

## Electronic structure of solid uranium tetrafluoride UF<sub>4</sub>

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X-ray photoelectron spectra (XPS) and conversion electron spectra of the outer (0–15 eV) and inner (15–40 eV) valence electrons for UF<sub>4</sub> were measured. Relativistic  $X_\alpha$  discrete variation ( $RX_\alpha$  DV) calculation data for the UF<sub>8</sub><sup>4-</sup> ( $C_{4v}$ ) cluster reflecting uranium close environment in solid UF<sub>4</sub> were used for the quantitative interpretation of the fine spectral structure. Quantitative agreement between the experimental and theoretical data was established. The U 5*f* electrons ( $\approx 1$  U 5*f* electron) were shown to participate directly in the chemical bond formation. This U 5*f* electron was shown to be delocalized within the outer valence molecular orbitals (OVMO) range (1–15 eV). The other U 5*f* electrons were shown to be localized and to participate weakly in the chemical bond formation. The XPS line associated with these electrons was observed at 3.8 eV. The vacant U 5*f* states are generally delocalized in the range of the low positive energies (0–7 eV). The contribution of the U 6*p* electronic density to the molecular orbitals of UF<sub>4</sub> was experimentally and theoretically evaluated. The U 6*p* electrons were experimentally shown to participate significantly (0.6 U 6*p* electrons) in the formation of the OVMO beside the formation of the inner valence molecular orbitals (IVMO). IVMO composition and sequence order in the binding energy range 15–40 eV in UF<sub>4</sub> were determined.

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### I. INTRODUCTION

X-ray photoelectron spectroscopy (XPS) studies of uranium, uranium fluorides and other compounds have revealed a complex fine structure in the low binding energy ( $E_b$ ) range.<sup>1–6</sup> For example, the XPS from solid uranium tetrafluoride UF<sub>4</sub> and uranyl fluoride UO<sub>2</sub>F<sub>2</sub> show a great difference in the structures of all the valence binding energy ( $E_b$ ) range 0–40 eV (Refs. 5 and 7) since uranium U<sup>6+</sup> ion's electronic configuration in UO<sub>2</sub>F<sub>2</sub> is {Rn}5*f*<sup>0</sup>, and that of U<sup>4+</sup> one in UF<sub>4</sub> is {Rn}5*f*<sup>2</sup>, where {Rn}—radon electronic configuration. Therefore, the XPS from UF<sub>4</sub> exhibits a relatively sharp (1.5 eV wide) peak at  $E_b=3.8$  eV attributed to the U 5*f* electrons weakly participating in the chemical bond. The XPS from UO<sub>2</sub>F<sub>2</sub> does not exhibit this peak.

The XPS peaks in the  $E_b$  range 0–40 eV from uranium fluorides and other compounds were observed<sup>7,8</sup> to be several eV wide. For example, for UF<sub>4</sub>, the F 1*s* peak's ( $E_b=685.3$  eV) full width at half-maximum (FWHM) is  $\Gamma=1.3$  eV, while the corresponding F 2*s* peak ( $E_b=29.9$  eV) is 3.7 eV wide and structured. The extra structure was observed on the both sides from the expected single F 2*s* peak. This F 2*s* widening (relative to the F 1*s* FWHM) contradicts the uncertainty ratio  $\Delta E \Delta \tau \sim h/2\pi$ , where  $\Delta E$ —natural width of the level from which an electron was removed,  $\Delta \tau$ —lifetime of a hole and  $h$ —Planck's constant. Since the hole lifetime ( $\Delta \tau$ ) decreases as the absolute energy of a level grows, the XPS atomic peaks are expected to narrow as the electron binding energy decreases. In case of UF<sub>4</sub> and UO<sub>2</sub>F<sub>2</sub> the picture is reversed. One of the reasons for the XPS peak widening in the binding energy range 0–40 eV was found to

be the formation of molecular orbitals (MOs).<sup>8–10</sup> These MOs form generally from the An 6*p* and L *ns* atomic shells of actinides and ligands (L). The formed MOs can be subdivided into the outer valence molecular orbitals (OVMO) (0–15 eV  $E_b$ ) and inner valence molecular orbitals (IVMO) (15–40 eV  $E_b$ ). Practically, the XPS spectra reflect the structure of the valence band (0–40 eV) and are observed as several eV wide bands. It was shown that under favorable conditions the IVMO could form in compounds of any elements of the periodic table.<sup>8–11</sup>

According to the earlier suggestions, the An 5*f* electrons are supposed to get promoted, for example, to the An 6*d* atomic orbitals before the chemical bond formation. However, the theoretical calculations show that the An 5*f* atomic shells can participate directly in the MO formation in actinide compounds.<sup>7,8</sup> Therefore, it is important to determine the experimental An 5*f* and An 6*p* partial electronic densities.

The qualitative identification of the XPS data for uranium tetrafluoride<sup>7</sup> allowed a qualitative interpretation of the conversion electron (CES) (Ref. 12) and x-ray O<sub>4,5</sub>(U) emission (XES) (Ref. 13) spectral structures. The absence of the relativistic electronic structure calculations did not allow a correct interpretation of these spectral structures. The calculations are too complicated because of the fact that uranium environment in solid tetrafluoride cluster (UF<sub>8</sub><sup>4-</sup> of symmetry group  $D_{4d}$ ) is too complex. The authors of Ref. 14 made an attempt to interpret qualitatively the electronic structure of solid UF<sub>4</sub> on the basis of the relativistic calculation results for the UF<sub>4</sub> ( $T_d$ ) cluster reflecting uranium environment in the gaseous tetrahedral molecule. Earlier the XPS, CES and

$O_{4,5}(U)$  XES results taking into account the relativistic calculations were used for the quantitative evaluation of the partial densities of the U  $6p$  and  $5f$  electronic states in the binding energy range 0–40 eV in  $UO_2$ ,<sup>15</sup>  $\gamma-UO_3$ ,<sup>16</sup> and  $UO_2F_2$ .<sup>17</sup>

This work analyzes the fine low binding energy (0–40 eV) high resolution XPS and CES structures from  $UF_4$  taking into account relativistic  $X_\alpha$  discrete variation method ( $RX_\alpha$  DVM) electronic structure calculation for the  $UF_8^{4-}$  ( $C_{4v}$ ) cluster reflecting uranium close environment in solid  $UF_4$ . As a result uranium electronic structure was interpreted quantitatively and the U  $6p$ ,  $5f$  partial electronic densities in uranium tetrafluoride were evaluated experimentally.

