Suppression of photocatalytic efficiency in highly N-doped anatase films

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We report the role of N in epitaxial films of anatase TiO₂. The films were artificially grown with a two-step temperature-tuned epitaxy which utilized the high-temperature cubic phase of LaAlO3 substrates. The preparation of highly crystallized anatase with various N concentrations ($C_N \leq 3.85$ at. %) allowed us to identify the optimum dopant concentration $(C_N = 1 - 2$ at. %). At higher doping levels, N is found to be difficult to substitute for O having been predicted to contribute to the band-gap narrowing, giving rise to the undesirable deep-level defects. In addition, a study by x-ray and Raman spectroscopy revealed that the growth of anatase became more difficult, and the stable phase was shifted to rutile at the higher N concentrations.

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I. INTRODUCTION

Photocatalysis is a generic term for surface chemical reactions resulting from photoinduced carriers. Among the well-known photocatalytic materials (e.g., ZnO, CdSe, $SrTiO₃$, and $TiO₂$), $TiO₂$ is the most promising photocatalyst with its chemical stability, high reactivity, and low environmental pollution. The crystal phase of $TiO₂$ is classified into three types, i.e., rutile, brookite, and anatase.^{1,2} The anatase presents the strongest photocatalytic activity among them and is widely used for photocatalyst; nevertheless, it is a low-temperature metastable phase and it is difficult to obtain a single crystal. The most significant limitation to the application of anatase is its optical band gap. Anatase has an indirect band gap of \sim 3.2 eV, and exhibits a photocatalytic activity only under ultraviolet light illumination (λ) $<$ 387 nm). This has prompted many groups to investigate the enhancement of its photocatalytic activity under visible light illumination. In the past decades, the doping of metal ions has been recognized as a promising solution.^{3–5} The improved photocatalytic performance is explained by the shallow charge-trapping sites produced by the cationic dopant which partially prevents the undesirable recombination of electron-hole pairs.

Sato reported sensitized photocatalytic activity under visible light illumination with a nitrogen dopant in 1986.⁶ In 2001, Asahi *et al.* explained the effect of N doping using spin-restricted local-density-functional approximation (LDA) calculations for anatase.⁷ They pointed out that the substitute N is effective for band-gap narrowing by mixing with the O 2*p* state in the anatase matrix, and experimentally verified the enhanced optical absorption at less than 500 nm. Subsequently, the superiority of N-doped $TiO₂$ is demonstrated from the viewpoint of photocatalytic activity, $8-10$ surface hydrophilicity, $1\overline{1}$ and photoinduced carrier response¹² under visible light illumination. However, samples with varying N-dopant concentrations have not been studied due to the difficulty of the chemically stable nitrogen doping into the metastable anatase. Therefore, there is a lack of correct understanding of the effects of N at high dopant levels. While Burda *et al.* reported the enhancement of N incorporation up to 8 at. $\%$ into anatase nanoparticles, degraded photocatalytic efficiencies at a high N concentration have been predicted by Lindgren *et al.*¹²

For the metal doping method, it is known that the photocatalytic activity strongly depends on the doping concentration, and there exists an optimum value for each dopant. Many hypotheses have been suggested to explain the mechanism for the drawbacks at the high doping levels: (1) The higher concentration traps more electrons and/or holes, lowering the carrier mobility;^{13,14} and (2) metal ions produce not only shallow traps but also deep-level defects which act as recombination centers.15 In this paper, we will report the growth of highly textured anatase films and demonstrate the optimum N concentration for the photocatalytic efficiency. The N-doped anatase crystal was artificially grown by employing two-step temperature-tuned epitaxy. The N concentration was controlled using homemade $TiO_{2-2x}N_x$ targets.

II. EXPERIMENTS

 N -doped TiO₂ films were grown by KrF excimer pulsedlaser deposition (PLD). The ablation targets were prepared by mixing rutile $TiO₂$ and TiN powders in various ratios. By changing the ratio, we obtained various N concentrations $(C_{\rm N})$ from 0 to 30 at. % in the targets. The films were deposited on (100) LaAlO₃ substrates. The substrates were heated by a standard resistive heater which was varied from 300 to 550 °C. The base pressure of the process chamber was evacuated to as low as 10^{-4} Pa, and then it was backfilled with oxygen to 10^{-1} Pa. A deposition run was performed for 30 min and a deposition rate of 20 nm min⁻¹ was achieved. To obtain information about the N states, x-ray photoelectron spectroscopy (XPS, 9000-MC, JEOL) was carried out. The structural characteristics of the thin films were determined by x-ray diffraction (XRD, RAD-C, Rigaku) and a Raman spectrometer (CCD-Raman-One, Chromex). The photocatalytic activity was evaluated by measuring the decomposition of a methylene blue solution (15 ppm).

