Pseudogap in underdoped $Ba_{1-x}K_xFe_2As_2$ as seen via optical conductivity

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We report the observation of a pseudogap in the *ab*-plane optical conductivity of underdoped $Ba_{1-x}K_xFe_2As_2$ ($x = 0.2$ and 0.12) single crystals. Both samples show prominent gaps opened by a spin density wave (SDW) order and superconductivity at the transition temperatures T_{SDW} and T_c, respectively. In addition, we observe an evident pseudogap below $T^* \sim 75$ K, a temperature much lower than T_{SDW} but much higher than T_c. A spectral weight analysis shows that SDW competes with superconductivity while the pseudogap is closely connected to the superconducting gap, indicating the possibility of its being a precursor of superconductivity. The doping dependence of the gaps is also supportive of such a scenario.

Among all the families of iron-pnictide superconductors discovered to date,[1–7] the $BaFe_2As_2$ (Ba122) family is one of the most studied. The parent $BaFe_2As_2$ composition is a poor Pauli-paramagnetic metal with a structural and magnetic phase transition at 140 K.[8] Superconductivity arises with the suppression of magnetism that can be achieved by applying pressure[2] or chemical substitution.[3–5] The substitution of Ba with K atoms yields hole doping[3] with a maximum $T_c \approx$ 39 K and the substitution of Fe atoms by Co or Ni atoms results in electron doping[4,5] with a maximum $T_c \approx 25$ K. Extensive studies have been carried out in the parent $BaFe_2As_2$,[9,10] electron-doped $BaFe_{2-x}A_xAs_2$ ($A =$ Co, Ni),[11–14] as well as optimally hole-doped $Ba_{0.6}K_{0.4}Fe_2As_2$[15–19] compounds. However, the hole-underdoped regime of the phase diagram is relatively unexplored.

This hole-underdoped region is of the utmost importance. First, the superconducting mechanism is deeply tied with magnetism. The interplay between magnetism and superconductivity is manifest in this regime, where SDW and superconductivity coexist over a large doping range.[20–24] Second, in cuprates, the most exciting, yet puzzling, physics takes place in the hole-underdoped regime. This regime is thus pivotal to the comparison between iron pnictides and cuprates.

Xu *et al.*'s angle-resolved photoemission (ARPES) measurements on underdoped $Ba_{1-x}K_xFe_2As_2$ (Ref. 25) showed a distinct pseudogap coexisting with the superconducting gap and suggested that both the pseudogap and superconductivity are driven by antiferromagnetic fluctuations. However, one key issue in understanding the origin of the pseudogap and, in particular, its relation to superconductivity, is the question of whether it shares electronic states with the superconducting condensate.[26] Infrared spectroscopy probes the charge dynamics of bulk materials and the spectral weight analysis is a powerful tool to address this issue.

We measured the broadband infrared response of two underdoped $Ba_{1-x}K_xFe_2As_2$ ($x = 0.2$ and 0.12) single crystals. In both samples, the opening of the SDW gap and the superconducting gap was clearly observed on the optical conductivity. In addition, another small gap opens below

$T^* \sim 75$ K, closely resembling the pseudogap in the hole-underdoped cuprates. We find that the SDW gap depletes the spectral weight available for the superconducting condensate, which suggests that the SDW order competes with superconductivity. On the other hand, the pseudogap shares electronic states with the superconducting gap. This is shown by the doping and temperature dependence of the optical conductivity spectral weight. This observation supports a scenario where the pseudogap is a precursor of superconductivity.

High quality $Ba_{1-x}K_xFe_2As_2$ single crystals were grown by the self-flux method using FeAs as the flux.[27] The left panel of Fig. 1 shows the temperature dependence of the dc resistivity [$\rho(T)$] for the $Ba_{1-x}K_xFe_2As_2$ ($x = 0.2$) sample. The $\rho(T)$ curve is characterized by a steep superconducting transition at $T_c = 19$ K. The inset shows the derivative of the resistivity $d\rho/dT$ as a function of temperature. The SDW transition manifests itself as a sharp peak in $d\rho/dT$ at $T_{SDW} = 104$ K, which corresponds to a small kink on the $\rho(T)$ curve. The right panel displays the same curves for the $x = 0.12$ sample, which has $T_c = 11$ K and $T_{SDW} = 121$ K.

