) and sample 2 (d) interpreted on the basis of calculated DQW magnetic levels and obtained free-charge densities. Red (black) magnetic levels are for the upper (lower) layer as in Fig. 11. (a) for $V_g \leq 0$; (b) for $V_g > 0$ in the range of multiple inversions; (c) in the vicinity of the charge-neutrality point. E_F is calculated the same way as for Fig. 3. The gaps for the same filling factors are painted the same colors.

point is the crossing between two valence subband levels around $B = 13$ T (the left corner of the actual $\nu = 1$ triangular gap), which appears due to that the hole level behaves in an

electron manner to the left from the field where it reaches the lateral maximum. This crossing point is responsible for the stop of moving in field of the $(i = 1) - (i = 2)$ PPT with

increased V_g . Another actual crossing point also appears in the calculations (which is more in the spirit of the model in Fig. 3): the crossing of the lowest-electron level with the valence subband edge between $B = 8$ and 9 T. This point correlates with the minimum formed in $\rho_{xy}(B)$ around 10 T, which sinks to negative values at almost constant fields manifesting an increase in contribution of the electron states to magnetotransport at E_F .

Similar behavior of quantum magnetotransport is observed at higher fields in sample 2 [Fig. 13(d)]. A distinct $i = +1$ QH plateau at $V_g \leq +2$ V (see also Fig. 2) correlates with location of E_F inside the triangular gap for $\nu = +1$. The curve for $V_g = -3$ V shows an inflection point at $\rho_{xy}(B) \approx \frac{1}{2}$ that correlates with the passage of E_F within the $\nu = +2$ gap, but closely to its upper edge, that is why this QH feature is weakly resolved. The $i = +1$ plateau is rather wide for $V_g = 0$. It corresponds to the passage of E_F through the two adjacent triangular $\nu = +1$ gaps crossing the neck between them. Surprisingly, this crossing is not reflected in the plateau shape. Probably there is some anticrossing within the neck that might smooth the expected feature. The $i = +2$ feature is rather blurred here probably because the $\nu = +2$ triangle gap is too narrow. E_F goes out into the $\nu = 0$ gap at high fields for $V_g = +3$ V in agreement with the observed turn of $\rho_{xy}(B)$ to negative values. The oscillatory behavior of the valence subband edge is clearly manifested here as E_F passes the positive $\nu = 1$ gap close to its round top between $B = 13$ and 21 T resulting in a tendency for $\rho_{xy}(B)$ to reach the $i = +1$ plateau and its further turn back to negative values at the weaker fields.

IV. CONCLUSIONS

In summary, we have found that the inverted energy spectrum of the quasi-2D HgTe is substantially modified in the DQW structure built of these layers even in the case of no interlayer tunneling. The overlap of the conduction and valence

subbands may be considerably increased in the DQW. Due to this fact, the critical field B_c for opening the gap in the energy spectrum shifts to much higher fields with respect to B_c in a single QW. As a result, the accompanying specific features in magnetotransport, such as multiple inversions in $\rho_{xy}(B)$, zero-filling-factor state with a concomitant manifestation of its insulator character in $\rho_{xx}(B)$, etc., are moved to the range of higher fields where the QH regime is well realized. This makes the quantum features more pronounced, with appeared superposed magnetoresistance oscillations, offering the possibility to analyze the observed peculiarities on the basis of calculated magnetic level patterns. We believe that the observed specificities are caused by the nonmonotonic course of the valence band magnetic levels resulted in triangular-shaped cyclotron gaps and an oscillating character of the valence subband top as a function of field. The effects of interlayer screening of the gate potential, which is never exhausted in the HgTe DQW with inverted spectrum, makes it possible to shift the energy spectra of the layers with respect to each other, thus making possible to regulate the overlap between the conduction and valence subbands by the gate voltage V_g . A huge difference in the density of states is found between the valence subband lateral maxima and the conduction subband minimum leading to that only the hole density in the upper layer is sensitive to V_g while the electron density remains almost constant.

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