# Quantitative size-dependent structure and strain determination of CdSe nanoparticles using atomic pair distribution function analysis

A. S. Masadeh, E. S. Božin, C. L. Farrow, G. Paglia, P. Juhas, and S. J. L. Billinge\* Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan 48824-1116, USA

A. Karkamkar and M. G. Kanatzidis

Department of Chemistry, Michigan State University, East Lansing, Michigan 48824-1116, USA (Received 9 April 2007; revised manuscript received 3 August 2007; published 12 September 2007)

The size-dependent structure of CdSe nanoparticles, with diameters ranging from 2 to 4 nm, has been studied using the atomic pair distribution function (PDF) method. The core structure of the measured CdSe nanoparticles can be described in terms of the wurtzite atomic structure with extensive stacking faults. The density of faults in the nanoparticles is  $\sim 50\%$ . The diameter of the core region was extracted directly from the PDF data and is in good agreement with the diameter obtained from standard characterization methods, suggesting that there is little surface amorphous region. A compressive strain was measured in the Cd–Se bond length that increases with decreasing particle size being 0.5% with respect to bulk CdSe for the 2 nm diameter particles. This study demonstrates the size-dependent quantitative structural information that can be obtained even from very small nanoparticles using the PDF approach.

DOI: 10.1103/PhysRevB.76.115413

PACS number(s): 61.46.Df, 61.10.-i, 78.66.Hf, 61.46.-w

# I. INTRODUCTION

Semiconductor nanoparticles are of increasing interest for both applied and fundamental research. Wurtzite-structured cadmium selenide is an important II-VI semiconducting compound for optoelectronics.<sup>1</sup> CdSe quantum dots are the most extensively studied quantum nanostructure because of their size-tunable properties, and they have been used as a model system for investigating a wide range of nanoscale electronic, optical, optoelectronic, and chemical processes.<sup>2</sup> CdSe also provided the first example of self-assembled semiconductor nanocrystal superlattices.<sup>3</sup> With a direct band gap of 1.8 eV, CdSe quantum dots have been used for laser diodes,<sup>4</sup> nanosensing,<sup>5</sup> and biomedical imaging.<sup>6</sup> In fundamental research, particles with a diameter in the 1-5 nm range are of particular importance since they cover the transition regime between the bulk and molecular domains where quantum size effects play an important role. Significant deviation from bulk properties are expected for particles with diameter below 5 nm and were observed in many cases<sup>6,7</sup> as well as in this study.

Accurate determination of atomic scale structure, homogeneous and inhomogeneous strain, structural defects, and geometrical particle parameters such as diameter and shape is important for understanding the fundamental mechanisms and processes in nanostructured materials. However, difficulties are experienced when standard methods are applied to small nanoparticles. In this domain, the presumption of a periodic solid, which is the basis of a crystallographic analysis, breaks down. Quantitative determinations of the nanoparticle structure require methods that go beyond crystallography. This was noted early on in a seminal study by Bawendi *et al.*<sup>8</sup> where they used the Debye equation, which is not based on a crystallographic assumption, to simulate semiquantitatively the scattering from some CdSe nanoparticles. However, despite the importance of knowing the nanoparticle structure quantitatively with high accuracy, this work has not been followed up with application of modern local structural methods<sup>9,10</sup> until recently.<sup>11–16</sup> In this study, we return to the archetypal CdSe nanoparticles to investigate the extent of information about the size-dependent structure of nanoparticles from the atomic pair distribution function (PDF) method. This is a local structural technique that yields quantitative structural information on the nanoscale from x-ray and neutron powder diffraction data.<sup>10</sup> Recent developments in both data collection<sup>17,18</sup> and modeling<sup>19,20</sup> make this a potentially powerful tool in the study of nanoparticles. Additional extensions to the modeling are necessary for nanoparticles, and some of these have been successfully demonstrated.<sup>11,12,21</sup>

In this paper, we present a detailed analysis of the structural information available from PDF data on (2-4 nm)CdSe nanoparticles. The PDF method is demonstrated here as a key tool that can yield precise structural information about the nanoparticles such as the atomic structure size of the core, the degree of crystallinity, local bonding, the degree of the internal disorder, and the atomic structure of the core region, as a function of the nanoparticle diameter. Three CdSe nanoparticle samples with different diameters that exhibit different optical spectra have been studied. The purpose of this paper is not only to explain the PDF data of CdSe nanoparticles through a modeling process but also to systematically investigate the sensitivity of the PDF data to subtle structural modifications in nanoparticles relative to the bulk material.

The measurement of the nanoparticle size can lead to significantly different results when performed by different methods, and there is no consensus as to which is the most reliable.<sup>22,23</sup> It is also not clear that a single diameter is sufficient to fully specify even a spherical particle since the presence of distinct crystalline core and disordered surface regions has been postulated.<sup>8</sup>

Powder diffraction is a well established method for structural and analytical studies of crystalline materials, but the applicability to such small particles of standard powder diffraction based on crystallographic methods is questionable and likely to be semiquantitative at best. Palosz *et al.*<sup>24</sup> have shown that the conventional tools developed for the elaboration of powder diffraction data are not directly applicable to nanocrystals.<sup>24</sup> There have been some reports<sup>8,23,25</sup> in the past few years extracting nanoparticle diameter from x-ray diffraction (XRD) using the Scherrer formula, which is a phenomenological approach that considers the finite size broadening of Bragg peaks.<sup>26</sup> This approach will decrease in accuracy with decreasing particle size, and for particle sizes in the range of a few nanometers, the notion of a Bragg peak becomes moot.<sup>24</sup> At this point, the Debye formula<sup>27</sup> becomes the more appropriate way to calculate the scattering.<sup>8</sup> The inconsistencies between the nanoparticle diameter determined from the standard characterization methods and the diameter obtained by applying the Scherrer formula have been observed by several authors.8,23,28

Previous studies of CdSe nanoparticle structure have demonstrated the sensitivity of the XRD pattern to the presence of planar disorder and thermal effects due to nanosize effects.<sup>8,29</sup> The diffraction patterns of CdSe nanoparticles smaller than 2.0 nm have been observed to appear markedly different from those of the larger diameters (see Ref. 29, Fig. 11), the large attenuation and broadening in the Bragg reflections in these small nanoparticles, making the distinction between wurtzite and zinc blende hard using conventional XRD methods. Murray et al.<sup>29</sup> reported that the combination of x-ray studies and transmission electron micrograph (TEM) imaging yields a description of the average CdSe nanoparticle structure. Strict classification of the CdSe nanoparticles structure as purely wurtzite or zinc blende is potentially misleading.<sup>29</sup> Bawendi et al.<sup>8</sup> reported that CdSe nanoparticles are best fitted by a mixture of crystalline structures intermediate between zinc blende and wurtzite. Here, we apply the PDF method to CdSe nanoparticles and refine quantitative structural parameters to a series of CdSe nanoparticles of different sizes.

