Unusual Nernst effect and spin density wave precursors in superconducting LaFeAs $O_{1-x}F_x$

Agnieszka Kondrat, Günter Behr, Bernd Büchner, and Christian Hess*

Leibniz-Institute for Solid State and Materials Research, IFW-Dresden, D-01171 Dresden, Germany (Received 3 June 2010; revised manuscript received 8 February 2011; published 28 March 2011)

The Nernst effect has recently been proven as a sensitive probe for detecting unusual normal state properties of unconventional superconductors. Here we present a systematic study of the Nernst effect of the iron pnictide superconductor LaFeAsO_{1-x}F_x with a particular focus on its evolution upon doping. For the parent compound we observe a huge negative Nernst coefficient accompanied with a severe violation of the Sondheimer cancellation in the spin density wave (SDW) ordered state. Surprisingly, an unusual and enhanced Nernst signal is also found at underdoping (x = 0.05) despite the presence of bulk superconductivity and the absence of static magnetic order, strongly suggestive of SDW precursors at $T \lesssim 150$ K. A more conventional normal state Nernst response is observed at optimal doping (x = 0.1) where it is rather featureless with a more complete Sondheimer cancellation.

DOI: 10.1103/PhysRevB.83.092507 PACS number(s): 74.70.Xa, 72.15.Jf

I. INTRODUCTION

The Nernst effect of unconventional superconductors has recently attracted a lot of attention for several reasons. 1-12 For type-II superconductors it is strongly enhanced by movement of magnetic flux lines (vortices). 13-16 Here the Nernst coefficient $\bar{\nu}$ is directly proportional to the drift velocity of the vortices, which has rendered this transport quantity a valuable tool for studying their dynamics. Based on this very fundamental property, the unusual enhancement of the Nernst coefficient in the normal state of cuprate high- T_c superconductors at temperatures much higher than the critical temperature T_c has been interpreted as the signature of vortex fluctuations. 1-3 More specifically, it was proposed that in the pseudogap phase above T_c long-range phase coherence of the superconducting order parameter is lost while the pair amplitude remains finite. As an alternative explanation, it has only recently been proposed that Fermi-surface distortions due to stripe or spin density wave (SDW) order could lead to an enhanced Nernst effect in the cuprates. 4-6,12 Very recently, a strong anisotropy of the Nernst coefficient arising from the broken rotation symmetry of electron-nematic order has been discussed both experimentally and theoretically.^{8,9}

The discovery of superconductivity in LaFeAsO_{1-x} F_x (Ref. 17) initiated a tremendous research effort which yielded soon after a large variety of new superconducting iron pnictide compounds with T_c up to 55 K.¹⁸ The parent compound LaFeAsO is a poor metal and exhibits an antiferromagnetic SDW ground state. The transition toward the SDW state occurs at $T_N = 137$ K and is accompanied by a structural tetragonal-to-orthorhombic transition at $T_s \approx 160 \text{ K.}^{19-21}$ Upon substituting fluorine for oxygen the SDW phase is destabilized, i.e., T_s and T_N gradually decrease and at some finite doping level ($x \leq 0.05$) superconductivity emerges. The actual nature of the doping-driven transition from SDW to superconductivity is much under debate. There is evidence that in LaFeAsO_{1-x}F_x the transition is abrupt and first-order-like toward a homogeneous superconducting state²² while in other systems (e.g., $SmFeAsO_{1-x}F_x$ or $BaFe_{2-x}Co_xAs_2$) experiments suggest a finite doping interval where superconductivity and static magnetism coexist. 23,24 This obvious proximity to antiferromagnetism suggests spin fluctuations being important for the mechanism of superconductivity with a respective impact on the normal-state properties, including the normal-state transport. ^{21,25–27}

However, little is known about the Nernst effect of the pnictide superconductors family. In a pioneering study Zhu $et\ al.$ reported an anomalous suppression of the off-diagonal thermoelectric current in optimally doped LaFeAsO_{1-x}F_x and suggested the presence of SDW fluctuations near the superconducting transition.²⁸ More recently, Matusiak $et\ al.$ observed a strong enhancement of the Nernst coefficient in the SDW state of the parent compound of CaFe_{2-x}Co_xAs₂, but did not find any particular anomaly in the Nernst effect of a superconducting doping level that could be attributed to either vortex flow or to SDW fluctuations.²⁹

In this Brief Report we systematically investigate the doping evolution of the Nernst effect in LaFeAsO $_{1-x}F_x$. For the parent compound we observe a huge negative Nernst coefficient accompanied with a severe violation of the so-called Sondheimer cancellation in the SDW state. Surprisingly, at underdoping (x=0.05) we find a similar enhanced ν at $T\lesssim 150$ K despite the absence of static magnetic order and the presence of bulk superconductivity, strongly suggestive of SDW fluctuations. More conventional transport is observed at optimal doping (x=0.1) where the normal-state Nernst signal is rather featureless with a more complete Sondheimer cancellation.