III. RESULTS AND DISCUSSION

A. Two-step temperature-tuned epitaxy of N-doped anatase films

It is well known that the crystal structure of $LaAlO₃$ undergoes crystallographic transitions by changing the tem-

FIG. 1. (Color online) $2\theta \cdot \theta$ XRD scans of TiO₂ films grown using an undoped TiO₂ target at (a) 500 °C and (b) 300 °C with a high-temperature buffer layer. The inset in (b) is the scan around the (004) anatase peak of the HTBL grown at 550 \degree C for 5 min.

perature. A gradual change from a rhombohedral to hightemperature cubic structure was observed by Geller and Bala.¹⁶ They pointed out that the critical temperature (T_c) was 435 °C. However, this T_c was not correctly measured. Nowadays, there are many reports concerning the T_c of $LaAlO₃$; however, this is beyond the current scope of this paper. The generally accepted transition temperature ranges between 530 and 550 $^{\circ}$ C.¹⁷ Since the lattice mismatch between the cubic $LaAlO₃$ and anatase is very small $(\sim 0.2\%)$, epitaxy in the high-temperature phase is absolutely necessary. It is expected that the crystallographic structure of the deposited films dramatically changes around T_c . In our study, while a clear anatase peak was easily obtained at 550 °C, the films grown below 400 °C did not indicate any XRD peaks.

Figure 1(a) shows the XRD spectrum of the film grown at 500 °C using an undoped $TiO₂$ target. The shallow anatase diffraction peak suggests the incomplete transition of LaAlO₃. At the same time, we must avoid the oxidation of the Ti-N bond by the *in situ* thermal annealing. Hsieh *et al.* observed the dependence of the film oxidation on the annealing temperature using TiN films grown on $SrTiO₃$ substrates.¹⁸ They found that the TiN was easily oxidized and fully disappeared at 600 $^{\circ}$ C. Figure 2(c) shows the N 1*s* core levels investigated with XPS for the film grown at 500 °C using a C_N =30 at. % target. As can be seen in Fig. 2(c), there is only a shallow peak from N (400 eV), suggesting that most of the TiN suffers from oxidation during the growth.

It is now clear that the anatase films are obtainable above the T_c at the expense of the N doping. To circumvent this dilemma, we developed the two-step temperature-tuned epitaxy, which used a high-temperature buffer layer (HTBL) as a nucleation layer of the anatase. The HTBL was prepared at 550 °C for 5 min. The substrate temperature was then decreased to the desired temperature. Figure $1(b)$ shows the

FIG. 2. (Color online) Differences in N 1s core-level XPS spectra of N-doped $TiO₂$ films grown at various conditions. (a) and (b) are grown at 300 °C with a HTBL using C_N = 30 at. % and 5 at. % target, respectively. (c) is grown at 500 °C using C_N = 30 at. % target without a HTBL.

XRD spectrum of the film grown at 300 °C with HTBL using an undoped $TiO₂$ target. A clear anatase peak is observable, which was not possible without the buffer layer at such a low temperature. The inset in Fig. $1(b)$ is a typical XRD spectrum of a buffer layer. As the nucleation layer is very thin, the diffraction peak is very shallow. Based on these figures, it is obvious that the diffraction peak seen in Fig. 1(b) is not from the buffer layer. It should be noted that even when we used a highly N-doped target (up to 30 at. $%$), we exclusively obtained a single oriented (004) anatase peak and no persisting TiN phases. The XPS spectrum of the N-doped TiO₂ film grown at 300 \degree C with the HTBL using a C_N =30 at. % target is illustrated in Fig. 2(a). It is clearly seen that the peak from N is dramatically increased compared to Fig. 2(c). The calculated C_N in the films and full width at half maximum (FWHM) of (004) anatase peaks are summarized in Table I.

