



Anisotropic H_{c2} , thermodynamic and transport measurements, and pressure dependence of T_c in $\text{K}_2\text{Cr}_3\text{As}_3$ single crystals

Tai Kong, Sergey L. Bud'ko, and Paul C. Canfield

Ames Laboratory, U.S. DOE, and Department of Physics and Astronomy, Iowa State University, Ames, Iowa 50011, USA

(Received 7 January 2015; revised manuscript received 20 January 2015; published 30 January 2015)

We present a detailed study of single crystalline $\text{K}_2\text{Cr}_3\text{As}_3$ and analyze its thermodynamic and transport properties, anisotropic $H_{c2}(T)$, and initial pressure dependence of T_c . In zero field, the temperature-dependent resistivity is metallic. Deviation from a linear temperature dependence is evident below 100 K and a T^3 dependence is roughly followed from just above T_c (~ 10 K) to ~ 40 K. Anisotropic $H_{c2}(T)$ data were measured up to 140 kOe with field applied along and perpendicular to the rodlike crystals. For the applied field perpendicular to the rod, $H_{c2}(T)$ is linear with a slope ~ -70 kOe/K. For field applied along the rod, the slope is about -120 kOe/K below 70 kOe. Above 70 kOe, the magnitude of the slope decreases to ~ -70 kOe/K. The electronic specific heat coefficient γ , just above T_c , is 73 mJ/mol K^2 ; the Debye temperature Θ_D is 220 K. The specific heat jump at the superconducting transition $\Delta C \sim 2.2\gamma T_c$. Finally, for hydrostatic pressures up to ~ 7 kbar, T_c decreases under pressure linearly at a rate of -0.034 K/kbar.

DOI: [10.1103/PhysRevB.91.020507](https://doi.org/10.1103/PhysRevB.91.020507)

PACS number(s): 74.70.Xa, 74.25.Op, 74.25.Bt, 74.62.Fj

$\text{K}_2\text{Cr}_3\text{As}_3$ has been recently discovered as a new superconducting material with a T_c of 6.1 K [1]. Having a structure that contains $(\text{Cr}_3\text{As}_3)^{2-}$ chains, it has quickly aroused interest as a potentially new quasi-one-dimensional (Q1D) superconductor. Related compounds, $\text{Rb}_2\text{Cr}_3\text{As}_3$ and $\text{Cs}_2\text{Cr}_3\text{As}_3$, with T_c values of about 4.8 and 2.2 K, respectively, were reported soon after the discovery of $\text{K}_2\text{Cr}_3\text{As}_3$ [2,3]. Band structure calculations were also conducted and compared with experimental results [4]. In the initial report [1], data were acquired primarily on polycrystalline samples of $\text{K}_2\text{Cr}_3\text{As}_3$. For Q1D superconductors, it is believed that with a magnetic field applied along the chain, the upper critical field H_{c2} should be much higher than that when the field is along other directions due to minimized orbital pair breaking. Experimentally, this was observed for Q1D systems such as $\text{Li}_{0.9}\text{Mo}_6\text{O}_{17}$ and $(\text{TMTSF})\text{ClO}_4$ [5,6]. In this Rapid Communication, we present detailed anisotropic $H_{c2}(T)$ data via resistance measurements on $\text{K}_2\text{Cr}_3\text{As}_3$ single crystals. In addition, with the recent discovery of the isostructural compound $\text{Rb}_2\text{Cr}_3\text{As}_3$, with a larger unit cell and a lower T_c value, it is interesting to study the pressure dependence of T_c of $\text{K}_2\text{Cr}_3\text{As}_3$. Results from low-field dc magnetization measurements with pressures up to 7 kbar will be presented.

$\text{K}_2\text{Cr}_3\text{As}_3$ single crystals were grown using a high-temperature solution growth method [7]. A schematic drawing of the Matryoshka-like ampoule assembly for crystal synthesis is shown in Fig. 1(a). Elemental K, Cr, and As [in bulk(lump) form] were packed in an alumina crucible following the ratio listed in Ref. [1] (K:Cr:As = 6:1:7). The crucibles and starting material were then welded into a tantalum tube and sealed in a silica ampoule under a partial argon atmosphere. To avoid a possible explosion due to the high vapor pressure of the unreacted elements, the whole ampoule was slowly heated up to 1000 °C over 2 days. It was then cooled down over ~ 100 h to 700 °C, at which temperature the single crystals and the remaining liquid were separated in a centrifuge by the alumina strainer that was placed in between the growth side crucible and an empty (catch) crucible. The crystals obtained are rodlike [Fig. 1(b)] and malleable. A rough powder x-ray diffraction