II. EXPERIMENT

Polycrystalline samples of LaFeAsO_{1-x}F_x (x = 0, 0.05, 0.1) were prepared and characterized by powder x-ray diffraction (XRD) and wavelength-dispersive x-ray spectroscopy (WDX,)²¹ and further investigated by resistivity,²⁷ magnetization,³⁰ nuclear quadrupole resonance (NQR),³¹ and muon spin rotation (μ SR) experiments.^{20,22} Measurements of the Nernst, Seebeck, and Hall coefficients were performed on cuboid-shaped samples of size \sim 0.6 \times 0.6 \times 1.5 mm³ in a homemade device. The Seebeck and Nernst coefficient were measured using a steady-state method with an surface mount resistor as a heater. The temperature gradient on the sample was measured using a Au-Chromel differential thermocouple.³² The Nernst coefficient was first determined as a function of temperature T at constant magnetic field B = 14 T with two

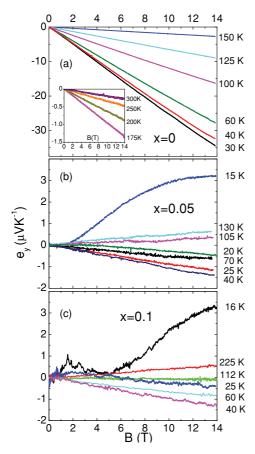


FIG. 1. (Color online) Nernst signal e_y of LaFeAsO_{1-x}F_x as a function of magnetic field for x = 0 (a), x = 0.05 (b), x = 0.1 (c).

opposite field polarizations, to compensate for the longitudinal thermal voltages and then subsequently measured at several temperature points as a function of magnetic field. Excellent agreement was obtained between the two methods. The Hall effect was measured using a standard four-point method by sweeping the magnetic field at stabilized temperatures.

III. RESULTS

Figure 1(a) presents the field dependence of the Nernst signal $e_y = E_y/\nabla T$ (with E_y the transverse electrical field) of the parent compound (x=0) for selected T. In the whole investigated temperature range e_y is negative and within the experiment's accuracy linearly field dependent. Between room temperature and 175 K the magnitude of the Nernst signal at 14 T is rather small with $|e_y| \lesssim 1.3 \ \mu\text{V/K}$. This changes drastically upon crossing T_S and T_N where a strong enhancement of the Nernst signal occurs. For example, we find a strikingly large magnitude of $|e_y| \approx 35 \ \mu\text{V/K}$ at 30 K and 14 T.

Significantly different is the behavior of the Nernst coefficient of the underdoped compound at x=0.05 shown at Fig. 1(b). At 15 K we observe a positive signal with the amplitude 3 μ V/K at 14 T and a nonlinear magnetic-field dependence which is indicative of a strong contribution caused by vortex flow. As T rises, $e_y(B)$ changes sign and becomes linear just above $T_c=20.6$ K. With further increasing T

the Nernst signal reaches a minimum at 40 K and becomes positive again above approximately 100 K. The absolute values are more than one order of magnitude lower than those of LaFeAsO.

At optimal doping (x = 0.1) the Nernst signal is at first glance very similar to that at x = 0.05: a strongly nonlinear and positive $e_y(B)$ due to vortex flow is observed below $T_c = 26.8$ K. As T rises, $e_y(B)$ becomes negative and linear with a minimum at about 40 K, which is followed by another sign change at further elevated T.

We turn now to the actual T dependence of the Nernst coefficient $v = e_v/B$ of the three samples shown in Fig. 2. For the undoped sample, $\nu(T)$ is negative over the whole T range. From $\nu = -0.02 \,\mu\text{VK}^{-1}\text{T}^{-1}$ at 300 K, the Nernst coefficient decreases moderately down to $\nu = -0.2 \,\mu\text{VK}^{-1}\text{T}^{-1}$ at about 150 K. At $T \leq 150$ K, i.e., in the SDW phase, a large negative contribution is visible in the curve—the slope changes strongly and $\nu(T)$ reaches its minimum at around 25 K with the "giant" value of $-2.5 \mu V K^{-1} T^{-1}$. While the enhancement of the Nernst coefficient qualitatively should be attributed to the Fermi-surface reconstruction that is associated with the SDW phase,⁵ its actual magnitude is quite surprising. The Nernst response is about one order of magnitude larger than that generated by vortex flow in the superconducting samples or in, e.g., cuprate superconductors, ^{3,16} which is often considered as a benchmark for a large Nernst effect. It is also