## II. EXPERIMENTAL

XPS spectra of the solid  $UF_4$  sample were measured with an electrostatic spectrometer HP 5950A Hewlett-Packard using monochromatized Al  $K_{\alpha 1,2}$  ( $h\nu=1486.6$  eV) radiation in a vacuum of  $1.3 \times 10^{-7}$  Pa at room temperature. The device resolution measured as full width ( $\Gamma$ , eV) on the half-maximum (FWHM) of the Au  $4f_{7/2}$  peak on the standard rectangular golden plate was 0.8 eV. The binding energies  $E_b$  (eV) were measured relative to the binding energy of the C  $1s$  electrons from hydrocarbons absorbed on the sample surface accepted to be equal to 285.0 eV. On the gold plate  $E_b(C 1s)=284.7$  eV at  $E_b(Au 4f_{7/2})=83.8$  eV. The FWHM were measured relative to the width of the C  $1s$  line of hydrocarbons accepted to be equal to  $\Gamma(C 1s)=1.3$  eV. The errors in determination of electron binding energies and the linewidths did not exceed 0.1 eV and that of the relative line intensities was less than 10%.

The  $UF_4$  sample for the XPS studies was prepared from the fine powder ground in the agate mortar as a thick dense flat layer pressed into In on a Ti substrate. The U  $4f_{7/2}$  binding energy was measured to be 482.7 eV. Oxygen concentration determined on the basis of O  $1s$  XPS peak was found to be less than 4 at. %, which testifies to the low surface oxidation. The CES spectrum was measured with the same spectrometer using an additional accelerating electronic system. A metallic copper substrate with evaporated  $UF_4$  layer for uranium implantation by the electrostatic collection of  $^{235m}U$  recoil atoms resulting from  $^{239}Pu$  decay in the inert atmosphere was used.<sup>12</sup> Peak identification and calibration of the CES was done using the XPS data for  $UF_4$ . Elastic scattering related background in the XPS was subtracted by Shirley.<sup>18</sup> For the CES the background was subtracted by Shirley and by exponent. The maximal difference in the relative intensities did not exceed 10%. This work gives the CES spectrum with the background subtracted by exponent (Fig. 1).

The  $UF_8^{4-}$  cluster of symmetry group  $C_{4v}$  reflecting uranium close environment in solid  $UF_4$  can be constructed if to rotate by  $45^\circ$  an upper cube face with the edge equal to 2.64 Å. Uranium ion located in the center of this cube is surrounded by fluorines located in the cube corners at  $R_{U-F}=2.29$  Å from uranium.<sup>19</sup> The calculation for this cluster was done in the  $X_\alpha$  DVM approximation<sup>20,21</sup> based on the Dirac-Slater equation for the four component spinors with

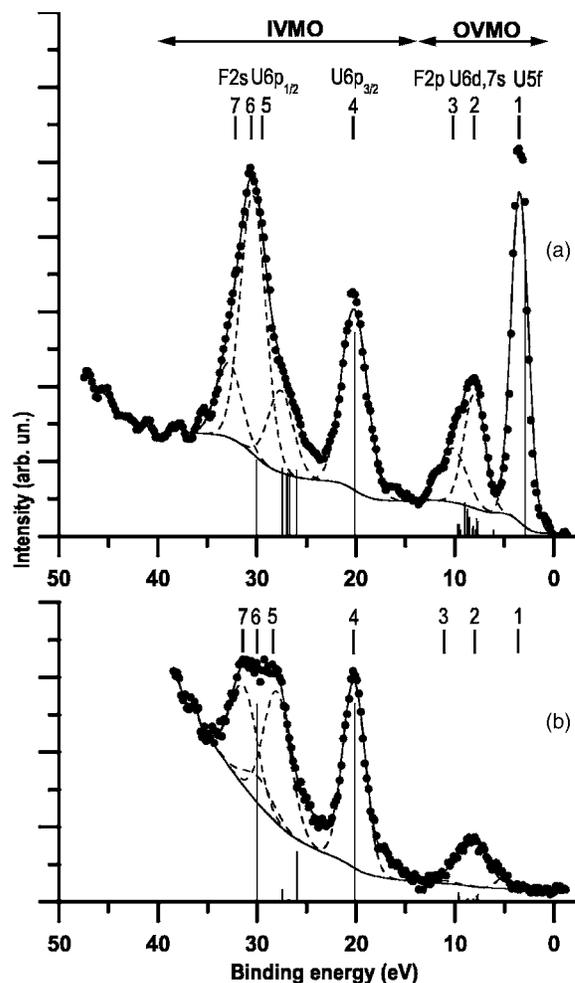


FIG. 1. XPS (a) and CES (b) from solid  $UF_4$ . The corresponding theoretical spectra are given under the experimental spectra as vertical bars. The shape of subtracted background and spectra decomposition are shown. The experimental spectral intensities are given in arbitrary units; the theoretical intensities are normalized in %.

exchange-correlation potential.<sup>22</sup> The extended basis of the numerical atomic orbitals (AO) from the solution of Dirac-Slater equation for the isolated atoms beside the filled included the vacant U  $7p_{1/2}$ ,  $7p_{3/2}$  states. The basis also took into account the cluster symmetry, i.e., by the technique of projecting operators<sup>20</sup> the regular AOs were reconstructed in the linear combinations converting by the irreducible representations of the binary group  $C_{4v}$ . For the relativistic basis calculations an original program of symmetrization was utilized. This program used the matrices of irreducible representations for the most part of binary groups obtained in Ref. 22 and transformation matrices given in Refs. 23 and 24. The numerical diophantine integration during the calculation of the secular equation matrix elements was done by the number of 22 000 points spread in the cluster space. It provided the convergence of MO energies of not worse than 0.1 eV. The local exchange-correlation potential was taken as  $X_\alpha$  with  $\alpha$  equal to the mean atomic value. Since the clusters were the fragments of the crystal, the ligand AO population renormalization during the self-consistent was done. It allowed an effective account of the stoichiometry and charge

redistribution between the ligands and surrounding crystal.