pattern was collected on ground/deformed crystals using a Rigaku Miniflex unit located in a N_2 glove box. The data, along with a fit to the crystallographic data given in Ref. [1], are shown in Fig. 1(d) and are consistent with the single crystals adopting the hexagonal, $a = 9.983$ Å, $c = 4.230$ Å unit cell. Given the unit cell dimensions and single crystal morphology, we identify the rod direction as being along the crystallographic c axis.

Resistance was measured using a standard four-probe technique. DuPont 4929N silver paint was used to attach platinum wires onto the sample in a N_2 glove box. Electric current was applied along the rod for anisotropic H_{c2} measurements [Fig. 1(c)]. Long, straight samples were used and supported by flat plastic pads for all resistivity measurements to avoid potential torque that could deform the crystals and thus change their alignment. The resistivity was estimated assuming the sample has a cylindrical shape. The absolute value of the resistivity is therefore only accurate to within a factor of 3. The temperature- and field-dependent resistance was measured in a Quantum Design (QD) physical property measurement system, PPMS-14 ($T = 1.8$ –305 K, $H = 0$ –140 kOe). Specific heat data were measured using a QD PPMS via relaxation method. A ^3He option was utilized to obtain specific heat data down to 0.4 K. Specific heat data were measured on a “raft” assembled out of several single crystals. Despite the fact that the samples’ shape and arrangement were not ideal for such measurements, the coupling constant was high, between 96% and 100%.

The pressure dependence of the superconducting transition temperature T_c was determined via low-field (20 Oe), zero-field-cooled (ZFC), dc magnetization measurements in a QD magnetic property measurement system (MPMS) using a commercial (HMD) Be-Cu piston cylinder pressure cell [8]. Daphne oil 7373 was used as the pressure medium and Pb was used as the manometer [9]. Since the samples are very air sensitive, exposure to air was limited as much as possible to avoid oxidation.

The temperature dependence of the resistivity of $\text{K}_2\text{Cr}_3\text{As}_3$ is metallic, as shown in Fig. 2. The residual resistance ratio

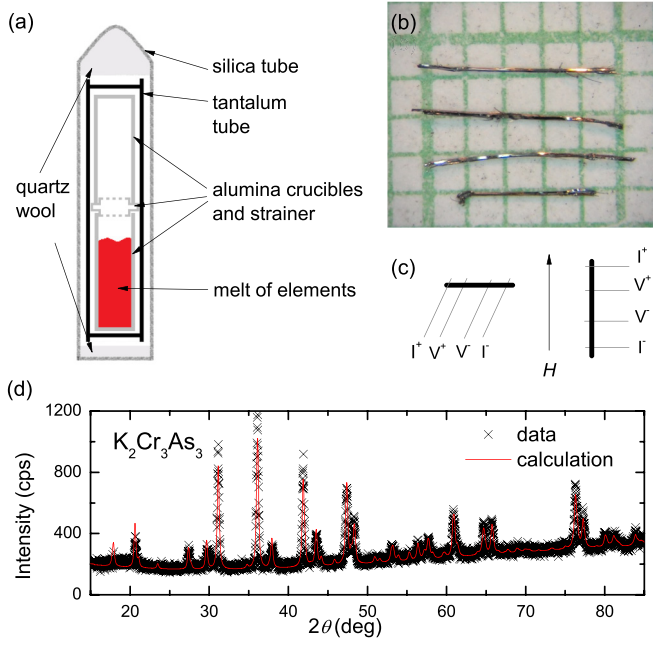


FIG. 1. (Color online) (a) A schematic drawing of the ampoule assembly for the material synthesis. (b) Typical habit of the single crystals of $\text{K}_2\text{Cr}_3\text{As}_3$ on a millimeter grid paper. (c) The contact arrangements and external field directions for the resistance measurements. (d) X-ray diffraction data on ground/deformed crystals. The black crosses represent the experimental data and the red solid line represents the calculation based on the crystallographic information. For details of the crystal structure, readers are directed to Ref. [1].

