Structural and magnetic properties of Mn-implanted Si

Shengqiang Zhou, K. Potzger, Gufei Zhang, A. Mücklich, F. Eichhorn, N. Schell, R. Grötzschel, B. Schmidt,

W. Skorupa, M. Helm, and J. Fassbender

Institute for Ion Beam Physics and Materials Research, Forschungszentrum Rossendorf, POB 510119, 01314 Dresden, Germany

D. Geiger

Institute of Structure Physics, Dresden University, Zellescher Weg 16, 01062 Dresden, Germany

(Received 6 January 2006; revised manuscript received 11 December 2006; published 7 February 2007)

Structural and magnetic properties in Mn-implanted, *p*-type Si were investigated. High resolution structural analysis techniques such as synchrotron x-ray diffraction revealed the formation of $MnSi_{1.7}$ nanoparticles already in the as-implanted samples. Depending on the Mn fluence, the size increases from 5 nm to 20 nm upon rapid thermal annealing. No significant evidence is found for Mn substituting Si sites either in the as-implanted or annealed samples. The observed ferromagnetism yields a saturation moment of $0.21\mu_B$ per implanted Mn at 10 K, which could be assigned to $MnSi_{1.7}$ nanoparticles as revealed by a temperature-dependent magnetization measurement.

DOI: 10.1103/PhysRevB.75.085203

PACS number(s): 75.50.Pp, 82.80.Yc, 68.37.Lp, 61.10.Nz

I. INTRODUCTION

Recently, diluted magnetic semiconductors (DMS) have attracted huge research attention.^{1,2} III-V and II-IV compound semiconductors doped with transition metals (TM) by ion implantation show ferromagnetism with a Curie temperature up to room temperature.^{3,4} However, Si-based DMS would be preferable to compound semiconductors due to the availability of high-quality Si in large sizes at relatively low costs. More importantly, the fabrication of Si-based DMS is compatible with the mature microelectronics technique. Based on the Zener model, Dietl et al.⁵ predicted a carriermediated ferromagnet if Si is doped with 5% Mn. Zhang et $al.^{6}$ reported on the crystalline Mn_{0.05}Si_{0.95} alloy with a Curie temperature up to 400 K. Also, epitaxially grown Mn_{0.07}Si_{0.93} thin films show ferromagnetic ordering up to 200 K, and at the same time, possess semiconducting properties.⁷ In addition to the aforementioned epitaxial techniques, ion implantation is another promising method for ferromagnetic doping of semiconductors, e.g., implantation of almost any TM element with an atomic concentration exceeding those possible in thermal equilibrium by far. Very recently. Bolduc *et al.*⁸ have reported room-temperature ferromagnetism of Mn-ion implantation Si at Mn fluences of 1×10^{15} cm⁻² and 1×10^{16} cm⁻². The Curie temperature has been evaluated to be above 400 K for both *n*-type and *p*-type Si. The strength of the ferromagnetism was affected by the Mn concentration, thermal annealing, and carrier type of the Si substrate. The ferromagnetic coupling was believed to be carrier mediated. However, the high-resolution structural characterization of the investigated Si:Mn systems (both epitaxially grown and ion implanted) was rather limited. The diffusion of Mn in Si upon annealing was only investigated by Rutherford backscattering (RBS) and secondary ion mass spectrometry.^{9,10} Generally Mn silicides are easily formed. For example, if Mn is deposited on Si, an epitaxial MnSi is formed.¹¹⁻¹³ Also transition metal (i.e., Co and Ni) implantation into Si can easily leads to silicides.^{14,15} Therefore a careful characterization of the Mn-doped Si samples has to be carried out in order to clarify the origin of the ferromagnetism, i.e., the formation of nanoparticles or a DMS system. For this purpose, in the present paper, RBS and/or channeling, transmission electron microscopy (TEM), and synchrotron radiation x-ray diffraction (SR-XRD), are used to clarify the possible origin of the ferromagnetism in Mn-implanted Si.