### III. RESULTS AND DISCUSSION

The low binding energy (0–40 eV) XPS from  $\text{UF}_4$  can be conditionally subdivided into the two ranges (Fig. 1). The first one 0–15 eV exhibits the structure attributed to the OVMO built mostly from the incompletely filled outer U  $5f$ ,  $6d$ ,  $7s$  and F  $2p$  AOs (Table I). The second one 15–40 eV shows the IVMO related fine structure. These IVMOs are built mostly from the completely filled inner valence U  $2p$  and F  $2s$  AOs. The fact that the IVMO XPS parameters correlate with uranium close environment structure in compounds encouraged this subdivision.<sup>8,9,11</sup> The OVMO XPS structure has typical features and can be subdivided into the three components. The IVMO spectral range exhibits pronounced peaks and can be subdivided into the four components (Fig. 1). The low intense shoulders on the lower binding energy side can be attributed to uranium oxide on the sample surface, as well as to the background subtraction imperfection. The areas under these peaks were taken into account. This subdivision allows the qualitative and quantitative comparisons between the XPS, CES and relativistic calculation results for the  $\text{UF}_8^{4-}$  ( $C_{4v}$ ) cluster.

Relativistic calculation results are given in Table I. Since the XPS reflect both ground and excited states of ions, the calculated binding energies for the transition (not ground) states must be used<sup>25</sup> for comparison between the theoretical and experimental data. However, it is known<sup>8,11</sup> that in a rough approximation one can suggest that for the valence region the binding energies for the transition state differ from those for the ground state by a constant shift toward the higher absolute energy. Therefore, the present work used the theoretical binding energies increased by 2.87 eV for comparison with the corresponding experimental values (Table II). Taking into account the MO compositions (Table I) and photoionization cross sections  $\sigma_i$  (Ref. 26) ( $\sigma_i$  for the U  $7p$  were calculated by Yarbemsky), the theoretical spectral intensities for the considered energy ranges were determined (Table II, Fig. 1). Experimental XPS intensities are given for comparison. A good qualitative agreement between the theoretical and experimental data was obtained (Table II). The worst agreement was reached for the middle ( $6\gamma_6-3\gamma_6$ ) IVMO region.

Earlier<sup>8,9</sup> the IVMO XPS structure of  $\text{UF}_4$  was interpreted on the basis of the binding energy differences between the core and valence electronic levels. It enabled us to identify qualitatively the fine spectral structure of  $\text{UF}_4$  in the binding energy range 15–40 eV and to attribute it to the IVMO electrons. The relativistic calculation results enabled us to interpret quantitatively the XPS fine structure in the whole range 0–40 eV.

Thus, the sharp peak at 3.8 eV is attributed to the U  $5f$  electrons weakly participating in the chemical bond, and the outer valence band—to the outer valence U  $5f$ ,  $6d$ ,  $7s$ ,  $7p$  and F  $2p$  AOs and to a lesser degree—to the U  $6p$  AO. Earlier<sup>16</sup> experimental evidence for the fact that the  $5f$  electrons can participate in the chemical bond in  $\gamma\text{-UO}_3$  without losing their  $f$  nature was established. However, for  $\text{UF}_4$  such

a strict conclusion could not be made. Thus, the experimental intensity ratios OVMO/IVMO with (and without) taking into account the U  $5f$  intensity at 3.8 eV are 0.60 (0.24). It differs slightly from the corresponding theoretical values 0.76 (0.42) (Table II). The theoretical value 0.42 is about twice higher than the corresponding experimental one 0.24 due to the contribution of the U  $5f$  electrons. Having attributed, for example, the intensity of the OVMO band only to the U  $6d^17s^25f^3$  and  $4F 2p^5$  electrons, and the IVMO band—to the U  $6p^6$  and  $4F 2s^2$  electrons in  $\text{UF}_4$ , the corresponding theoretical values are 0.78 (0.38). It agrees with the theoretical values 0.76 (0.42), but more than the corresponding experimental values 0.60 (0.24). However, if to attribute the OVMO band to the U  $6d^27s^25f^2$  and  $4F 2p^5$  electrons, and the IVMO band—to the U  $6p^6$  and  $4F 2s^2$  electrons, the corresponding theoretical values are 0.62 (0.22). It is in a good agreement with the experimental data. It indicates that the U  $5f$  electron involved in the chemical bond is either strongly delocalized or loses its  $f$  nature. To the greatest degree, the direct participation of the U  $5f$  electrons in the chemical bond was observed in  $\gamma\text{-UO}_3$ ,<sup>16</sup> and to a lesser degree—in  $\text{UO}_2$ .<sup>15</sup> Unlike  $\gamma\text{-UO}_3$  and  $\text{UO}_2$ , in the case of  $\text{UF}_4$  this participation can be much lower because of the higher fluorine electronegativity and higher uranium—fluorine interatomic distance. It can result in the delocalization of the U  $5f$  electrons and decrease of the U  $5f$  photoionization cross section.

In the IVMO XPS range the best agreement in the binding energy was reached only for the  $5\gamma_7$ ,  $7\gamma_6(4)$ , and  $2\gamma_6(7)$  orbitals determining the spectral width. The experimental intensities of peaks and the IVMO group  $6\gamma_6-3\gamma_6(6)$  in some cases are much higher than the theoretical values (Table II). Thus, the  $6\gamma_6(5)$  and  $2\gamma_6(7)$  experimental IVMO intensities are 2.1 and 1.5, respectively. These data do not allow a correct experimental evaluation of the participation degree of these electronic shells in the IVMO formation, since the U  $6p_{1/2}$  and F  $2s$  photoionization cross section are comparable (Table I).