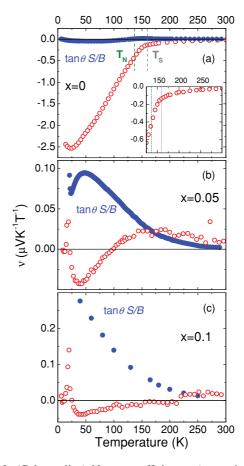


FIG. 2. (Color online) Nernst coefficient ν (open circles) and $\tan \theta S/B$ (full circles) of LaFeAsO_{1-x}F_x as a function of temperature for x=0 (a), x=0.05 (b), x=0.1 (c).

one order of magnitude larger than the presumably SDW-related enhancement of the Nernst coefficient in electron-doped cuprates. 3,5,10,11 Instead, its size rather compares with the so-called giant Nernst effect of URu_2Si_2 and Bechgaard salts. $^{7,33-35}$ However, the absolute magnitude of ν is almost one order of magnitude larger in our data for LaFeAsO. In order to judge to what extent ν is unusual in terms of its magnitude, we consider the expression 2,36

$$\nu = \left(\frac{\alpha_{xy}}{\sigma} - S \tan \theta\right) \frac{1}{B}.\tag{1}$$

Here S is the Seebeck coefficient, $\tan\theta$ is the Hall angle, σ is the electrical conductivity, and α_{xy} is the off-diagonal Peltier conductivity. In a one-band metal, the two terms in Eq. (1) cancel exactly if the Hall angle is independent of energy (Sondheimer cancellation). Above However, in a multiband material such as LaFeAsO_{1-x}F_x under scrutiny here, this cancellation is no longer valid. The degree of its violation can be determined experimentally by comparing the measured ν with the term $S \tan\theta/B$, which can be easily calculated from our thermopower and Hall data displayed in Fig. 3 and the resistivity. The direct comparison of both quantities shown in Fig. 2(a) reveals clear $|\nu| \gg |S \tan\theta|/B$, i.e., a severe violation of the Sondheimer cancellation in the SDW phase.

The Nernst coefficient of the underdoped sample [x=0.05, Fig. 2(b)] exhibits a strong positive contribution arising from vortex motion in the superconducting state which extends up to a remarkably large $T\approx 40$ K. Surprisingly, at higher temperature $\nu(T)$ resembles that of the SDW ordering as seen in the parent compound: Between 300 K and about 150 K we observe a rather flat $\nu(T)$ with a weak negative slope. However, at around 150 K the T dependence changes and a sizable negative contribution becomes apparent in ν which leads to a sign change at ~ 100 K and a minimum at ~ 40 K where the vortex contribution sets in. Note that in this regime

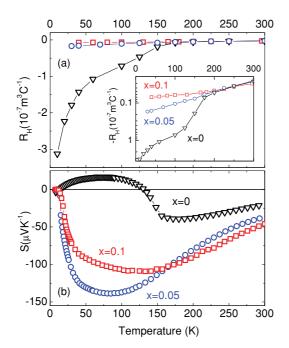


FIG. 3. (Color online) Hall coefficient R_H (a) and thermopower S (b) as a function of temperature for x = 0, 0.05, 0.1.

 $|\nu| \approx |S \tan \theta|/B$, i.e., a significant violation of the Sondheimer cancellation is still obvious, and the negative contribution between 40 and 150 K is of similar size as the vortex contribution at low T. The presence of the SDW-like signature at $T \lesssim 150$ K is completely unexpected since our sample exhibits bulk superconductivity where μ SR and Mössbauer spectroscopy show no trace of magnetic ordering in this T regime. We therefore conclude that instead of true SDW order, a form of a respective precursor, such as fluctuations or possibly nematic phases, give rise to the enhanced Nernst response. In fact, we observe an enhancement of ν with a similar magnitude also for the parent compound at $T \gtrsim T_N$ [cf. Fig. 1(a)], i.e., in a T range where SDW precursors are truly present, 37,38 which corroborates this conclusion.

Despite a very similar behavior in the vicinity of T_c and similar magnitude the Nernst coefficient of the optimally doped compound $[x=0.1, {\rm Fig.~2(c)}]$ displays a completely different normal-state behavior as compared to the underdoped compound. At $T\gtrsim 40~{\rm K}$, i.e., in the whole normal state, $\nu(T)$ is rather featureless with a weak positive slope. In particular, no anomaly that could be related to SDW fluctuations is present and, moreover, a more complete Sondheimer cancellation, i.e., $|\nu| \ll |S \tan \theta|/B$ is found at low T. Altogether this suggests a more conventional transport behavior for this doping level.

