

Ground-state phase diagram of the spin- $\frac{1}{2}$ square-lattice J_1 - J_2 model with plaquette structure

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Using the coupled-cluster method for high orders of approximation and Lanczos exact diagonalization we study the ground-state phase diagram of a quantum spin- $\frac{1}{2}$ J_1 - J_2 model on the square lattice with plaquette structure. We consider antiferromagnetic ($J_1 > 0$) as well as ferromagnetic ($J_1 < 0$) nearest-neighbor interactions together with frustrating antiferromagnetic next-nearest-neighbor interaction $J_2 > 0$. The strength of interplaquette interaction λ varies between $\lambda = 1$ (which corresponds to the uniform J_1 - J_2 model) and $\lambda = 0$ (which corresponds to isolated frustrated 4-spin plaquettes). While on the classical level ($s \rightarrow \infty$) both versions of models (i.e., with ferro- and antiferromagnetic J_1) exhibit the same ground-state behavior, the ground-state phase diagram differs basically for the quantum case $s = 1/2$. For the antiferromagnetic case ($J_1 > 0$) Néel antiferromagnetic long-range order at small J_2/J_1 and $\lambda \gtrsim 0.47$ as well as collinear striped antiferromagnetic long-range order at large J_2/J_1 and $\lambda \gtrsim 0.30$ appear which correspond to their classical counterparts. Both semiclassical magnetic phases are separated by a nonmagnetic quantum paramagnetic phase. The parameter region, where this nonmagnetic phase exists, increases with decreasing λ . For the ferromagnetic case ($J_1 < 0$) we have the trivial ferromagnetic ground state at small $J_2/|J_1|$. By increasing J_2 this classical phase gives way for a semiclassical plaquette phase, where the plaquette block spins of length $s = 2$ are antiferromagnetically long-range ordered. Further increasing of J_2 then yields collinear striped antiferromagnetic long-range order for $\lambda \gtrsim 0.38$, but a nonmagnetic quantum paramagnetic phase $\lambda \lesssim 0.38$.

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I. INTRODUCTION

The spin- $\frac{1}{2}$ quantum Heisenberg antiferromagnet with nearest-neighbor (NN), $J_1 > 0$, and next-nearest-neighbor (NNN), $J_2 \geq 0$, bonds on the square lattice has attracted much interest (see, e.g., Refs. 1–26) as a canonical model to study the interplay between frustration and quantum fluctuations. In particular, the quantum phase transitions inherent in this model as well as the nature of its quantum paramagnetic phase in the region $0.4 \lesssim J_2/J_1 \lesssim 0.6$ is a matter of intensive debate. The Néel antiferromagnetic (NAF) long-range order (LRO) at small J_2/J_1 and the collinear striped antiferromagnetic (CAF) LRO at large J_2/J_1 correspond to their classical counterparts; however, the sublattice magnetization is reduced by quantum fluctuations.

Motivated by recent investigations on quasi-two-dimensional frustrated magnetic materials with ferromagnetic (FM) NN bonds, for example, $\text{Pb}_2\text{VO}(\text{PO}_4)_2$,^{27–31} $\text{SrZnVO}(\text{PO}_4)_2$,^{30,32–34} and $\text{BaCdVO}(\text{PO}_4)_2$,^{29,32,35} the FM J_1 - J_2 model ($J_1 < 0, J_2 \geq 0$) has recently been studied.^{14,36–46} Although on the classical level the ground-state (GS) phase diagrams of both variants of the J_1 - J_2 model are quite similar, the GS properties of the quantum model with $J_1 < 0$ are basically different from those for $J_1 > 0$. Currently, the GS properties of the FM model around $J_2/|J_1| = 0.5$ are under controversial debate. On the one hand, in Refs. 37 and 45 arguments for a nematic phase separating the classical FM and the semiclassical CAF phases are given; on the other hand, in Refs. 44 and 46 an intermediate quantum paramagnetic GS phase is not found (or it exists in a very small parameter region around $J_2 \sim 0.4|J_1|$ only). Hence, a direct first-order transition between the FM GS and the GS with CAF LRO at $J_2 \sim 0.4|J_1|$ could take place.^{44,46}

Quantum phase transitions can also occur by competition between antiferromagnetic (AFM) NN bonds of different strength, that is, without frustration. One example is the local singlet formation in dimerized Heisenberg models (see, e.g., Ref. 47 and references therein). As a certain extension to dimerized models the square-lattice Heisenberg model with a plaquette structure has been considered,^{9,48–53} where local quadrumer singlet formation can destroy NAF LRO.

However, for the description of real materials a modification of the J_1 - J_2 model might be necessary. For example, for $\text{PbZnVO}(\text{PO}_4)_2$ a spatially anisotropic model has been derived,⁵⁴ whereas for $(\text{CuCl})\text{LaNb}_2\text{O}_7$ ⁵⁵ a J_1 - J_2 model with additional plaquette structure was proposed in Ref. 49. This plaquette model, however, has been questioned recently by Rosner and co-workers.⁵⁶

In this paper we consider the J_1 - J_2 model on the square lattice with plaquette structure (see Fig. 1) merging this way the two mechanisms to destroy magnetic LRO, namely frustration and local singlet formation. Moreover, the model we consider corresponds to that proposed in Ref. 49 for $(\text{CuCl})\text{LaNb}_2\text{O}_7$. The Hamiltonian of our model is given by

$$H = \sum_{a=0}^1 (\delta_{0a} + \delta_{1a}\lambda) \left(J_1 \sum_{\langle ij \rangle_a} \mathbf{s}_i \mathbf{s}_j + J_2 \sum_{\langle\langle ij \rangle\rangle_a} \mathbf{s}_i \mathbf{s}_j \right), \quad (1)$$

where the sum is taken over the nearest ($\langle \cdot \cdot \rangle$) and next-nearest ($\langle\langle \cdot \cdot \rangle\rangle$) neighbors, and δ_{ab} is the Kronecker symbol. We consider FM and AFM NN bonds $J_1 = \pm 1$ and AFM (i.e., frustrating) NNN bonds $J_2 \geq 0$. The intra- ($a = 0$) and interplaquette ($a = 1$) interactions differ by the factor λ . The case $\lambda = 1$ corresponds to the standard uniform J_1 - J_2 model on the square lattice, whereas the limit $\lambda = 0$ describes unconnected 4-spin J_1 - J_2 plaquettes. For the parameter λ