Taking into account the calculations and experimental core–valence levels binding energy differences for metallic uranium<sup>8</sup> and  $\text{UF}_4$ , a MO schematic diagram can be built (Fig. 2). This diagram was built in the MO LCAO (molecular orbitals as linear combinations of atomic orbitals) approximation. It enables us to understand the real XPS structure of  $\text{UF}_4$ . In this approximation one can separate formally the antibonding  $5\gamma_7$ ,  $7\gamma_6(4)$ , and  $6\gamma_6(5)$  and bonding  $1\gamma_7$ ,  $3\gamma_6(6)$ , and  $2\gamma_6(7)$  IVMOs, as well as the quasiautomatic (in a certain approximation)  $4\gamma_7$ ,  $3\gamma_7$ ,  $2\gamma_7$ ,  $5\gamma_6$ , and  $4\gamma_6(6)$  IVMOs attributed to the F  $2s$  electrons (Fig. 2, Tables I and II). The XPS data show that the binding energies of quasiautomatic IVMOs must be close to the magnitude. Indeed, the F  $1s$  peak is symmetric, and its FWHM is  $\Gamma=1.3$  eV. The F  $2s$  binding energy must be about 29.9 eV (Fig. 2). This value is the difference between  $E_b=685.3$  eV and  $\Delta E_F=655.4$  eV, where  $E_b$  is the F  $1s$  binding energy in  $\text{UF}_4$ , and  $\Delta E_F$  is the difference between the F  $1s$  and F  $2s$  binding energies for  $\text{MnF}_2$  (Ref. 11) (Fig. 2). Theoretical results agree qualitatively with these data (Table II). Since  $\Delta E_U=360.6$  eV,  $\Delta E_1=362.6$  eV, one can find  $\Delta_1=2.0$  eV (Fig. 2). Since the

TABLE I. MO composition and energies  $E_0^a$  (eV) for the  $\text{UF}_8^{4-}$  ( $C_{4v}$ ) at  $\text{RU}_{-F}=2.29 \text{ \AA}$  ( $\text{RX}_\alpha$  DVM), photoionization cross sections  $\sigma_i^b$  and conversion one-electron partial probabilities  $\alpha_i^c$ .

MO <sup>d</sup>	Q	$-E_0$ , eV	MO composition												
			U										F		
			$6s$	$6p_{1/2}$	$6p_{3/2}$	$6d_{3/2}$	$6d_{5/2}$	$7s$	$5f_{5/2}$	$5f_{7/2}$	$7p_{1/2}$	$7p_{3/2}$	$2s$	$2p_{1/2}$	$2p_{3/2}$
		$\sigma_i$	1.14	0.89	1.29	0.61	0.55	0.12	3.67	3.48	0.07	0.10	1.44	0.13	0.13
		$\alpha_i$	0.07	43.38	23.55	6.55	7.71	0.01	0.07	0.04	8.23	4.39			
						0.10	0.75		0.01				0.03	0.07	0.04
						0.43	0.41			0.01			0.03	0.02	0.10
							0.83		0.01	0.01			0.02	0.06	0.07
								0.03		0.01	0.91		0.02	0.01	0.02
								0.89			0.03		0.03	0.02	0.03
						0.71	0.10			0.04			0.01	0.02	0.12
					0.01							0.96	0.01		0.02
									0.01			0.96	0.01		0.02
						0.41	0.44							0.04	0.11
						0.04			0.01	0.86				0.06	0.03
O							0.02		0.03	0.87				0.02	0.06
										0.93	0.01			0.04	0.02
						0.01				0.93				0.02	0.04
V							0.01		0.86	0.05				0.01	0.07
									0.92					0.01	0.07
M							0.01		0.91	0.01				0.01	0.06
														0.30	0.70
						0.06				0.01				0.21	0.72
O														0.24	0.76
						0.05			0.02					0.35	0.58
														0.11	0.89
														0.54	0.46
														0.23	0.77
														0.44	0.56
						0.01				0.01	0.01	0.01		0.58	0.38
					0.01							0.03		0.25	0.71
									0.02	0.02		0.02		0.27	0.67
					0.01				0.01	0.01	0.02			0.38	0.57
										0.06				0.36	0.58
									0.06	0.01				0.24	0.69
									0.02	0.04				0.45	0.49
					0.01				0.03	0.02				0.11	0.83
									0.07	0.02			0.01	0.10	0.80
									0.01	0.07				0.55	0.37
					0.01			0.04					0.01	0.33	0.61
						0.02	0.12						0.01	0.32	0.53
							0.11						0.01	0.52	0.36
							0.12						0.01	0.46	0.41
						0.11	0.01						0.01	0.05	0.82
						0.13								0.15	0.72
												0.01	0.11	0.02	0.03

TABLE I. (Continued.)

MO <sup>d</sup>	Q	-E <sub>0</sub> , eV	MO composition													
			U										F			
			6s	6p <sub>1/2</sub>	6p <sub>3/2</sub>	6d <sub>3/2</sub>	6d <sub>5/2</sub>	7s	5f <sub>5/2</sub>	5f <sub>7/2</sub>	7p <sub>1/2</sub>	7p <sub>3/2</sub>	2s	2p <sub>1/2</sub>	2p <sub>3/2</sub>	
		σ <sub>i</sub>	1.14	0.89	1.29	0.61	0.55	0.12	3.67	3.48	0.07	0.10	1.44	0.13	0.13	
		α <sub>i</sub>	0.07	43.38	23.55	6.55	7.71	0.01	0.07	0.04	8.23	4.39				
	7γ <sub>6</sub>	2	17.23			0.83							0.01	0.11	0.01	0.04
	6γ <sub>6</sub>	2	23.09		0.20								0.02	0.78		
I	4γ <sub>7</sub>	2	23.81				0.02		0.01					0.97		
	3γ <sub>7</sub>	2	23.81			0.01	0.01			0.01				0.97		
V	2γ <sub>7</sub>	2	24.00			0.01	0.03							0.96		
	5γ <sub>6</sub>	2	24.01			0.02	0.02							0.96		
M	4γ <sub>6</sub>	2	24.14	0.01				0.04						0.95		
	1γ <sub>7</sub>	2	24.53			0.10							0.01	0.88	0.01	
O	3γ <sub>6</sub>	2	24.58			0.10							0.01	0.88		0.01
	2γ <sub>6</sub>	2	27.14		0.79									0.19	0.01	0.01
	1γ <sub>6</sub>	2	43.29	0.98										0.02		

<sup>a</sup>Calculated energies are shifted down toward the negative values (down) by 1.8 eV of absolute scale.