IV. DISCUSSION

Our findings provide fresh information about the doping evolution of the normal-state properties of LaFeAsO_{1-x} F_x , in particular concerning the interesting nature of the normal state close to the boundary to the SDW phase. We point out that the onset of anomalous Nernst response at $T \lesssim$ 150 K for x = 0.05 coincides with a resistivity anomaly at about the same temperature, which has been interpreted as a signature of SDW fluctuations as well.²⁷ Furthermore, recent NMR and NQR studies provide clear-cut evidence for the slowdown of spin fluctuations in this T regime³⁹ and for nanoscale electronic order,³¹ which further underpins our conclusion. It thus seems that in contrast to other pnictide systems such as $BaFe_{2-x}Co_xAs_2$ where strong evidence for the coexistence of static magnetic order and superconductivity exists, 40 LaFeAsO_{1-x}F_x represents one rare example among the iron pnictides where both order parameters do not coexist but truly compete with each other. However, one might speculate that rather small perturbations of the system, e.g., the disorder that Co doping generates in the Fe₂As₂ plane in $BaFe_{2-x}Co_xAs_2$, might be sufficient to stabilize the ubiquitous SDW fluctuations and drive the system toward the coexistence of static magnetism and superconductivity. 41,42

It seems noteworthy that our results provide one more example of an unconventional superconductor where an unusually large normal-state Nernst effect is observed. The Nernst signal related to the SDW precursors clearly exceeds that of recently reported stripe and/or nematic phases observed recently in cuprate superconductors. Furthermore, the negative sign of the unusual contribution to the Nernst coefficient unambiguously rules out vortex fluctuations ^{1–3} as a possible origin since these should give rise to a positive Nernst response.

V. SUMMARY

In summary, an unusually enhanced Nernst effect accompanied with a considerable violation of the Sondheimer cancellation is observed in both undoped as well as in underdoped LaFeAsO_{1-x}F_x at $T \lesssim 150$ K. While the enhancement in the parent compound is huge (giant Nernst effect) and coincides with the onset of SDW order, it is about one order of magnitude smaller in the underdoped, bulk superconducting sample. Since no trace of magnetic order is observed, our results suggest the presence of SDW fluctuations at this

doping level. At optimal doping a more conventional, rather featureless Nernst response with a more complete Sondheimer cancellation is observed.

ACKNOWLEDGMENTS

This work was supported by the Deutsche Forschungsgemeinschaft through Grant No. BE1749/12, the Research Unit FOR538 (Grant No. BU887/4), and the Priority Programme SPP1458 (Grant No. GR3330/2).

- *c.hess@ifw-dresden.de
- ¹Z. A. Xu, N. P. Ong, Y. Wang, T. Kakeshita, and S. Uchida, Nature (London) **406**, 486 (2000).
- ²Y. Wang, Z. A. Xu, T. Kakeshita, S. Uchida, S. Ono, Y. Ando, and N. P. Ong, Phys. Rev. B **64**, 224519 (2001).
- ³Y. Wang, L. Li, and N. P. Ong, Phys. Rev. B **73**, 024510 (2006).
- ⁴O. Cyr-Choiniere et al., Nature (London) **458**, 743 (2009).
- ⁵A. Hackl and S. Sachdev, Phys. Rev. B **79**, 235124 (2009).
- ⁶A. Hackl, M. Vojta, and S. Sachdev, Phys. Rev. B **81**, 045102 (2010).
- ⁷K. Behnia, J. Phys.: Condens. Matter **21**, 113101 (2009).
- ⁸R. Daou et al., Nature (London) **463**, 519 (2010).
- ⁹A. Hackl and M. Vojta, Phys. Rev. B **80**, 220514 (2009).
- ¹⁰P. Fournier, X. Jiang, W. Jiang, S. N. Mao, T. Venkatesan, C. J. Lobb, and R. L. Greene, Phys. Rev. B **56**, 14149 (1997).
- ¹¹P. Li and R. L. Greene, Phys. Rev. B **76**, 174512 (2007).
- ¹²C. Hess, E. Ahmed, U. Ammerahl, A. Revcolevschi, and B. Büchner, Eur. Phys. J. Special Topics 188, 103 (2010).
- ¹³R. P. Huebener and A. Seher, Phys. Rev. **181**, 701 (1969).
- ¹⁴F. A. Otter and P. R. Solomon, Phys. Rev. Lett. **16**, 681 (1966).
- ¹⁵S. J. Hagen, C. J. Lobb, R. L. Greene, M. G. Forrester, and J. Talvacchio, Phys. Rev. B 42, 6777 (1990).
- ¹⁶H.-C. Ri, R. Gross, F. Gollnik, A. Beck, R. P. Huebener, P. Wagner, and H. Adrian, Phys. Rev. B **50**, 3312 (1994).
- ¹⁷Y. Kamihara, T. Watanabe, M. Hirano, and H. Hosono, J. Am. Chem. Soc. **130**, 3296 (2008).
- ¹⁸Z.-A. Ren et al., Chin. Phys. Lett. **25**, 2215 (2008).
- ¹⁹C. de la Cruz *et al.*, Nature (London) **453**, 899 (2008).
- ²⁰H.-H. Klauss et al., Phys. Rev. Lett. **101**, 077005 (2008).
- ²¹A. Kondrat *et al.*, Eur. Phys. J. B **70**, 461 (2009).
- ²²H. Luetkens, H.-H. Klauss, M. Kraken, F. J. Litterst, T. Dellmann, R. Klingeler, C. Hess, R. Khasanov, A. Amato, C. Baines, M. Kosmala, O. J. Schumann, M. Braden, J. Hamann-Borrero, N. Leps, A. Kondrat, G. Behr, J. Werner, and B. Büchner, Nat. Mater. 8, 305 (2009).
- ²³A. J. Drew *et al.*, Nat. Mater. **8**, 310 (2009).