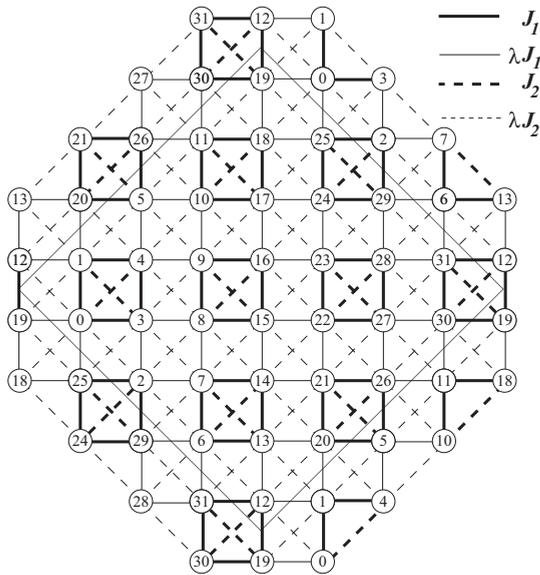


FIG. 1. (Color online) The two-dimensional spin- $\frac{1}{2}$ J_1 - J_2 finite square lattice of $N = 32$ sites (periodic boundary conditions) with plaquette structure [see Eq. (1)]. Solid lines, NN bonds J_1 (thick) and λJ_1 (thin); dashed lines, NNN bonds J_2 (thick) and λJ_2 (thin).

we consider the interval $\lambda \in [0,1]$, this way interpolating between a system of isolated plaquettes and the uniform J_1 - J_2 model.

The model (1) has been studied previously by series expansion⁹ for AFM $J_1 > 0$, and by bond-operator mean-field theory as well as a second-order perturbation theory in λ (Ref. 49) for (FM) $J_1 < 0$. Bearing in mind the intensive work on the standard J_1 - J_2 model (see Refs. 1–26 and references therein), it seems to be desirable to discuss the much less studied plaquette J_1 - J_2 model by methods alternative to those used in Refs. 9 and 49, thereby confirming or questioning the findings of those papers.

In this paper we consider both the AFM ($J_1 > 0$) and the FM ($J_1 < 0$) version of the model (1), and we derive the GS quantum phase diagram for both versions of the model. We use the Lanczos exact diagonalization (see Sec. III A) and the coupled cluster method (CCM) (see Sec. III B) to analyze the GS of the model. Both methods are powerful general many-body methods and, particularly, the combination of both can lead to a consistent description of various aspects of quantum spin systems. We note that another important method for quantum spin systems, the quantum Monte Carlo (QMC) method, cannot be used for the considered frustrated model since it suffers from the minus-sign problem.

II. CLASSICAL GROUND-STATE PHASES

It is well known that the GS for the classical case ($s \rightarrow \infty$) of the uniform J_1 - J_2 model (i.e., $\lambda = 1$) has three phases: the NAF phase for $J_2 < J_1/2$ and $J_1 > 0$, the FM phase for $J_2 < -J_1/2$ and $J_1 < 0$, and a phase consisting of two interpenetrating Néel-ordered square lattices. The angle between the directions of the two interpenetrating Néel states is arbitrary in the classical limit, whereas, due to the *order-from-*

disorder effect, collinear states are favored by fluctuations; that is, the CAF is realized. By contrast to the uniform model, two additional plaquette phases appear for $\lambda < 1$ (see Fig. 2), namely the plaquette antiferromagnetic (PAF) phase separating the NAF from the CAF phase for $J_1 > 0$ and the plaquette ferromagnetic (PFM) between the FM and the CAF phases for $J_1 < 0$. These plaquette phases have some relation to the CAF phase: They also can be understood as a system of two interpenetrating Néel-ordered square lattices, where again the angle between the directions of the two interpenetrating Néel states is arbitrary in the classical limit. However, as an elementary unit (a “lattice site”) of the two interpenetrating square lattices does not act a single spin (site), but rather a 4-spin-(4-site-)plaquette carrying strong bonds. Again, due to the *order-from-disorder* effect collinear states are favored by fluctuations; that is, the collinear PAF or PFM phases are realized, respectively. A graphical illustration of these classical phases is given in Fig. 2. The energy of the classical plaquette phases reads $E_{\text{PAF,PFM}} = -[\pm \frac{1}{4} J_1 - \frac{1}{8} J_2] - \frac{1}{8} \lambda J_2$, where the upper (lower) sign corresponds to PAF (PFM). Obviously, the interplaquette interaction is due to λJ_2 only. The transition line between the CAF and the PAF (PFM) phases is given by $J_2 = J_1/(1 + \lambda)$ [$J_2 = -J_1/(1 + \lambda)$], the transition between the NAF and the PAF phases as well as the FM and the PFM phases are horizontal lines $J_2 = J_1/2$ and $J_2 = -J_1/2$, respectively. Besides the trivial degeneracies, the CAF, the PAF, and the PFM states are twofold degenerate, since the stripes of parallel spins (CAF) or parallel plaquettes (PAF, PFM) can be arranged along either the horizontal or the vertical direction.

III. METHODS

A. Exact diagonalization of finite lattices

The Lanczos exact diagonalization method was successfully used to discuss the GS phases of the uniform $s = 1/2$ J_1 - J_2 model (i.e., $\lambda = 1$) using finite lattices of $N = 16, 20, 32, 36$, and 40 sites.^{2–5,21,44} However, the new classical phases, PAF and PFM, appearing due to the plaquette structure do not fit to the periodic boundary conditions of the finite lattices of $N = 20, 36$, and 40 sites, and, therefore, we do not have the possibility to perform a finite-size extrapolation used previously for the uniform model.^{3,21,44} Using Jörg Schulenburg’s SPINPACK⁵⁷ we have calculated the GS focusing on the finite lattice of $N = 32$ sites shown in Fig. 1 for a large set of λ and J_2 values fixing J_1 either to $+1$ or to -1 .

To analyze the GS magnetic ordering we have calculated the spin-spin correlation function $\langle s_0 s_i \rangle$ as well as a general finite-size order parameter,^{58,59}

$$m^2 = \frac{1}{(N-4)^2} \sum_{i,j \neq P}^{N-4} |\langle s_i s_j \rangle|, \quad (2)$$

that adds the total strength of the spin-spin correlation functions. Bearing in mind the strong intraplaquette correlation functions for $\lambda < 1$ we have excluded these correlations in the sum, this way increasing the weight of distant correlation functions.

