Formation process and superparamagnetic properties of (Mn,Ga)As nanocrystals in GaAs fabricated by annealing of (Ga,Mn)As layers with low Mn content

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X-ray diffraction, transmission electron microscopy, and magnetization measurements are employed to study the structural and magnetic properties of Mn-rich (Mn,Ga)As nanocrystals embedded in GaAs. These nanocomposites are obtained by moderate-temperature (400 °C) and high-temperature (560 °C and 630 °C) annealing of (Ga,Mn)As layers with Mn concentrations between 0.1% and 2%, grown by molecular beam epitaxy at 270 °C. Decomposition of (Ga,Mn)As is already observed at the lowest annealing temperature of 400 °C for layers with initial Mn content of 1% and 2%. Both cubic and hexagonal (Mn,Ga)As nanocrystals, with similar diameters of 7–10 nm, are observed to coexist in layers with an initial Mn content of 0.5% and 2% after higher-temperature annealing. Measurements of magnetization relaxation in the time span 0.1–10 000 s provide evidence for superparamagnetic properties of the (Mn,Ga)As nanocrystals, as well as for the absence of spin-glass dynamics. These findings point to weak coupling between nanocrystals even in layers with the highest nanocrystal density.

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I. INTRODUCTION

Prospects for fabricating application-ready spintronic devices by using dilute magnetic semiconductors (DMSs) are facing difficulties due to the lack of a suitable material that maintains ferromagnetic properties up to room temperature. 1,2 As a result, there is renewed interest in nanocomposite systems consisting of ferromagnetic nanocrystals embedded in a semiconductor matrix, ^{2–6} for which a number of functionalities have been predicted. 4,5 Ferromagnetic nanocrystals can be formed by crystallographic phase separation, i.e., by the precipitation of either a transition metal (TM) compound (to form a condensed magnetic semiconductor) or a TM.^{2,3} Interestingly, in many cases chemical phase separation (often referred to as spinodal decomposition⁴) occurs, forming nanoscale regions that are rich in TM cations without any difference in crystallographic structure from the surrounding semiconductor matrix and are thus difficult to detect.^{2–5} The efficiency of the decomposition may depend on the position of the Fermi level in the semiconducting host. ^{7,8} The occurrence of a decomposition can also be difficult to identify. A combination of several experimental techniques is necessary to address this question.⁹

In (Ga,Mn)As-one of the most comprehensively studied DMS materials with hole-induced ferromagnetism¹-phase separation results in the formation of Mn-rich nanocrystals that have high-temperature (HT) ferromagnetic properties (with critical temperatures in the range of 300–350 K). In this system, phase separation can be achieved in a controlled manner by the HT annealing of (Ga,Mn)As ternary alloy layers grown by molecular beam epitaxy (MBE) at low temperature (LT). LT MBE growth of chemically uniform (Ga,Mn)As layers takes place at substrate temperatures in the range 170 °C–300 °C. ^{10,11} HT postgrowth annealing, leading

to detectable phase separation, is then performed at temperatures of 400 °C–700 °C. After HT annealing, Mn-rich nanocrystals inside the GaAs matrix can be identified by transmission electron microscopy (TEM) techniques. The physical parameters of the nanocrystals, e.g., their sizes and densities, can be controlled by the annealing temperature used and by the initial Mn content. Thus, HT-annealed (Ga,Mn)As can be used as a model system for studying phase separation in TM-doped DMS materials.

In this paper, we study the properties of (Mn,Ga)As:GaAs phase-separated material produced by moderate-temperature to HT annealing of LT MBE-grown (Ga,Mn)As ternary alloy with low Mn content (in the range of 0.1%–2%). Although many reports on this system have been published, both prior to¹² and after¹³⁻²⁰ the first successful growth of (Ga,Mn)As ferromagnetic semiconductor,²¹ there are still unexplained issues concerning both the properties of the (Mn,Ga)As:GaAs granular system and the detailed mechanisms of the formation of (Mn,Ga)As nanocrystals. Even basic characteristics of the (Mn,Ga)As inclusions, e.g., their compositions, are not vet sufficiently well understood. 12-20 We use the notation (Mn,Ga)As, since both Ga and Mn atoms can occur in the volume of an individual inclusion. Moreover, correct estimation of the exact composition of individual nanocrystals is difficult due to the small nanocrystal size and the presence of surrounding GaAs matrix. It is likely that the nanoinclusions consist of MnAs with some admixture of Ga; however, their magnetic properties, e.g., Curie temperatures in the case of a nanocrystal with ferromagnetic properties and coercive field values, are rather close to those of MnAs.

To answer some of these questions, we carried out the following experiments: (1) HT annealing of a set of very

diluted (Ga,Mn)As samples that had starting Mn content of 0.1%–2%, along with subsequent investigation of their structural and magnetic properties, and (2) HT annealing of (Ga,Mn)As with a starting Mn content of 0.5% in the TEM column for *in situ* observation of the phase separation process.

In this paper, we concentrate on the first set of experiments. We investigate the structural and magnetic properties of (Ga,Mn)As samples with Mn content increasing from 0.1%, or paramagnetic (Ga,Mn)As, to 2%, or ferromagnetic (Ga,Mn)As.

