

## Observation of room-temperature high-energy resonant excitonic effects in graphene

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Using a combination of ultraviolet-vacuum ultraviolet reflectivity and spectroscopic ellipsometry, we observe a resonant exciton at an unusually high energy of 6.3 eV in epitaxial graphene. Surprisingly, the resonant exciton occurs at room temperature and for a very large number of graphene layers  $N \approx 75$ , thus suggesting a poor screening in graphene. The optical conductivity ( $\sigma_1$ ) of a resonant exciton scales linearly with the number of graphene layers (up to at least 8 layers), implying the quantum character of electrons in graphene. Furthermore, a prominent excitation at 5.4 eV, which is a mixture of interband transitions from  $\pi$  to  $\pi^*$  at the M point and a  $\pi$  plasmonic excitation, is observed. In contrast, for graphite the resonant exciton is not observable but strong interband transitions are seen instead. Supported by theoretical calculations, for  $N \leq 28$  the  $\sigma_1$  is dominated by the resonant exciton, while for  $N > 28$  it is a mixture between excitonic and interband transitions. The latter is characteristic for graphite, indicating a crossover in the electronic structure. Our study shows that important elementary excitations in graphene occur at high binding energies and elucidate the differences in the way electrons interact in graphene and graphite.

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Graphene, a one-layer-thick carbon honeycomb structure, has recently attracted a lot of attention due its exotic quantum properties.<sup>1</sup> In addition to the well-known electronic properties, such as ballistic electron transport,<sup>1-3</sup> quantum Hall effect,<sup>4</sup> tunable band gap,<sup>5,6</sup> and physics driven by many-body interactions,<sup>7</sup> graphene also displays very interesting optical properties. For instance, due to its low-energy excitation graphene becomes highly transparent in the visible spectral range and its infrared optical conductivity can be tuned using a gate voltage.<sup>8</sup> These particular properties of graphene potentially result in exciting optoelectronic applications.

Recent theoretical studies based on the *ab initio* GW and Bethe-Salpeter equation (BSE) approach by Yang *et al.*<sup>9</sup> and Trevisanutto *et al.*<sup>10</sup> have predicted the existence and have highlighted the importance of resonant excitonic effects in the optical absorption of graphene. However, there is disagreement as to the origin and position of the exciton. In Ref. 9 the calculations were done up to 7 eV and the resonant exciton was predicted to occur at 4.6 eV due to electron-hole interaction in the  $\pi^*$  and  $\pi$  bands at the M point. In Ref. 10, however, the calculations were done at a much higher energy of 22 eV and the resonant exciton was predicted to appear at 8.3 eV due to the background single-particle continuum of dipole forbidden transition at the  $\Gamma$  point. Despite their disagreement, both have agreed that resonant excitons play important role in elementary excitations in graphene and thus its understanding is crucial.

A direct way to probe a resonant exciton in graphene is to measure its complex dielectric response in wide energy

range.<sup>10</sup> Despite reports on the optical properties of graphene over the visible energy range<sup>7,11-16</sup> there has been no reliable optical conductivity data at higher energies. Therefore, it is crucial to study the optical properties of graphene in an unprecedented high-energy range, not only to provide a direct evidence of the resonant excitonic effects *per se*, but also for our understanding of electronic structure of graphene.

Herewith, we report the optical conductivity ( $\sigma_1$ ) of epitaxial graphene on a 6H-SiC(0001)/buffer layer substrate in an unprecedented wide photon energy range from 0.5 to 35 eV using a combination of spectroscopic ellipsometry and ultraviolet-vacuum ultraviolet (UV-VUV) reflectometry.<sup>17</sup> It has been recently shown that the combination of these two techniques enables a stabilized Kramers-Kronig transformation,<sup>17</sup> which is crucial to resolve precisely the dielectric function. We study the evolution of  $\sigma_1$  as a function of the number of layers  $N$  (henceforth  $\sigma_{1,N}$ ) including graphite.