Strain in nanosystems has been observed before in different studies, as well as in this study. Using combined PDF and extended x-ray-absorbtion fine structure (EXAFS) methods, Gilbert et al.<sup>11</sup> observed a compressive strain compared to the bulk in ZnS nanocrystals. Using an electric-field-induced resonance method. Chen *et al.*<sup>30</sup> detected the enhancement of Young's modulus of ZnO nanowires along the axial direction when the diameters are decreased. Very recently, Ouyang et  $al.^{31}$  developed an analytical model for the size-induced strain and stiffness of a nanocrystal from the perspective of thermodynamics and a continuum medium approach. It was found theoretically that the elastic modulus increases with the inverse of crystal size and the vibration frequency is higher than that of the bulk.<sup>31</sup> Experimentally, the CdQ (Q =S,Se,Te) first-neighbor distances have been studied using both XRD and EXAFS methods.<sup>32</sup> The distances were found smaller than those in the bulk compounds by less than 1.0%. Herron et al.<sup>33</sup> studied CdS nanocrystals and showed a bond contraction of ~0.5% compared to the bulk. Carter *et al.*<sup>34</sup> studied a series of CdSe nanoparticles using the EXAFS method. In the first shell around both the Se and Cd atoms, they found essentially no change in the first-neighbor distance. Chaure *et al.*<sup>35</sup> studied the strain in nanocrystalline CdSe thin films, using Raman scattering, and observed a peak shift with a decrease in particle size, which was attributed to the increase in stress with decreasing particle size.<sup>35</sup>

Local structural deviations or disorder mainly affect the diffuse scattering background. The XRD experiments probe for the presence of periodic structure, which are reflected in the Bragg peaks. In order to have information about both long-range order and local-structure disorder, a technique that takes both the Bragg and the diffuse scattering need to be used, such as the PDF technique. Here, we apply the PDF method to study the structure, size, and strain in CdSe nanoparticles as a function of nanoparticle diameter. The core structure of the CdSe nanoparticles can be described by a mixture of crystalline structures intermediate between zinc blende and wurtzite, which is wurtzite containing a stacking fault density (SFD) of up to  $\sim$ 50%, with no clear evidence of a disordered surface region, certainly down to 3 nm diameter. The structural parameters are reported quantitatively. We measure a size-dependent strain on the Cd-Se bond which reaches 0.5% at the smallest particle size. The size of the well-ordered core extracted directly from the data agrees with the size determined from other methods.

## **II. EXPERIMENTAL DETAILS**

#### A. Sample preparation

CdSe nanoparticles were synthesized from cadmium acetate, selenium, trioctyl phosphine, and trioctyl phosphine oxide. 64 g of trioctylphosphine oxide (TOPO) containing cadmium acetate was heated to 360 °C under flowing argon. Cold stock solution (38.4 ml) of (Se:trioctylphosphine =2:100 by mass) was quickly injected into the rapidly stirred, hot TOPO solution. The temperature was lowered to 300 °C by the injection. At various time intervals, 5–10 ml aliquots of the reaction mixture were removed and precipitated in 10 ml of methanol. The color of the sample changed from bright yellow to orange to red to brown with a time interval variation from 20 to 1200 s. Three nanoparticle sizes, CdSe I (small), CdSe II (medium), and CdSe III (large), were used for this study, as well as a bulk CdSe sample for reference.

The samples were further purified by dissolving and centrifuging in methanol to remove excess TOPO. This process also resulted in a narrower particle size distribution. The TEM images (Fig. 1) show uniformly sized nanoparticles with no signs of aggregation. The ultraviolet visible (UV-vis) absorption and photoluminescence (PL) spectra of the aliquots were recorded by redissolving the nanocrystals in toluene. The spectra are shown in Fig. 2.

The band-gap values obtained for the measured samples can be correlated with the diameter of the nanoparticles based on the table provided in the supplementary information of Peng *et al.*<sup>36</sup> using the data on exciton peaks measured with UV-vis light absorption and photoluminescence peaks. The particle sizes were measured by TEM as well. The measured values of particle diameter using these various methods are summarized in Table I.



FIG. 1. TEM image of CdSe nanocrystal prepared using the method described in the text. CdSe obtained by 1200 s (top) and 15 s (bottom) nucleation. The line bar is 10 nm in size in both images.

#### B. Atomic Pair distribution function method

The atomic PDF analysis of x-ray and neutron powder diffraction data is a powerful method for studying the structure of nanostructured materials.<sup>9,10,37–40</sup> Recently, it has been explicitly applied to study the structure of discrete nanoparticles.<sup>11,12,40–42</sup> The PDF method can yield precise structural and size information, provided that special care is applied to the measurement and to the method used for analyzing the data. The atomic PDF G(r) is defined as

$$G(r) = 4\pi r [\rho(r) - \rho_0],$$
 (1)

where  $\rho(r)$  is the atomic pair density,  $\rho_0$  is the average atomic number density, and *r* is the radial distance.<sup>43</sup> The



FIG. 2. (Color online) (a) Room temperature UV-vis absorption and (b) photoluminescence spectra from the sample of CdSe nanocrystals: ( $\bigcirc$ ) CdSe I, ( $\triangle$ ) CdSe II, and ( $\square$ ) CdSe III.

PDF yields the probability of finding pairs of atoms separated by a distance r. It is obtained by a sine Fourier transformation of the reciprocal space total scattering structure function S(Q), according to

$$G(r) = \frac{2}{\pi} \int_0^\infty Q[S(Q) - 1] \sin Qr dQ, \qquad (2)$$

where S(Q) is obtained from a diffraction experiment. This approach is widely used for studying liquids and amorphous and crystalline materials but has recently also been successfully applied to nanocrystalline materials.<sup>10</sup>

#### C. High-energy x-ray diffraction experiments

X-ray powder diffraction experiments to obtain the PDF were performed at the 6IDD beamline at the Advanced Photon Source at Argonne National Laboratory. Diffraction data were collected using the recently developed rapid acquisition pair distribution function (RAPDF) technique<sup>17</sup> that benefits from two-dimensional (2D) data collection. Unlike TEM,

TABLE I. CdSe nanoparticle diameter as determined using various methods.