The *ab*-plane reflectivity [$R(\omega)$] was measured at a near-normal angle of incidence on Bruker IFS113v and IFS66v spectrometers. An *in situ* gold overfilling technique[28] was used to obtain the absolute reflectivity of the samples. Data from 20 to 12 000 cm^{-1} were collected at 18 different temperatures from 5 to 300 K on freshly cleaved surfaces. In order to use a Kramers-Kronig analysis, we extended the data to the visible and UV range (10 000–55 000 cm^{-1}) at room temperature with an AvaSpec-2048 $\times$14 model fiber optic spectrometer.

Figure 2 shows the infrared reflectivity at selected temperatures for both samples up to 1200 cm^{-1}. The inset in each panel displays the reflectivity for the full measured range at 300 K. For the $x = 0.2$ sample, shown in the top panel, the reflectivity exhibits a metallic response and approaches unity at zero frequency. Below $T_{SDW} = 104$ K, a substantial suppression of $R(\omega)$ at about 650 cm^{-1} sets in and intensifies with decreasing temperature. Simultaneously, the low frequency reflectivity continues increasing towards unity. This is a signature of
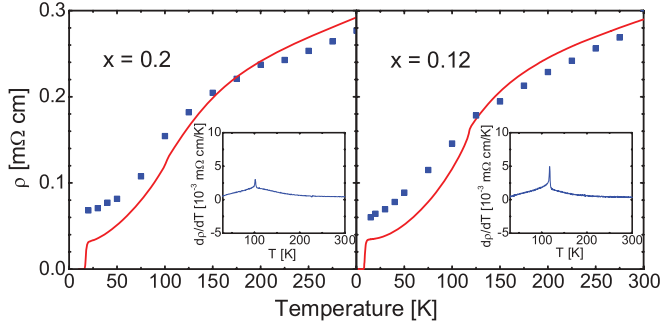
FIG. 1. (Color online) Left: Temperature dependence of the resistivity of the $Ba_{1-x}K_xFe_2As_2$ ($x = 0.2$) single crystal (red solid line). A steep superconducting transition can be seen at $T_c = 19$ K. The blue solid squares are values from the zero frequency extrapolation of the optical conductivity. The inset shows the derivative of the resistivity $d\rho/dT$ as a function of temperature. The sharp peak at 104 K in $d\rho/dT$ is associated with the SDW transition. Right: Same curves for the $x = 0.12$ sample with $T_c = 11$ K and $T_{SDW} = 121$ K.

a partial SDW gap on the Fermi surface. Below 75 K, defined as T^* here, another suppression of $R(\omega)$ appears in a lower-energy scale (~ 150 cm^{-1}), signaling the opening of a second partial gap (pseudogap) with a smaller value. Upon crossing the superconducting transition, the reflectivity below ~ 150 cm^{-1} increases, indicating the opening of a superconducting gap. Similar features are observed on $R(\omega)$

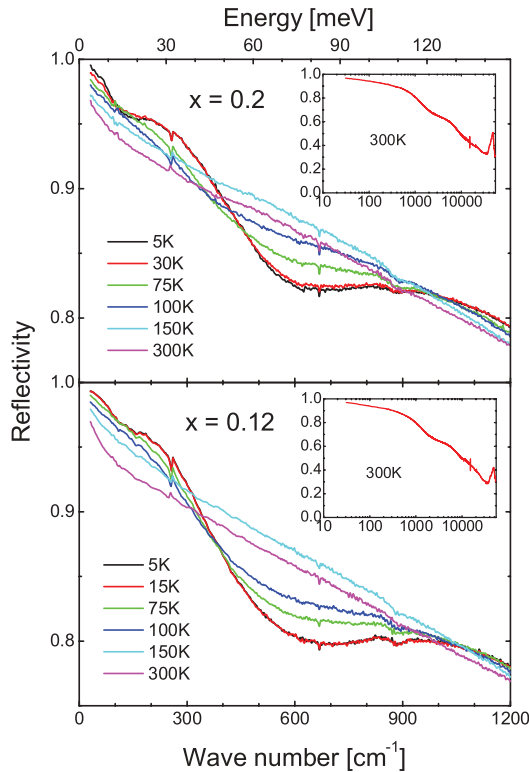FIG. 2. (Color online) Reflectivity of $Ba_{1-x}K_xFe_2As_2$ single crystals below 1200cm^{-1} at various temperatures. Top panel: $x = 0.2$; Bottom panel: $x = 0.12$. Inset: Reflectivity of full measured range at 300 K.
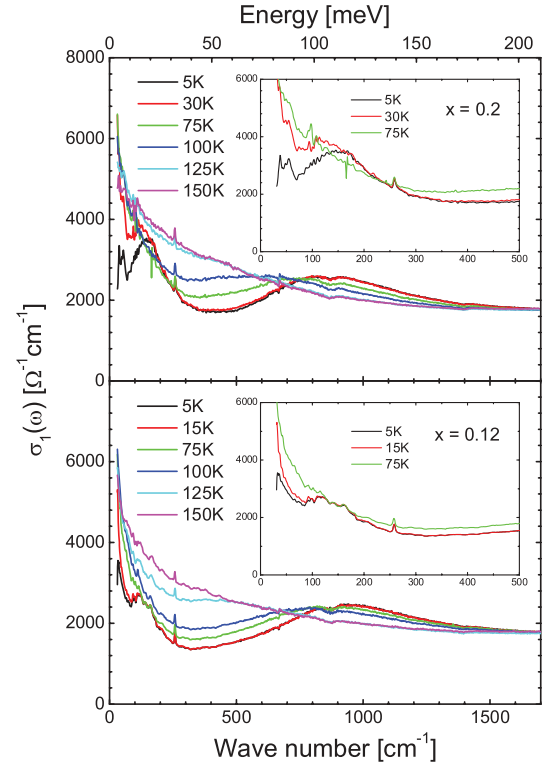