Epitaxial graphene samples on 6H-SiC(0001) were prepared by a Si desorption process in Si flux as described in previous studies.<sup>18,19</sup> This process was found to lead to multilayer graphene with ordered stacking.<sup>19</sup> The thickness was determined using scanning tunneling microscopy and Raman spectroscopy.<sup>19,20</sup> For reference, we have used a cleaved highly ordered pyrolytic graphite. The reproducibility was checked by fabricating and repeating individual measurements on three different samples that were grown under identical conditions, with at least three different locations on each sample yielding reproducible results.<sup>21</sup>

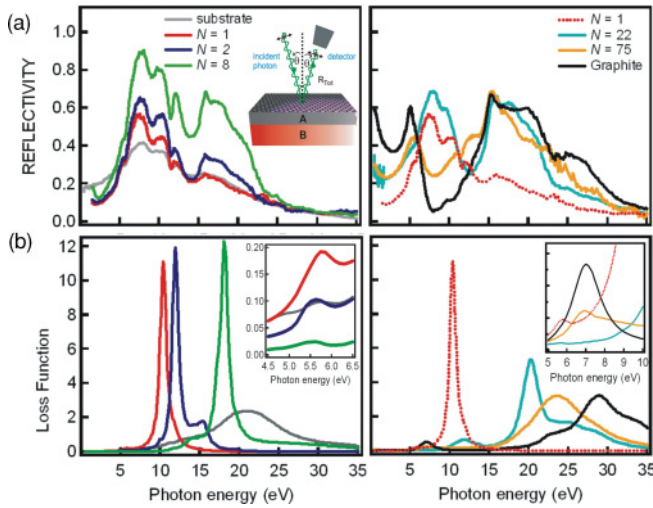


FIG. 1. (Color online) Room-temperature experimental results of (a) reflectivity and (b) loss function,  $\text{Im}(\epsilon^{-1})$ . The inset of (a) shows the experimental geometry while the insets of (b) shows  $(\epsilon^{-1})$  on an expanded scale from 4.5 to 6.5 eV and from 5 to 10 eV.

Figure 1(a) shows room-temperature reflectivity data for  $N$ -dependent epitaxial graphene, graphite, and the substrate. Interestingly, one can see that the reflectivity of epitaxial graphene is very much dependent on  $N$  and very much distinguished from that of graphite. The reflectivity has rich and distinct structures especially in the range from 5 to 8 eV for  $N = 1, 2$ , and 8 while additional structures occur around 14 to 20 eV for  $N = 22$  and 75. On the other hand, the reflectivity of graphite has shown strong structures below 0.5 eV and  $\sim 5$  eV and around 14 to 20 eV which are similar to published result.<sup>22</sup> Further rigorous discussion will be achieved by analyzing the  $\sigma_1$ .<sup>23</sup> The  $\sigma_{1,N}$  shows striking results [Figs. 2(a) and 2(b)]. For  $N = 1$ , the  $\sigma_{1,1}$  is dominated by two well-defined peaks, very pronounced at 6.3 eV (peak B) and less pronounced at 5.4 eV (peak A), with almost equal full width at half maximum (FWHM) of  $\sim 0.8$  eV. These structures, especially peak B, are considerably sharp for this high-energy range.

Now we focus our discussion on peak B as it shows an interesting dependence on  $N$  [Fig. 2(c)]. The  $\sigma_{1,N}(B)$  increases rapidly toward  $N = 8$ , while its FWHM and peak position remain independent of  $N$ . For  $N > 8$ ,  $\sigma_{1,N}(B)$  decreases while the peak position shifts toward higher energy (6.7 and 11 eV for  $N = 22$  and  $N = 75$ , respectively). For graphite, peak B disappears and our result is similar to the published data.<sup>22</sup>

To find out the origin of peak B, we have directly compared our experimental data with theoretical calculations.<sup>9,10</sup> As shown in inset 1 of Fig. 2(a), the line shape and the  $\sigma_{1,1}(B)$  between experimental data and theoretical calculations from Ref. 10 are surprisingly very similar. This comparison shows decisive evidence of the high-energy resonant exciton in epitaxial graphene. The resonant exciton arises from dipole transitions of the single-particle continuum. However, the observed resonant exciton peak occurs *redshifted* compared to the calculations. Our calculations show that the redshifted exciton could result from film-substrate interactions, which are not included in the previous calculations (see below for the details). Thus, the origin of peak B is the high-energy resonant exciton as predicted.