	CdSe III	CdSe II	CdSe I
Nucleation time (s)	1200	630	15
Diameter (nm)			
TEM	3.5(2)	2.7(2)	2.0(2)
UV-vis	3.5(4)	2.9(3)	≤1.90
PL	3.6(4)	2.9(3)	≤2.1
PDF	3.7(1)	3.1(1)	2.2(2)



FIG. 3. (Color online) Two-dimensional XRD raw data collected using the image plate detector from (a) CdSe bulk and (b) nanoparticle CdSe III samples.

XRD probes a large number of crystallites that are randomly oriented. The powder samples were packed in a flat plate with thickness of 1.0 mm sealed between Kapton tapes. Data were collected at room temperature with an x-ray energy of 87.005 keV ( $\lambda$ =0.142 48 Å). An image plate camera (Mar345) with a diameter of 345 mm was mounted orthogonally to the beam path with a sample to detector distance of 208.857 mm, as calibrated by using a silicon standard sample.<sup>17</sup> The image plate was exposed for 10 s and this was repeated five times for a total data collection time of 50 s. The RAPDF approach avoids detector saturation while allowing sufficient statistics to be obtained. This approach also avoids sample degradation in the beam that was observed for the TOPO coated nanoparticles during longer exposures, on the scale of hours, that were required using conventional point-detector approaches. To reduce the background scattering, lead shielding was placed before the sample with a small opening for the incident beam.

Examples of the raw 2D data are shown in Fig. 3. These data were integrated and converted to intensity versus  $2\theta$  using the software FIT2D,<sup>44</sup> where  $2\theta$  is the angle between the incident and scattered x-ray beams. The integrated data were normalized by the average monitor counts. The data were corrected and normalized<sup>9</sup> using the program PDFGETX2 (Ref.



FIG. 4. (Color online) (a) The experimental reduced structure function F(Q) of CdSe nanoparticle with different diameters and (b) the corresponding PDF G(r) obtained by Fourier transformation of the data in (a) with  $Q_{max}$ =25.0 Å<sup>-1</sup>, from top to bottom: bulk, CdSe III, CdSe II, and CdSe I.

45) to obtain the total scattering structure function S(Q) and the PDF G(r), which are shown in Figs. 4(a) and 4(b), respectively. The scattering signal from the surfactant (TOPO) was measured independently and subtracted as a background in the data reduction.

In the Fourier transform step to get from S(Q) to the PDF G(r), the data are truncated at a finite maximum value of the momentum transfer,  $Q=Q_{max}$ . Different values of  $Q_{max}$  may be chosen. Here, a  $Q_{max}=25.0$  Å<sup>-1</sup> was found to be optimal.  $Q_{max}$  is optimized such as to avoid large termination effects and to reasonably minimize the introduced noise level as the signal to noise ratio decreases with the Q value.

Structural information was extracted from the PDFs using a full-profile real-space local-structure refinement method<sup>46</sup> analogous to Rietveld refinement.<sup>47</sup> We used an updated version<sup>48</sup> of the program PDFFIT (Ref. 19) to fit the experimental PDFs. Starting from a given structure model and given a set of parameters to be refined, PDFFIT searches for the best structure that is consistent with the experimental PDF data. The residual function ( $R_w$ ) is used to quantify the agreement of the calculated PDF from model to experimental data,

$$R_{w} = \sqrt{\frac{\sum_{i=1}^{N} \omega(r_{i}) [G_{obs}(r_{i}) - G_{calc}(r_{i})]^{2}}{\sum_{i=1}^{N} \omega(r_{i}) G_{obs}^{2}(r_{i})}}.$$
 (3)

Here, the weight  $\omega(r_i)$  is set to unity which is justified because in G(r), the statistical uncertainty on each point is approximately equal.<sup>49,50</sup>

The structural parameters of the model were unit cell parameters, anisotropic atomic displacement parameters (ADPs), and the fractional coordinate z of the Se and/or Cd atom. Nonstructural parameters that were refined were a correction for the finite instrumental resolution ( $\sigma_Q$ ), low-r correlated motion peak sharpening factor ( $\delta$ ),<sup>51,52</sup> and scale factor. When estimating the particle size, a new version of the fitting program with particle size effects included as a refinable parameter<sup>53</sup> was used. The sample resolution broadening was determined from a refinement to the crystalline CdSe and the silicon standard sample and fixed, and the particle diameter refined, as described below. Good agreement between these results was obtained.

# **III. RESULTS AND DISCUSSION**

The reduced structure functions for the bulk and nanocrystalline samples are shown plotted over a wide range of Qin Fig 4(a). All of the patterns show significant intensity up to the highest values of Q, highlighting the value of measured data over such a wide Q range. All of the diffraction patterns have peaks in similar positions reflecting the similarity of the basic structures, but as the nanoparticles get smaller, the diffraction features become broadened out due to finite size effects.<sup>26</sup>

The PDFs are shown in Fig. 4(b). What is apparent is that, in real space, the PDF features at low r are comparably sharp in all the samples. The finite size effects do not broaden features in real space. The finite particle size is evident in a falloff in the intensity of structural features with increasing r. Later, we will use this to extract the average particle size in the material. The structure apparent in the G(r) function comes from the atomic order within the nanoparticle. The value of r where these ripples disappear indicates the particle core region diameter, or at least the diameter of any coherent structural core of the nanoparticle. By direct observation (Fig. 9), we can put a lower limit on the particle diameters to be 3.6, 2.8, and 1.6 nm for CdSe III, II and I, respectively, where the ripples can be seen to die out by visual inspection. These numbers will be quantified more accurately later.

#### A. Nanoparticle structure

Features in the PDF at low r reflect the internal structure of the nanoparticles. The nanoparticle PDFs have almost the same features as in the bulk in the region below 8.0 Å, reflecting the fact that they share a similar atomic structure on



FIG. 5. (Color online) Fragments from the (a) wurtzite structure, space group ( $P6_3mc$ ) and (b) zinc-blende structure, space group ( $F\overline{4}3m$ ).

average. In the finite nano-size regime, local structural deviations from the average bulk structure are expected.