To our knowledge, the results of a systematic investigation of samples produced by HT annealing of (Ga,Mn)As with such a low Mn content have not been reported previously. In the literature concerning the (Mn,Ga)As:GaAs granular system, usually the initial GaMnAs samples (subsequently used for HT annealing, leading to the formation of (Mn,Ga)As nanocrystals in GaAs matrix) had much higher Mn content than our samples described in this paper. For example Moreno et al. used HT-annealed GaMnAs with Mn content in the range of 3%–6%. 14–16 We have published elsewhere the results from HT-annealed GaMnAs samples with Mn content between 2% and 6%.¹⁷ Recently, results concerning HT annealing of Be-codoped GaMnAs with Mn content of 5%-8% have been reported. 18 To our knowledge, results from HT-annealed GaMnAs with Mn content lower than 2% have not been published so far.

There is also renewed interest in the (Mn,Ga)As:GaAs composite system due to a recently discovered property of (Mn,Ga)As inclusions embedded in semiconducting GaAs: an electromotive force can be induced by the reorientation of the magnetic moments in the (Mn,Ga)As nanocrystals by an external magnetic field.⁶

The results of our annealing experiments performed inside the TEM have been published elsewhere.²²

II. EXPERIMENTAL DETAILS

Samples studied were been deposited in a Kryovak MBE system dedicated to (Ga,Mn)As. An epitaxy-ready GaAs (100) substrate was glued using In to the Mo holders, which ensures good thermal contact and lateral temperature uniformity during MBE growth. The substrate temperature was monitored using an infrared pyrometer, was the same for all samples studied, and was chosen to be \sim 270 °C. The initial (Ga,Mn)As layers were grown with an As₂ flux generated by a valve cracker effusion cell, with an As/Ga flux ratio of ~2. The Mn content for compositions of 1% or higher was estimated from the increase in growth rate with respect to that of LT GaAs layers grown prior to (Ga,Mn)As deposition by using the period of reflection high-energy electron diffraction intensity oscillations.²³ For lower Mn compositions, this parameter was estimated by extrapolation of the temperature dependence of the Mn flux. The thickness of the (Ga,Mn)As layers was 1 μ m for Mn content of 0.1% and 0.3%, 0.7 μ m for Mn content of 0.5%, and $0.4~\mu m$ for Mn content of 1% and 2%. After MBE growth and removal from the vacuum system, the samples were each divided into four pieces for annealing at different temperatures, with one piece left in the as-grown state. The samples were annealed at three temperatures: 400 °C, 560 °C, and 630 °C, with the annealing temperature controlled by an infrared pyrometer. For each temperature, the samples were annealed simultaneously on the same Mo holder, which was mounted in the MBE growth chamber. Annealing at the highest temperature of 630 $^{\circ}$ C was carried out in the presence of an As₂ flux to prevent the surfaces of the annealed samples from degrading due to As desorption.²⁴ The annealing ran for 40 min at 400 $^{\circ}$ C and 1 h at 560 $^{\circ}$ C and 630 $^{\circ}$ C.

The lattice constant in the growth direction (the perpendicular lattice parameter) was measured using high-resolution x-ray diffraction in a Philips X'pert diffractometer. Structural characterization and chemical analysis were carried out on cross-sectional TEM specimens prepared using conventional mechanical polishing and Ar ion milling. The TEM specimens were finished at low ion energies (<1 keV) to minimize ion beam—induced sample preparation artifacts. Both image and probe aberration-corrected TEM and scanning transmission electron microscopy (STEM) studies were carried out using FEI Titan microscopes operated at 300 kV.

Temperature-dependent magnetization measurements (shown in Sec. IV) were performed using a commercial superconducting quantum interference device (SQUID) magnetometer from Quantum Design. Magnetic relaxation studies were carried out using a noncommercial low-field SQUID magnetometer.²⁵ A small superconducting magnet with a time constant of ~ 1 ms delivered small magnetic fields and was employed to measure the magnetization magnitude for 0.2-0.3 s after field switching. The initial cooling rate to reach the measurement temperature for the relaxation experiments was \sim 5 K/min. To improve the thermal contact between the thermometer and the top surface, rather than the backside of the GaAs substrate, the (Ga,Mn)As layer was glued to a sapphire rod connected to the thermometer using a silver paste.

III. STRUCTURAL PROPERTIES

The in-plane lattice parameter of the layers is identical to that of the GaAs(001) substrate, since layers are coherently strained to the GaAs substrate. Figure 1 shows the $2\theta/\omega$ X-ray diffractometer scans for 006 symmetric reflections for samples with 0.1%, 0.3%, 0.5%, 1%, and 2% Mn.

The scans are sensitive to strain perpendicular to the sample surface. The 006 reflection was chosen because the angular position of the 004 reflection from (Ga,Mn)As layers with low Mn content is too close to the 004 diffraction peak from the GaAs substrate. Figure 1 shows the angular positions of the 006 diffraction peaks of the (Ga,Mn)As layers measured before and after annealing at temperatures of 400 °C, 560 °C, and 630 °C. The intensities were normalized to that of the 006 reflection of the GaAs substrate (the highest intensity peak at an angular position 2θ of \sim 109.65°).