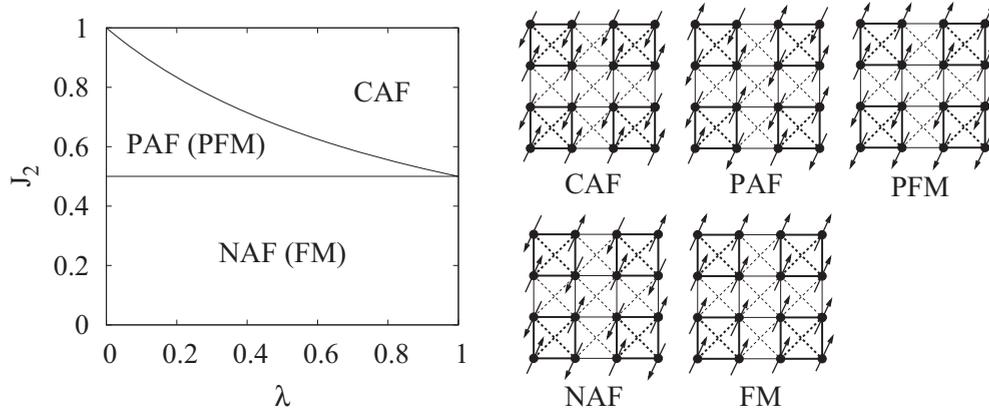


FIG. 2. (Left) The classical phase diagram of the two-dimensional J_1 - J_2 square-lattice Heisenberg model with plaquette structure, see Eq. (1). For the AFM model ($J_1 = +1$) the GS phases CAF, NAF, and PAF exist. For the FM model ($J_1 = -1$) the CAF GS phase also exists, but the NAF and the PAF phases are replaced by the FM and the PFM phases, respectively. (Right) Illustration of the classical phases (CAF, collinear striped antiferromagnetic order; PAF, plaquette antiferromagnetic order; NAF, Néel antiferromagnetic order; FM, ferromagnetic order; PFM, plaquette ferromagnetic order).

B. Coupled-cluster method

We first mention that the CCM yields results directly in the thermodynamic limit $N \rightarrow \infty$. The CCM has been previously reviewed in several articles (e.g., for the AFM^{8,15–17} and for the FM J_1 - J_2 model,^{44,60} as well as for dimerized models^{61–64}) and will not be repeated here in detail. For more general information on the methodology of the CCM, see, for example, Refs. 65–68.

The CCM is a quantum many-body method and is defined by a reference (or model) state $|\Phi\rangle$ and a complete set of mutually commuting many-body creation operators C_I^+ . For our model we choose the classical GSs (see Fig. 2) as reference states. It is convenient to perform an appropriate rotation of the local axis of the spins such that in the rotated coordinate frame the reference state is a tensor product of spin-down states $|\Phi\rangle = |\downarrow\rangle|\downarrow\rangle|\downarrow\rangle \dots$. The creation operators are then the multispin creation operators $C_I^+ = s_i^+, s_i^+ s_j^+, s_i^+ s_j^+ s_k^+, \dots$, where the indices i, j, k, \dots denote arbitrary lattice sites.

The CCM parametrizations of the ket and bra GSs are given by [with $C_I^- = (C_I^+)^+$]

$$|\Psi\rangle = e^S |\Phi\rangle, \quad S = \sum_{I \neq 0} S_I C_I^+; \quad (3)$$

$$\langle \tilde{\Psi} | = \langle \Phi | \tilde{S} e^{-S}, \quad \tilde{S} = 1 + \sum_{I \neq 0} \tilde{S}_I C_I^-.$$

Using $\langle \Phi | C_I^+ = 0 = C_I |\Phi\rangle \quad \forall I \neq 0, \quad C_0^+ \equiv 1$, the orthonormality condition $\langle \Phi | C_I C_J^+ |\Phi\rangle = \delta_{IJ}$, and the completeness relation $\sum_I C_I^+ |\Phi\rangle \langle \Phi | C_I = 1 = |\Phi\rangle \langle \Phi | + \sum_{I \neq 0} C_I^+ |\Phi\rangle \langle \Phi | C_I$ we get a set of nonlinear and linear equations for the correlation coefficients S_I and \tilde{S}_I , respectively. The order parameter (sublattice magnetization) in the rotated coordinate frame is given by

$$M = -\frac{1}{N} \sum_i \langle \tilde{\Psi} | s_i^z | \Psi \rangle. \quad (4)$$

In the CCM the only approximation is the truncation of the expansion of the correlation operator S and \tilde{S} . We use

the well-established LSUB m scheme, where all multispin correlations on the lattice with m or fewer contiguous sites are taken into account. The number of these configurations can be reduced, using lattice symmetry and conservation laws, but it is increasing very rapidly with m . In the highest order of approximation considered here, LSUB10, we have, for example, 180 957 configurations for the collinear stripe reference state (CAF) and 219 446 for the antiferromagnetic plaquette state (PAF), that is, finally 180 957 or 219 446 corresponding coupled nonlinear equations which have to be solved numerically.

Since the LSUB m approximation scheme becomes exact for $m \rightarrow \infty$, we can improve our results by extrapolating the “raw” LSUB m data to $m \rightarrow \infty$. There is ample empirical experience regarding how one should extrapolate the magnetic order parameter $M(m)$ for systems with quantum phase transition between magnetically ordered and disordered GS phases. Following Refs. 15–17 and 44, we use $M(m) = b_0 + b_1(1/m)^{1/2} + b_2(1/m)^{3/2}$ to extrapolate to $m \rightarrow \infty$. As a rule,^{15–17,44} the lowest level of approximation, LSUB2, is excluded from extrapolation. For the particular 4-spin plaquette model considered here, we expect LSUB2 to be especially poor, because its cluster size (2 spins) is smaller than the unit cell (4 spins) of the model. Hence, we use LSUB4, LSUB6, LSUB8, and LSUB10 data for the extrapolations.