Figures $2(a)$ and $2(b)$ show the N states in the lattice with different dopant concentrations prepared using $C_N = 30$ and 5 at. % targets. These spectra represent two differential peaks of N centered at 396 and 400 eV. The lower-energy peak corresponds to Ti-N bonding, as observed by Saha and Tompkins.19 They also assigned the peaks of 400 and 402 eV to molecularly chemisorbed N_2 . Indeed, the clear peaks at 402 and 400 eV were observed from TiON synthesized through TiN oxidation or TiO₂ nitridation.^{7,19} However, the 402 eV peak was hardly seen from the TiON films grown by sputtering or PLD, not only in this study but in previous reports. $9,11$ Here, we believe that the 402 eV peak results from moleculelike N_2 , which is confined under the oxide layer or literally adsorbed on the surface (or just below) during the TiN oxidation or $TiO₂$ nitridation process. Also, it is reasonable to consider that $N₂$ evaporates during thin-film growth processes, especially when they are performed in a vacuum chamber at high temperatures. Taking into account that the TiAlN film XPS peaks appeared at 403.2 eV and no other peak was observed above 398 eV,20 the N 1*s* peak at 400 eV is likely newly created by oxygen, and it should be interpreted as the Ti-O-N oxynitride. This discussion is well consistent with recent silicon oxynitride studies.^{21,22} But,

	$C_{\rm N}$ in film	FWHM of XRD for anatase	Redshift	Remaining C_{MB} after the irradiation (ppm)	
$C_{\rm N}$ in target $(at, \%)$	$(at, \%)$	(deg)	denoted by $\Delta E_{\rm g}$ (eV)	u v	VIS
$\overline{0}$		0.167		3.01	7.42
2	\leq 1	0.171	0.072	1.77	7.28
5	1.28	0.178	0.094	1.21	6.41
10	1.69	0.180	0.090	0.907	6.55
20	2.55	0.179	0.099	2.06	6.59
30	3.85	0.374	0.107	3.03	6.43

TABLE I. Dependence of N-dopant concentration (C_N) on the film characteristics for anatase samples grown on HTBL/(100)LaAlO₃ substrates at 300 °C.

even if we cannot detect a N_2 signal, it does not mean there is no moleculelike N_2 , since a small quantity of it may be embedded below the XPS detection limit.²¹ Indeed, interstitially doped N (N_i) can make a N_2 bond with substitutional $N(N_0)$ at high doping levels in a N-doped metal oxide.²³

According to Ref. 7, the replacement of the O atom with N is effective for band-gap narrowing, while the interstitial ones in the anatase matrix produce deep-level defects in the band gap. Therefore, the increment of N*ⁱ* is unfavorable. Provided that N atoms substitute at O sites, both 397 eV(Ti-N) and 400 eV(Ti-O-N) peaks should simultaneously increase. However, the peak at 400 eV is highlighted along with the C_N , while the 397 eV peak remains unchanged, as can be seen in Fig. 2. This implies that the doping into the desired site (N_O) is saturated at low dopant levels, and the excess of N is doped interstitially within the matrix. In addition, it is worth pointing out that the anatase growth at a low temperature has a significant importance for the substitutionally doped impurity due to the structural instabilities at high temperatures.²⁴ In the case of N, the formation of a Ti-O bond is energetically favorable compared with that of Ti-N. The film grown at 500 $^{\circ}$ C [Fig. 2(c)] does not indicate the spectrum of the Ti-N bond, suggesting the formation of N*ⁱ* defects resulting from O atoms which fill the appropriate sites by pushing aside the N_O during the annealing process at such a high temperature. This result is reasonable considering the differential formation enthalpy of ΔH_f $=-938.7$ kJ mol⁻¹ for anatase and -337.9 kJ mol⁻¹ for TiN.

B. Properties of the N-doped anatase films

Figure 3(a) shows the photographs of as-grown N-doped anatase epilayers with various N concentrations. The color of the film changed from transparent to yellowish along with the change in the N concentration, indicating that the absorption of visible light is enhanced, which corresponds well with other reports. $6-12$ However, the effect of the N dopant should be determined with caution. Nevertheless, heavily N-doped films certainly show highly enhanced optical absorption; the redshifts in the absorption edge (ΔE_{g}) are less affected by the increased C_N . The absorption edge of the N-doped anatase films is presented in Fig. 3(b). These spectra are significantly affected by the Burstein shift and bandtailing effect especially when N is heavily doped. The optical transmission of the N-doped film grown from a C_N $= 30$ at. % target [lower right in Fig. 3(a)] gradually increases up to 500– 550 nm; however, it is still less than half of the undoped $TiO₂$ film in the whole visible spectrum. The remarkable feature seen in Fig. $3(b)$ is that the visible light is transmitted when it is longer than 380 nm, indicating that the band-gap narrowing has hardly occurred. The inset in Fig. 3(b) shows the estimated optical band gap from these curves. The observed redshifts of the absorption edges from the undoped film are also presented in Table I. The small quantity of N seems to narrow the band gap, even though it is only by 0.07 eV. However, with an increase of C_N , it is clearly seen that the band gap hardly became any narrower. The highly doped film $(C_N = 3.85$ at. %) has a downward shift of only