II. EXPERIMENT

The Mn-implanted Si samples were prepared from commercially available, Czochralski-grown single-crystal Si(001) wafers, which were *p*-type doped with a B concentration of 1.2×10^{19} cm⁻³. Mn⁺ ions were implanted at an energy of 300 keV with fluences of 1×10^{15} cm⁻², 1×10^{16} cm⁻², and 5×10^{16} cm⁻², which corresponds to Mn peak concentrations of 0.08, 0.8, and 4 at. %, respectively, with a projected range (R_p) of 258 ± 82 nm. The samples were held at 350 °C during implantation to avoid amor-



FIG. 1. RBS random and channeling along Si[001] spectra of as-implanted Si samples. Random (\bigcirc), channeling (+) for the fluence of $5 \times 10^{16} \text{ cm}^{-2}$; channeling (\blacksquare) for the fluence of $1 \times 10^{16} \text{ cm}^{-2}$; channeling (\triangle) for the fluence of $1 \times 10^{15} \text{ cm}^{-2}$; and channeling for virgin Si (—) which is slightly below triangle symbols.

Fluence and process	$1 \times 10^{15} \text{ cm}^{-2}$ as implanted	$1 \times 10^{15} \text{ cm}^{-2},$ RTA ^a	$1 \times 10^{16} \text{ cm}^{-2}$ as implanted	$1 \times 10^{16} \text{ cm}^{-2},$ RTA ^a	$5 \times 10^{16} \text{ cm}^{-2}$ as implanted	$5 \times 10^{16} \text{ cm}^{-2},$ RTA ^a
$\chi_{\rm min}$ -Si (RBS/C)	7.4%	6.4%	51%	41%	82%	70%
χ_{\min} -Mn (RBS/C) ^b	_	_	91%	89%	96%	98%
Conv. XRD ^c	No	No	No	No	No	No
TEM ^c	No	Yes	No	Yes	Yes	Yes
SR-XRD ^c	No	No	No	Yes	Yes	Yes
				(11 nm) ^d	(6 nm) ^d	(15 nm) ^d
Ferromagnetism	No	No	No	Strong	Weak	Weak
(10 K)				$(0.21 \mu_{\rm B}/{\rm Mn})$	$(0.04 \mu_{\rm B}/{\rm Mn})$	$(0.02\mu_{\rm B}/{\rm Mn})$

TABLE I. Map of structural and magnetic properties of all samples.

^aRTA: rapid thermal annealing in a forming gas of N_2 at 800 °C for 5 min.

^bFor the Mn signal, a very weak channeling effect was observed, which reveals that the majority of implanted Mn ions are not substituted on Si sites.

c"Yes" or "No" to answer if precipitates of Mn silicides are detectable by corresponding technique.

^dAverage grain size.

phization. In order to reduce channeling effects, the angle between the sample surface normal and the incident beam was set to 7°. After implantation, rapid thermal annealing (RTA) was performed at a temperature of 800 °C for 5 min in a forming gas of N₂. The RBS spectra were taken with a collimated 1.7 MeV He⁺ beam with a surface barrier detector at 170°. TEM specimens were conventionally prepared for cross-section inspection and examined on a Philips CM 30 FEG microscope. Conventional XRD (conv. XRD) measurements were performed on a Siemens D5005 diffractometer using Cu irradiation. SR-XRD was performed at Rossendorf beamline (ROBL) at the European Synchrotron Radiation Facility with monochromatic x rays of 0.1541-nm wavelength. Magnetic properties were analyzed using a superconducting quantum interference device (SQUID) magnetometer (Quantum Design MPMS).

III. RESULTS AND DISCUSSION

A. Mn lattice location

After Mn-ion implantation, the samples were analyzed by RBS and/or channeling to check the implantation-induced damage and the lattice location of Mn with respect to the Si lattice. Figure 1 shows the RBS and/or channeling spectra for all samples. The Mn signals display a Gaussian distribution which corresponds to an R_p of around 258 nm. For samples with a fluence of 1×10^{16} cm⁻² and 5×10^{16} cm⁻², the R_p and the amount of Mn agree well with SRIM simulation¹⁶ and designed values. However, for the lowest fluence sample $(1 \times 10^{15} \text{ cm}^{-2})$, the concentration of Mn is in the order of the detection limit. All χ_{\min} values, the ratio of the backscattering yield at channeling condition to that for a random beam incidence,¹⁷ which label the degree of lattice disordering, for Si and Mn signals are listed in Table I. As expected, a higher fluence implantation introduces a larger lattice disordering (a larger χ_{\min}) in the silicon substrates. For the Mn signal, χ_{\min} amounts to more than 90% for both samples with the fluence of $1 \times 10^{16} \text{ cm}^{-2}$ and 5 $\times 10^{16}$ cm⁻², which proves that the majority of the implanted Mn atoms are not substituted to Si sites with respect to the [001] crystallographic axis.