IV. RESULTS FOR THE QUANTUM $s = 1/2$ MODEL

A. Antiferromagnetic nearest-neighbor exchange J_1

We first consider the AFM case, set $J_1 = +1$, and start with the discussion of ED data (see Figs. 3 and 4). Corresponding to the three classical phases there are three regimes in the quantum model. The spin-spin correlation functions $\langle s_0 s_i \rangle$ presented in Fig. 3 for $\lambda = 0.5$ and $\lambda = 0.7$ are quite strong for small and large J_2 , thus indicating semiclassical magnetic LRO, where the signs of $\langle s_0 s_i \rangle$ fit to the classical phases NAF and CAF. In an intermediate regime, around $J_2 = 0.5$, the interplaquette spin-spin correlations are weak. These three different regimes are also well seen in Fig. 4, where the

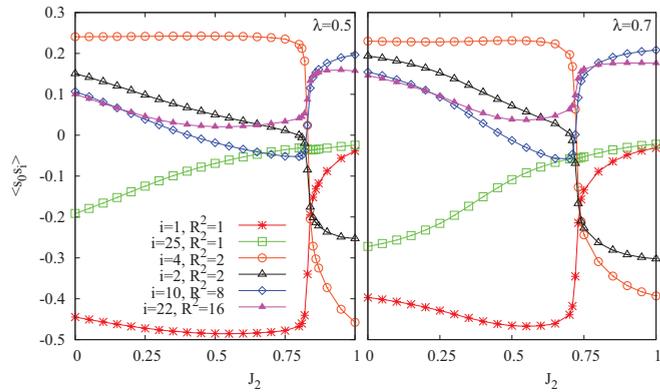


FIG. 3. (Color online) ED data for the spin-spin correlation functions $\langle s_0 s_i \rangle$ versus J_2 (AFM $J_1 = +1$) for two values of λ for the finite lattice of $N = 32$ sites. Except the site indices 0 and i corresponding to Fig. 1 we also give the square of separation R^2 of the sites 0 and i . Note that $\langle s_0 s_1 \rangle$ and $\langle s_0 s_4 \rangle$ are intraplaquette correlation functions, whereas the other ones are interplaquette correlation functions.

finite-size order parameter m^2 [see Eq. (2)] is shown. In particular, the m^2 data yield clear evidence for regions with weak magnetic order. The widths of those regions increase drastically with decreasing λ . Bearing in mind that for $\lambda = 1$ the region of magnetic disorder is $0.4 \lesssim J_2 \lesssim 0.6$ (see, e.g., Refs. 12,17,18,21,22, and 24), our ED data for $N = 32$ suggest that there is no semiclassical antiferromagnetic LRO of plaquette type (PAF). Moreover, for small values of $\lambda \lesssim 0.5$ the finite-size order parameter is small in the whole range of J_2 values, indicating the absence of any magnetic LRO.

Now we discuss the CCM results ($N \rightarrow \infty$). Figure 5 provides the CCM GS energy per site e_0 for $\lambda = 0.5$ and $\lambda = 0.7$ compared with the corresponding ED data. Obviously, the CCM and the ED agree well. To discuss magnetic LRO we consider the CCM-LSUB m magnetic order parameter M (sublattice magnetization) extrapolated to $m \rightarrow \infty$ [see

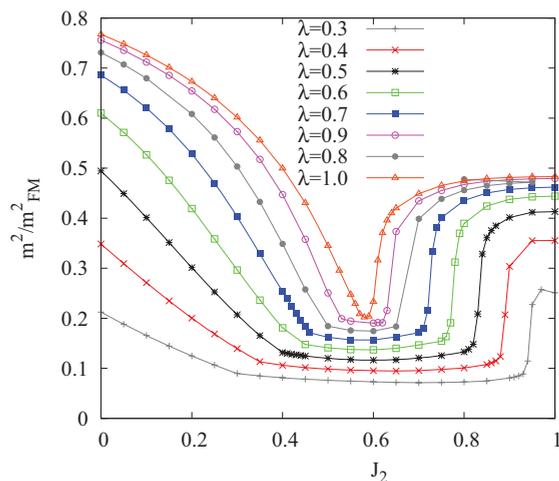


FIG. 4. (Color online) ED data for the finite-size order parameter m^2/m_{FM}^2 as defined in Eq. (2) versus J_2 (AFM $J_1 = +1$) for various values of λ for the finite lattice of $N = 32$ sites. For convenience m^2 is scaled by its value for the FM state $m_{\text{FM}}^2 = 0.25$.

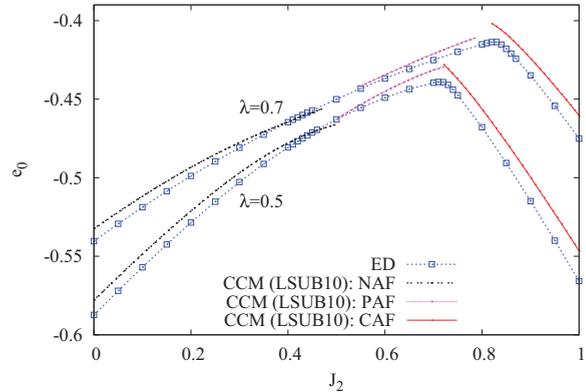


FIG. 5. (Color online) CCM LSUB10 data for the GS energy e_0 (AFM $J_1 = +1$) for $\lambda = 0.5$ and $\lambda = 0.7$ compared with corresponding ED data ($N = 32$). The CCM results correspond to the NAF reference state (small J_2), the PAF reference state (intermediate J_2), and the CAF reference state (large J_2).

Eq. (4)] that is depicted in Fig. 6. To illustrate the quality of the used extrapolation for the order parameter M of the “raw” LSUB m data to the limit $m \rightarrow \infty$ we also show, as an example, corresponding plots for the $\lambda = 0.7$ in Fig. 7. It is obvious that the LSUB m data are well fitted by the applied extrapolation function.

First we notice that the extrapolated order parameter using the PAF reference state for all relevant parameter sets of λ and J_2 vanishes (see also Fig. 7); that is, there is no PAF magnetic LRO in the quantum model. This finding is in agreement with the ED data for m^2 discussed above (and see again Fig. 4). The range of stability of the semiclassical NAF and CAF phases is visible in Fig. 6. We define the quantum critical points $J_2^{c1}(\lambda)$ and $J_2^{c2}(\lambda)$ as those points, where the extrapolated CCM magnetic order parameter M vanishes. For $\lambda = 1$ (uniform model) we find $J_2^{c1} = 0.447J_1$ and $J_2^{c2} = 0.586J_1$, which is in agreement with previous CCM predictions.^{15–17} As expected,

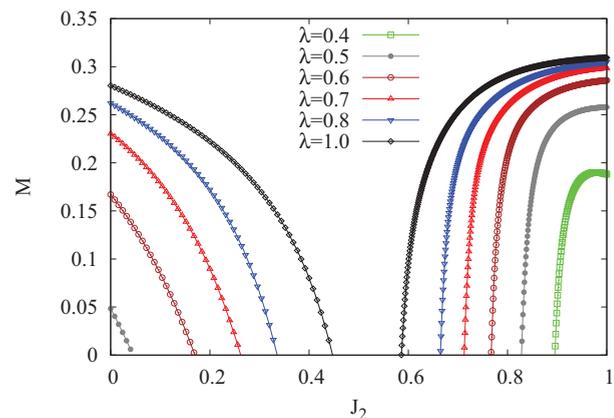


FIG. 6. (Color online) Extrapolated CCM data for the magnetic order parameter (sublattice magnetization) M for the AFM case ($J_1 = +1$) and for various values of λ . The CCM reference state is the NAF state (left-hand side, small J_2) and the CAF state (right-hand side, large J_2). The extrapolated CCM data are obtained using LSUB m with $m = 4, 6, 8, 10$ and the extrapolation rule $M(m) = b_0 + b_1(1/m)^{1/2} + b_2(1/m)^{3/2}$.