FIG. 3. (Color online) (a) Photographs of N-doped anatase epilayers. The C_N of ablation targets used are indicated below the respective specimens. (b) C_N dependence of absorption edges of N-doped anatase epilayers. The curves from left to right correspond to those of the samples grown from $C_N=0$, 2, 5, 10, 20, and 30 at. % targets. The redshift of the band gap (E_{g}) is plotted in the inset.

 \sim 0.04 eV compared with the film containing the lowest C_N . Therefore, the absorption and photocatalytic activity produced by the visible light $(\lambda = 400 - 550 \text{ nm})$ likely originate from the N-related visible light absorptive defects, rather than a band-to-band transition within a narrowed band gap. Such discussion is consistent with recent papers.^{10–12,25}

Taking into account the XPS results, the saturation of band-gap shift and the film coloring likely originate from N_O and N_i , respectively. The absorption of the visible light (λ) >400 nm) is merely caused by the deep-level defects, so that the optimum C_N should be more pronounced. Recently, Nakano *et al.* reported the defect created by nitrogen in N-doped anatase film deposited on n^+ -GaN/Al₂O₃ substrates.26 They point out that a deep-level defect located at 2.48 eV below the conduction band is introduced by N. Although it is hard to conclude that this deep-level defect contributes to the band-gap narrowing according to the above discussion, it is very important to determine the effect of this defect. To investigate the photocatalytic activity of N-doped films, we have used both uv and visible light illumination. The uv light source applied is a commercial black light (EFD15BLB, Toshiba) which emits long-wavelength uv rays centered on 352 nm, where wavelengths below 320 nm and visible light from the Hg lamp are cut by the filter. The visible light (hereafter VIS) source is a blue-light-emitting diode (LED E1L55-7B0A, Toyoda Gosei) centered on $475 \text{ nm}(\sim 2.6 \text{ eV})$ with a photon energy slightly higher than the N-related defects (2.48 eV) and well lower than the optical band gap of N-doped anatase.

The concentration of methylene blue (C_{MB}) after a 36-h uv illumination is summarized in Table I. As can be clearly seen in Table I, there is a strong correlation with C_N . We found that the photocatalytic activity was enhanced under ultraviolet illumination owing to the N doping at low C_N , and it reached the highest decomposition efficiency at C_N $= 1.69$ at. %. The enhancement is explained by the small band-gap shift (resulting from N_O) and charge-trapping sites (N_i) , which raise the efficiency of electron-hole generation and prevent undesirable recombination. The degraded photocatalytic efficiency at the high dopant levels is reasonable considering the interfusions of excessive N*ⁱ* defects, which can lower the carrier mobility and act as recombination centers. The C_{MB} after a 72-h VIS illumination is summarized in Table I, as well. Here, we also find that the decomposition rate of the N-doped $TiO₂$ is enhanced. The improved efficiency originates from visible light absorptive defects considering the photon energy of the LED light. The decomposition from the undoped $TiO₂$ sample is likely caused by defect levels within the band gap or direct photodecomposition of methylene blue. Here, we have to note the degree of enhancement through N doping is not significant compared to the uv irradiation experiment. More importantly, the decomposition rate is not strongly affected by the C_N in the film when it is larger than 1 at. %. The improvement up to 1 at. % is reasonable, since the optical absorption is enhanced by N, as can be seen in Fig. $3(a)$. Considering the use of N-doped $TiO₂$ under solar light or fluorescence light, where both contain a small amount of uv light, N doping is still very attractive, since it can benefit from the visible and

FIG. 4. C_N dependence of the lattice distortion represented by the *c*-axis length. The dashed line is a guide indicating the possible behavior as a function of the C_N .

uv spectra. Here, we point out that C_N of 1–2 % is the optimum concentration for practical applications, and more N is not necessary; more N turns colorless $TiO₂$ yellowish and degrades its efficiency.