Figure 2 shows the RBS spectra for the sample implanted with 5×10^{16} cm⁻² Mn ions after RTA annealing at 800 °C for 5 min. The changes of χ_{min} values are listed in Table I. The annealing process removed only a very small fraction of the implantation-induced damage (a slight decrease of χ_{min}) and did not change the Mn element distribution, i.e., Mn atoms did not diffuse to deep or shallow places, nor to substitutional Si sites. The other samples show similar behavior after annealing (see Table I).

B. Observation of crystalline MnSi_x nanoparticles

In order to confirm the formation of Mn-silicides nanoparticles, high resolution TEM (HRTEM) was performed.



FIG. 2. RBS random and channeling along Si[001] spectra for sample implanted with a fluence of 5×10^{16} cm⁻². Random (\bigcirc) and channeling (+) for as-implanted sample, channeling (—) for RTA sample, and channeling (\triangle) for virgin Si. Only a small fraction of damage was recovered by annealing. For Mn signals, no channeling was observed either for as-implanted or annealed samples, which means Mn does not substitute Si lattice site.



FIG. 3. TEM-overview images show the implantation induced damage and nanosized precipitates after RTA: (a) 1×10^{15} cm⁻², (b) 1×10^{16} cm⁻², and (c) 5×10^{16} cm⁻². The nanosized precipitates are visible as spotlike contrast, and they are growing with increasing fluence. The arrow marks the direction towards the sample surface.

For the highest fluence $(5 \times 10^{16} \text{ cm}^{-2})$, some small nanoparticles with diameters of ~ 5 nm are detected already in the as-implanted samples (not shown). Figure 3 shows typical TEM-overview images at lower magnification of samples after RTA treatment. Nanosized precipitates are unambiguously found in all annealed samples with different fluence. They exhibit mean diameters of 5 nm for a fluence of 1 $\times 10^{15}$ cm⁻², 10 nm for a fluence of 1×10^{16} cm⁻², and 20 nm for a fluence of 5×10^{16} cm⁻². Figures 4(a) and 4(b) show HRTEM images for the samples of 1E16 RTA and 5E16 RTA (see Table I for the sample abbrevation). Most of those nanoparticles are spherical-like, their diameters increase from around 10 nm to 20 nm with increasing fluence. However, considering the rich varieties of Mn silicides.¹³ it is extremely difficult to identify these nanoparticles even by HRTEM.¹⁸ Table II lists the structural information of some higher manganese silicides with Si composition of around 1.7 after Ref. 13. All different compositions are named as MnSi₁₇ throughout the rest of the article. The crystallite area of some particles was Fourier-analysed as shown in Fig. 4(c). The reflections in the diffractogram indicated by arrows reveal a lattice spacing of around 1.7 nm, which can be attributed to $MnSi_{1.7}$ (see Table II).

In addition, XRD was performed in order to confirm the formation of Mn-silicide nanoparticles. Since a conventional XRD diffractometer (Siemens D5005) failed to detect any crystalline nanoparticles in both as-implanted and annealed samples $(5 \times 10^{16} \text{ cm}^{-2})$ due to the small amount of nanoparticles and the limited intensity of Cu-target x-ray source, we performed XRD measurements at a synchrotron radiation

