

In-plane magnetic field dependence of cyclotron relaxation time in a Si two-dimensional electron system

Tasuku Chiba,¹ Ryuichi Masutomi,¹ Kentarou Sawano,² Yasuhiro Shiraki,² and Tohru Okamoto¹

¹*Department of Physics, University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-0033, Japan*

²*Research Center for Silicon Nano-Science, Advanced Research Laboratories, Tokyo City University, 8-15-1 Todoroki, Setagaya-ku, Tokyo 158-0082, Japan*

(Received 15 February 2012; published 13 July 2012)

Cyclotron resonance of two-dimensional electrons is studied for a high-mobility Si/SiGe quantum well in the presence of an in-plane magnetic field, which induces spin polarization. The relaxation time τ_{CR} shows a negative in-plane magnetic-field dependence, which is similar to that of the transport scattering time τ_t obtained from dc resistivity. The resonance magnetic field shows an unexpected negative shift with increasing in-plane magnetic field.

DOI: [10.1103/PhysRevB.86.045310](https://doi.org/10.1103/PhysRevB.86.045310)

PACS number(s): 71.30.+h, 76.40.+b, 73.40.Lq

Low-density and strongly correlated two-dimensional (2D) systems have attracted much attention.¹⁻³ Metallic temperature dependence of resistivity $\rho(T)$ has been observed in 2D systems where a Wigner-Seitz radius r_s is much larger than unity. Extensive experimental and theoretical studies have been carried out, as this metallic behavior contradicts the scaling theory of localization in two dimensions.⁴ However, the origin of the metallic behavior remains unclear and controversial. Another intriguing feature of low-density 2D systems is a dramatic response to a magnetic field parallel to the 2D plane. Strong positive magnetoresistance has been reported for various 2D carrier systems such as Si-metal oxide semiconductor field-effect transistors (MOSFETs),^{5,6} *n*-Si quantum wells (QWs),⁷⁻⁹ *p*-GaAs/AlGaAs,^{10,11} *n*-GaAs/AlGaAs heterojunctions,^{12,13} *n*-AlAs,^{14,15} *p*-GaAs,¹⁶ and *p*-SiGe QWs.¹⁷ It is related to the spin polarization P since an in-plane magnetic field B_{\parallel} does not couple to the 2D motion of carriers.⁶ The positive dependence of ρ on P is also a subject of discussion.

Recently, Masutomi *et al.* have performed cyclotron resonance (CR) measurements on high-mobility Si 2D electron systems (2DESs).¹⁸ The relaxation time τ_{CR} , obtained from the linewidth, was found to be comparable to the transport scattering time τ_t . It increases with decreasing temperature in a fashion similar to τ_t . The results indicate that the scattering time has the metallic T dependence over a wide frequency range. In this work, we study τ_{CR} in the presence of an in-plane magnetic field. It decreases as B_{\parallel} increases. The B_{\parallel} dependence is also similar to that of τ_t , which corresponds to the positive magnetoresistance.

The sample was fabricated from the same wafer as the one studied in Ref. 18. It is a Si/SiGe heterostructure with a 20-nm-thick strained Si QW sandwiched between relaxed Si_{0.8}Ge_{0.2} layers.¹⁹ The electrons are provided by a Sb- δ -doped layer 20 nm above the channel. The 2D electron concentration N_s was adjusted to $1.13 \times 10^{15} \text{ m}^{-2}$ at 20 K with bias voltage $V_{\text{BG}} = -5.5 \text{ V}$ of a *p*-type Si(001) substrate 2.1 μm below the channel. A two-axis vector magnet system was used to apply independently B_{\parallel} and the perpendicular magnetic field B_{\perp} for CR measurements. Instead of using a bolometer,¹⁸ we observe electron heating in the 2DES under excitation at 100 GHz. A Hall bar sample, whose channel width is 200 μm , was mounted

inside an oversized waveguide with an 8-mm bore inserted into a liquid-helium cryostat.