<sup>b</sup>Photoionization cross section σ<sub>i</sub> (Barn per electron) from Ref. 26.

<sup>c</sup>α<sub>i</sub>, relative partial conversion probability per one electron (%) of E3 multipole of <sup>235</sup>U nucleus with participation of electrons from the *nlj* shells, determined on the basis of the data from Ref. 27.

<sup>d</sup>To simplify the data, all the sequence numbers are shifted like 27γ<sub>6</sub> → 1γ<sub>6</sub> and 22γ<sub>7</sub> → 1γ<sub>7</sub>, respectively.

<sup>e</sup>Upper filled orbital 18γ<sub>7</sub> (two electrons), filling number for *n*γ<sub>6</sub><sup>±</sup> and *n*γ<sub>7</sub><sup>±</sup> MO is 2.

U 6*p* atomic spin-orbit splitting according to the calculation<sup>28</sup> and experimental<sup>29</sup> data is Δ*E*<sub>*sl*</sub>(U 6*p*) = 10.0 eV, the U 6*p*<sub>3/2</sub> and F 2*s* binding energies are comparable by the magnitude. The 6γ<sub>6</sub>(5) and 2γ<sub>6</sub>(7) IVMO binding energy difference is 4.2 eV. In this case the values characterizing the antibonding Δ<sub>2</sub> and bonding Δ<sub>3</sub> are approximately equal and they are ≈2.1 eV (Fig. 2). It agrees with the data on the FWHMs and intensities. Indeed, the intensities of these lines are comparable and FWHM of the antibonding 6γ<sub>6</sub>(5) IVMO is equal to that of the bonding 2γ<sub>6</sub>(7) IVMO. It indicates the compensation of the antibonding (Fig. 2, Table II). Since in the U 6*p*<sub>1/2</sub>-F 2*s* binding energy range the structure is complicated due to the IVMO overlapping, it is very difficult to identify correctly all the IVMOs and to make a conclusion about the contributions of the IVMO electrons to the covalent component of the chemical bond in UF<sub>4</sub>.

The conversion transition of the E3 multipole for the <sup>235m</sup>U nucleus (*T*<sub>1/2</sub> ≈ 26 min) from the first excited state (spin *I*<sub>1</sub> = 1/2<sup>+</sup>, *E*<sub>1</sub> = 76.5 ± 0.4 eV) to the ground nucleus state (spin *I*<sub>0</sub> = 7/2<sup>-</sup>, *E*<sub>0</sub> = 0 eV) is accompanied by the low energy electron photoemission. The conversion process is energetically permitted for the U 6*s*<sup>2</sup>6*p*<sup>6</sup>5*f*<sup>3</sup>6*d*<sup>1</sup>7*s*<sup>2</sup>7*p*<sup>0</sup> electrons. These shells can participate effectively in the OVMO and IVMO formation in uranium compounds. In this case the partial conversion probability per one electron α<sub>*i*</sub>(*EL*, *I*<sub>1</sub> → *I*<sub>0</sub>, *n*<sub>*i*</sub>*l*<sub>*i*</sub>*j*<sub>*i*</sub> → ε<sub>*i*</sub>) of ejection of (*n*<sub>*i*</sub>*l*<sub>*i*</sub>*j*<sub>*i*</sub>) electron is proportional to the electronic factor *w*<sub>*e*</sub>(E3, *n*<sub>*i*</sub>*l*<sub>*i*</sub>*j*<sub>*i*</sub>, *hω*), where ε<sub>*i*</sub>, electron kinetic energy; *n*<sub>*i*</sub>, principal quantum number; *l*<sub>*i*</sub> and *j*<sub>*i*</sub>, orbital and total angular moments; and *hω*, nucleus excitation energy.<sup>27</sup>

The present work built the theoretical CES spectrum on the basis of the relativistic calculations and relative one-electron partial conversion probabilities<sup>27</sup> [Tables I and II, Fig. 1(b)]. Since the U 6*p* conversion cross section is much higher than that for the other electronic shells (Table I), the CES spectrum from UF<sub>4</sub> reflects rather the U 6*p* partial electronic density. One can see a qualitative agreement between experimental XPS and CES. Also a good agreement between the theoretical and experimental CES spectra was found. Comparison of the XPS and CES data gives three important conclusions. First, the U 6*p* shell participates effectively in the IVMO formation. Second, the U 6*p* shell participates significantly in the OVMO formation. Third, the U 5*f* electrons from the 18γ<sub>7</sub>(1) OVMO and electrons from the quasiatomic 4γ<sub>7</sub>, 3γ<sub>7</sub>, 2γ<sub>7</sub>, 5γ<sub>6</sub>, and 4γ<sub>6</sub>(6) IVMOs, as expected, practically are not observed at 3.8 and 29.9 eV. The 6γ<sub>6</sub>(5), 1γ<sub>7</sub>, 3γ<sub>6</sub>(6), and 2γ<sub>6</sub>(7) IVMO energies differ slightly from the corresponding theoretical values and agree with the experimental XPS parameters (Figs. 1 and 2, Table II).

For comparative quantitative analysis of the experimental and theoretical intensities the considered spectra were decomposed with the data of Fig. 2 in mind. The diagram of Fig. 2 was built on the basis of the experimental binding energy differences and theoretical intensities. The identification of the XPS and CES structures is reflected in Table II and Fig. 3. These data show that the experimental XPS and CES binding energies practically coincide to within the measurement errors. However, the experimental IVMO intensities often differ from the corresponding theoretical ones. The best agreement was reached for the 5γ<sub>7</sub>, 7γ<sub>6</sub>(4) IVMOs. Taking into account conversion cross sections and experimental

TABLE II. XPS and CES parameters for the  $\text{UF}_8^{4-}$  ( $C_{4v}$ ) cluster at  $R_{\text{U-F}}=2.29 \text{ \AA}$  ( $RX_\alpha$  DVM), and the experimental density  $\rho_i(e^-)$  of U  $6p$  electronic states in  $\text{UF}_4$ .