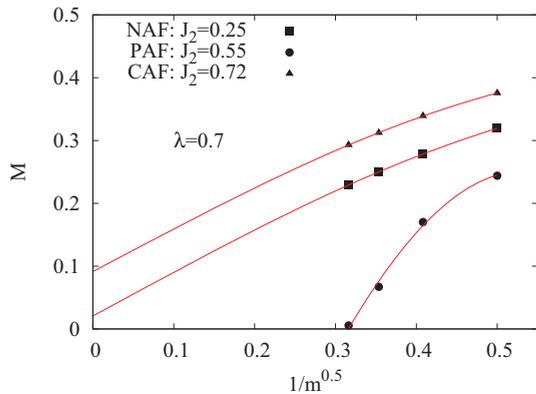


FIG. 7. (Color online) Illustration of the extrapolation (solid lines) of the CCM-LSUB m data (symbols) for the magnetic order parameter M for $\lambda = 0.7$, $J_1 = +1$ and specific values of J_2 corresponding to the NAF, PAF, CAF reference states. For the extrapolations to $m \rightarrow \infty$ according to $M(m) = b_0 + b_1(1/m)^{1/2} + b_2(1/m)^{3/2}$ we have used LSUB m data for $m = 4, 6, 8, 10$.

with decreasing λ the region of J_2 values without magnetic LRO increases (see Fig. 6). Collecting the data for the quantum critical points J_2^{c1} and J_2^{c2} for various λ we get the GS phase diagram of the quantum model as shown in Fig. 8. Our phase diagram is in good agreement with the only available corresponding one of Ref. 9. Hence, one can argue that the phase diagram is basically correct.

Concerning the order of the quantum phase transitions at J_2^{c1} and J_2^{c2} our data are in favor of a continuous transition at J_2^{c1} and a first-order transition at J_2^{c2} as it was discussed for the uniform model.^{3,9,11,21} The scenario of a first-order transition between the QPM and the CAF phases is supported by (i) the kinklike behavior in e_0 (see Fig. 5), (ii) the steep fall in the CCM order parameter M near J_2^{c2} (see Fig. 6), and (iii) the jumplike behavior of the spin-spin correlations functions (see Fig. 3) and of the finite-size order parameter (see Fig. 4).

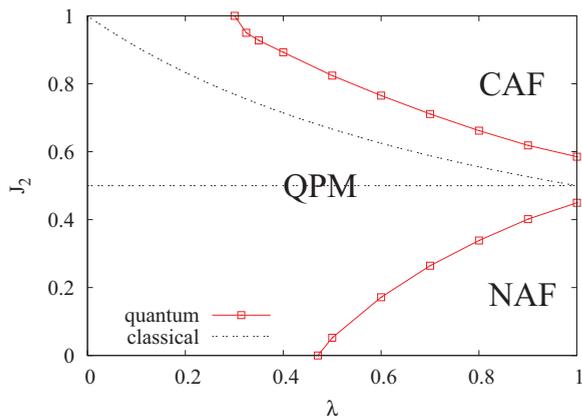


FIG. 8. (Color online) The CCM GS phase diagram of the model (1) for the AFM case ($J_1 = +1$). NAF denotes the semiclassical phase with Néel LRO, CAF the semiclassical phase with CAF LRO, and QPM the magnetically disordered quantum paramagnetic phase. For comparison we also show the classical transition lines (thin dashed lines) (cf. Fig. 2).

For the unfrustrated case $J_2 = 0$ we can compare our result $\lambda_c = 0.47$ (and see Fig. 8) for the critical λ , where the NAF LRO breaks down with several previous results, namely $\lambda_c = 0.555$ (Ising series expansion⁹), $\lambda_c = 0.6$ (exact diagonalization⁴⁸), $\lambda_c = 0.5485$ (QMC^{50,52}), $\lambda_c = 0.5491, \dots, 0.5513$ (contractor renormalization expansion⁵⁰), and $\lambda_c = 0.4822$ (real space renormalization group approach⁵³). Likely, the QMC result is most accurate. Our result is in reasonable agreement with that result, but slightly overestimates the stability region of NAF LRO.

Let us briefly discuss another limit of the model, namely the limit of large J_2 . As discussed above in that limit the system splits into two interpenetrating square-lattice Heisenberg anti-ferromagnets (where J_2 plays the role of the AFM NN bond), which are Néel-ordered for $\lambda = 1$. However, for $\lambda < 1$ each of the two interpenetrating square lattices carries now a staggered arrangement of J_2 and λJ_2 bonds, which corresponds precisely to the so-called J - J' model discussed in Refs. 51,52,61,69–71. The QMC estimate of the critical λ is $\lambda_c = 0.397$ (see Ref. 51). For $J_2 = 1$ we find from our CCM data a critical value of $\lambda_c = 0.301$ (see Fig. 8). Increasing J_2 beyond $J_2 = 1$ (not shown in Fig. 8) yields a steep increase of the critical line as already indicated by the last two data points near $J_2 = 1$ in Fig. 8. In the limit of large J_2 the critical λ_c is even smaller as for $J_2 = 1$. For example, at $J_2 = 5$ we obtain $\lambda_c = 0.375$ that is in good agreement with the QMC result for the critical λ_c of the J - J' model. Note that our value of λ_c deviates from an early CCM result $\lambda_c = 0.316$.⁶¹ The difference is related to the fact that we use here (i) a higher approximation level (LSUB10) and (ii) an improved extrapolation in comparison with Ref. 61.

B. Ferromagnetic nearest-neighbor exchange J_1

We now consider the FM case and set $J_1 = -1$. Although, the classical phase diagrams for $J_1 = -1$ and $J_1 = +1$ are identical (see Fig. 2), we know from the uniform model (i.e., $\lambda = 1$) that the phase diagrams in the quantum case $s = 1/2$ are basically different.^{14,36–46} In particular, it was found for $J_1 = -1$ that the region of semiclassical CAF LRO extends up to much smaller values of J_2 compared with the case of $J_1 = +1$.^{44,46} Hence, we expect also for $\lambda < 1$ basic differences between $J_1 = +1$ and $J_1 = -1$.