Figure 1 shows B_{\parallel} dependence of ρ at $B_{\perp} = 0$ for different T without millimeter-wave radiation. The in-plane magnetic field is oriented along the Hall bar direction. The critical magnetic field for the full spin polarization ($P = 1$) at $T = 0$ is estimated to be 5.0 T at this density.⁸ Arrows indicate calculated values of B_{\parallel} for $P = 0.5$ at each temperature.²⁰ While P decreases with increasing T for a fixed B_{\parallel} , the reduction of P is not significant in this temperature range. In contrast to the case of Si-MOSFETs^{5,6} and *p*-GaAs/AlGaAs heterojunctions,^{10,11} high-mobility Si 2DESs exhibit apparent metallic behavior even in the spin-polarized regime.⁷⁻⁹ In Ref. 8, Okamoto *et al.* proposed a schematic phase diagram and pointed out the importance of low disorder and the valley degree of freedom. The essential role of the valley degree of freedom for the metallic behavior is also reported for an AlAs 2DES.¹⁵

Figure 2(a) shows B_{\perp} dependence of the longitudinal resistivity ρ_{xx} at $B_{\parallel} = 3.0 \text{ T}$. The lattice or bath temperature, T_L , was kept at 1.7 K and the magnetoresistance curves were obtained with and without electromagnetic-wave excitation. We also measured the magnetoresistance at higher temperatures and confirmed that T dependence of ρ_{xx} remains metallic in the measurement region. The radiation-induced increase in ρ_{xx} can be attributed to electron heating. The electron temperature T_e is evaluated from the data obtained without radiation for $T_e = T_L > 1.7 \text{ K}$. Electron cooling to the lattice is expected to occur via electron-phonon coupling. Heat transfer rate can be obtained experimentally from the dc current-voltage characteristics.^{21,22} Using a sample fabricated from the same wafer, Toyama *et al.* found a T^5 power law in the range from 0.6 to 8 K.²³ The relaxation time is calculated to be 2 ns at 1.7 K, which is much longer than τ_t , τ_{CR} , the period of the electromagnetic wave ($= 10 \text{ ps}$) and the dephasing time ($\sim \hbar/k_B T = 4 \text{ ps}$). The weakness of the electron-phonon coupling ensures well-defined steady-state temperature T_e of the electron system. In Fig. 2(b), the CR absorption is shown assuming that it is proportional to $T_e^5 - T_L^5$. Since the radiation power was kept low so that temperature difference $\Delta T = T_e - T_L$ is much smaller than T_L , the CR absorption is almost proportional to ΔT .

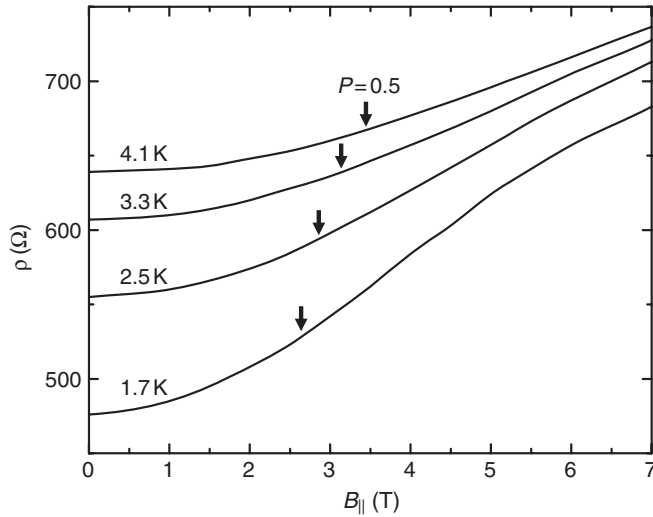


FIG. 1. Resistivity ρ as a function of $B_{||}$ at $B_{\perp} = 0$ for different temperatures. Calculated values of $B_{||}$ for $P = 0.5$ are indicated by arrows.