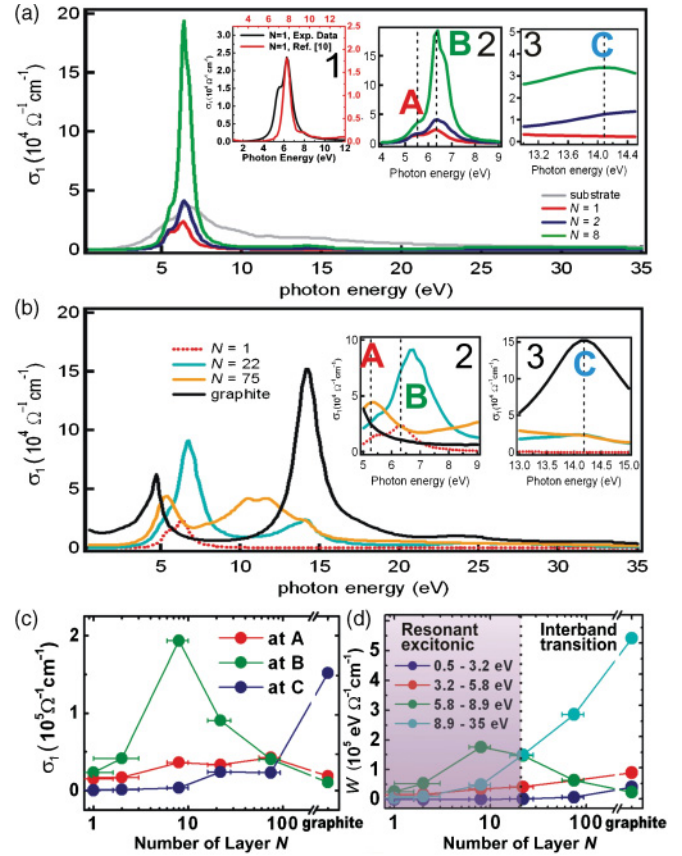


FIG. 2. (Color online) The optical conductivity ( $\sigma_1$ ) of (a) substrate and graphene ( $N = 1, 2, 8$ ) and (b) graphene ( $N = 22, 75$ ) and graphite shows three peaks at 5.4 eV (label A), 6.3 eV (label B), and 14.1 eV (label C). Inset 1 shows the comparison of  $\sigma_{1,1}$  between experimental data and theoretical calculations. Insets 2 and 3 show the  $\sigma_1$  on an expanded scale at various energy ranges. (c) The value of  $\sigma_1$  at A, B, and C as a function of  $N$ . (d) The partial spectral weight ( $W$ ) as a function of  $N$ .

Furthermore, our detail study rules out the interband transitions and plasmonic excitations as the origin of peak B. Based on density functional theory (DFT), we have calculated optical conductivity for  $N = 1$  and 2 and found that peak B does not originate from interband transitions [see Fig. 3(a) and discussion below]. Second, one may argue that we have to consider contributions from a plasmonic excitation. To address this issue, we have studied in detail an energy-loss function  $\text{Im}(\epsilon^{-1})$  which can reveal collective excitations such as plasmonic excitations.<sup>24,25</sup> As shown in Fig. 1(b), there is no feature that can be attributed to a plasmonic excitation. For  $N = 1$ ,  $\text{Im}(\epsilon^{-1})$  is dominated by a strong structure at  $\sim 10$  eV and a weak structure at 5.5 eV. These structures are from  $\pi + \sigma$  and  $\pi$  in-plane plasmon modes, respectively.<sup>25</sup> Based on the symmetry, the 5.5-eV plasmon peak structure is only visible for light polarization parallel to the  $c$  axis ( $\mathbf{E} \parallel c$ ). In our measurement, the  $\mathbf{E}$  was mixed between  $\mathbf{E} \parallel c$  and  $\mathbf{E} \perp c$  with the main contribution from  $\mathbf{E} \perp c$ ; thus, the observed  $\pi$  plasmon is very weak. On the other hand, the  $\pi + \sigma$  plasmon mode is very strong in our geometry, which is consistent with electron-energy-loss spectroscopy measurements.<sup>25,26</sup> As  $N$  increases, the  $\pi + \sigma$  plasmon mode shifts toward higher