A large number of semiconductor alloys, especially some sulfides and selenides, do not crystallize in the cubic zincblende structure but in the hexagonal wurtzite structure.<sup>54</sup> Both wurtzite and zinc-blende structures are based on the stacking of identical two-dimensional planar units translated with respect to each other, in which each atom is tetrahedrally coordinated with four nearest neighbors. The layer stacking is described as *ABABAB*... along the [001] axis for wurtzite and as *ABCABC*... along the [111] axis for zinc blende. As can be seen in Fig. 5, each cadmium and selenium are tetrahedrally coordinated in both structures. However, the next nearest and more distant coordination sequences are different in the two structures.

The largest changes in structure are expected in the smallest nanoparticles. In these small nanoparticles, the proportion of atoms on the surface is large, making the notion of a well-ordered crystalline core moot. The fractions of atoms involved in the surface atoms were estimated as 0.6, 0.45, and 0.35 for 2, 3, and 4 nm nanoparticle diameters, respectively. This was estimated by taking different spherical cuts from the bulk structure, then counting the atom with coordination number 4 as the core atom and the one with less than 4 as the surface atom. For the smallest particles, the small number of atoms in the core makes it difficult to define a core crystal structure, making the distinction between wurtzite and zinc blende difficult using the conventional XRD methods as the nanoparticle size decreases.<sup>29</sup> The principle difference between these structures is the topology of the CdSe<sub>4</sub> connections, which may also be becoming defective in the small nanoparticles.

Two structure models, wurtzite (space group  $P6_3mc$ ) and zinc blende (space group  $F\overline{4}3m$ ), were fit to the PDF data. The results of the full-profile fitting to the PDF data are shown in Fig. 6. In this figure, we compare fits to the (a) wurtzite and (b) zinc-blende structure models using isotropic atomic displacement factors ( $U_{iso}$ ) in both models. The wurtzite structure gives superior fits for the bulk structure. However, for all the nanoparticle sizes, the fits of wurtzite and zinc blende are comparable as evident from the difference curves in Fig. 6 and the  $R_w$  values reported in Table II. This indicates that classification of the CdSe nanoparticle structure as purely wurtzite or zinc blende is misleading<sup>29</sup> and it is better described as being intermediate between the two structures, as has been reported earlier.<sup>8</sup>



FIG. 6. (Color online) The experimental PDF G(r) with  $Q_{max} = 19.0 \text{ Å}^{-1}$  (blue solid dots) and the calculated PDF from the refined structural model (red solid line), with the difference curve offset below (black solid line). PDF data are fitted using (a) the wurtzite structure model with space group  $P6_{3}mc$  and (b) the zinc-blende model with space group  $F\overline{4}3m$ . In both models, isotropic atomic displacement factors ( $U_{iso}$ ) are used.

Introducing anisotropic ADPs  $(U_{11}=U_{22} \neq U_{33})$  into the wurtzite model resulted in better fits to the data. The refined parameters are reproduced in Table III and the fits are shown in Fig. 7(a). The values for the nanoparticles are rather close to the values in the bulk wurtzite structure. The model with anisotropic ADPs resulted in lower  $R_w$ . There is a general increase in the ADPs with decreasing particle size. This re-

TABLE II. The refined residual  $(R_w)$  values obtained from PDF analysis assuming the wurtzite and zinc blende structure models with space groups  $P6_{3}mc$  and  $F\overline{4}3m$ , respectively. In both models, isotropic atomic displacement factors  $(U_{iso})$  are used.

	CdSe bulk	CdSe III	CdSe II	CdSe I
Wurtzite $(R_w)$	0.16	0.31	0.28	0.31
Zinc blende $(R_w)$	0.52	0.32	0.30	0.35

flects the inhomogeneous strain accommodation in the nanoparticles, as we will discuss below. However, the values of the ADPs along the z direction for Se atoms  $(U_{33})$  are four times larger in the nanoparticles compared with the bulk where they are already unphysically large. The fact that this parameter is large on the Se site and small on the Cd site is not significant since we can change the origin of the unit cell to place a Cd ion at the (1/3, 2/3, z) position and the enlarged  $U_{33}$  shifts to the Cd site in this case.

The unphysically large  $U_{33}$  value on the Se site is likely to be due to the presence of faults in the basal plane stacking. For example, similar unphysical enlargements of perpendicular thermal factors in PDF measurements are explained by the presence of turbostratic disorder in layered carbons,<sup>55</sup> which is a similar effect to faults in the *ABABAB* wurtzite stacking. Also, the presence of stacking faults in the nanoparticles has been noted previously.<sup>8</sup> It is noteworthy that this parameter is enlarged in EXAFS analyses of bulk wurtzite structures, probably for the same reason.<sup>32,56,57</sup> We suspect that the enlargement in this parameter ( $U_{33}$ ) is related to the stacking fault density present in the bulk and that is increasing in the nanoparticles.

To test this idea, we simulated PDF data using the wurtzite structure containing different stacking fault densities. The stacking faults were simulated for different densities (0.167, 0.25, 0.333, and 0.5) by creating wurtzite superlattices with different stacking sequences along the *C* axis. The program DISCUS (Ref. 58) was used to create the stacking fault models and PDFGUI (Ref. 48) was used to generate the corresponding PDFs. The PDFs were simulated with all the ADPs fixed at  $U_{ii}$ =0.0133 Å<sup>2</sup>, the value observed in the experimental bulk data collected at room temperature (see Table III).

To see if this results in enlarged  $U_{33}$  values, we refined the simulated data containing stacking faults using the wurtzite model without any stacking faults. Indeed, the refined Se site  $U_{33}$  increased monotonically with increasing stacking fault density. The results are plotted in Fig. 8.

Figure 8 can be considered as a calibration curve of stacking fault density in the wurtzite structure, based on the enlargement in the ADPs along the z direction  $U_{33}$ . From this, we can estimate a stacking fault density of ~35% for our bulk CdSe sample and ~50% for each of the nanoparticles.