FIG. 3. (Color online) Top panel: Optical conductivity of $Ba_{1-x}K_xFe_2As_2$ ($x = 0.2$) at selected temperatures. The inset shows the enlarged view of the optical conductivity at low frequencies. The bottom panel displays the same spectra for the $x = 0.12$ sample.

for the $x = 0.12$ sample, as shown in the bottom panel of Fig. 2.

The real part of the optical conductivity $\sigma_1(\omega)$ was determined by Kramers-Kronig analysis of the measured reflectivity. Figure 3 shows $\sigma_1(\omega)$ at different temperatures for the two samples. The zero frequency extrapolations of $\sigma_1(\omega)$ represent the inverse dc resistivity of the sample, shown as blue solid squares in Fig. 1, which are in good agreement with the transport measurement. The top panel of Fig. 3 shows $\sigma_1(\omega)$ for $Ba_{1-x}K_xFe_2As_2$ ($x = 0.2$) below 1700 cm^{-1}. At 150 and 125 K, hence above T_{SDW}, a Drude-like metallic response dominates the low frequency optical conductivity. Below T_{SDW}, $\sigma_1(\omega)$ below about 650 cm^{-1} is severely suppressed. Meanwhile, it increases in a higher-energy scale from 650 to 1700 cm^{-1}. The optical conductivity for the normal state and that for the SDW state just below T_{SDW} show an intersection point at about 650 cm^{-1}. As the temperature decreases, both the low-energy spectral suppression and the high-energy bulge become stronger, resulting in a prominent peak at about 840 cm^{-1}. This spectral evolution manifests the behavior of the SDW gap in this material: transfer of low frequency spectral weight to high frequencies. If we take the peak position as an estimate of the gap value, we can see that the gap increases with decreasing temperature. The effect of the SDW transition on optical conductivity has been well established in the parent $BaFe_2As_2$,[9,29,30] which is characterized by two distinct SDW gaps, associated with different Fermi surfaces, at 360 and 890 cm^{-1}, respectively.

Below $T^* \sim 75$ K, a second suppression in the optical conductivity sets in below roughly 110 cm^{-1} with a bulge extending from about 110 to 250 cm^{-1}, implying the opening of the pseudogap on the Fermi surface. The inset of Fig. 3 shows the enlarged view of the low temperature optical conductivity at low frequencies, where the pseudogap is seen more clearly. This pseudogap is unlikely due to the SDW transition as (i) it opens at a temperature between 50 and 75 K, well below T_{SDW}, and (ii) it redistributes spectral weight at a different, much lower-energy scale, which is about the energy scale of the superconducting gap in the optimally doped $Ba_{0.6}K_{0.4}Fe_2As_2$.[15–17]

The superconducting transition at $T_c = 19$ K implies the opening of a superconducting gap. As shown in the inset of Fig. 3, this leads to the reduction of the optical conductivity at low frequencies between 30 and 5 K. The spectral weight lost in the transition is recovered by the $\delta(\omega)$ function at zero frequency representing the infinite dc conductivity in the superconducting state. This $\delta(\omega)$ function is not visible in the $\sigma_1(\omega)$ spectra, because only finite frequencies are experimentally measured. Nevertheless, its weight can be calculated from the imaginary part of the optical conductivity.[31,32] Note that, the spectral depletion in $\sigma_1(\omega)$ due to the superconducting condensate extends up to 180 cm^{-1}. This is the same energy scale of the pseudogap, hinting that the superconducting gap and the pseudogap share the same electronic states, and may have the same origin.