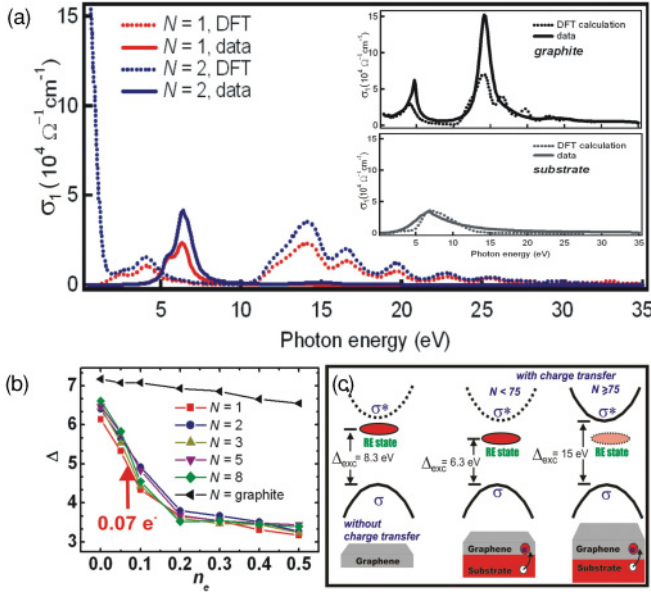


FIG. 3. (Color online) (a) Comparison between experimental data and calculated  $\sigma_1$ . (Top inset)  $\sigma_1$  spectra of graphite. (Bottom inset)  $\sigma_1$  spectra of the substrate. (b) The background single-particle continuum transitions ( $\Delta$ ) at the  $\Gamma$  point as a function of nominal charge transfer ( $n_e$ ) from substrate to graphene for various graphene layer ( $N$ ). The red arrow shows the observed charge transfer in our sample. (c) A proposed model of optical absorption of exciton in graphene and graphene on substrate as function of  $N$ . RE stands for the resonant exciton and ( $\Delta_{\text{exc}}$ ) is excitation energy.

energy and gets broader, while the structure at 5.5 eV is nearly  $N$  dependent. This *blueshift* may be due to strong effects of the interlayer Coulomb interaction on the total plasmon.<sup>24</sup>

The exciton energy of 6.3 eV suggests an untypically high binding energy which can also be made responsible for the presence of the exciton up to room temperature. If one considers the angular-resolved photoemission data, the distance of the  $\sigma$  band at the  $\Gamma$  point to the Fermi level is around 4.5 eV.<sup>27</sup> By assuming the distance of the  $\sigma^*$  band to the Fermi level to be similar to that of the  $\sigma$  band one will end up with 9 eV as the separation between the  $\sigma$  and  $\sigma^*$  bands. Based on these values the binding energy of the resonant exciton is 2.7 eV, which is 108 times larger than the thermal energy. This binding energy is significantly larger compared to that of ZnO, which is around 60 meV (2.4 times thermal energy)<sup>28</sup> and is very clearly seen at room temperature while the binding energy of the excitation for GaAs is about 3.4 meV and can only be seen at low temperatures.<sup>29</sup> Such a very high-binding-energy exciton observed on graphene is indeed remarkable and a direct consequence of its low dimensionality in conjunction with its intriguing electronic properties.