For the as-grown samples, the observed lattice expansion of the (Ga,Mn)As ternary alloy, with respect to GaAs, results from Mn partially replacing Ga in the GaAs host lattice. ^{21,26} However, a strong contribution from Mn atoms located at interstitial positions has also been predicted theoretically and observed experimentally. ^{26–29} After HT annealing, the difference between the lattice constant of the layer originally constituting the (Ga,Mn)As ternary alloy and the lattice constant of the GaAs substrate is due *only* to strain in the GaAs

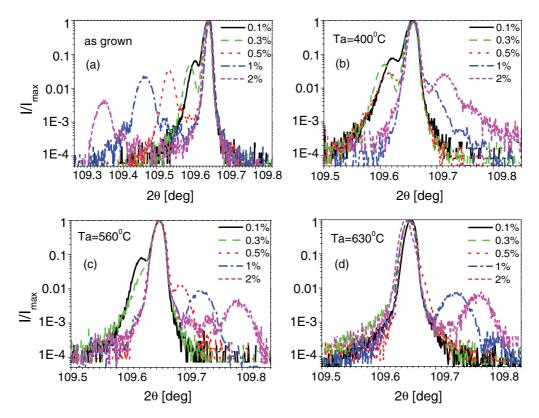


FIG. 1. (Color online) 006 reflections in $2\theta/\omega$ x-ray diffractometer scans, measured for (a) as-grown and (b) $400\,^{\circ}$ C-, (c) $560\,^{\circ}$ C-, and (d) $630\,^{\circ}$ C-annealed (Ga,Mn)As layers containing 0.1%, 0.3%, 0.5%, 1%, and 2% Mn.

host lattice caused by the presence of Mn-rich clusters. ^{15,16,29} The clusters themselves do not contribute to the diffraction peaks shown in Fig. 1.

It is known from previous reports that Mn-rich nanocrystals can exert a compressive strain on a surrounding GaAs matrix. ^{15,16,30} Since the in-plane lattice parameter of the layer is fixed to that of the GaAs substrate and the nanocrystals cause the local contraction of the surrounding GaAs crystal lattice, the formation of clusters causes a reversal in the sign of the strain from negative (compressive) to positive (tensile). This transition can be seen in Fig. 2, which shows the strain calculated from the lattice constants obtained from angular positions of diffraction peaks shown in Fig. 1. The relaxed lattice constant values used for strain evaluation were calculated from the measured (perpendicular), strained values, taking the in-plane lattice parameters of the layers as identical to that of the GaAs substrate. ²⁶ Strain is defined here as

$$\varepsilon = (a_{\rm s} - a_{\rm l})/a_{\rm l},\tag{1}$$

where a_s is the lattice constant of the GaAs substrate and a_l is the relaxed lattice constant of the layer.

Values of a_1 were calculated from the measured values of a_{\perp} , the lattice parameter of the layer in the direction perpendicular to the substrate, by using the formula

$$a_1 = (a_{\perp} + a_{||} \cdot 2b)/(1 + 2b),$$
 (2)

where

$$a_{||} = a_{s}$$
$$b = C_{12}/C_{11}$$

and C_{11} and C_{12} are the elastic constants of the layer material, which are assumed to be the same as for GaAs and take the following values:²⁶ $C_{11} = 11.82 \times 10^{10}$ Pa and $C_{12} = 5.326 \times 10^{10}$ Pa.

For (Ga,Mn)As layers containing 1% and 2% Mn, the reversal in the sign of the strain occurs after annealing at the lowest temperature ($400\,^{\circ}$ C). Thus, this relatively low

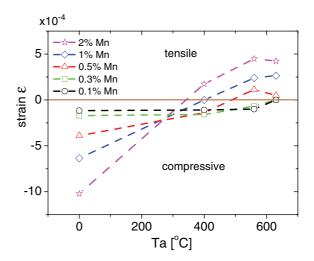


FIG. 2. (Color online) Evolution of strain in the as-grown and annealed (Ga,Mn)As layers. Annealing temperatures (T_a) were set to $400\,^{\circ}\text{C}$, $560\,^{\circ}\text{C}$, and $630\,^{\circ}\text{C}$. The dashed lines joining the experimental points are visual guides and do not necessarily follow the true values of strain in layers annealed at intermediate temperatures.

TABLE I. Parameters of individual (Mn,Ga)As nanocrystals in HT-annealed (Ga,Mn)As layers with Mn content of 2% and 0.5%, estimated from TEM studies.

	Average size (nm)		Density $(\times 10^{-6} \text{ nm}^3)$		Voids		As		Cubic (Mn,Ga)As		Hexagonal (Mn,Ga)As	
$%Mn/T_a$	560 °C	630 °C	560 °C	630 °C	560 °C	630 °C	560 °C	630 °C	560 °C	630 °C	560 °C	630 °C
0.5%	9.8	10.8	8	23	Yes	Yes	Yes	No	Yes	Yes	Yes	Yes
2%	3.9	8.8	44	27	Yes	Yes	No	No	Yes	Yes	Yes	Yes

 $^{{}^{}a}T_{a}$, annealing temperature.

annealing temperature is already thought to result in phase separation and in formation of (Mn,Ga)As clusters. For (Ga,Mn)As with a Mn content of 0.5%, strain reversal occurs at a higher annealing temperature (560 °C) and was not observed for samples with 0.1% and 0.3% Mn, presumably due to the much lower internal stress in the GaAs matrix associated with the lower density of clusters. The magnitude of the strain for 1% and 2% Mn is reduced by \sim 50% upon HT annealing (Fig. 2). For 0.1% and 0.3% Mn, the angular positions of the diffraction peaks from the tensile-strained layers, after annealing at 630 °C, may be too close to the position of the 006 peak of the GaAs substrate to be distinguishable. The apparent tendency for the temperature for strain reversal to increase with decreasing Mn content in HT-annealed (Ga,Mn)As, which can be seen in Fig. 2, may be caused by the higher thermal stability of (Ga,Mn)As with a lower Mn content or by the contribution of As antisite defects to the (Ga,Mn)As lattice constant, noticeable for Mn content lower than $\sim 1\%$.²⁸