source. However, also at synchrotron XRD, in a symmetric beam geometry (Bragg-Brentano) one fails to detect any Mn silicides in all samples (not shown). Therefore, a grazing incidence geometry was used during the measurement. Grazing incidence x-ray diffraction (GIXRD) refers to a method where the incident x-ray beam is aligned at a small (here (0.4°) angle to the sample surface. This has the advantage to limit the penetration depth of the x rays into the sample with consequently low background scattering from the substrate. Figure 5 shows GIXRD patterns of all samples. The diffraction peaks at around 42° and 46.3° cannot be attributed to the Si substrate, but to MnSi_{1.7} (see Table II). For the highest fluence sample $(5 \times 10^{16} \text{ cm}^{-2})$, nanoparticles were observed already in the as-implanted state, and they grew from around 5 nm to 15 nm in grain size after annealing. MnSi_{1.7} was also found for a fluence of 1×10^{16} cm⁻² after annealing, while there is no indication of any crystalline nanoparticles in the lowest fluence sample $(1 \times 10^{15} \text{ cm}^{-2})$, even after annealing. Using the Scherrer formula,¹⁹ the average size of these nanoparticles is calculated, and listed in Table I.

C. Magnetic properties

Figure 6 shows the magnetization versus field curves (*M-H* loops) recorded at 10 K. The sample of 1E16 RTA shows a clear hysteretic behavior with a saturation of $0.21 \mu_{\rm B}$ /Mn and a coercivity of 275 ± 25 Oe. In contrast, both samples of 5E16 as implanted and 5E16 RTA show very weak hysteretic behaviors. All other samples (1E16 as implanted, 1E15 as implanted, and 1E15 RTA) did not show

TABLE II. Composition and lattice parameters of tetragonal MnSi_{1.7} phases (after Ref. 13).

Phase	Composition	<i>a</i> (nm)	<i>c</i> (nm)	Space group	PDF number ^a
Mn ₄ Si ₇	1.750	0.5525	1.7463	<i>P</i> -4 <i>c</i> 2	72-2069
Mn ₁₁ Si ₁₉	1.727	0.5518	4.8136	P - 4n2	_
Mn ₁₅ Si ₂₆	1.733	0.5525	6.5550	I-42d	20-724
Mn ₂₇ Si ₄₇	1.741	0.5530	11.7940	P - 4n2	26-1251

^aInorganic crystal structure database (ICSD)



FIG. 4. High-resolution TEM images for representative $MnSi_x$ nanoparticles in (a) 1E16 RTA, (b) 5E16 RTA, and (c) the diffractogram of one nanoparticle in (b). The reflection spots marked by the arrows in (c) reveal a lattice spacing of around 1.7 nm, which can be attributed to $MnSi_{1.7}$.

any detectable magnetism (not shown here). In order to clarify the origin of the ferromagnetism of the sample 1E16 RTA, zero-field cooled (ZFC) and field cooled (FC) magnetization curves were measured in a field of 50 Oe from



FIG. 5. (Color online) XRD grazing incidence scans of all investigated samples implanted with Mn. In some samples two peaks arise, which can be attributed to Mn-silicide nanoparticles (MnSi_{1,7}).

10 to 300 K (the inset of Figure 6). The ZFC curve shows a gradual increase at low temperature, and reaches a maximum at a temperature of T_B (around 23 K), while the FC curve continues to decrease with increasing temperature. At a higher temperature than T_B , the FC curve merges together with the ZFC curve. Figure 6 also shows a *M*-*H* loop taken at 100 K. Below T_B (10 K), these nanoparticles show ferromagnetic behavior, while above T_B (100 K), the *M*-*H* loop shows neither cleary remanence nor coercivity, and these nanoparticles are superparamagnetic. The above-mentioned magnetic behavior is typical for systems containing magnetic nanoparticles.^{20–22} Moreover, the wasp-waist shape of the loop is associated with magnetic phases with different coer-



FIG. 6. (Color online) M-H hysteresis loops for Mn-implanted Si samples. Only the sample of 1E16 RTA exhibits a clear hysteretic behavior. Inset: ZFC (open circle) and FC (solid line) magnetization curves reveal a typical characteristic of a magnetic nanoparticle system for the sample of 1E16 RTA. Above the blocking temperature, the sample of 1E16 RTA shows superparamagnetic properties, where neither coercivity nor remanence was observed at 100 K. None of the other samples (1E15 RTA, 1E15 as implanted, and 1E16 as implanted) exhibit ferromagnetism, even at temperatures down to 10 K (not shown).

civities, which could be attributed to the size distribution of nanoparticles as indicated by HRTEM results.^{23–25} Although the sample preparation in our work is identical with Ref. 8, the magnetic properties are drastically different. There could be some other parameters, or technical tricks in sample preparation to result in the discrepancy.