We next discuss peak A and its origin. As shown in Fig. 2(c),  $\sigma_{1,N}(A)$  increases monotonically as a function of  $N$ . Interestingly, for  $N > 8$  as well as for graphite the position of peak A shifts toward lower energies, as opposed to peak B, while the line shape is rather symmetric and the FWHM is broader for larger  $N$ . Based on electronic band structure calculations,<sup>14</sup> peak A is the result of interband transitions from the  $\pi$  to the  $\pi^*$  bands at the M point, where the van Hove singularity occurs. Furthermore, from analysis of energy-loss

function [Fig. 1(b)] the  $\text{Im}(\epsilon^{-1})$  shows a weak  $\pi$ -plasmon contribution. Thus, one can conclude that peak A is a mixture of interband transitions and a plasmonic excitation.

Another notable observation is a broad structure at 14.1 eV (peak C). The peak C is absent for  $N = 1$  and  $\sigma_{1,N}(C)$  increases dramatically for  $N \geq 2$  [Fig. 2(c)]. Interestingly, for graphite the  $\sigma_1$  is dominated by peak C with symmetric lineshape. Based on electronic band structure calculations,<sup>10,14,30</sup> this structure is interband transitions from  $\sigma$  to  $\sigma^*$  bands at the  $\Gamma$  point. Peak C clearly has an intimate relationship with optical conductivity in graphite while peak B is a unique characteristic of graphene.

We next discussed partial spectral weight integral ( $W$ ) because it describes the effective number of electrons excited by photons of respective energy. The  $\sigma_1$  is restricted by the  $f$ -sum rule:  $\int_0^\infty \sigma_1(E) dE = \frac{\pi n e^2}{2m^*}$ , where  $n$  is the electron density,  $e$  is elementary charge and  $m^*$  is effective electron mass. Hence, one can extract the  $W \equiv \int_{E_1}^{E_2} \sigma_1(E) dE$  for various energy ranges. Due to the  $f$ -sum rule and charge conservation, the  $W$  is constant, and thus one can study the spectral weight transfer and reveal interactions as well as the effect of  $N$  in the broad energy range of 0.5 to 35 eV. Figure 2(d) shows  $W$  for the different energy regions: 0.5 to 3.2 eV (region I,  $W_I$ ), 3.2 to 5.8 eV (region II,  $W_{II}$ ), 5.8 to 8.9 eV (region III,  $W_{III}$ ), and 8.9 to 35.0 eV (region IV,  $W_{IV}$ ).

The  $W_I$  is mainly governed by the transition around the Dirac cone, that is,  $\pi$  to  $\pi^*$  around the  $K$  point in the Brillouin zone, which is consistent with previous publications.<sup>7,13</sup> The  $W_{II}$ ,  $W_{III}$ , and  $W_{IV}$  show the main contribution from peaks A, B, and C, respectively. Interestingly,  $W_I$  and  $W_{II}$  show an almost similar trend in which the  $W$  increases monotonically as function of  $N$ , while the  $W_{III}$  and  $W_{IV}$  show completely different fashion. For  $N \leq 8$  the  $W_{III}$  increases as  $N$  increases, while for  $N > 8$   $W_{III}$  decreases as  $N$  increases. On the other hand,  $W_{IV}$  increases rapidly for  $N \geq 8$ . In graphite,  $W_{IV}$  is maximum while  $W_{III}$  obtains its smallest value. We find a crossover between  $W_{III}$  and  $W_{IV}$  at  $N \sim 28$  [see Fig. 2(d)]. Thus, we propose that for  $N < 28$  the optical conductivity of epitaxial multilayer graphene is dominated by high-energy resonant excitonic effects, while for  $N > 28$ , the optical conductivity is dominated by interband transitions.