It is then possible to carry out a refinement using a structural model that contains an appropriate stacking fault density. The PDF data of bulk CdSe were therefore fitted with a wurtzite model with a 33% density, and the nanoparticle PDF fitted with a model with 50% of stacking faults. The refinements give excellent fits, as is evident in Fig. 7(b). The results are presented in Table III. The enlarged  $U_{33}$  parameter on the Se site is no longer present and it is now possible to refine physically reasonable values for that parameter. As well as resulting in physically reasonable ADPs, the quality of the fits to the data is excellent, though the  $R_w$  value is slightly larger in the nanoparticles.

Attempts to characterize the structure changes using direct measurements such as TEM technique for such small CdSe nanoparticles<sup>59</sup> were unsuccessful due to the poor contrast. However, in the present study, we were successful in exploring the local atomic structure for CdSe nanoparticles, in real space, at different length scales. The PDF fits clearly

	CdSe bulk		CdSe III		CdSe II		CdSe I	
Stacking fault density (%)	0.0	33.0	0.0	50.0	0.0	50.0	0.0	50.0
a (Å)	4.3014(4)	4.3012(4)	4.2997(9)	4.2987(9)	4.3028(9)	4.3015(9)	4.2930(9)	4.2930(8)
<i>c</i> (Å)	7.0146(9)	7.0123(9)	7.0145(4)	7.0123(4)	6.9987(9)	6.9975(9)	6.9405(9)	6.9405(7)
Se Z-frac.	0.3774(3)	0.3771(3)	0.3761(9)	0.3759(9)	0.3751(6)	0.3747(6)	0.3685(9)	0.3694(9)
Cd $U_{11} = U_{22} (Å^2)$	0.0108(2)	0.0102(2)	0.0146(7)	0.0149(7)	0.0149(6)	0.0112(5)	0.0237(9)	0.0213(8)
$U_{33}$ (Å <sup>2</sup> )	0.0113(3)	0.0112(3)	0.0262(9)	0.0241(9)	0.0274(9)	0.0271(9)	0.0261(9)	0.0281(9)
Se $U_{11} = U_{22} (Å^2)$	0.0109(9)	0.0102(9)	0.0077(7)	0.0138(7)	0.0083(7)	0.0121(7)	0.0110(9)	0.0191(9)
$U_{33}$ (Å <sup>2</sup> )	0.0462(9)	0.0115(9)	0.1501(9)	0.02301(9)	0.1628(9)	0.0265(9)	0.1765(9)	0.0311(9)
NP <sup>a</sup> diameter (nm)	$\infty$	$\infty$	3.7(1)	3.7(1)	3.1(1)	3.1(1)	2.4(2)	2.2(2)
$R_w$	0.12	0.09	0.20	0.14	0.18	0.15	0.27	0.21

TABLE III. The refined parameter values obtained from PDF analysis assuming the wurtzite structure, space group  $P6_{3}mc$ , with different stacking fault densities (SFDs).

<sup>a</sup>NP refers to nanoparticle.

indicate that the structure can be described in terms of locally distorted wurtzite structure containing  $\sim$ 50% stacking fault density (i.e., intermediate between wurtzite and zinc blende) even for the 2 nm diameter particles, Fig. 7.

Interestingly, there is little evidence in our data for a significant surface modified region. This surface region is sometimes thought of as being an amorphouslike region. Amorphous structures appear in the PDF with sharp first-neighbor peaks but rapidly diminish and broaden higher neighbor peaks. Thus, in the presence of a surface amorphous region, we might expect to see extra intensity at the first-peak position when the wurtzite model is scaled to fit the higher-rfeatures coming just from the crystalline core. As evident in Fig. 7, this is not observed. Furthermore, as we will describe below, the diameter of the crystalline core that we refine from the PDF agrees well with other estimates of nanoparticle size, suggesting that there is no surface amorphous region in these nanoparticles. The good agreement in the intensity of the first PDF peak also presents a puzzle in the opposite direction since we might expect surface atoms to be undercoordinated, which would result in a decrease in the intensity of this peak. It is possible that the competing effects of surface amorphous behavior and surface under coordination perfectly balance each other out, and this cannot be ruled out, though it seems unlikely that it would work perfectly at all nanoparticle diameters. This is also not supported by the nanoparticle size determinations described below.

#### **B.** Nanoparticle size

We describe here how we extracted more accurate nanoparticle diameters. This determination is important since the physical properties are size dependent. It is also important to use complementary techniques to determine particle size as different techniques are more dependent on different aspects of the nanoparticle structure, for example, whether or not the technique is sensitive to any amorphous surface layer on the nanoparticle. More challenges are expected in accurate size determination as the nanoparticle diameter decreases, due to poor contrast near the surface of the nanoparticle. In the literature, CdSe nanoparticles with a diameter of 2.0 nm have been considered to be an especially stable size with an associated band edge absorption centered at 414 nm,<sup>60</sup> that size was observed earlier<sup>29,61</sup> with an estimated diameter of  $\leq 2.0$  nm. There are some reported difficulties in determining the diameter of such small CdSe nanoparticles. Attempts to characterize the structure changes by TEM and x-ray diffraction techniques<sup>59</sup> were unsuccessful due to the small diameter of the particles relative to the capping material.

If we assume the nanoparticle to have a spherical shape (a reasonable approximation based on the TEM in Fig. 1) cut from the bulk, then the measured PDF will look like the PDF of the bulk material that has been attenuated by an envelope function given by the PDF of a homogeneous sphere as follows:<sup>62</sup>

$$G(r,d)_s = G(r)f(r,d), \tag{4}$$

where G(r) is given in Eq. (1) and f(r,d) is a sphere envelope function given by

$$f(r,d) = \left[1 - \frac{3r}{2d} + \frac{1}{2}\left(\frac{r}{d}\right)^3\right]\Theta(d-r),$$
 (5)

where *d* is the diameter of the homogeneous sphere and  $\Theta(x)$  is the Heaviside step function, which is equal to 0 for negative *x* and 1 for positive.