D. Discussion

In Table I, all structural and magnetic properties of the investigated samples are listed. Obviously the dependence of magnetization on the fluence and annealing is rather anomalous. In order to explain the magnetization, the magnetic properties for Mn silicides have to be reviewed. Unfortunately, very limited work has been done on the magnetism of $MnSi_{1,7}$ ²⁶ In Ref. 26, single crystalline Mn_4Si_7 (MnSi_{1,75}) were reported to exhibit weak itinerant magnetism with an ordering temperature of 47 K and with a very low saturation moment of $0.012\mu_{\rm B}/{\rm Mn}$. For the nanoparticles embedded in Si, i.e., with different surface and/or volume ratio and undergoing different pressure, the magnetism can be much different from bulk crystals. During intense investigations on MnSi (a weak itinerant magnetic material), the ordering temperature has been found to strongly depend on lattice strain, induced, e.g., by hydrostatic pressure or by the lattice mismatch with substrate.^{27–30} If the exchange interaction is changed for instance by pressure or by alloying, the system no longer orders even at T=0 K.²⁷ Considering the different grain size of MnSi1.7 in our sample, i.e., different stress provided by the surrounding lattice, we could explain the dependence of the magnetization on the Mn fluence and postimplantation annealing as follows. Those $MnSi_{1.7}$ nanoparticles in the sample (1E16 RTA) have a particular medium grain size, i.e., particular lattice deformation or strain, therefore a higher ordering temperature and even a much higher moment (0.21 μ_B/Mn) than the bulk crystal (0.012 μ_B/Mn), while the bigger nanoparticles in the sample of 5E16 RTA are more bulklike and have very weak ferromagnetism. Those smaller MnSi_{1.7} nanoparticles in other samples (5E16 as implanted and 1E15 RTA) could be in a rather worse crystalline quality or different strain status, and therefore only show weak or nondetectable ferromagnetism. However a nonambiguous reason for the anomalous dependence of the magnetization on the particle size is not yet known.

IV. CONCLUSION

By investigating Mn-ion implanted Si, we have found a nonsystematic dependence of the induced ferromagnetization on the implantation fluence. There was no significant indication for Mn substituting Si either in the as-implanted or the rapidly annealed samples. The observed ferromagnetization is maximized at a certain medium grain size of $MnSi_{1.7}$ nanoparticles, which could result in a particular strain status, consequently a stronger ferromagnetic coupling. Moreover, high resolution analysis techniques are necessary in order to identify those nanoparticles and to clarify the origin of the observed ferromagnetism.

ACKNOWLEDGMENT

We thank S. Gemming for valuable discussion.

- ¹S. J. Pearton, C. R. Abernathy, D. P. Norton, A. F. Hebard, Y. D. Park, L. A. Boatner, and J. D. Budai, Mater. Sci. Eng., R. **40**, 137 (2003).
- ² A. H. Macdonald, P. Schiffer, and N. Samarth, Nat. Mater. **4**, 195 (2005).
- ³N. Theodoropoulou, A. F. Hebard, M. E. Overberg, C. R. Abernathy, S. J. Pearton, S. N. G. Chu, and R. G. Wilson, Appl. Phys. Lett. **78**, 3475 (2001).
- ⁴M. Venkatesan, C. B. Fitzgerald, J. G. Lunney, and J. M. D. Coey, Phys. Rev. Lett. **93**, 177206 (2004).
- ⁵T. Dietl, H. Ohno, F. Matsukura, J. Cibert, and D. Ferrand, Science **287**, 1019 (2000).
- ⁶F. M. Zhang, X. C. Liu, J. Gao, X. S. Wu, Y. W. Du, H. Zhu, J. Q. Xiao, and P. Chen, Appl. Phys. Lett. **85**, 786 (2004).
- ⁷H. M. Kim, T. W. Kang, and K. S. Chung, J. Ceram. Proc. Res. **5**, 238 (2004).
- ⁸M. Bolduc, C. Awo-Affouda, A. Stollenwerk, M. B. Huang, F. G. Ramos, G. Agnello, and V. P. LaBella, Phys. Rev. B **71**, 033302 (2005).
- ⁹B. E. Egamberdiev and M. Yu. Adylov, Tech. Phys. Lett. 27, 168 (2001).
- ¹⁰H. Francois-Saint-Cyr, E. Anoshkina, F. Stevie, L. Chow, K. Richardson, and D. Zhou, J. Vac. Sci. Technol. B **19**, 1769 (2001).
- ¹¹Y. C. Lian and L. J. Chen, Appl. Phys. Lett. 48, 359 (1986).