We start again with ED results for the spin-spin correlation functions $\langle s_0 s_i \rangle$ (see Fig. 9) and the finite-size order parameter m^2 defined in Eq. (2) (see Fig. 10). The three classical phases lead again to three different regimes in the quantum model. The trivial FM state at small J_2 is, of course, also present in the quantum model. It gives way at $J_2^{c1}(\lambda)$ for an intermediate singlet state. This transition at J_2^{c1} is clearly seen by jumps in $\langle s_0 s_i \rangle$ and m^2 . J_2^{c1} is smaller than the classical value $J_{2, \text{clas}}^{c1} = 0.5$, and it depends on λ for the quantum model. Moreover, there is a second jumplike behavior of the spin-spin correlation functions indicating the transition from the intermediate regime to the CAF regime. We can use the positions of these jumps to extract the transition points J_2^{c1} and J_2^{c2} between the regimes from our ED calculations, see the discussion of Fig. 12 given below.

The main difference in comparison to the AFM model ($J_1 = +1$) concerns the intermediate regime. From Fig. 9 it

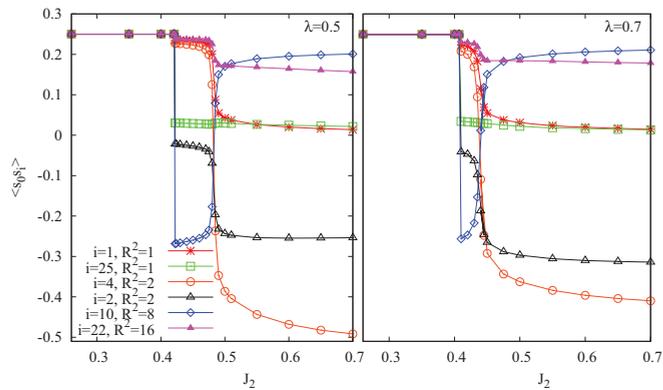


FIG. 9. (Color online) ED data for the spin-spin correlation functions $\langle s_0 s_i \rangle$ versus J_2 (FM $J_1 = -1$) for two values of λ for the finite lattice of $N = 32$ sites. Except the site indices 0 and i corresponding to Fig. 1 we also give the square of separation R^2 of the sites 0 and i . Note that $\langle s_0 s_1 \rangle$ and $\langle s_0 s_4 \rangle$ are intraplaquette correlation functions, whereas the other ones are interplaquette correlation functions.

is evident that the interplaquette correlation functions $\langle s_0 s_{10} \rangle$ and $\langle s_0 s_{22} \rangle$ (see Fig. 1) are not small. Hence, the intermediate phase might be long-range ordered for FM J_1 . Indeed, in the finite-size order parameter m^2 shown in Fig. 10 an intermediate regime is clearly visible for $\lambda \lesssim 0.5$, where the finite-size order parameter m^2 in this regime is much larger than that for $J_1 = +1$ (cf. Fig. 4). The possible appearance of magnetic LRO can be understood on the basis of the classical PFM state (see also Sec. II): The elementary unit is the 4-spin-plaquette (with strong bonds). Due to the strong FM intraplaquette bond J_1 the plaquette carries an effective block (composite) spin $s = 2$. These block spins interact via λJ_2 and form effectively two interpenetrating square-lattice $s = 2$ Heisenberg antiferromagnets.

Now we pass to the CCM results ($N \rightarrow \infty$). Figure 11 provides the CCM GS energy per site e_0 for $\lambda = 0.5$ and $\lambda = 0.7$ in LSUB6 and LSUB10 approximation compared with the corresponding ED data. The agreement between the

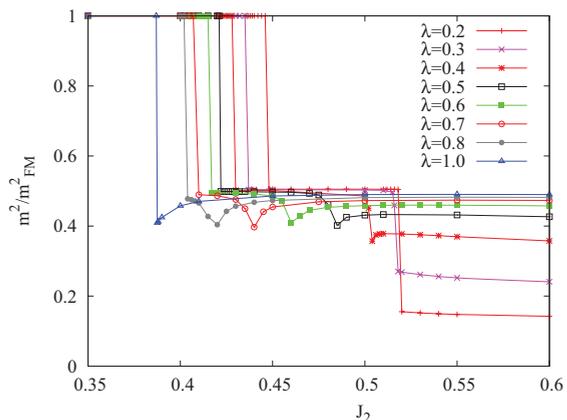


FIG. 10. (Color online) ED data for the finite-size order parameter m^2/m_{FM}^2 as defined in Eq. (2) versus J_2 (FM $J_1 = -1$) for various values of λ for the finite lattice of $N = 32$ sites. For convenience m^2 is scaled by its value for the FM state $m_{\text{FM}}^2 = 0.25$.

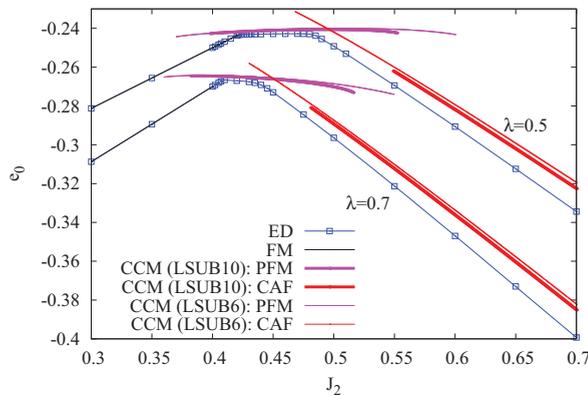


FIG. 11. (Color online) CCM-LSUB10 (thick lines) and LSUB6 (thin lines) data for the GS energy e_0 ($J_1 = -1$) for $\lambda = 0.5$ and $\lambda = 0.7$ compared with corresponding ED data ($N = 32$). The CCM results correspond to the PFM reference state (intermediate J_2) and the CAF reference state (large J_2). The black solid line is the exact FM GS energy.