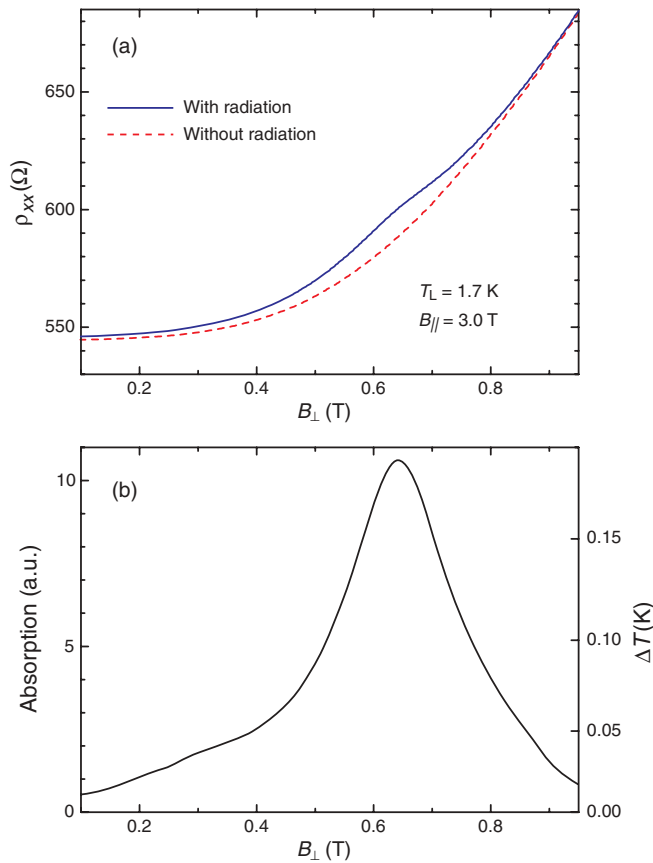


FIG. 2. (Color online) (a) B_{\perp} dependence of the longitudinal resistivity ρ_{xx} , with (solid curve) and without (dashed curve) electromagnetic-wave excitation. The in-plane magnetic field and the lattice temperature are fixed at $B_{||} = 3.0$ T and $T_L = 1.7$ K, respectively. (b) CR trace deduced from the electron temperature T_e . The energy absorption per unit area is assumed to be proportional to $T_e^5 - T_L^5$. The corresponding temperature difference, $\Delta T = T_e - T_L$, is indicated on the right axis.