(RRR) is about 50, a factor of 5 better than the $\text{RRR} \sim 10$ found for the polycrystalline data [1]. A clear, sharp transition at 6.1 K indicates the superconducting transition. At room

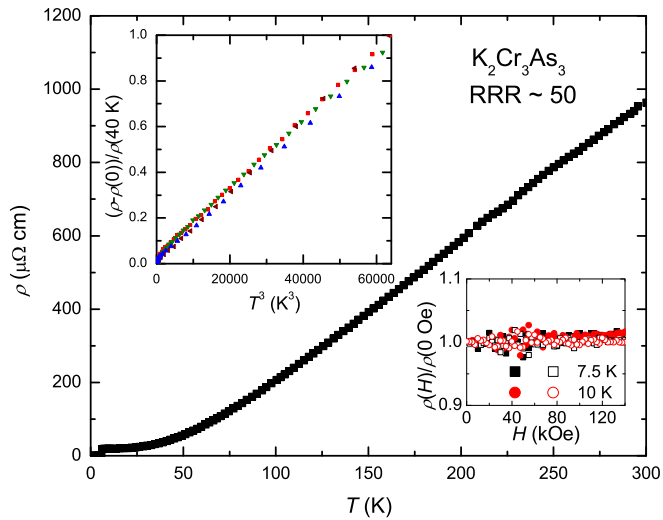


FIG. 2. (Color online) Temperature-dependent resistivity of $\text{K}_2\text{Cr}_3\text{As}_3$. The left inset shows the normalized resistances vs T^3 below 40 K for four different samples. The right inset shows the magnetoresistance measured at 7.5 and 10 K. Solid and open symbols represent transverse and longitudinal magnetoresistance, respectively.

temperature, the resistivity is of the same order of magnitude as the reported value [1]. However, in contrast to the reported linear temperature dependence from 7 to 300 K in Ref. [1], we see a clear deviation from linearity below ~ 100 K. For the four samples we have measured, we get slopes of 3.0, 2.8, 3.1, and 3.1 from a $\log(\rho - \rho_0)$ vs $\log T$ plot between 10 and 40 K. Therefore, in the left inset of Fig. 2, we plot the resistance data from several samples normalized to their value at 40 K as a function of T^3 . It is clear that a T^3 power law describes these data well over the temperature region that is shown. The slight deviation below 10 K might be due to the proximity to the superconducting transition. Although we did not see a temperature region where Fermi-liquid behavior is dominant, this could be due to the fact that the measurement temperature is not low enough as compared with its Debye temperature. As will be discussed below, due to the extremely large H_{c2} values associated with this compound, measuring normal state resistivity to low temperatures will require extremely high magnetic fields.

In the normal state, despite a relatively large RRR value, virtually no magnetoresistance was observed. The right inset of Fig. 2 shows the normalized resistance as a function of applied field for both transverse ($H \perp I$) and longitudinal ($H \parallel I$) directions of the applied field. Right above the T_c , the resistance stays close to constant up to 140 kOe. Given this null response, an upper limit of 2% (at 140 kOe) can be set on the low-temperature magnetoresistance.

Figures 3(a) and 3(b) show the temperature-dependent resistance $R(T)$ measured under different magnetic fields for $H \parallel \text{rod}$ and $H \perp \text{rod}$, respectively. Figure 3(c) shows the anisotropic $H_{c2}(T)$ data for $\text{K}_2\text{Cr}_3\text{As}_3$ from several samples. The superconducting transition is very sharp in our resistivity measurements. The inset to Fig. 3(c) shows that, at over 125 kOe, the onsets and offsets of the transition differ by less than 0.25 K as opposed to an over 0.6 K width at 80 kOe for polycrystalline data [1]. An offset criterion [10] was used to determine T_c [see the inset of Fig. 3(c)]. For field applied perpendicular to the rod, the T_c decreases almost linearly with increasing magnetic field. The slope is ~ -70 kOe/K. For field applied along the rod, the initial slope is roughly -120 kOe/K, and above 70 kOe the slope is close to -70 kOe/K, which is similar to that when field is applied perpendicular to the rod. Different batches of samples as well as field-sweep data at fixed temperature give consistent results. In Fig. 3(c), data from Ref. [1] are plotted as blue solid diamonds. It is close to what we obtained for field perpendicular to the rod and the slope is consistent with our current results.