CCM and the ED is good. The three regimes, FM, PFM, and CAF, are clearly seen in the energy data. For the CCM estimate for the transition point J_2^{c1} between the FM state and the PFM regime we use the intersection points between the FM energy and the CCM-LSUB10 as well as the LSUB6 energy calculated with the PFM reference state. Both results for J_2^{c1} almost coincide (see Fig. 12). Moreover, we find an excellent agreement between the ED and CCM results for J_2^{c1} . Concerning the order of the phase transition at J_2^{c1} we have clear evidence for a first-order transition as indicated by (i) the jumplike behavior of the spin-spin correlations functions (see Fig. 9) and of the finite-size order parameter (see Fig. 10), (ii)

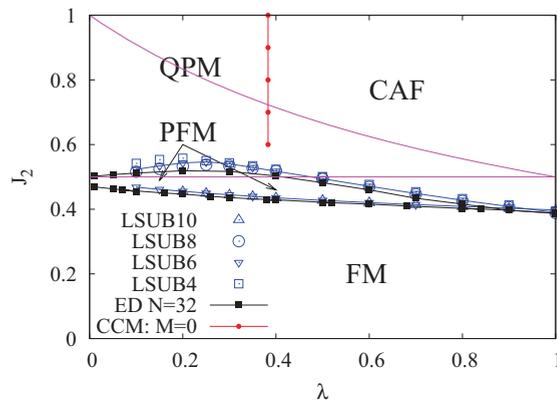


FIG. 12. (Color online) GS phase diagram of the plaquette model of Eq. (1) for the FM case ($J_1 = -1$). The transition lines between the FM and the PFM phases as well as the PFM and the CAF/QPM phases are obtained by (i) ED by using the jumps in the ED data (black lines) (see Figs. 9 and 10) and (ii) by CCM from the intersection points between the GS energies using various reference states for the CCM calculation (blue lines) (see text). The transition lines between the CAF and the QPM phase (red line) correspond to those parameter values (λ, J_2), where the extrapolated CCM order parameter vanishes (and see Fig. 14). For comparison we also show the classical transition lines (magenta line) (cf. Fig. 2).

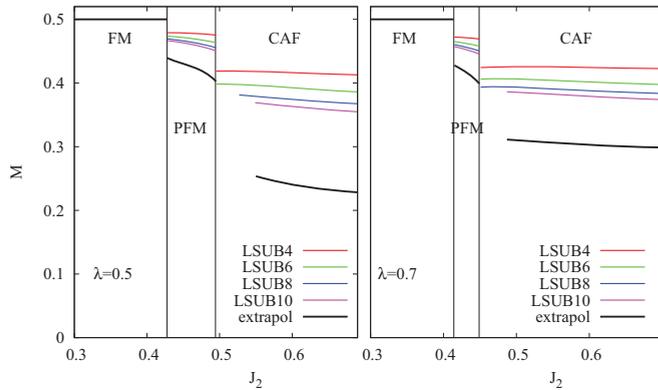


FIG. 13. (Color online) CCM magnetic order parameter M for the FM case ($J_1 = -1$) versus J_2 using the PFM and the CAF states as CCM reference states for $\lambda = 0.5$ and $\lambda = 0.7$. The extrapolated CCM results are obtained using LSUB m with $m = 4, 6, 8, 10$ and the extrapolation rule $M(m) = b_0 + b_1(1/m)^{1/2} + b_2(1/m)^{3/2}$.

the kinklike behavior in e_0 (see Fig. 11), and (iii) the jump in the CCM order parameter M at J_2^{c1} (see Fig. 13).

To determine the transition point J_2^{c2} between the PFM and the CAF regimes we are faced with the problem that near this transition no CCM-LSUB10 solutions for the CAF reference state could be found. This observation, that for high orders of CCM approximation in the vicinity of an intersection point of two CCM energy curves belonging to two different GS phases no solution of the large set of coupled nonlinear ket equations can be obtained, is often found when dealing with strongly frustrated systems (see, e.g., Refs. 13 and 44). However, from Fig. 11 it is obvious that the parameter range, where solutions for lower levels m of LSUB m approximations can be found, is much larger than for LSUB10. Moreover, the LSUB m data for e_0 for various m are very close to each other. Hence, we can use the intersection points for lower levels of CCM-LSUB m approximation to determine the second transition point J_2^{c2} . Thus, we find intersection points for LSUB8, LSUB6, and LSUB4 for $\lambda \geq 0.7$, $\lambda \geq 0.5$, and $\lambda \geq 0.25$, respectively. In addition, we can also take benefit from the almost straight behavior of $e_0(J_2)$ curves (see Fig. 11), which allows a reliable extrapolation of $e_0(J_2)$ until a hypothetical intersection point (cf. also Refs. 44 and 46). This gives finally various sets of $J_2^{c2}(\lambda)$ data, which, however, agree well with each other (cf. Fig. 12). Only for $\lambda \lesssim 0.3$ a slight difference is visible. Nevertheless, it is necessary to mention that the CCM estimate for $J_2^{c2}(\lambda)$ becomes less reliable for smaller values of λ due to the increasing distance between the hypothetical intersection points and the last data points for which LSUB m solutions for the CAF reference state could be found. The results for $J_2^{c1}(\lambda)$ and $J_2^{c2}(\lambda)$ obtained by ED and CCM are collected in the phase diagram presented in Fig. 12. It is obvious that both ED and CCM yield very similar values for J_2^{c1} and J_2^{c2} .

The question about magnetic LRO in the various regimes we address next by analyzing the CCM magnetic order parameter M for the PFM and CAF regimes. (Remember that in the trivial FM phase we have always $M = 1/2$.) In Fig. 13 we show M versus J_2 for $\lambda = 0.5$ and $\lambda = 0.7$. Obviously, the order parameter in the PFM regime is nonzero and even larger than in the CAF regime. This observation is in agreement with

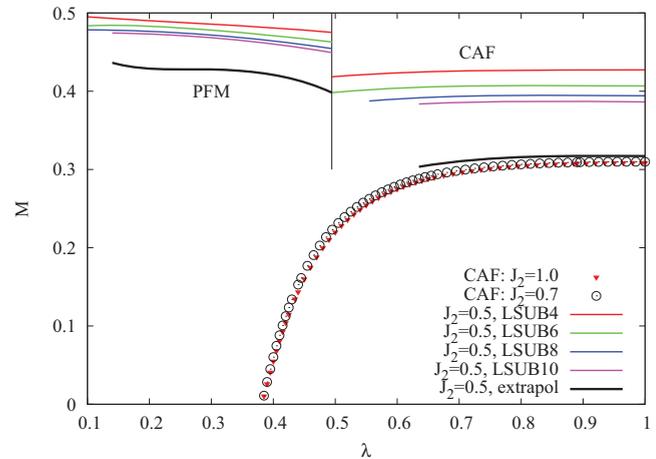


FIG. 14. (Color online) Magnetic order parameter M for the FM case ($J_1 = -1$) versus λ for $J_2 = 1.0$ (solid triangles), 0.7 (circles), and 0.5 (lines) using the PFM and the CAF states as CCM reference states. The extrapolated CCM results are obtained using LSUB m with $m = 4, 6, 8, 10$ and the extrapolation rule $M(m) = b_0 + b_1(1/m)^{1/2} + b_2(1/m)^{3/2}$.

our ED results (cf. Fig. 10) and can be related to the effective block-spin $s = 2$ model discussed above. The order parameter in the CAF is only weakly dependent on J_2 , but its magnitude depends on λ . As discussed above we do not get CAF CCM solutions for LSUB m for higher values of m near the transition to the PFM phase. Nevertheless, the curves presented in Fig. 13 indicate that there is likely a direct first-order transition at J_2^{c2} between semiclassical phases with PFM and CAF magnetic LRO for $\lambda \gtrsim 0.4$.