MO	$-E^a$ , eV	XPS			CES			$\rho_i$ , $e^-$ units		
		Energy <sup>b</sup> , eV	Intensity, %		Energy <sup>b</sup> , eV	Intensity, %		U $6p_{3/2}$	U $6p_{1/2}$	
			Experiment	Theory		Experiment	Experiment			Theory
	18 $\gamma_7^c$	2.87	3.8(1.5) <sup>d</sup>	20.9	22.4		0.1			
	19 $\gamma_6$	6.07		0.4		5.3(1.4)		0.8		
	17 $\gamma_7$	7.67		1.0			1.4			
	18 $\gamma_6$	7.72		0.4						
	17 $\gamma_6$	7.82		1.2			1.1			
	16 $\gamma_7$	7.82		0.4						
	16 $\gamma_6$	7.85		0.4						
	15 $\gamma_7$	7.90		0.4						
	14 $\gamma_7$	7.92		0.4						
O	15 $\gamma_6$	8.14	8.0(2.5)	0.7	9.3		0.4			
	13 $\gamma_7$	8.19		0.5		8.2(3.7)	0.4	10.1		
V	14 $\gamma_6$	8.53		1.3			0.1			
	13 $\gamma_6$	8.67		0.9			0.6			
M	12 $\gamma_6$	8.71		1.6						
	12 $\gamma_7$	8.72		1.9						
O	11 $\gamma_7$	8.75		1.7						
	11 $\gamma_6$	8.78		1.5			0.5			
	10 $\gamma_7$	8.98		2.3						
	9 $\gamma_7$	9.00		2.0						
	10 $\gamma_6$	9.43		0.4						
	9 $\gamma_6$	9.57		0.8			1.0			
	8 $\gamma_7$	9.63		0.7			0.8			
	7 $\gamma_7$	9.63		0.8			0.9			
	8 $\gamma_6$	9.67		0.8			0.8			
	6 $\gamma_7$	9.69	10.2(3.2)	0.8	5.6	11.1(2.4)	0.9	0.4		
	$\Sigma I_i^e$			44.2	37.3		9.0	11.3	0.5	0.1
	5 $\gamma_7$	20.10	20.1(2.5)	7.1	9.8	20.2(2.7)	19.0	16.9		
	7 $\gamma_6$	20.10	20.1(2.5)	7.1	9.9	20.2(2.7)	19.0	16.9	2.9	
I	6 $\gamma_6$	25.96	27.6(3.0)	4.6	9.5	28.0(3.2)	9.6	25.7		1.0
	4 $\gamma_7$	26.68		4.6			0.2			
V	3 $\gamma_7$	26.68		4.6			0.1			
	2 $\gamma_7$	26.87		4.4			0.3			
M	5 $\gamma_6$	26.88	30.4(3.0)	4.4	25.4		0.3			
	4 $\gamma_6$	27.01		4.3						
O	1 $\gamma_7$	27.40		4.7		29.8(3.5)	2.3	3.3		
	3 $\gamma_6$	27.45		4.7		29.8(3.5)	2.3	3.3	0.6	
	2 $\gamma_6$	30.01	31.8(3.0)	5.3	8.1	31.4(3.2)	37.9	23.2		0.9
	$\Sigma I_i^e$			55.8	62.7		91.0	88.7	3.5	1.9
	1 $\gamma_6$	46.16	48.3(6.0)	$\approx 7.1$			$\approx 0.1$			

<sup>a</sup>Calculated energies are shifted down toward the negative energies by 2.87 eV so that the calculated energy of the 5 $\gamma_7$  IVMO would be 20.1 eV.

<sup>b</sup>FWHMs in eV are given in the parentheses.

<sup>c</sup>Upper filled MO 18 $\gamma_7$  (two electrons), filling number for the  $n\gamma_6$  and  $n\gamma_7$  MOs is 2.

<sup>d</sup>FWHM is given relative to the  $\Gamma(\text{C } 1s)=1.3 \text{ eV}$ .

<sup>e</sup>Total intensities and the U  $6p$  electronic state densities.

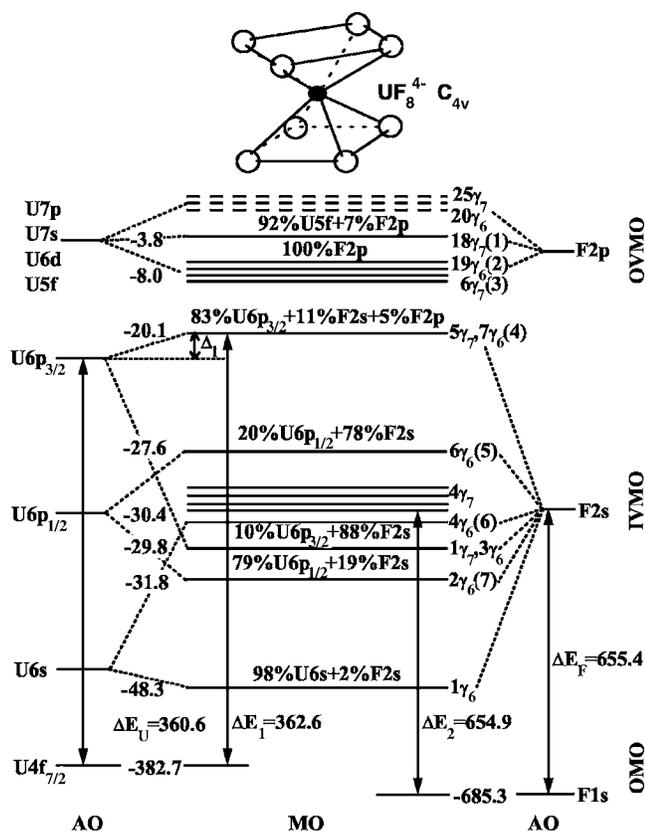


FIG. 2. MO schematic diagram for the  $UF_8^{4-}$  ( $C_{4v}$ ) cluster built taking into account the theoretical and experimental data. The chemical shift during the cluster formation is not shown. The arrows indicate some experimentally measurable binding energy differences. The experimental binding energies (eV) are given to the left side. Energetic scale is not kept.

intensities, the partial U  $6p_{3/2,1/2}$  electronic density in solid  $UF_4$  was evaluated (Table II).