From one-band BCS theory, for an s -wave, isotropic material, in the clean limit, the orbital $H_{c2}^{\text{orb}} = -0.73|dH_{c2}/dT|_{T_c} T_c$ [11]. In our case, for field perpendicular to the rod, $H_{c2}^{\text{orb}} \sim 312$ kOe. The Pauli limit for a simple BCS superconductor is $H^p = 1.84T_c = 110$ kOe. The H_{c2} of real systems will be influenced by both H_{c2}^{orb} and H^p . The Maki parameter [12] that describes the relative importance of the two critical field limits, $\alpha = \sqrt{2}H_{c2}^{\text{orb}}/H^p$, is equal to 4. This suggests that at low temperature, the H_{c2} might be Pauli limited. With the current data up to 140 kOe that exceed the one-band BCS estimated Pauli limit, we see no clear sign of saturation of H_{c2} . Note

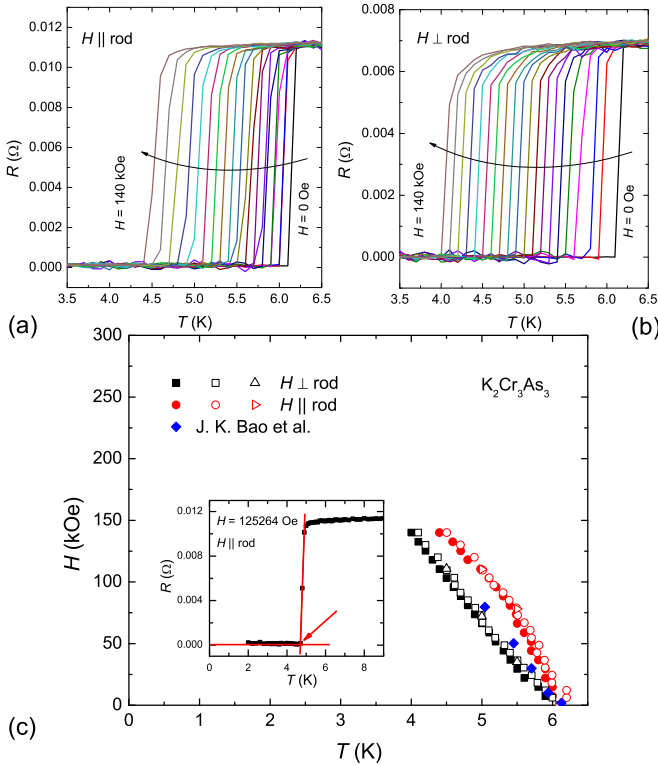


FIG. 3. (Color online) $R(T)$ data measured for (a) $H \parallel \text{rod}$ and (b) $H \perp \text{rod}$ from 0 Oe up to 140 kOe in 20 uniform field steps ($\Delta H = 7368.4$ Oe). (c) Anisotropic H_{c2} of $\text{K}_2\text{Cr}_3\text{As}_3$. The black and red symbols represent data for $H \parallel \text{rod}$ and $H \perp \text{rod}$. The blue points are from Ref. [1]. The solid and open symbols show data from two different batches of samples. The triangles are data from field sweeps at fixed temperatures. The squares and circles are data from temperature sweeps at fixed field. The inset shows a typical superconducting transition measured at 125 264 Oe, with red solid lines and the arrow indicating the criteria for determining the transition temperature (see more in text).

that, with multiband as well as varying coupling strength, the Pauli limit can be different from a one-band BCS prediction. For example, in many Fe-based superconductors, H^p values are significantly enhanced [13].

The specific heat data for $\text{K}_2\text{Cr}_3\text{As}_3$ are shown in Fig. 4. A clear jump in the specific heat at around 6 K corresponds to the superconducting transition. In the normal state, $C = \gamma T + \beta T^3$ fits the data quite well. In the range of 7–10 K, we obtained a $\gamma \sim 73$ mJ/mol K² and the Debye temperature Θ_D , derived from β , is around 220 K. These data are consistent with recently reported values [1]. At the superconducting transition, the specific heat jump is roughly $2.2\gamma T_c$. This is larger than the simple s -wave BCS prediction and comparable to the value obtained in Ref. [1]. Possibly, strong coupling is involved [14]. Below 1.5 K, we observe a clear upturn in C/T . Assuming that the normal state values of γ and β stay constant below T_c , this upturn exacerbates the difficulty of having entropy conserved, which was already hinted at in the previously reported data down to 2 K [1]. Additional, nonsuperconducting contributions, such as impurities or a nuclear Schottky anomaly, may be responsible for this low-