Next we fix J_2 and consider the behavior of the order parameter M in dependence on λ (see Fig. 14). For values of $J_2 \gtrsim 0.55$ the PFM regime is not relevant. As explained in Sec. IV A in the limit of large J_2 our plaquette model corresponds to the J - J' model, where at $\lambda = \lambda_c = 0.397$ (QMC result, see Ref. 51) a second-order transition to a QPM phase takes place. This behavior is clearly seen in our CCM data for $M(\lambda)$ for $J_2 = 1.0$ and $J_2 = 0.7$ shown in Fig. 14. We find that for $J_2 \gtrsim 0.6$ the extrapolated CCM order parameter M vanishes continuously at a critical value $\lambda = \lambda_c(J_2)$ defining a second-order transition between a semiclassical phase with CAF magnetic LRO and the QPM phase, which is depicted as a red line in Fig. 12. Interestingly, λ_c practically does not depend on J_2 and we have $\lambda_c \approx 0.383$ for $0.6 \lesssim J_2 \leq 1$. This value is close to the QMC value⁵¹ $\lambda_c = 0.397$ valid for $J_2 \rightarrow \infty$ and also to $\lambda_c = 0.375$ obtained for the AFM model ($J_1 = +1$) for $J_2 = 5$ (see Sec. IV A).

For $J_2 = 0.5$ the situation is different, since both, the PFM and CAF regimes, play a role. For $\lambda < 0.494$ we are inside the PFM regime. Obviously, the PFM order parameter is large and the variation with λ is weak. Within the CAF regime we are again faced with the problem that the CAF LSUB m solutions for higher m terminate before meeting the corresponding PFM LSUB m solutions. Hence, the critical line $\lambda_c(J_2)$ (red vertical line in Fig. 12) terminates before meeting the transition line $J_2^{c2}(\lambda)$, and we cannot give a secure statement on the continuation of the critical line $\lambda_c(J_2)$ toward lower values of J_2 .

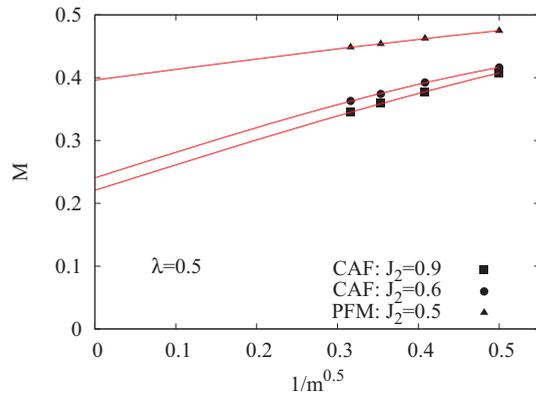


FIG. 15. (Color online) Illustration of the extrapolation (solid lines) of the CCM-LSUB m data (symbols) for the magnetic order parameter M for $\lambda = 0.5$, $J_1 = -1$, and specific values of J_2 corresponding to the CAF and PFM reference states. For the extrapolations to $m \rightarrow \infty$ according to $M(m) = b_0 + b_1(1/m)^{1/2} + b_2(1/m)^{3/2}$ we have used LSUB m data for $m = 4, 6, 8, 10$.

Finally, we illustrate the quality of the used extrapolation for the order parameter M of the “raw” LSUB m data to the limit $m \rightarrow \infty$ for $\lambda = 0.7$ in Fig. 15. It is obvious that the LSUB m data are well fitted by the applied extrapolation function.

Let us briefly discuss the relation of our results to the schematic phase diagram previously presented in Ref. 49. As stated in Ref. 49 the approximations used there (bond-operator mean-field theory as well as a second-order perturbation theory in λ) may be not reliable for large $\lambda \sim 1$ and large $J_2 \sim 1$. The main features of our phase diagram agree well with those presented in Ref. 49. For $\lambda \rightarrow 0$ we obtain the same transition point $J_2^c = 0.5$ between the QPM and the PFM phases. However, for $\lambda \sim 1$ and $J_2 \sim 1$ some differences appear. For instance, we do not find a nematic phase for $\lambda \lesssim 1$, as it was discussed in Ref. 49. Note that the absence of the nematic phase

is in agreement with the recent findings for the uniform model, that is, at $\lambda = 1$.^{44,46} Moreover, the second-order transition between the QPM and CAF phases for large J_2 is found in Ref. 49 at $\lambda \approx 0.5$, but it should be at $\lambda \approx 0.4$.

V. SUMMARY

Inspired by a recent investigation of a frustrated two-dimensional Heisenberg model proposed to describe the magnetic properties of (CuCl)LaNb₂O₇ (see Refs. 55 and 49), we investigate the GS phase diagram of the spin- $\frac{1}{2}$ J_1 - J_2 Heisenberg model on the square lattice with plaquette structure. The 4-site plaquettes carrying the (strong) intraplaquette bonds J_1 and J_2 are coupled to each other by (weaker) interplaquette bonds λJ_1 and λJ_2 , $0 \leq \lambda \leq 1$. The parameter λ can also be thought of modeling a distortion of the underlying square lattice.

We consider AFM ($J_1 > 0$) as well as FM nearest-neighbor exchange coupling ($J_1 < 0$). Except the phases with magnetic LRO (FM, Néel, and collinear striped AFM) and the nonmagnetic quantum paramagnetic phase known from the standard spin- $\frac{1}{2}$ J_1 - J_2 model we find for FM J_1 a new plaquette phase showing AFM long-range of $s = 2$ block spins associated with 4-spin plaquettes. For the AFM model ($J_1 > 0$) the region of the quantum paramagnetic phase is significantly increased for $\lambda < 1$ compared to the standard model, thus increasing the prospects of finding a magnetically disordered low-temperature phase in real magnetic materials, where the exchange pattern may deviate from the standard J_1 - J_2 model (see, e.g., Refs. 54, 72, and 73).

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