- ²⁴N. Ni, S. L. Bud'ko, A. Kreyssig, S. Nandi, G. E. Rustan, A. I. Goldman, S. Gupta, J. D. Corbett, A. Kracher, and P. C. Canfield, Phys. Rev. B 78, 014507 (2008).
- ²⁵P. Prelovšek, I. Sega, and T. Tohyama, Phys. Rev. B 80, 014517 (2009).
- ²⁶P. Prelovšek and I. Sega, Phys. Rev. B **81**, 115121 (2010).
- ²⁷C. Hess, A. Kondrat, A. Narduzzo, J. E. Hamann-Borrero, R. Klingeler, J. Werner, G. Behr, and B. Büchner, Europhys. Lett. **87**, 17005 (2009).
- ²⁸Z. W. Zhu, Z. A. Xu, X. Lin, G. H. Cao, C. M. Feng, G. F. Chen, Z. Li, J. L. Luo, and N. L. Wang, New J. Phys. **10**, 063021 (2008).
- ²⁹M. Matusiak, Z. Bukowski, and J. Karpinski, Phys. Rev. B 81, 020510 (2010).
- ³⁰R. Klingeler *et al.*, Phys. Rev. B **81**, 024506 (2010).
- ³¹G. Lang, H.-J. Grafe, D. Paar, F. Hammerath, K. Manthey, G. Behr, J. Werner, and B. Büchner, Phys. Rev. Lett. **104**, 097001 (2010).
- ³²C. Hess, B. Büchner, U. Ammerahl, and A. Revcolevschi, Phys. Rev. B 68, 184517 (2003).
- ³³R. Bel, H. Jin, K. Behnia, J. Flouquet, and P. Lejay, Phys. Rev. B 70, 220501 (2004).
- ³⁴M.-S. Nam, A. Ardavan, W. Wu, and P. M. Chaikin, Phys. Rev. B 74, 073105 (2006).
- ³⁵A qualitatively similar impact of SDW order on the Nernst effect has been observed also in CaFe₂As₂ by Matusiak *et al.* (Ref. 29). However, the absolute magnitude of ν is almost one order of magnitude larger in our data for LaFeAsO.
- ³⁶E. H. Sondheimer, Proc. R. Soc. London, Ser. A **193**, 484 (1948).
- ³⁷Y. Nakai, K. Ishida, Y. Kamihara, M. Hirano, and H. Hosono, J. Phys. Soc. Jpn. **77**, 073701 (2008).
- ³⁸L. Wang, U. Köhler, N. Leps, A. Kondrat, M. Nale, A. Gasparini, A. de Visser, G. Behr, C. Hess, R. Klingeler, and B. Büchner, Phys. Rev. B 80, 094512 (2009).
- ³⁹C. Hess, H. J. Grafe, and B. Büchner (unpublished).
- ⁴⁰Y. Laplace, J. Bobroff, F. Rullier-Albenque, D. Colson, and A. Forget, Phys. Rev. B 80, 140501 (2009).
- ⁴¹R. M. Fernandes et al., Phys. Rev. B **81**, 140501 (2010).
- ⁴²R. M. Fernandes and J. Schmalian, Phys. Rev. B **82**, 014521 (2010).