The approach is as follows. First, we refine the bulk CdSe data using PDFFIT. This gives us a measure of the PDF intensity falloff due to the finite resolution of the measurement.<sup>9</sup> Then, the measured value of the finite resolution was kept as an unrefined parameter after that, while all the other structural and nonstructural parameters were refined. To measure the PDF intensity falloff due to the finite particle size, the refined PDF is attenuated, during the refinement, by the envelope function [Eq. (5)], which has one refined parameter, the particle diameter. The fitted results are shown in Fig. 9 and the resulting values of particle diameter from the PDF refinement are recorded in Table I. The insets show the calculated and measured PDFs on an expanded scale. The ac-



FIG. 7. (Color online) The experimental PDF G(r) with  $Q_{max} = 19.0 \text{ Å}^{-1}$  (blue solid dots) and the calculated PDF from the refined structural model (red solid line), with the difference curve offset below (black solid line). PDF data are fitted using the wurtzite structure model (a) with no stacking fault and (b) with 33% stacking fault density for bulk and 50% for all nanoparticle sizes. In both cases, anisotropic atomic displacement factors ( $U_{aniso}$ ) are used.

curacy of determining the nanoparticle size can be evaluated directly from this figure. Features in the measured PDFs that correspond to the wurtzite structure are clearly seen disappearing smoothly attenuated by the spherical PDF envelope function. The procedure is least successful in the smallest nanoparticles, where the spherical particle approximation on the model results in features that extend beyond those in the data. In this case, the spherical approximation may not be working so well.

The particle diameters determined from the PDF are consistent with those obtained from TEM, UV-vis, and photoluminescence measurements. In particular, an accurate determination of the average diameter of the smallest particles is possible in the region where UV-vis and photoluminescence measurements lose their sensitivity.<sup>23</sup> In this analysis, we have not considered particle size distributions, which are small in these materials. The good agreement between the data and the fits justify this, though some of the differences



FIG. 8. (Color online) The enlargement in the ADPs along the z direction for Se site  $U_{33}$ , as a function of the stacking fault density.

at high r may result from this and could contribute an error to the particle size. Several additional fits to the data were performed to test the sphericity of the nanoparticles. Attempts were made to fit the PDF with oblate and prolate spheroid nanoparticle form factors. These fits resulted in ellipticities very close to 1, and large uncertainties in the refined ellipticity and particle diameters, which suggests that the fits are over parametrized. Another series of fits attempted to profile the PDF with a log-normal distribution of spherical nanoparticles. Allowing the mean nanoparticle diameter and log-normal width to vary resulted in nonconvergent fits, which implies that the particle sizes are not lognormal distributed. Therefore, there appears to be little evidence for significant ellipticity, nor a significant particle size distribution, as fits assuming undistributed spherical particles give the best results.

The simple fitting of a wurtzite structure with  $\sim$  50% SFD to the data will result in an estimate of the coherent structural core of the nanoparticle that has a structure which can be described by a mixture of crystalline structures intermediate between zinc blende and wurtzite. Comparing the nanoparticle core diameter extracted from PDF analysis with the diameter determined from the standard characterization methods yields information about the existence of a surface amorphous region. The agreement between the core diameter extracted from PDF and that determined from the standard methods (Table I) indicates that within our measurement uncertainties, there is no significant heavily disordered surface region in these nanoparticles, even at the smallest diameter of 2 nm (Fig. 9). In contrast with ZnS nanoparticles,<sup>11</sup> the heavily disordered surface region is about 40% of the nanoparticle diameter for a diameter of 3.4 nm, the surface region thickness being around 1.4 nm.11

## C. Internal strain

The local bonding of the tetrahedral Cd-Se building unit was investigated vs nanoparticle diameter. The nearest neighbor peaks at r=2.6353(3) Å come from covalently bonded



FIG. 9. (Color online) The experimental PDF G(r) shown as solid dots. Sphere envelope function [Eq. (5)] is used to transform the calculated PDF of bulk CdSe, using wurtzite structure containing 50% stacking fault density, to give a best fit replication of the PDF of CdSe nanoparticles (red solid line). The inset shows an expanded scale for the high-*r* region of experimental G(r) on the top of simulated PDF data for different diameters of CdSe nanoparticles (solid line): (a) CdSe III, (b) CdSe II, (c) CdSe I. Dashed lines are guides for the eyes.

Cd-Se pairs. The positions and the width of these peaks have been determined by fitting a Gaussian [Fig. 10(a)] and the results presented in Table IV. The results indicate that there



FIG. 10. (Color online) (a) The first PDF peak, ( $\bigcirc$ ) bulk, ( $\bigcirc$ ) CdSe III, ( $\square$ ) CdSe II, and ( $\triangle$ ) CdSe I fitted with one Gaussian (—). Dashed line represents the position of first PDF peak in the bulk data. (b)( $\blacktriangle$ ) The first PDF peak width vs nanoparticle size obtained from one Gaussian fit. Dashed line represents the width of first PDF peak in the bulk data. (c) Strain in Cd-Se bond ( $\Delta r/r$ ) (%) vs nanoparticle size. ( $\blacksquare$ ) Bond values obtained from the localstructure fitting and ( $\bigcirc$ ) obtained from one Gaussian fit to the first PDF peak. Dotted curves are guides for the eyes.

is a significant compressive strain on this near-neighbor bond length, and it is possible to measure it with the PDF with high accuracy. The bond length of Cd-Se pairs shorten as the nanoparticle diameter decreases, suggesting the presence of an internal stress in the nanoparticles. The Cd-Se bond lengths extracted from the PDF structural refinement are also in good agreement with those obtained from the first-peak Gaussian fit, as shown in Fig. 10(c). Thus, we have a model independent and a model dependent estimate of the strain that are in quantitative agreement. The widths of the first PDF peaks have also been extracted vs nanoparticle diameter from the Gaussian fits (Table IV). They remain comparably sharp as the nanoparticles get smaller, as shown in Fig. 10(b). Apparently, there is no size-dependent inhomogeneous strain measurable on the first peak. However, peaks at higher r do indicate significant broadening [Fig. 4(b)], suggesting that there is some relaxation taking place through bond bending. This is reflected in enlarged thermal factors that are refined in the nanoparticle samples. This is similar to what is observed in semiconductor alloys where most of the structural relaxation takes place in relatively lower energy bondbending distortions.<sup>63,64</sup>

TABLE IV. The first PDF peak position (FPP) and width (FPW) for different CdSe nanoparticle sizes and the bulk.

	CdSe bulk	CdSe III	CdSe II	CdSe I
PDF FPP (Å)	2.6353(3)	2.6281(3)	2.6262(3)	2.6233(3)
PDF FPW (Å)	0.1985(09)	0.1990(19)	0.2021(25)	0.2032(25)

# **IV. CONCLUSION**

The PDF is used to address the size and structural characterization of a series of CdSe nanoparticles prepared by the method mentioned in the text. The core structure of the measured CdSe nanoparticles was found to possess a welldefined atomic arrangement that can be described in terms of locally disordered wurtzite structure that contains  $\sim 50\%$ stacking fault density, and quantitative structural parameters are presented.