- ¹²S. Teichert, S. Schwendler, D. K. Sarkar, A. Mogilatenko, M. Falke, G. Beddies, C. Kleint, and H.-J. Hinneberg, J. Cryst. Growth **227-228**, 882 (2001).
- ¹³A. Mogilatenko, Ph.D thesis, Technical University Chemnitz, 2002.
- ¹⁴A. Vantomme, M. F. Wu, I. Dézsi, G. Langouche, K. Maex, and J. Vanhellemont, Mater. Sci. Eng., B 4, 157 (1989).
- ¹⁵ M. F. Wu, J. De Wachter, A.-M. Van Bavel, R. Moons, A. Vantomme, H. Pattyn, G. Langouche, H. Bender, J. Vanhellemont, K. Temst, and Y. Bruynseraede, J. Appl. Phys. **78**, 1707 (1995).
- ¹⁶J. F. Ziegler, J. P. Biersack, and U. Littmark, *The Stopping and Range of Ions in Solids* Stopping and Ranges of Ions in Matter Vol. 1 (Pergamon Press, New York 1984).
- ¹⁷W. K. Chu, J. W. Mayer, and M. A. Nicolet, *Backscattering Spectrometry* (Academic, New York, 1978).
- ¹⁸M. Eizenberg and K. N. Tu, J. Appl. Phys. **53**, 6885 (1982).
- ¹⁹B. D. Cullity, *Elements of X-ray Diffractions* (Addison-Wesley, Reading MA, 1978), p. 102.
- ²⁰ M. Respaud, J. M. Broto, H. Rakoto, A. R. Fert, L. Thomas, B. Barbara, M. Verelst, E. Snoeck, P. Lecante, A. Mosset, J. Osuna, T. O. Ely, C. Amiens, and B. Chaudret, Phys. Rev. B 57, 2925 (1998).
- ²¹Y. Sun, M. B. Salamon, K. Garnier, and R. S. Averback, Phys. Rev. Lett. **91**, 167206 (2003).

- ²²P. Dutta, A. Manivannan, M. S. Seehra, N. Shah, and G. P. Huffman, Phys. Rev. B **70**, 174428 (2004).
- ²³L. H. Bennett and E. D. Torre, J. Appl. Phys. **97**, 10E502 (2005).
- ²⁴F. Brem, L. Tiefenauer, A. Fink, J. Dobson, and A. M. Hirt, Phys. Rev. B **73**, 224427 (2006).
- ²⁵ Shengqiang Zhou, K. Potzger, Gufei Zhang, F. Eichhorn, W. Skorupa, M. Helm, and J. Fassbender, J. Appl. Phys. **100**, 114304 (2006).
- ²⁶U. Gottlieb, A. Sulpice, B. Lambert-Andron, and O. Laborde, J. Alloys Compd. **361**, 13 (2003).
- ²⁷C. Pfleiderer, S. R. Julian, and G. G. Lonzarich, Nature (London) 414, 427 (2001).
- ²⁸K. Schwinge, C. Mueller, A. Mogilatenko, J. J. Paggel, and P. Fumagalli, J. Appl. Phys. **97**, 103913 (2005).
- ²⁹K. Koyama, T. Goto, T. Kanomata, and R. Note, Phys. Rev. B 62, 986 (2000).
- ³⁰W. Yu, F. Zamborszky, J. D. Thompson, J. L. Sarrao, M. E. Torelli, Z. Fisk, and S. E. Brown, Phys. Rev. Lett. **92**, 086403 (2004).