To gain further insight, we have calculated the optical conductivity ( $\sigma_{1,\text{calc}}$ ) of graphene, graphite, and substrate using DFT. We have studied the role of interband transitions and have compared them with experimental results [see Fig. 3(a)]. Noting that our calculations do not include electron-hole (e-h) interactions, ( $\sigma_{1,\text{calc}}$ ) is mainly driven by interband transitions and serves as a reference for the noncorrelated case.<sup>31</sup> Based on the DFT calculation, ( $\sigma_{1,\text{calc}}$ ) mimics reasonable well the ( $\sigma_1$ ) for graphite as well as for the 6H-SiC(0001)/buffer layer. For graphite, the peak at 4 eV is dominated by transitions from  $\pi \rightarrow \pi^*$  bands while the peak at 14.1 eV is dominated by transitions from  $\sigma \rightarrow \sigma^*$  bands. This is consistent with previous theoretical study.<sup>30</sup>

On the other hand, the DFT calculations for graphene show that its optical conductivity cannot be explained with interband transitions. As shown in Fig. 3(a), the calculated optical conductivity for graphene shows a completely different result than the experimental data and thus our attempt to mimic  $\sigma_1$  for  $N = 1$  and 2 failed. Calculations based on DFT show that the  $\sigma_{1,\text{calc}}$  for graphene ( $N = 1, 2$ ) above 3 eV is



very similar to that seen with graphite. This is in contrast to our experimental results. This further supports that the  $\sigma_{1,N}$  is mainly driven by strong e-h interactions which lead to a high-energy resonant exciton and form an unusual electronic band structure while the optical conductivity of graphite is mainly driven by interband transitions. It is shown that for  $N = 1$ , once one turns on the e-h interactions, the peak around 4 eV and the peak around 14.1 eV [inset of Fig. 3(a)] vanish, resulting in a new and very strong peak in between.<sup>10</sup>

To find out the origin of the redshift of the excitonic excitation, we have calculated the variation of the  $\sigma$  bands and the background single-particle continuum transitions ( $\Delta$ ) as functions of charge transfer ( $n_e$ ) and strain. We find out that while the  $\sigma$  bands are nearly independent from these two effects, the background single-particle continuum transitions very much depend on the charge transfer  $n_e$  from substrate to the graphene layer [Fig. 3(b)]. Based on our DFT calculations, the charge transfer from substrate to graphene  $n_e$  is about  $\sim 0.07e$  (per graphene unit) and thus the background single-particle continuum transitions reduce by  $\sim 1.4$  eV compared to that of for  $n_e = 0$ . The strain in graphene due to the lattice mismatch with the substrate reduces the background single-particle continuum transitions by 0.6 eV. These two effects altogether decrease the single-particle continuum transitions to about 2.0 eV in total. Thus, one may expect to see the exciton at a lower energy of  $\sim 6.3$  eV. Our result may suggest that energy excitation of the exciton depends on the  $n_e$  and the strain.

Interestingly, our calculation shows that  $\Delta$  does not depend on  $N$  significantly [Fig. 3(b)]. This is, in fact, consistent with our experimental result. In contrast, the theoretical predictions in Ref. 10 show an energy shift between the exciton in graphene and bilayer graphene which is attributed to enhanced screening. Our experimental results show, however, that the position of the exciton does not depend on the  $N$  value until about  $N = 22$ . This can be reconciled with Ref. 10 only if one assumes that screening effects are significantly reduced. Thus, by comparing the experimental results and theoretical calculations, one may conclude that screening effects in graphene are much weaker than one would expect.

Finally, based on our experimental results of high-energy optical conductivity, DFT calculations, and recent *ab initio* GW-BSE calculations,<sup>10</sup> we have proposed the following phase diagram [Fig. 2(c)–2(d)] and a model for optical absorption which is proportional to  $\sigma_1$  [Fig. 3(c)]. For free-standing graphene of  $N = 1$ , the optical absorption is dominated by high-energy resonant exciton effects which occur around 8.3 eV. However, in our case due to the charge transfer and lattice mismatch between graphene and the substrate, the resonant exciton reduces to 6.3 eV. For  $N < 28$ , the optical absorption shows very strong resonant excitoniclike structure and weak interband transitions. For  $28 < N < 75$  the optical absorption shows a mixture between exciton and interband transitions [Figs. 2(d) and 3(c)]. For  $N > 75$  as well as for graphite, the optical absorption is dominated by interband transitions.