It was demonstrated previously^{30,31} that strain in the GaAs matrix surrounding (Mn,Ga)As inclusions is larger for zincblende clusters than for hexagonal ones. Since the angular positions of the 006 diffraction peaks after annealing at 560 °C, are higher than the angular positions of the same peaks after annealing at 630 °C [Fig. 1(c) and 1(d)], we can infer that the 560 °C-annealed samples contain a greater proportion of cubic Mn-rich clusters than the samples annealed at 630 °C.

More detailed information concerning the local structure of three samples containing 0.1%, 0.5%, and 2% Mn and annealed at 400 $^{\circ}$ C, 560 $^{\circ}$ C, and 630 $^{\circ}$ C have been obtained from TEM investigations. Structural parameters of the annealed samples are summarized in Table I.

Figure 3 shows representative TEM images of (Ga,Mn)As layers with an initial Mn content of 0.5% that had been annealed at different temperatures. In the TEM image shown in Fig. 3(a), the structure of the layer annealed at 400 °C appears to be homogenous. However, nanometer-sized regions with dark contrast are present, which may indicate the start of a phase separation process. The dark areas may also have resulted from the cross-sectional sample preparation process, even though the specimen was finished at low ion energy. Figure 3(b) shows that HT annealing induces nanocrystal formation. The image was recorded slightly underfocus to enhance the contrast of the edges of the particles. Interestingly, regions of bright contrast are visible adjacent to many of the nanocrystals. Our detailed electron microscopy analysis²² reveals that these features are voids. Similar regions can be observed in bright-field images of LT-grown and HT-annealed Mn-doped GaAs layers recorded by other groups; 12,16,20

however, the voids have not been addressed. The structures of the nanocrystals were determined using nanobeam electron diffraction to be cubic (zinc blende, ZnS type) and hexagonal (NiAs type), with two types of crystal structure coexisting after annealing for 1 h both at 560 °C and at 630 °C. This situation is different from that typically observed in HT-annealed (Ga,Mn)As with a higher Mn content (5% and above), in which zinc-blende clusters are usually identified after annealing at lower temperatures (550 °C and below) and hexagonal clusters are observed after annealing at higher temperatures (600 °C and above).

In addition, we identified rhombohedral (space group 166, symbol *R-3m*) and orthorhombic (space group 64, symbol *Bmab*) As nanocrystals in annealed samples that had been doped with less than 1% Mn. Figure 3(c) shows an aberration-corrected high-resolution annular dark-field (ADF) STEM image of a void, a hexagonal (Mn,Ga)As, and a rhombohedral As nanocrystal embedded in GaAs in a Ga_{0.995}Mn_{0.005} As layer that was annealed at 560 °C. The

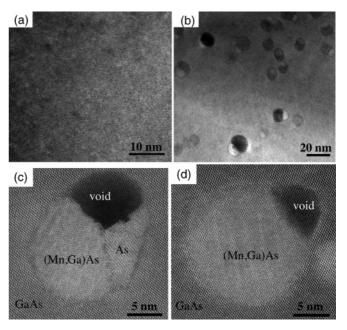


FIG. 3. Cross-sectional bright-field TEM images of $Ga_{0.995}Mn_{0.005}As$ layers annealed at (a) $400\,^{\circ}C$ and (b) $630\,^{\circ}C$. (c) Aberration-corrected ADF STEM image of a void, a hexagonal (Mn,Ga)As nanocrystal, and a rhombohedral As nanocrystal in a sample annealed at $560\,^{\circ}C$. (d) Aberration-corrected ADF STEM image of a void and a cubic (Mn,Ga)As nanocrystal in a GaAs host. The ADF inner detector semiangle used was $47.4\,$ mrad.

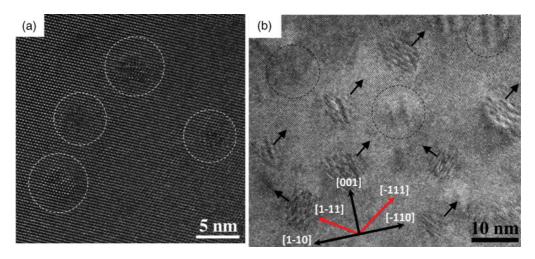


FIG. 4. (Color online) High-resolution aberration-corrected TEM images of hexagonal and cubic (Mn,Ga)As nanocrystals in $Ga_{0.98}Mn_{0.02}As$ layers annealed at (a) 560 °C and (b) 630 °C. (a) Dashed circles mark the location of the nanocrystals. (b) Dashed circles mark cubic (Mn,Ga)As nanocrystals, while arrows indicate the *c*-axes of hexagonal crystals that are parallel with {111} orientation of the GaAs host.