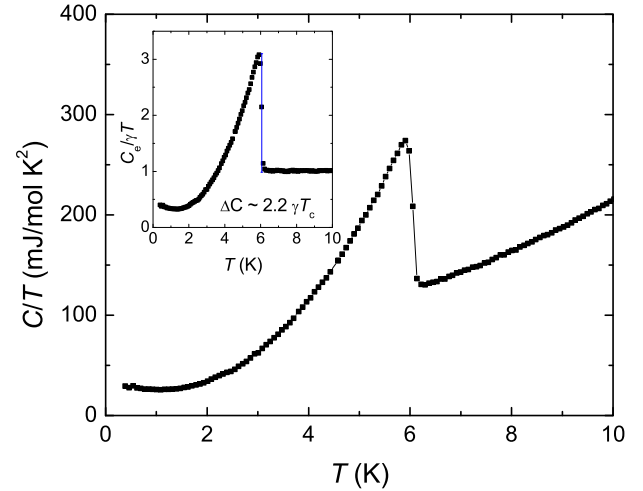


FIG. 4. (Color online) C/T plotted as a function of T down to 0.4 K. The inset shows the $C_e/\gamma T$ as a function of T . The electronic part of the specific heat C_e was obtained by subtracting βT^3 from the total specific heat.

temperature rise. But whatever its origin, it makes fitting of the data to specific models problematic.

At ambient pressure, the normal state dc magnetic susceptibility ($\chi = M/H$) is $\sim 0.9 \times 10^{-3}$ emu/mol and decreases by more than 30% from 7 to 300 K (Fig. 5). If we fit these data to a Curie-Weiss law, we find a temperature-independent susceptibility $\chi_0 \sim 0.7 \times 10^{-3}$ emu/mol, a small effective moment of $0.36\mu_B/\text{f.u.}$ (or $0.21\mu_B/\text{Cr}$), and a Curie-Weiss temperature $\Theta \sim -40$ K. At this point in time, it is not clear if such a local-moment fit is appropriate or intrinsic for this material. Given that $\gamma = 9$ mJ/mol atom K² is relatively high, it is appropriate to contrast it with the magnetic susceptibility. Assuming spin = 1/2, the estimated Wilson ratio [15] R_W is in a range of 0.9–1.1 by taking the low-temperature magnetic susceptibility value. If we consider the Curie tail as a result of an impurity contribution and thus take the value of temperature-independent χ_0 , R_W is about 0.7. These values

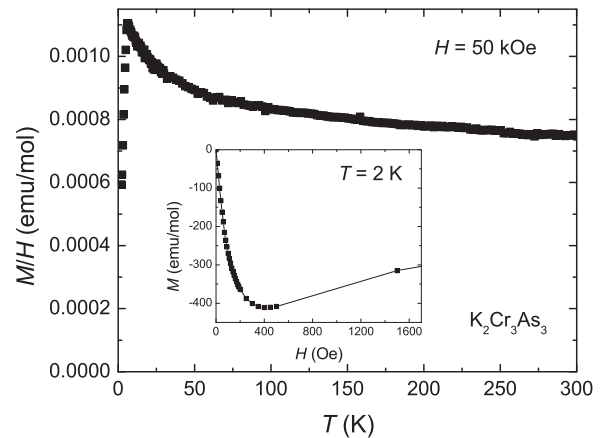


FIG. 5. The magnetic susceptibility measured at 50 kOe on a randomly oriented collection of single crystals. The inset shows the magnetization isotherm measured at 2 K up to 1500 Oe.