It is also noteworthy that, in the superconducting state, large residual low frequency conductivity, representing unpaired quasiparticles, is present, even at the lowest measured temperature 5 K. This is in sharp contrast to the optimally doped $Ba_{0.6}K_{0.4}Fe_2As_2$, where the low frequency optical conductivity is dominated by a fully open gap.[15–17] Here, the residual low frequency conductivity observed in the underdoped $Ba_{1-x}K_xFe_2As_2$ might be due to the well-known phase separation[20,21,33] or nodes in the superconducting gap.[34,35]

In the $x = 0.12$ sample, very similar features are observed, as shown in the bottom panel of Fig. 3, but remarkable differences exist: (i) The SDW gap opens at a higher temperature ($T_{SDW} = 121$ K) and the gap value shifts to a higher-energy scale (~ 890 cm^{-1}); (ii) the low frequency spectral weight suppression due to the SDW gap is stronger, indicating that a larger part of the Fermi surface is removed in the SDW state; and (iii) in contrast to the SDW gap, both the pseudogap and the superconducting gap features are weaker in the more underdoped sample. The evolution of the three gaps (SDW, pseudogap, and superconducting) with doping also suggests that the pseudogap and the superconducting gap may have a common origin while the SDW is a competitive order to superconductivity.

In order to investigate the origin of the pseudogap and the relationship among all gaps, we analyzed the data utilizing a restricted spectral weight, defined as

$$W_{\omega_a}^{\omega_b} = \int_{\omega_a}^{\omega_b} \sigma_1(\omega)d\omega, \qquad (1)$$

where ω_a and ω_b are lower and upper cutoff frequencies, respectively. By choosing appropriate values for ω_a and ω_b, one can study the relations among different phase transitions.[26]

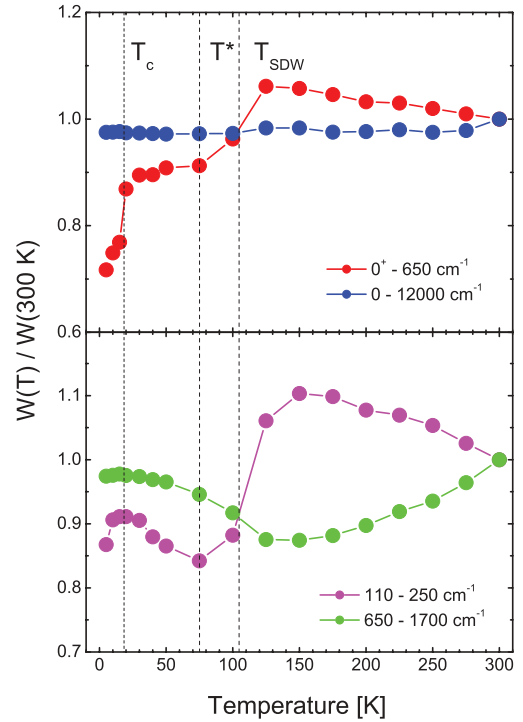
FIG. 4. (Color online) Temperature dependence of the spectral weight, $W_{\omega_a}^{\omega_b} = \int_{\omega_a}^{\omega_b} \sigma_1(\omega)d\omega$, between different lower and upper cutoff frequencies for the $x = 0.2$ sample. The vertical dashed lines denote T_c, T^* and T_{SDW}.

When replacing ω_a by 0 and ω_b by ∞, we fall back to the standard f-sum rule and the spectral weight is conserved.

Figure 4 shows the temperature dependence of the $x = 0.2$ sample spectral weight, normalized by its value at 300 K, at different cutoff frequencies. The vertical dashed lines denote T_c, T^*, and T_{SDW}. The blue solid circles in the top panel of Fig. 4 are the normalized spectral weight with cutoff frequencies $\omega_a = 0$ and $\omega_b = 12\,000$ cm^{-1} as a function of temperature. Here, the weight of the zero frequency δ function is included below T_c. Moreover, since the optical conductivity is measured only down to 20 cm^{-1}, we estimate the spectral weight below that energy by fitting the low frequency normal-state optical conductivity to a Drude model. The upper cutoff frequency ($\omega_b = 12\,000$ cm^{-1}) is high enough to cover the whole spectrum responsible for the phase transitions in this material. Hence the blue solid circles form a flat line at about unity, indicating that the spectral weight is conserved.