nanocrystal is associated with the void that exhibited dark contrast in the ADF STEM image. The As phases were identified using a combination of TEM, STEM images, and energy dispersive x-ray spectroscopy signals. Fewer than 5% of the precipitate complexes contain such As nanocrystals. The formation of As precipitate in LT-grown GaAs layers during HT annealing is well known and has been reported, e.g., in Ref. 32 and in the references therein. Our TEM results suggest that As nanocrystals could form in a low-doped (<1% Mn) (Ga,Mn)As layers as a function of the annealing temperature and Mn concentration. Interestingly, As nanocrystals were found only adjacent to hexagonal (Mn,Ga)As nanocrystals and voids. Figure 3(d) shows a void and cubic (Mn,Ga)As nanocrystal in a Ga_{0.995}Mn_{0.005}As layer that were annealed at 630 °C. Vertical Moiré fringes are visible at the position of the nanocrystal due to the lattice parameters of the cubic (Mn,Ga)As, which are different from those of GaAs. Interestingly, despite the lattice mismatch expected between the cubic nanocrystal and the GaAs host, no misfit dislocation formation was observed, as shown in Fig. 3(d).

The annealing of $Ga_{0.98}Mn_{0.02}As$ samples at $560\,^{\circ}C$ and $630\,^{\circ}C$ induces the formation of hexagonal and cubic (Mn,Ga)As nanocrystals, as shown in Fig. 4. At $560\,^{\circ}C$, the average size of the crystals is lower than in $Ga_{0.995}Mn_{0.005}As$ annealed at the same temperature, as presented in Table I. Arsenic precipitates were not observed in this sample. In the case of the hexagonal (Mn,Ga)As nanocrystals, the c-axis is parallel to one of the $\{111\}$ planes of GaAs, while the cubic nanocrystals are always coherent with the surrounding GaAs matrix.

IV. MAGNETIC PROPERTIES AND DISCUSSION

Figure 5 shows the temperature dependence of magnetization for the different (Ga,Mn)As samples, recorded for the as-grown [Fig. 5(a)] and the annealed [Figs. 5(b)–5(d)] layers. The temperature dependence of the magnetization of the as-grown samples with Mn content of 0.5% and higher is typical for that expected for thick (Ga,Mn)As layers, with a clear onset of ferromagnetism.³³ No significant magnetic signal is observed for the as-grown layers with 0.1% and

0.3% Mn. A weak magnetic signal on the order of 0.1 kA/m or below was observed for all layers after annealing at $400\,^{\circ}$ C, while some pieces that were annealed at higher temperatures exhibited larger magnetic signals.

In Fig. 5(b), the layers with the highest Mn content have similar magnetic responses, with a maximum in zero-field cooled (ZFC) magnetization around T = 15 K and a

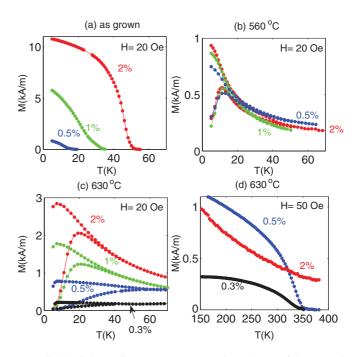


FIG. 5. (Color online) Temperature T dependence of the (a) and (d) FC and (b) and (c) ZFC/FC magnetization M of all layers measured (a) before and (b)—(d) after heat treatment. (d) Temperature dependence of the FC magnetization at higher temperatures for the layers annealed at 630 °C. No significant magnetic signal is observed for the as-grown or annealed layers with 0.1% Mn. In both the as-grown and the annealed cases, M is converted to units of kiloamps per meter by considering the initial volumes of the (Ga,Mn)As layers.

paramagnetic-like field-cooled (FC) magnetization, suggesting a superparamagnetic behavior. Similar curves are observed in Fig. 5(c) for the 1% and 2% Mn samples annealed at the highest temperature (630 °C), albeit with a maximum near 20 K. If we consider this temperature as a blocking temperature $T_{\rm b}$, we can estimate the anisotropy constant of the particles: Assuming that the Arrhenius law: $\ln(\tau/\tau_0) = KV/k_{\rm B}T$ is followed with τ of ~30 s (i.e., the time scale of a typical temperature-dependent magnetization measurement) and τ_0 of ~10⁻¹⁰ s (i.e., the typical value for nanoparticles of several nanometers; see, e.g., Ref. 34), we obtain that $KV/k_{\rm B}T_{\rm b}=26.43$. With $T_{\rm b}=20$ K and a diameter of 10 nm for the particles, we obtain $K=13\,930$ J/m³ (or 139 300 erg/cm³), which is in agreement with earlier determinations.

The LT FC magnetization of both layers is flatter than that measured for 560 °C-annealed samples, suggesting increased magnetic interaction and possible spin-glass behavior.³⁴ The FC magnetization curve of the layer with 0.5% Mn is even flatter. The layer with 0.3% Mn, which did not displayed any sizable magnetism in the as-grown state and after annealing at 560 °C, was found to exhibit a magnetic response similar to that of the layer with 0.5% Mn after annealing at 630 °C. A step decrease in magnetization at LTs (<6-7 K) can be observed in these measurements for the layers annealed at the highest temperature. Rather than the effect of magnetic (dipolar) interaction, we believe that the reduction of the magnetization is associated with the annealing-induced diffusion of the In used to fix the GaAs substrate to its holder in the MBE chamber. We have performed magnetization measurements from lower temperatures (2 K, not shown), which show that the magnetization below 6-7 K decreases and becomes negative, continuing to decrease until it becomes temperature independent. This behavior is typical for diamagnets, and we believe that it reflects the superconducting behavior associated with the diffusion of In, as discussed in Ref. 36.