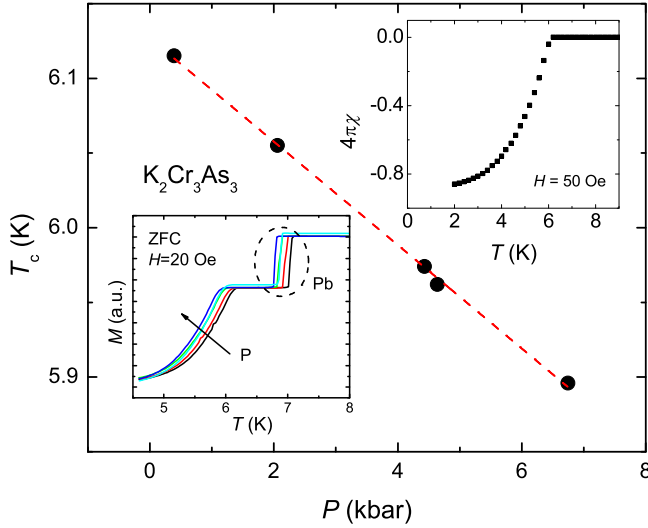


FIG. 6. (Color online) Pressure dependence of T_c . The dotted line is a guide for the eye. The left inset shows the raw data with Pb serving as a manometer. The right inset shows the zero-field-cooled magnetization data under ambient pressure measured at 50 Oe.

are close to what one would expect for a Fermi-liquid system ($R_W = 1$).

Zero-field-cooled dc magnetization data are shown in the right inset of Fig. 6. The superconducting transition appears sharply at 6.1 K, consistent with both resistance and specific heat measurements. To avoid oxidation, which could affect both sample mass and superconducting volume fraction, the sample mass was measured in a glove box and exposure to air was minimized.

At 2 K, the low-field $M(H)$ data deviate from the linear field dependence at around 70 Oe (see the inset of Fig. 5). The minimum in magnetization appears at around 400 Oe. Taking $H_{c1} = 70$ Oe and $H_{c2} \sim H_{c2}^{\text{orb}} = 312$ kOe and using $H_{c1}/H_{c2} = \ln(\kappa)/2\kappa^2$, the Ginzburg-Landau (GL) parameter κ is ~ 100 . From the specific heat jump at T_c using the Rutgers relation $\Delta C/T_c = (1/8\pi\kappa^2)(dH_{c2}/dT)|_{T_c}$ [16], we obtain a similar κ value of 116. This high value of κ suggests that $\text{K}_2\text{Cr}_3\text{As}_3$ is deep in the type II regime.

The pressure dependence of T_c is plotted in Fig. 6. The superconducting transition of Pb was used to determine the pressure within the pressure cell [9]. As shown in the left inset of Fig. 6, both superconducting transitions from the Pb and the sample are clear. Up to 7 kbar, the T_c value for $\text{K}_2\text{Cr}_3\text{As}_3$ decreases linearly with increasing pressure at a rate of -0.034 K/kbar. Given that the larger-unit-celled $\text{Rb}_2\text{Cr}_3\text{As}_3$ undergoes a superconducting transition at a slightly lower temperature, 4.8 K, the effects of chemical and physical pressure are clearly not equivalent in this material.

In conclusion, we grew single crystals of $\text{K}_2\text{Cr}_3\text{As}_3$ using a high-temperature solution growth method. Being able to

decant the molten solution at high temperature allows us to get relatively large, free-standing, high-quality single crystals, which also enables us to perform all of our measurements on single crystals or their arrays. In contrast to previously reported linear-temperature-dependent resistivity from 300 down to 7 K, the resistivity of a single crystal with a RRR of 50 (RRR ~ 10 for polycrystalline samples [1]) deviates from the linear temperature dependence below 100 K and roughly follows T^3 over the range of 10–40 K. Whereas the differences between the single crystalline and polycrystalline sample transports are easily explained by grain-boundary scattering and potential resistive anisotropies, the rather clear T^3 power law that we observe will require more experimental and theoretical investigations. Despite the fact that the temperature dependence is different, if the RRR values reflect the intrinsic quality of the samples, and are not due to grain-boundary scattering, the independence of T_c on the value of RRR, or impurity scattering, suggests a conventional mechanism of its superconductivity. Anisotropic $H_{c2}(T)$ data up to 140 kOe were obtained from resistivity measurements. Up to 140 kOe, for field applied perpendicular to the rod, $H_{c2}(T)$ is linear with a slope of -70 kOe/K. For field along the rod, $dH_{c2}/dT|_{T_c}$ is about -120 kOe/K. Above 70 kOe, the slope decreases to around -70 kOe/K. The T_c as well as the slope of H_{c2} is close to previously reported data [1]. The anisotropy in H_{c2} is not as large as one would generally expect for a Q1D superconductor. Virtually zero magnetoresistance was observed at temperatures right above the T_c . The electronic specific heat γ is 73 mJ/mol K² or 9 mJ/mol atom K². Although this is a relatively large value, the equally enhanced magnetic susceptibility leads to a Wilson ratio that is close to 1 and suggests Fermi-liquid-like properties. At ambient pressure, H_{c1} is close to 70 Oe, which results in a GL parameter that is ~ 100 , taking 312 kOe as the estimated H_{c2} . For pressures up to 7 kbar, the superconducting transition temperature T_c decreases linearly with a rate of -0.034 K/kbar. Comparing to the newly discovered $\text{Rb}_2\text{Cr}_3\text{As}_3$, it seems to suggest that physical pressure and chemical pressure have different effects on the superconductivity in these compounds.