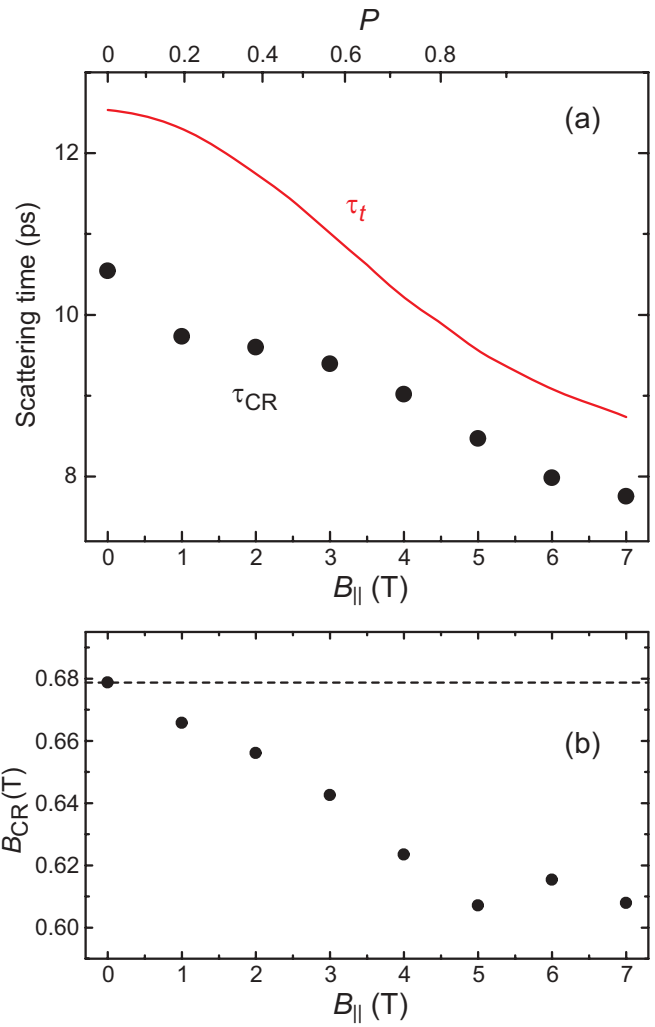


FIG. 3. (Color online) (a) Scattering times τ_t and τ_{CR} at $T = 1.7$ K as a function of $B_{||}$. The corresponding spin polarization is indicated on the upper axis. (b) The resonance magnetic field B_{CR} vs $B_{||}$. The dotted line represents $B_{CR} = 0.19m_e\omega/e$.

Figure 3(a) shows $B_{||}$ dependence of τ_t and τ_{CR} at $T = 1.7$ K. The transport scattering time $\tau_t = m^*/e^2 N_s \rho$ is determined from the data shown in Fig. 1 with the effective mass $m^* = 0.19m_e$. From the half width at half maximum ΔB of the CR absorption line, the relaxation time is obtained as $\tau_{CR} = B_{CR}/(\omega\Delta B)$. Here B_{CR} is the resonance magnetic field and ω is the microwave frequency ($\omega/2\pi = 100$ GHz). Although the data of Fig. 3(a) were obtained for a constant T_L , the difference in T_e for different $B_{||}$ is small since ΔT was kept low and about 0.2 K at the peak for all $B_{||}$. The obtained τ_{CR} exhibits a negative dependence on $B_{||}$. It is similar to that of τ_t , corresponding to the positive magnetoresistance. This suggests that the scattering time has a negative dependence on the spin polarization over a very wide frequency range from dc to 100 GHz. We believe that the present results, together with those of Ref. 18, will provide a strong constraint on theoretical models.

In Fig. 3(b), B_{CR} is plotted as a function of $B_{||}$. Unexpectedly, B_{CR} deviates from $0.19m_e\omega/e$ and decreases as $B_{||}$ increases. Electron-spin-resonance measurements demonstrate

that the spin-orbit interactions are very small in the present system.²⁴ An in-plane magnetic field can modify the wave function in the confinement direction and cause a distortion of the 2D Fermi lines.²⁵ However, this effect increases B_{CR} . The enhancement of the cyclotron mass is estimated to be only about 1% for 7 T in the present 2DES.²⁶

In summary, we have performed the cyclotron resonance measurements on a high-mobility Si 2DES in the presence of an in-plane magnetic field B_{\parallel} . The relaxation time τ_{CR} ,

obtained from the linewidth, was found to have a negative B_{\parallel} dependence, which is similar to that of the transport scattering time τ_t . The resonance peak shifts unexpectedly toward lower B_{\perp} as B_{\parallel} increases.

Acknowledgment. This work was partly supported by a Grant-in-Aid for Scientific Research (A) (Grant No. 21244047) from the Ministry of Education, Culture, Sports, Science, and Technology, Japan.