In Fig. 5(d), which shows FC magnetization at higher temperatures, the layers that showed superparamagnetic-like behavior (1% and 2% Mn) do not show further features at higher temperatures, while those that showed "flat FC magnetization" (0.3% and 0.5% Mn) appear to undergo magnetic phase transitions near T = 350 K. Similar HT ferromagnetic phase transitions were previously observed in HT-annealed (Ga,Mn)As systems and were then associated with Mn-rich nanocrystals with cubic (zinc blende) structures, which have a higher Curie temperature (T_C) than the MnAs in its natural hexagonal phase.³¹ Our results suggest that HT annealing yields microscopic configurations and crystal structures of (Mn,Ga)As inclusions for layers with more than 1% Mn that are different from those in layers with less than 1% Mn. After annealing at 560 °C, the samples with 0.5%, 1%, and 2% Mn exhibit almost the same M(T) dependencies, i.e., superparamagnetic behavior with a similar T_b close to 15 K. We assume that the formation of (Mn,Ga)As precipitates during HT annealing starts with the nucleation of small cubic clusters. Upon increasing the annealing temperature, the small cubic clusters coalesce into larger ones. Then, depending on the Mn content, either the size increases further or the clusters undergo transitions from the cubic to the hexagonal phase. Since bulk, pure MnAs does not exist in a zinc-blende structure, above a certain critical size the clusters adopt only the hexagonal phase,

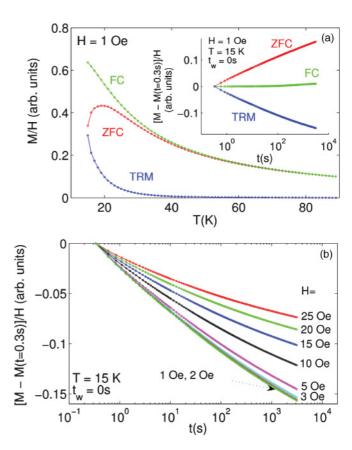


FIG. 6. (Color online) (a) Temperature (main frame) and time (inset) dependence of the ZFC, FC, and TRM magnetization of the layer with 2% Mn, plotted as M/H and [M-M(t=0.3 s)]/H, respectively. (b) TRM relaxation recorded at measurement temperature $T_{\rm m}=15$ K for different magnetic fields.

provided that they are sufficiently Mn rich. This critical size is close to 15 nm (Fig. 3). Zinc-blende (Mn,Ga)As nanocrystals with larger sizes were not observed either in our TEM images or in TEM results published by other groups. 9,15,16,31

The magnetic properties of granular system with two kinds of nanoparticles are not trivial to understand. In ferrofluids composed of single-domain magnetic nanoparticles separated by surfactants, superparamagnetism is usually observed in systems in which the nanoparticles do not interact magnetically.³⁴ If the nanoparticles interact because of their sizes or spacings, then they may instead display spin-glass behavior.³⁴ In the present study, the most concentrated systems (1% and 2% Mn) show superparamagnetic behavior up to HTs, while the layers that contain less Mn appear to exhibit HT magnetic transitions.

The dynamic (time dependent) magnetic properties of layers with 2% Mn that had been annealed at higher temperature were investigated in more detail. The results of these measurements are shown in Figs. 6 and 7.

In addition to usual temperature-dependent ZFC, FC, and thermoremanent magnetization (TRM) measurements, Fig. 6(a) shows time-dependent ZFC, FC, and TRM relaxation measurements performed at constant temperature. If smallenough magnetic fields are used, then the response of the system is linear and the system is only probed by the magnetic field. Fig. 6(b) suggests that a linear response is achieved only when the magnetic field is lower than $H \approx 3$ Oe.

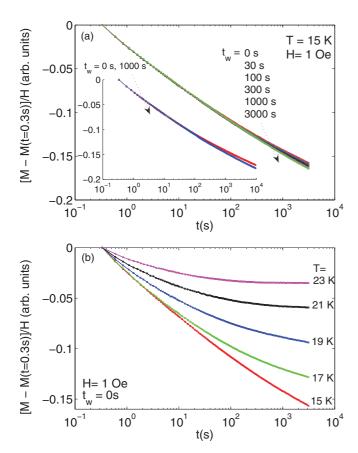


FIG. 7. (Color online) Time t dependence of the TRM of the layer with 2% Mn, plotted as $[M-M(t=0.3~{\rm s})]/H$. The sample was cooled rapidly from a reference temperature $T_{\rm ref}=120~{\rm K}$ to the measurement temperature $T_{\rm m}$ in a magnetic field $H=1~{\rm Oe}$. After a waiting time $t_{\rm w}$, the magnetic field was switched off and the magnetization was recorded as a function of time while keeping the temperature constant. (a) The TRM magnetization was recorded at $T_{\rm m}=15~{\rm K}$ for different waiting times (the inset shows examples of measurements recorded over longer time scales). (b) The TRM magnetization was recorded at different temperatures without a waiting time, i.e., for $t_{\rm w}=0~{\rm s}$.