The diameter of the CdSe nanoparticles was extracted from the PDF data and is in good agreement with the diameter obtained from standard characterization methods, indicating that within our measurement uncertainties, there is no significant heavily disordered surface region in these nanoparticles, even at the smallest diameter of 2 nm. In contrast with ZnS nanoparticles,<sup>11</sup> the heavily disordered surface region is about 40% of the nanoparticle diameter for a diameter of 3.4 nm, the surface region thickness being around 1.4 nm.<sup>11</sup>

Compared with the bulk PDF, the nanoparticle PDF peaks are broader in the high-*r* region due to strain and structural defects in the nanoparticles. The nearest neighbor peaks at r=2.6353(3) Å which come from covalently bonded Cd-Se pairs shorten as the nanoparticle diameter decreases, resulting in a size-dependent strain on the Cd-Se bond that reaches 0.5% at the smallest particle size.

### ACKNOWLEDGMENTS

We would like to acknowledge help from Didier Wermeille, Doug Robinson, Mouath Shatnawi, Moneeb Shatnawi, and He Lin for help in collecting data. We are grateful to Christos Malliakas for the valuable assistance with the transmission electron microscopy. Many thanks to Hyun-Jeong Kim for useful discussion. We are grateful to Reinhard Neder for the valuable help with the stacking fault simulation. This work was supported in part by National Science Foundation (NSF) Grant No. DMR-0304391. Data were collected at the 6IDD beamline of the MUCAT sector at the Advanced Photon Source (APS). Use of the APS is supported by the U.S. DOE, Office of Science, Office of Basic Energy Sciences, under Contract No. W-31-109-Eng-38. The MUCAT sector at the APS is supported by the U.S. DOE, Office of Science, Office of Basic Energy Sciences, through the Ames Laboratory under Contract No. W-7405-Eng-82.

\*billinge@pa.msu.edu

- <sup>1</sup>G. Hodes, A. Albu-Yaron, F. Decker, and P. Motisuke, Phys. Rev. B **36**, 4215 (1987).
- <sup>2</sup>S. H. Tolbert and P. A. Alivisatos, Science **265**, 373 (1994).
- <sup>3</sup>C. B. Murray, C. R. Kagan, and M. G. Bawendi, Science **270**, 1335 (1995).
- <sup>4</sup>V. L. Colvin, M. C. Schlamp, and A. P. Alivisatos, Nature (London) **370**, 354 (1994).
- <sup>5</sup> P. T. Tran, E. R. Goldman, G. P. Anderson, J. M. Mauro, and H. Mattoussi, Phys. Status Solidi B **229**, 427 (2002).
- <sup>6</sup>M. Bruchez, M. Moronne, P. Gin, S. Weiss, and A. P. Alivisatos, Science **281**, 2013 (1998).
- <sup>7</sup>L. S. Li and A. P. Alivisatos, Phys. Rev. Lett. **90**, 097402 (2003).
- <sup>8</sup>M. G. Bawendi, A. R. Kortan, M. L. Steigerwald, and L. E. Brus, J. Chem. Phys. **91**, 7282 (1989).
- <sup>9</sup>T. Egami and S. J. L. Billinge, Underneath the Bragg Peaks: Structural Analysis of Complex Materials (Pergamon, New York/Elsevier, Oxford, 2003).
- <sup>10</sup>S. J. L. Billinge and M. G. Kanatzidis, Chem. Commun. (Cambridge) 2004, 749.
- <sup>11</sup>B. Gilbert, F. Huang, H. Zhang, G. A. Waychunas, and J. F. Banfield, Science **305**, 651 (2004).
- <sup>12</sup>R. B. Neder and V. I. Korsunskiy, J. Phys.: Condens. Matter 17, S125 (2005).
- <sup>13</sup>R. C. Howell, T. Proffen, and S. D. Conradson, Phys. Rev. B 73, 094107 (2006).
- <sup>14</sup> K. Page, T. Proffen, H. Terrones, M. Terrones, L. Lee, Y. Yang, S. Stemmer, R. Seshadri, and A. K. Cheetham, Chem. Phys. Lett. **393**, 385 (2004).
- <sup>15</sup>V. Petkov, M. Gateshki, J. Choi, E. G. Gillan, and Y. Ren, J. Mater. Chem. **15**, 4654 (2005).
- <sup>16</sup>M. Gateshki, V. Petkov, G. Williams, S. K. Pradhan, and Y. Ren,

Phys. Rev. B 71, 224107 (2005).

- <sup>17</sup> P. J. Chupas, X. Qiu, J. C. Hanson, P. L. Lee, C. P. Grey, and S. J. L. Billinge, J. Appl. Crystallogr. **36**, 1342 (2003).
- <sup>18</sup>T. Proffen, T. Egami, S. J. L. Billinge, A. K. Cheetham, D. Louca, and J. B. Parise, Appl. Phys. A: Mater. Sci. Process. **74**, s163 (2002).
- <sup>19</sup>T. Proffen and S. J. L. Billinge, J. Appl. Crystallogr. **32**, 572 (1999).
- <sup>20</sup>M. G. Tucker, M. P. Squires, M. T. Dove, and D. A. Keen, J. Phys.: Condens. Matter **13**, 403 (2001).
- <sup>21</sup>C. Kumpf et al., J. Chem. Phys. **123**, 224707 (2005).
- <sup>22</sup>J. N. Wickham, A. B. Herhold, and A. P. Alivisatos, Phys. Rev. Lett. **84**, 4515 (2000).
- <sup>23</sup>J. Nanda, S. Sapra, D. D. Sarma, N. Chandrasekharan, and G. Hodes, Chem. Mater. **12**, 1018 (2000).
- <sup>24</sup>B. Palosz, E. Grzanka, S. Gierlotka, S. Stel'Makh, R. Pielaszek, U. Bismayer, J. Neuefeind, H. P. Weber, and W. Palosz, Acta Phys. Pol. A **102**, 57 (2002).
- <sup>25</sup>H. Borchert, E. V. Shevchenko, A. Robert, I. Mekis, A. Kornowski, G. Grubel, and H. Weller, Langmuir **21**, 1931 (2005).
- <sup>26</sup>A. Guinier, X-ray Diffraction in Crystals, Imperfect Crystals, and Amorphous Bodies (Freeman, San Francisco, 1963).
- <sup>27</sup> P. Debye, Ann. Phys. **46**, 809 (1915).
- <sup>28</sup>Z. Kaszkur, Z. Kristallogr. 23, 147 (2006).
- <sup>29</sup>C. B. Murray, D. J. Norris, and M. G. Bawendi, J. Am. Chem. Soc. **115**, 8706 (1993).
- <sup>30</sup>C. Q. Chen, Y. Shi, Y. S. Zhang, J. Zhu, and Y. J. Yan, Phys. Rev. Lett. **96**, 075505 (2006).
- <sup>31</sup>G. Ouyang, X. L. Li, X. Tan, and G. W. Yang, Appl. Phys. Lett. 89, 031904 (2006).
- <sup>32</sup>M. A. Marcus, L. E. Brus, C. Murray, M. G. Bawendi, A. Prasad, and A. P. Alivisatos, Nanostruct. Mater. 1, 323 (1992).