Evaluation of the number of the U  $6p_{3/2,1/2}$  electrons participating in the chemical bond in  $UF_4$  was done in the following approximation. The  $5\gamma_7$ ,  $7\gamma_6(4)$  and  $1\gamma_7$ ,  $3\gamma_6(6)$  IVMO CES intensity was suggested to be formed only by the U  $6p_{3/2}$  electrons, while the  $6\gamma_6(5)$  and  $2\gamma_6(7)$  one—only by the U  $6p_{1/2}$  electrons. This suggestion is well grounded (see Table I). While considering the theoretical intensities ratio first it was suggested that the  $5\gamma_7$ ,  $7\gamma_6(4)$  IVMOs are populated by the 3.32 U  $6p_{3/2}$  electrons and the  $6\gamma_6(5)$  and  $2\gamma_6(7)$  IVMOs—by the 2.0 U  $6p_{1/2}$  electrons (Table I). After that, the changes in the intensities were renormalized taking into account the experimental data. Also, on the basis of the analysis of the experimental and theoretical data it was suggested that the relative CES intensity at 8.2 eV, being 11.3 (Table II) had a 6.35 contribution from the U  $6p$  electrons [see Tables I and II, Fig. 1(b)]. As a result it was found that the OVMOs include 0.6 U  $6p$  electrons (Table II). It is more than the corresponding theoretical value 0.3 electrons (Table I). In the beginning it was also suggested that mainly the U  $6p_{3/2}$  electrons participate in the OVMO formation. Table I shows that the  $6p_{1/2}$  contribution to the OVMO is about 6 times lower. Therefore, 0.1 U  $6p_{1/2}$  electrons were suggested to be spread around among the OVMOs. The remaining

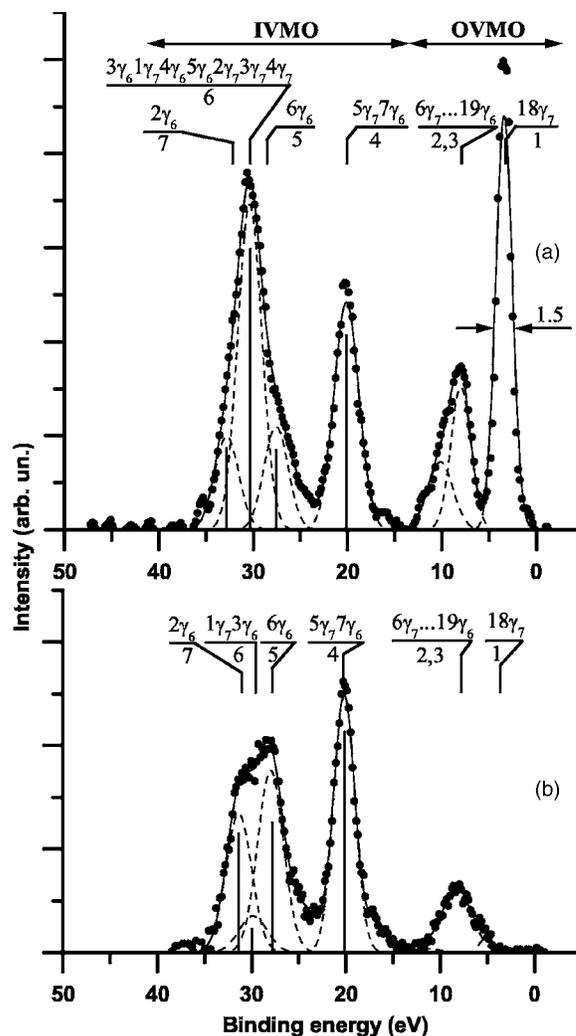


FIG. 3. XPS (a) and CES (b) from  $UF_4$  with subtracted background. The corresponding expected spectra obtained on the basis of the theoretical and experimental data are given under the experimental spectra as vertical bars. The spectral intensities are given in arbitrary units; the theoretical intensities are normalized in %.

1.9 U  $6p_{1/2}$  electrons were suggested to be spread among the  $6\gamma_6(5)$  and  $2\gamma_6(7)$  taking into account their intensities.

In the IVMO energy range a significant difference between the theoretical and experimental data was observed. For example, the antibonding  $5\gamma_7$ ,  $7\gamma_6(4)$  IVMO contains 2.9 U  $6p_{3/2}$  electrons, which is less than the corresponding calculated value 3.32 (Table I). The difference was also observed for the bonding  $1\gamma_7$ ,  $3\gamma_6(6)$  IVMO: 0.6 electrons, experiment (Table II); and 0.4, theory (Table I). For the antibonding  $6\gamma_6(5)$  and bonding  $2\gamma_6(7)$  the corresponding experimental and theoretical values are 1.0 and 0.4 electrons for the  $6\gamma_6(5)$ ; 0.9 and 1.58 for the  $2\gamma_6(7)$ . These data show that the IVMO formation (AO mixing) in reality takes place in a much higher scale (13%–60%) than it was predicted by the theory.

The quantum mechanics can predict qualitatively the alterations of the considered spectral intensities depending on the binding energy. Therefore, one can vary the initial data for the initial clusters for the better agreement between the

theoretical and experimental data. As it was shown in Ref. 9, the wider a MO peak is, the more its electrons participate (bonding, antibonding) in the chemical bond. The lines  $5\gamma_7$ ,  $7\gamma_6(4)$  corresponding to the antibonding IVMOs in the CES are observed narrower than the corresponding bonding  $1\gamma_6$ ,  $3\gamma_6(6)$  IVMO (Table II). It must be noted that the  $5\gamma_7$ ,  $7\gamma_6(4)$  XPS peak was calibrated, which cannot be done for the CES peak. The observed relative decrease of the  $5\gamma_7$ ,  $7\gamma_6(4)$  IVMO FWHM comparing to the corresponding bonding  $1\gamma_6$ ,  $3\gamma_6(6)$  IVMO FWHM can be explained by the contribution from, for example, the F  $2p$  and U  $7p$  AOs in the  $5\gamma_7$ ,  $7\gamma_6(4)$  IVMO. As it was shown earlier for the XPS analysis, it leads to the loss of the antibonding nature of this IVMO. Comparison of the experimental (16.9%) and theoretical (19.0%)  $5\gamma_7$ ,  $7\gamma_6(4)$  IVMO intensity shows that the contribution from the U  $6p_{3/2}$  AO to this IVMO was theoretically overestimated. The experimental (25.7%) contribution from the U  $6p_{1/2}$  AO to the  $6\gamma_6(5)$  IVMO is significantly higher than the theoretical one (9.6%), and the contribution (23.2%) of this AO to the  $2\gamma_6(7)$  IVMO is significantly lower than the theoretical one (37.9%). It indicates that in reality the degree of participation of the U  $6p_{1/2}$  AO in the IVMO formation is about 2 times higher as it follows from the theoretical results. As it follows from the XPS and CES data,  $6\gamma_6(5)$  and  $2\gamma_6(7)$  IVMO peaks are about equal by the intensity. It indicates about equal contribution from the  $6p_{1/2}$  and O  $2s$  AOs to this IVMO.