The red solid circles in the top panel show the temperature dependence of the normalized spectral weight with cutoff frequencies $\omega_a = 0^+$ and $\omega_b = 650$ cm^{-1}. Here 0^+ means that the superfluid weight is not included. Above T_{SDW}, the continuous increase of the normalized value of $W_{0^+}^{650}$ with decreasing T is related to the narrowing of the Drude band. This is the typical optical response of a metallic material. A strong spectral weight suppression occurs at T_{SDW}, which is the consequence of the SDW gap opening. At T_c, another sharp drop of the spectral weight breaks in, indicating the superconducting gap opening.

The temperature dependence of the normalized value of W_{650}^{1700}, shown as green solid circles in the bottom panel

of Fig. 4, provides clues about the relation between the superconducting and the SDW gaps. Above T_{SDW}, the material shows a metallic response which can be described by a Drude peak centered at zero frequency. With the temperature decrease, the dc conductivity increases and the scattering rate reduces. The continuous narrowing of the Drude band induces a transfer of spectral weight from the midinfrared to the far infrared, resulting in the continuous decrease of the spectral weight observed in the 650–1700 cm^{-1} range. Below T_{SDW}, the opposite behavior dominates the optical conductivity. The SDW gap depletes the spectral weight below 650 cm^{-1} and transfers it to the 650–1700 cm^{-1} range, leading to the continuous increase of W_{650}^{1700} with decreasing T. This behavior continues into the superconducting state and does not show any feature at T_c. These observations indicate that the SDW and superconducting gaps are separate and even act as competitive orders in this material.

If a partial gap is due to a precursor order of superconductivity, for example performed pairs without phase coherence, once the long range superconductivity is established, a significant part of the spectral weight transferred to high frequencies by the partial gap should be transferred back to low energies and join the superconducting condensate.[36,37] Conversely, a partial gap due to a competitive order to superconductivity depletes the low-energy spectral weight and holds it in a high-energy scale without transferring it back to the superfluid weight below T_c.[26] From the normalized value of W_{650}^{1700} vs T curve (green solid circles) we note that no loss of spectral weight is observed at T_c. This means that the spectral weight transferred to high frequencies by the SDW gap remains in the high frequency scale and does not contribute to the superconducting condensate. Therefore, the SDW acts as a competitive order to superconductivity in this material.

Along these lines, the origin of the pseudogap and its relationship to superconductivity can be revealed by a close inspection of the temperature dependence of the normalized value of W_{110}^{250}, shown as purple solid circles in the bottom panel. Above T^*, this curve shows the same feature as the normalized value of $W_{0^+}^{650}$ vs T curve, i.e., a continuous increase upon cooling down followed by a suppression at T_{SDW} due to the SDW gap opening. At T^*, the spectral weight in the 110–250 cm^{-1} range reaches a minimum and starts to increase with decreasing temperature. This is due to the opening of the pseudogap. The pseudogap, opening at T^*,

depletes the spectral weight below 110 cm^{-1} and retrieves it in the 110–250 cm^{-1} frequency range, leading to the increase of W_{110}^{250} below T^*. An interesting phenomenon happens to the pseudogap when the material undergoes the superconducting transition. In contrast to the case of the SDW gap, a significant loss of spectral weight in the 110–250 cm^{-1} frequency range is observed below T_c. This observation indicates that the spectral weight transferred to the 110–250 cm^{-1} range by the pseudogap joins the superconducting condensate when superconductivity is established. Hence, the pseudogap is likely a precursor order with respect to superconductivity.

A recent ultrafast pump-probe study of underdoped $Ba_{1-x}K_xFe_2As_2$ revealed a normal-state precursor order to superconductivity at $T^* \sim 60$ K.[24] The precursor scenario is also supported by ^{75}As nuclear-quadrupole-resonance (NQR) measurements on $Ca(Fe_{1-x}Co_x)_2As_2$ showing a pseudogap-like phase as a precursor state for the coherent superconducting phase[38] and Andreev reflection studies in $Ba(Fe_{1-x}Co_x)_2As_2$ which offered evidence for phase-incoherent superconducting pairs in the normal state.[39]

In summary, we measured the optical conductivity of two underdoped $Ba_{1-x}K_xFe_2As_2$ ($x = 0.2$ and 0.12) single crystals. In both samples, besides the SDW gap and superconducting gap, the optical conductivity reveals another small partial gap (pseudogap) opening below $T^* \sim 75$ K, an intermediate temperature between T_{SDW} and T_c. A spectral weight analysis shows that the SDW gap diminishes the low-energy spectral weight available for the superconducting condensate while the pseudogap shares the same electronic states with the superconducting gap. These observations, together with the doping dependence of these gaps, suggest the SDW as a competitive order and the pseudogap as a precursor to superconductivity.

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