-
- ¹E. Abrahams, S. V. Kravchrenko, and M. P. Sarachik, *Rev. Mod. Phys.* **73**, 251 (2001).
- ²S. Das Sarma and E. H. Hwang, *Solid State Commun.* **135**, 579 (2005).
- ³B. Spivak, S. V. Kravchrenko, S. A. Kivelson, and X. P. A. Gao, *Rev. Mod. Phys.* **82**, 1743 (2010).
- ⁴E. Abrahams, P. W. Anderson, D. C. Licciardello, and T. V. Ramakrishnan, *Phys. Rev. Lett.* **42**, 673 (1979).
- ⁵D. Simonian, S. V. Kravchenko, M. P. Sarachik, and V. M. Pudalov, *Phys. Rev. Lett.* **79**, 2304 (1997).
- ⁶T. Okamoto, K. Hosoya, S. Kawaji, and A. Yagi, *Phys. Rev. Lett.* **82**, 3875 (1999).
- ⁷T. Okamoto, K. Hosoya, S. Kawaji, A. Yagi, A. Yutani, and Y. Shiraki, *Physica (Amsterdam)* **6E**, 260 (2000).
- ⁸T. Okamoto, M. Ooya, K. Hosoya, and S. Kawaji, *Phys. Rev. B* **69**, 041202(R) (2004).
- ⁹K. Lai, W. Pan, D. C. Tsui, S. A. Lyon, M. Mühlberger, and F. Schäffler, *Phys. Rev. B* **72**, 081313(R) (2005).
- ¹⁰J. Yoon, C. C. Li, D. Shahar, D. C. Tsui, and M. Shayegan, *Phys. Rev. Lett.* **84**, 4421 (2000).
- ¹¹E. Tutuc, E. P. De Poortere, S. J. Papadakis, and M. Shayegan, *Phys. Rev. Lett.* **86**, 2858 (2001).
- ¹²E. Tutuc, S. Melinte, and M. Shayegan, *Phys. Rev. Lett.* **88**, 036805 (2002).
- ¹³J. Zhu, H. L. Stormer, L. N. Pfeiffer, K. W. Baldwin, and K. W. West, *Phys. Rev. Lett.* **90**, 056805 (2003).
- ¹⁴E. P. De Poortere, E. Tutuc, Y. P. Shkolnikov, K. Vakili, and M. Shayegan, *Phys. Rev. B* **66**, 161308(R) (2002).
- ¹⁵O. Gunawan, T. Gokmen, K. Vakili, M. Padmanabhan, E. P. De Poortere, and M. Shayegan, *Nature Physics* **3**, 388 (2007).
- ¹⁶X. P. A. Gao, G. S. Boebinger, A. P. Mills, A. P. Ramirez, L. N. Pfeiffer, and K. W. West, *Phys. Rev. B* **73**, 241315(R) (2006).
- ¹⁷I. L. Drichko, I. Yu. Smirnov, A. V. Suslov, O. A. Mironov, and D. R. Leadley, *Phys. Rev. B* **79**, 205310 (2009).
- ¹⁸R. Masutomi, K. Sasaki, I. Yasuda, A. Sekine, K. Sawano, Y. Shiraki, and T. Okamoto, *Phys. Rev. Lett.* **106**, 196404 (2011).
- ¹⁹A. Yutani and Y. Shiraki, *Semicond. Sci. Technol.* **11**, 1009 (1996); *J. Cryst. Growth* **175–176**, 504 (1997).
- ²⁰Assuming an energy-independent density of states, the Pauli paramagnetic spin polarization P is calculated as a function of B_{\parallel} and T . Here we take into account the effects of electron-electron interactions on the g factor and the effective mass according to V. M. Pudalov, M. E. Gershenson, H. Kojima, N. Butch, E. M. Dizhur, G. Brunthaler, A. Prinz, and G. Bauer, *Phys. Rev. Lett.* **88**, 196404 (2002).
- ²¹O. Prus, M. Reznikov, U. Sivan, and V. Pudalov, *Phys. Rev. Lett.* **88**, 016801 (2001).
- ²²X. P. A. Gao, G. S. Boebinger, A. P. Mills Jr., A. P. Ramirez, L. N. Pfeiffer, and K. W. West, *Phys. Rev. Lett.* **94**, 086402 (2005).
- ²³K. Toyama, T. Nishioka, T. Okamoto, K. Sawano, and Y. Shiraki (unpublished). From the analysis of the $I - V$ characteristics, the relationship between T_e and the Joule heating power Q was found to be $Q \propto T_e^5 - T_L^5$ in the temperature range from 0.6 to 8 K.
- ²⁴J. Matsunami, M. Ooya, and T. Okamoto, *Phys. Rev. Lett.* **97**, 066602 (2006).
- ²⁵L. Smrčka and T. Jungwirth, *J. Phys.: Condens. Matter* **6**, 55 (1994).
- ²⁶Due to a one-side doping and applied gate voltage, the confining potential is regarded as a triangular well rather than a square well. The average distance of electrons from the interface is calculated to be 4 nm, which is smaller than the magnetic length $l_B = (\hbar/eB)^{1/2}$ (10 nm at 7 T).