The obtained data for the first time allowed an experimental determination of the IVMO composition in  $UF_4$ . Thus, the bonding  $2\gamma_6(7)$  and the corresponding antibonding  $6\gamma_6(5)$  IVMOs were found to form mostly from the U  $6p_{1/2}$  and F  $2s$  AOs. The calculated compositions of these shells (79% of the U  $6p_{1/2}$  and 19% of the F  $2s$  AOs and 20% of the U  $6p_{1/2}$  and 78% of the F  $2s$  AOs) differ significantly from the experimental compositions (45% of the U  $6p_{1/2}$  and 55% of the F  $2s$  AOs and 50% of the U  $6p_{1/2}$  and 50% of the F  $2s$  AOs). The bonding  $1\gamma_6$ ,  $3\gamma_6(6)$  and the corresponding antibonding  $5\gamma_7$ ,  $7\gamma_6(4)$  IVMOs were found to form mostly from the U  $6p_{3/2}$  and F  $2s$  AOs, and their calculated compositions (10% of the U  $6p_{3/2}$  and 88% of the F  $2s$  AOs and 83% of the U  $6p_{3/2}$  and 11% of the F  $2s$  AOs) differ less from the experimental compositions (15% of the U  $6p_{3/2}$  and 85% of the F  $2s$  AOs and 73% of the U  $6p_{3/2}$  and 27% of the F  $2s$  AOs) (Tables I and II).

The most ambiguous for interpretation is the  $4\gamma_7-3\gamma_6(6)$  IVMO XPS region. Comparison of the XPS and CES data shows that the structure in this range in general is formed from the F  $2s$  electrons. It agrees with the calculation results. However, the U  $6p$ -F  $2s$  AO overlap during formation of  $5\gamma_7$ ,  $7\gamma_6(4)$  and  $1\gamma_6$ ,  $3\gamma_6(6)$  IVMOs was underestimated, but the IVMO sequence order was determined correctly. The knowledge on the correct IVMO sequence order is critical for understanding of the IVMO contributions to the covalent component of the chemical bond in  $UF_4$ .

In conclusion we would like to note that despite the approximation imperfections, the calculation results for the

$UF_8^{4-}$  ( $C_{4v}$ ) cluster reflecting uranium close environment in solid  $UF_4$  are in a satisfactory agreement with the experimental data. It allowed for a reliable identification of the peaks of the U  $5f$  electrons weakly participating in the chemical bond, as well as the IVMO electrons. These results can be also used for the interpretation of other x-ray spectra (Auger, emission, absorption, etc.<sup>8</sup>) of uranium compounds.

#### IV. CONCLUSIONS

Low binding energy (0–40 eV) x-ray photoelectron and conversion electron spectra of  $UF_4$  were measured and interpreted in the relativistic  $X_\alpha$  discrete variation approximation for the  $UF_8^{4-}$  ( $C_{4v}$ ) cluster reflecting uranium close environment in solid  $UF_4$ . It yielded a satisfactory qualitative and in some cases quantitative agreement between the theoretical and experimental data.

It is theoretically shown that in  $UF_4 \approx 1$  U  $5f$  electron can directly participate in formation of chemical bond. The experimental evaluation, however, has shown that the U  $5f$  electron in  $UF_4$  is more delocalized in the OVMO binding energy range from –5 to –11 eV and less maintains the  $f$  nature than in  $\gamma$ - $UO_3$  and  $UO_2$ . About 2 U  $5f$  electrons weakly participating in the chemical bond are localized at –3.8 eV, and the vacant U  $5f$  electronic states are generally delocalized in the low positive energy range (0–7 eV).

The U  $6p$  electrons (0.6 U  $6p$  electrons) were experimentally shown to participate significantly in the OVMO formation beside the IVMO formation. It agrees with the theoretical data.

The peaks in the binding energy range 0–40 eV in the studied spectra were identified, the IVMO sequence order and quantitative experimental compositions were determined. The bonding  $2\gamma_6(7)$  and corresponding antibonding  $6\gamma_6(5)$  IVMO were determined to form mostly from the U  $6p_{1/2}$  and F  $2s$  atomic shells, and their calculated compositions (79% of the U  $6p_{1/2}$  and 19% of the F  $2s$  AOs and 20% of the U  $6p_{1/2}$  and 78% of the F  $2s$  AOs) were found to differ significantly from the experimental compositions (45% of the U  $6p_{1/2}$  and 55% of the F  $2s$  AOs and 50% of the U  $6p_{1/2}$  and 50% of the F  $2s$  AOs). The bonding  $1\gamma_6$ ,  $3\gamma_6(6)$  and the corresponding antibonding  $5\gamma_7$ ,  $7\gamma_6(4)$  IVMOs were found to form mostly from the U  $6p_{3/2}$  and F  $2s$  AOs, and their calculated compositions (10% of the U  $6p_{3/2}$  and 88% of the F  $2s$  AOs and 83% of the U  $6p_{3/2}$  and 11% of the F  $2s$  AOs) differ less from the experimental compositions (15% of the U  $6p_{3/2}$  and 85% of the F  $2s$  AOs and 73% of the U  $6p_{3/2}$  and 27% of the F  $2s$  AOs).

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