At lower fields, M/H is approximately independent of the magnetic field. We therefore chose H=1 Oe to perform our experiments. The inset to Fig. 6(a) shows that the relaxation curves essentially obey the superposition relation: $M_{\rm ZFC} \sim M_{\rm FC} - M_{\rm TRM}$. We therefore considered the relaxation of the TRM magnetization instead of the ZFC magnetization, as is more usual in studies of, e.g., spin-glass systems, to record the magnetization in zero magnetic field and thus limit the contribution from the diamagnetic substrate.

Figure 7 shows relaxation curves of TRM, recorded after rapidly cooling the layer with 2% Mn to LT in a small magnetic field and recording the evolution of the magnetization with time in a zero magnetic field while keeping the temperature constant. For spin glasses, such relaxation curves depend on the history of the system and on how long we wait before switching off the magnetic field and recording the magnitude of the magnetization.³⁸ During such a waiting time, the magnetic configuration of the spin glass is rearranged toward its equilibrium configuration without ever reaching it; the spin

glass ages. Such an aging phenomenon is not observed in the case of a superparamagnet, whose magnetic relaxation is related mainly to thermally activated processes associated with the magnetic anisotropies of the individual particles.³⁴

Figure 7(a) shows that TRM relaxation curves recorded at LT are essentially waiting-time independent, confirming superparamagnetic behavior. Spin-glass memory experiments³⁸ were also performed and did not reveal aging, memory, or rejuvenation effects that would be typical for spin glasses. Figure 7(b) shows that superparamagnetic relaxation is maintained up to a temperature just above that at which a cusp is observed in the ZFC magnetization.

V. CONCLUSIONS

MBE-grown layers of a (Ga,Mn)As ternary alloy with Mn content of 0.1%, 0.3%, 0.5%, 1%, and 2% have been subjected to HT postgrowth annealing at temperatures of 400 °C, 560 °C, and 630 °C. Annealing of layers with such low Mn content leads to the formation of a phase-separated nanocomposite system in which Mn-rich nanocrystals are buried in the GaAs matrix, with a much lower density than that reported previously. X-ray diffraction, TEM, and SQUID magnetometry have been used to show that the phase separation process is already initiated at a temperature as low as 400 °C by the formation of nanoscale (nanometer sized) Mn-rich nanocrystals, which have a zinc-blende structure and coalesce into larger (5-15 nm) crystals at higher annealing temperatures. In HT-annealed (Ga,Mn)As with a low Mn content (0.5%), both cubic (zinc blende) and hexagonal (NiAs type) crystals coexist, even after annealing at temperatures as high as 630 °C. The maximum size of the cubic nanocrystals is limited to ~15 nm. Moreover, the annealing of (Ga,Mn)As layers with low Mn content can result in a more complex structure consisting of hexagonal (Mn,Ga)As, As nanocrystals, and voids. The minimum Mn concentration that results in detectable phase separation and in ferromagnetic or superparamagnetic properties in the HT-annealed (Ga,Mn)As is close to 0.3%. For the samples with Mn content of below 1% that contain both cubic and hexagonal crystals, besides the signature of a superparamagnetic phase, a ferromagnetic phase transition with a $T_{\rm C}$ of \sim 350 K is observed. In HT-annealed (Ga,Mn)As with a higher Mn content (1% and 2%), only superparamagnetic properties are observed. Measurements of magnetization relaxation over the time spans of 0.1–10000 s corroborate the observation of superparamagnetic behavior of the (Mn,Ga)As nanocrystals, as well as the absence of spin-glass dynamics.

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- ¹H. Ohno, Nat. Mater. **9**, 952 (2010).
- ²T. Dietl, Nat. Mater. **9**, 965 (2010).
- ³A. Bonanni and T. Dietl, Chem. Soc. Rev. **39**, 528 (2010).
- ⁴H. Katayama-Yoshida, K. Sato, T. Fukushima, M. Toyoda, H. Kizaki, V. A. Dinh, and P. H. Dederichs, Phys. Status Solidi A **204**, 15 (2007).
- ⁵T. Dietl, J. Appl. Phys. **103**, 07D111 (2008).
- ⁶D. C. Ralph, Nature **474**, E6 (2011).
- ⁷S. Kuroda, N. Nishizawa, K. Takita, M. Mitome, Y. Bando, K. Osuch, and T. Dietl, Nat. Mater. **6**, 440 (2007).
- ⁸A. Bonanni, A. Navarro-Quezada, T. Li, M. Wegscheider, Z. Matej, V. Holy, R. T. Lechner, G. Bauer, M. Rovezzi, F. D'Acapito, M. Kiecana, M. Sawicki, and T. Dietl, Phys. Rev. Lett. 101, 135502 (2008).
- ⁹M. Kaminska, A. Twardowski, and D. Wasik, J. Mater. Sci. Mater. Electron. 19, 828 (2008).
- ¹⁰R. P. Campion, K. W. Edmonds, L. X. Zhao, K. Y. Wang, C. T. Foxon, B. L. Gallagher, and C. R. Staddon, J. Cryst. Growth 247, 42 (2003).
- ¹¹D. Chiba, Y. Nishitani, F. Matsukura, and H. Ohno, Appl. Phys. Lett. **90**, 122503 (2007).
- ¹²J. De Boeck, R. Oesterholt, A. Van Esch, H. Bender, C. Bruynseraede, C. Van Hoof, and G. Borghs, Appl. Phys. Lett. 68, 2744 (1996).
- ¹³H. Akinaga, S. Miyanishi, K. Tanaka, W. Van Roy, and K. Onodera, Appl. Phys. Lett. **76**, 97 (2000).
- ¹⁴M. Moreno, A. Trampert, B. Jenichen, L. Däweritz, and K. H. Ploog, J. Appl. Phys. **92**, 4672 (2002).
- ¹⁵M. Moreno, B. Jenichen, L. Däweritz, and K. H. Ploog, J. Vac. Sci. Technol. B 23, 1700 (2005).
- ¹⁶M. Moreno, V. M. Kaganer, B. Jenichen, A. Trampert, L. Däweritz, and K. H. Ploog, Phys. Rev. B **72**, 115206 (2005).
- ¹⁷ A. Kwiatkowski, D. Wasik, M. Kamińska, R. Bożek, J. Szczytko, A. Twardowski, J. Borysiuk, J. Sadowski, and J. Gosk, J. Appl. Phys. 101, 113912 (2007).
- ¹⁸D. W. Rench, P. Schiffer, and N. Samarth, Phys. Rev. B **84**, 094434 (2011).
- ¹⁹K. Lawniczak-Jablonska, J. Libera, A. Wolska, M. T. Klepka, P. Dłużewski, J. Sadowski, D. Wasik, A. Twardowski, A. Kwiatkowski, and K. Sato, Phys. Status Solidi RRL 5, 62 (2011).
- ²⁰K. Lawniczak-Jablonska, J. Bak-Misiuk, E. Dynowska, P. Romanowski, J. Z. Domagala, J. Libera, A. Wolska, M. T. Klepka, P. Dluzewski, J. Sadowski, A. Barcz, D. Wasik, A. Twardowski, and A. Kwiatkowski, J. Solid State Chem. 184, 153 (2011).