Despite the fact that $\text{K}_2\text{Cr}_3\text{As}_3$ can be considered to be close to Q1D in a crystallographic sense, the “small” anisotropy in H_{c2} appears to suggest a rather three-dimensional nature. In addition, the exotic pairing symmetry requires more detailed study.

We would like to thank U. S. Kaluarachchi, A. E. Böhmer, D. K. Finnemore, V. G. Kogan, and V. Taufour for useful discussions. This work was supported by the US Department of Energy (DOE), Office of Science, Basic Energy Sciences, Materials Science and Engineering Division. The research was performed at the Ames Laboratory, which is operated for the US DOE by Iowa State University under Contract NO. DE-AC02-07CH11358.

[1] J.-K. Bao, J.-Y. Liu, C.-W. Ma, Z.-H. Meng, Z.-T. Tang, Y.-L. Sun, H.-F. Zhai, H. Jiang, H. Bai, C.-M. Feng, Z.-A. Xu, and G.-H. Cao, [arXiv:1412.0067](https://arxiv.org/abs/1412.0067).

[2] Z.-T. Tang, J.-K. Bao, Y. Liu, Y.-L. Sun, A. Ablimit, H.-F. Zhai, H. Jiang, C.-M. Feng, Z.-A. Xu, and G.-H. Cao, *Phys. Rev. B* **91**, 020506 (2015).

- [3] Z.-T. Tang, J.-K. Bao, Z. Wang, H. Bai, H. Jiang, Y. Liu, H.-F. Zhai, C.-M. Feng, Z.-A. Xu, and G.-H. Cao, *Science China Materials* **58**, 16 (2015).
- [4] H. Jiang, G. Cao, and C. Cao, [arXiv:1412.1309](https://arxiv.org/abs/1412.1309).
- [5] J.-F. Mercure, A. F. Bangura, X. Xu, N. Wakeham, A. Carrington, P. Walmsley, M. Greenblatt, and N. E. Hussey, *Phys. Rev. Lett.* **108**, 187003 (2012).
- [6] S. Yonezawa, S. Kusaba, Y. Maeno, P. Auban-Senzier, C. Pasquier, K. Bechgaard, and D. Jérôme, *Phys. Rev. Lett.* **100**, 117002 (2008).
- [7] P. C. Canfield and Z. Fisk, *Philos. Mag. B* **65**, 1117 (1992).
- [8] www.qdusa.com/products/high-pressure-cell-mpms.html.
- [9] A. Eiling and J. S. Schilling, *J. Phys. F: Met. Phys.* **11**, 623 (1981).
- [10] N. Ni, M. E. Tillman, J.-Q. Yan, A. Kracher, S. T. Hannahs, S. L. Bud'ko, and P. C. Canfield, *Phys. Rev. B* **78**, 214515 (2008).
- [11] M. Tinkham, *Introduction to Superconductivity* (Courier Dover, New York, 2012).
- [12] K. Maki, *Phys. Rev.* **148**, 362 (1966).
- [13] S. Khim, B. Lee, J. W. Kim, E. S. Choi, G. R. Stewart, and K. H. Kim, *Phys. Rev. B* **84**, 104502 (2011).
- [14] J. P. Carbotte, *Rev. Mod. Phys.* **62**, 1027 (1990).
- [15] G. R. Stewart, *Rev. Mod. Phys.* **56**, 755 (1984).
- [16] U. Welp, W. K. Kwok, G. W. Crabtree, K. G. Vandervoort, and J. Z. Liu, *Phys. Rev. Lett.* **62**, 1908 (1989).