- <sup>34</sup>A. C. Carter, C. E. Bouldin, K. M. Kemmer, M. I. Bell, J. C. Woicik, and S. A. Majetich, Phys. Rev. B 55, 13822 (1997).
- <sup>35</sup>S. Chaure, N. B. Chaure, and R. K. Pandey, Physica E (Amsterdam) **28**, 439 (2005).
- <sup>36</sup>X. G. Peng, J. Wickham, and A. P. Alivisatos, J. Am. Chem. Soc. 120, 5343 (1998).
- <sup>37</sup> V. Petkov, S. J. L. Billinge, J. Heising, and M. G. Kanatzidis, J. Am. Chem. Soc. **122**, 11571 (2000).
- <sup>38</sup> V. Petkov, P. N. Trikalitis, E. S. Božin, S. J. L. Billinge, T. Vogt, and M. G. Kanatzidis, J. Am. Chem. Soc. **124**, 10157 (2002).
- <sup>39</sup>V. Petkov, P. Y. Zavalij, S. Lutta, M. S. Whittingham, V. Parvanov, and S. Shastri, Phys. Rev. B **69**, 085410 (2004).
- <sup>40</sup> P. Juhas, D. M. Cherba, P. M. Duxbury, W. F. Punch, and S. J. L. Billinge, Nature (London) **440**, 655 (2006).
- <sup>41</sup>H. Z. Zhang, B. Gilbert, F. Huang, and J. F. Banfield, Nature (London) **424**, 1025 (2003).
- <sup>42</sup> V. Petkov, Y. Peng, G. Williams, B. Huang, D. Tomalia, and Y. Ren, Phys. Rev. B **72**, 195402 (2005).
- <sup>43</sup>B. E. Warren, X-ray Diffraction (Dover, New York, 1990).
- <sup>44</sup>A. P. Hammersley, S. O. Svenson, M. Hanfland, and D. Hauserman, High Press. Res. 14, 235 (1996).
- <sup>45</sup>X. Qiu, J. W. Thompson, and S. J. L. Billinge, J. Appl. Crystallogr. **37**, 678 (2004).
- <sup>46</sup>S. J. L. Billinge, in *Local Structure From Diffraction*, edited by S. J. L. Billinge and M. F. Thorpe (Plenum, New York, 1998), p. 137.
- <sup>47</sup>H. M. Rietveld, J. Appl. Crystallogr. 2, 65 (1969).
- <sup>48</sup>C. L. Farrow, P. Juhas, J. W. Liu, D. Bryndin, E. S. Božin, J. Bloch, T. Proffen, and S. J. L. Billinge, J. Phys.: Condens. Mat-

ter 19, 335219 (2007).

- <sup>49</sup>B. H. Toby and S. J. L. Billinge, Acta Crystallogr., Sect. A: Found. Crystallogr. **60**, 315 (2004).
- <sup>50</sup>B. H. Toby and T. Egami, Acta Crystallogr., Sect. A: Found. Crystallogr. 48, 336 (1992).
- <sup>51</sup>I.-K. Jeong, T. Proffen, F. Mohiuddin-Jacobs, and S. J. L. Billinge, J. Phys. Chem. A **103**, 921 (1999).
- <sup>52</sup>I. K. Jeong, R. H. Heffner, M. J. Graf, and S. J. L. Billinge, Phys. Rev. B 67, 104301 (2003).
- <sup>53</sup>C. Farrow (unpublished).
- <sup>54</sup>I. Kikuma and M. Furukoshi, J. Cryst. Growth 71, 136 (1985).
- <sup>55</sup> V. Petkov, R. G. DiFrancesco, S. J. L. Billinge, M. Acharya, and H. C. Foley, Philos. Mag. B **79**, 1519 (1999).
- <sup>56</sup>M. A. Marcus, W. Flood, M. Stiegerwald, L. Brus, and M. Bawendi, J. Phys. Chem. **95**, 1572 (1991).
- <sup>57</sup>J. Rockenberger, L. Troger, A. L. Rogach, M. Tischer, M. Grundmann, H. Weller, and A. Eychmuller, Phys. Chem. Chem. Phys. **102**, 1561 (1998).
- <sup>58</sup>Th. Proffen and R. B. Neder, J. Appl. Crystallogr. **30**, 171 (1997).
- <sup>59</sup>C. Landes, M. Braun, C. Burda, and M. El-Sayed, Nano Lett. 1, 667 (2001).
- <sup>60</sup>X. B. Chen, A. C. Samia, Y. B. Lou, and C. Burda, J. Am. Chem. Soc. **127**, 4372 (2005).
- <sup>61</sup>L. Qu, Z. Peng, and X. Peng, Nano Lett. 1, 333 (2001).
- <sup>62</sup>K. Kodama, S. Iikubo, T. Taguchi, and S. Shamoto, Acta Crystallogr., Sect. A: Found. Crystallogr. **62**, 444 (2006).
- <sup>63</sup> V. Petkov, I.-K. Jeong, J. S. Chung, M. F. Thorpe, S. Kycia, and S. J. L. Billinge, Phys. Rev. Lett. **83**, 4089 (1999).
- <sup>64</sup>I.-K. Jeong, F. Mohiuddin-Jacobs, V. Petkov, S. J. L. Billinge, and S. Kycia, Phys. Rev. B 63, 205202 (2001).