In summary, we have observed a high-energy room-temperature stable resonant exciton at  $\sim 6.3$  eV with a large binding energy of  $\sim 2.7$  eV in the optical conductivity of multilayer epitaxial graphene revealing strong collective e-h interactions. The resonant exciton persists for very large  $N$  and thus dominates the electronic properties owing to the poor screening in graphene. Furthermore, the mixture of interband transition from  $\pi$  to  $\pi^*$  at the M point and weak  $\pi$  plasmonic excitation gives rise to peak in the optical conductivity at 5.4 eV. These findings demonstrate the importance of high-energy optical conductivity and have strong implications on the understanding of the electronic structure of epitaxial graphene.

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<sup>1</sup>A. H. Castro Neto *et al.*, *Rev. Mod. Phys.* **81**, 109 (2009).

<sup>2</sup>K. S. Novoselov *et al.*, *Science* **306**, 666 (2004).

<sup>3</sup>K. S. Novoselov *et al.*, *Nature (London)* **438**, 197 (2005).

<sup>4</sup>Y. Zhang *et al.*, *Nature (London)* **438**, 201 (2005).

<sup>5</sup>T. Ohta *et al.*, *Science* **313**, 951 (2006).

<sup>6</sup>Y. Zhang *et al.*, *Nature (London)* **459**, 820 (2009).

<sup>7</sup>A. Bostwick *et al.*, *Nat. Phys.* **3**, 36 (2007).

<sup>8</sup>Z. Li *et al.*, *Nat. Phys.* **4**, 532 (2008).

<sup>9</sup>L. Yang *et al.*, *Phys. Rev. Lett.* **103**, 186802 (2009).

<sup>10</sup>P. E. Trevisanutto *et al.*, *Phys. Rev. B* **81**, 121405(R) (2010).

<sup>11</sup>R. R. Nair *et al.*, *Science* **320**, 1308 (2008).

<sup>12</sup>A. B. Kuzmenko *et al.*, *Phys. Rev. Lett.* **100**, 117401 (2008).

<sup>13</sup>K. F. Mak *et al.*, *Phys. Rev. Lett.* **101**, 196405 (2008).

<sup>14</sup>S. Latil and L. Henrard, *Phys. Rev. Lett.* **97**, 036803 (2006).

<sup>15</sup>T. Ohta *et al.*, *Phys. Rev. Lett.* **98**, 206802 (2007).

<sup>16</sup>V. G. Kravets *et al.*, *Phys. Rev. B* **81**, 155413 (2010).

<sup>17</sup>A. Rusydi *et al.*, *Phys. Rev. B* **78**, 125110 (2008).

<sup>18</sup>W. Chen *et al.*, *Surface Science* **596**, 176 (2005).

<sup>19</sup>H. Huang *et al.*, *ACS Nano* **2**, 2513 (2008).

<sup>20</sup>S. Shivaraman *et al.*, *J. Electron. Mater.* **38**, 725 (2009).

<sup>21</sup>See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.84.081403> for details of the samples.

<sup>22</sup>E. A. Taft and H. R. Philipp, *Phys. Rev. A* **138**, 197 (1965).

<sup>23</sup>See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.84.081403> for details of the analysis of reflectivity, complex dielectric function, and optical conductivity.

<sup>24</sup>A. G. Marinopolous *et al.*, *Phys. Rev. Lett.* **89**, 076402 (2002).

<sup>25</sup>T. Eberlein *et al.*, *Phys. Rev. B* **77**, 233406 (2008).

<sup>26</sup>J. Lu *et al.*, *Phys. Rev. B* **80**, 113410 (2009).

<sup>27</sup>K. V. Emtsev *et al.*, *Phys. Rev. B* **77**, 155303 (2008).

<sup>28</sup>D. G. Thomas, *J. Phys. Chem. Solids* **15**, 86 (1960).

<sup>29</sup>J. Hegarty and M. D. Sturge, *J. Opt. Soc. Am. B* **2**, 1143 (1985).

<sup>30</sup>R. Ahuja *et al.*, *Phys. Rev. B* **55**, 4999 (1997).

<sup>31</sup>See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.84.081403> for the details of the calculation.