- ²¹H. Ohno, A. Shen, F. Matsukura, A. Oiwa, A. Endo, S. Katsumoto, and Y. Iye, Appl. Phys. Lett. 69, 363 (1996).
- ²²A. Kovács, J. Sadowski, T. Kasama, J. Z. Domagala, R. Mathieu, T. Dietl, and R. E. Dunin-Borkowski, J. Appl. Phys. **109**, 083546 (2011)
- ²³ J. Sadowski, J. Z. Domagala, J. Bak-Misiuk, S. Kolesnik, M. Sawicki, K. Swiatek, J. Kanski, L. Ilver, and V. Ström, J. Vac. Sci. Technol. B 18, 1697 (2000).
- ²⁴Z. Y. Zhou, C. X. Zheng, W. X. Tang, D. E. Jesson, and J. Tersoff, Appl. Phys. Lett. **97**, 121912 (2010).
- ²⁵J. Magnusson, C. Djurberg, P. Granberg, and P. Nordblad, Rev. Sci. Instrum. 68, 3761 (1997).
- ²⁶I. Kuryliszyn-Kudelska, J. Z. Domagala, T. Wojtowicz, X. Liu, E. Łusakowska, W. Dobrowolski, and J. K. Furdyna, J. Appl. Phys. 95, 603 (2004).
- ²⁷J. Masek, J. Kudrnovsky, and F. Maca, Phys. Rev. B 67, 153203 (2003).
- ²⁸J. Sadowski and J. Z. Domagala, Phys. Rev. B **69**, 075206 (2004).
- ²⁹L. X. Zhao, R. P. Campion, P. F. Fewster, R. W. Martin, B. Ya Ber, A. P. Kovarsky, C. R. Staddon, K. Y. Wang, K. W. Edmonds, C. T. Foxon, and B. L. Gallagher, Semicond. Sci. Technol. 20, 369 (2005).
- ³⁰J. Bak-Misiuk, J. Z. Domagala, P. Romanowski, E. Dynowska, E. Lusakowska, A. Misiuk, W. Paszkowicz, J. Sadowski, A. Barcz, and W. Caliebe, Radiat. Phys. Chem. 78, S116 (2009).
- ³¹M. Yokoyama, H. Yamaguchi, T. Ogawa, and M. Tanaka, J. Appl. Phys. **97**, 10D317 (2005).
- ³²M. Toufella, P. Puech, R. Carles, E. Bedel, C. Fontaine, A. Claverie, and G. Benassayag, J. Appl. Phys. 85, 2929 (1999).
- ³³J. Sadowski, R. Mathieu, P. Svedlindh, J. Z. Domagala, J. Bak Misiuk, K. Swiatek, M. Karlsteen, J. Kanski, L. Ilver, H. Åsklund, and U. Södervall, Appl. Phys. Lett. 78, 3271 (2001).
- ³⁴P. E. Jönsson, Adv. Chem. Phys. **128**, 191 (2004).
- ³⁵R. S. DiPietro, H. G. Johnson, S. P. Bennett, T. J. Nummy, L. H. Lewis, and D. Heiman, Appl. Phys. Lett. 96, 222506 (2010).
- ³⁶Y. K. Li, Y. Huang, Z. Fan, C. Jiang, X. B. Mei, B. Yin, J. M. Zhou, J. C. Mao, J. S. Fu, and E. Wu, J. Appl. Phys. **71**, 2018 (1992).
- ³⁷P. Nordblad, L. Lundgren, and L. Sandlund, J. Magn. Magn. Mater. 54–57, 185 (1986).
- ³⁸P. E. Jönsson, R. Mathieu, P. Nordblad, H. Yoshino, H. Aruga Katori, and A. Ito, Phys. Rev. B 70, 174402 (2004).