Another finding of the adverse impact of N*ⁱ* is on the crystallographic structure. For the anatase polymorph, the structural variation following the replacement of one O atom with N in the anatase is less pronounced, because Ti-N bond lengths $(1.964$ and 2.081 Å) are only slightly longer than the Ti-O ones $(1.942$ and 2.002 Å) in the matrix.²⁵ Therefore, N incorporation at an O site may not degrade the crystallinity. A small concentration of N even promotes the growth of anatase, as reported by Lindgren *et al.*¹² However, it is controversial at high dopant levels. Figure 4 shows the distortion of the lattice constant represented by the *c*-axis lengths as a function of C_N in the ablation target. The errors correspond to the dispersion of the data for different films grown under the same conditions or to the experimental error in the determination of the lattice constant by XRD. The lattice was expanded by N incorporation at $C_N = 2$ at. %; however, it decreases with further increments of C_N . The first expansion originates from the replacement of the Ti-O bonds with Ti -N ones. The following shrinkage behavior suggests that the amount of N_i is increasing, as the lattice is compressed by the interstitial dopant. At the same time, the N*ⁱ* inside the anatase elevates the density of the $TiO₂$ matrix by their volume, which may act as a driving force of the structural phase transformations.27 For instance, the density of rutile (4.25 g cm^{-3}) is higher than that of anatase (3.89 g cm^{-3}) .

To elucidate the dominant growth mode with various N concentrations, we have grown films on amorphous $SiO₂$ substrates at 300 °C. There is no epitaxial relationship here, and films can grow free from the influence of substrates. Figure 5 shows the Raman shift of the films, where C_N $=(a)$ 2, (b) 10, and (c) 30 at. % targets were used, respectively. In Fig. 5, the peaks of the 442 and 610 cm⁻¹ bands marked with *R* indicate the existence of the rutile phase, while the 399, 520, and 643 cm−1 bands marked with *A* are the anatase ones. 28 The film deposited at a low dopant concentration clearly exhibits the anatase spectra partially mixed with the rutile modes. With increasing C_N , the characteristic band of anatase at 520 cm^{-1} is weakened. Also, the anatase peaks of 399 and 643 cm−1 are gradually shifting toward the

FIG. 5. (Color online) Raman spectra of the as-grown N-doped TiO₂ films on amorphous SiO_2 substrates with various C_N : (a) 2, (b) 10, and (c) 30 at. %. The Raman shifts from rutile and anatase are indicated by *R* and *A*, respectively.

rutile peaks of 442 and 610 cm⁻¹, respectively. These results reveal that the amount of the rutile phase increases at the higher dopant levels, indicating that the growth of anatase becomes more difficult. This is consistent with the broadened FWHM of the XRD spectra of the anatase as shown in Table I. The superiority of the photocatalytic activity using the anatase phase is claimed elsewhere. Moreover, the blueshift in the absorption edge is pointed out in the N-doped rutile $TiO₂$.^{25,29} This observation also explains the origin of drawbacks, in addition to the N-related defect formation at high N-doping levels.

The difficulty imposed upon realizing higher photocatalytic efficiency and/or band-gap narrowing in anatase is the existence of large amounts of N_i , because the N that is effective for them is not N_i but N_O . Actually, at lower concentrations, the substitute N does not cause band-gap narrowing, and the N 2*p* states are localized just above the O 2*p* valence band as recently proven by Valentin *et al.*²⁵ In this regard, only a small redshift of E_g obtained in this work is reasonable. The theoretical prediction for the band-gap narrowing calculated by Asahi *et al.* was modeled by replacing an O atom with N in the 16 atom of O in the unit cell of anatase, producing an \sim 4.2 at. % N_O doping;⁷ however, such a high substitutional doping level is unrealistic in our work. This observation leads to the question as to whether a further

band-gap narrowing occurs. Our results clearly indicate that N_O is saturated at low dopant levels, and there exists a large amount of N_i at high dopant levels. However, the possibility remains that one could obtain a further redshift of the absorption edge, if the amount of N_O were successfully increased, for instance, through low-temperature nonequilibrium growth conditions.

IV. SUMMARY

We have given microscopic insights into the nature of the photocatalytic N-doped anatase films. The N-doped epitaxial anatase films were artificially prepared by employing a hightemperature buffer layer as nucleation layer. This technique enabled anatase growth even at a temperature lower than the crystallographic transition of LaAlO₃, and led to the enrichment of N up to 3.85 at. %. The transparent film color gradually turns to yellowish as C_N increases, indicating that the optical absorption is enhanced. However, the optical bandgap narrowing effect anticipated by N doping did not occur with the increment of C_N . This is explained by the saturation of N doping into O sites, and the excessive N is interstitially doped within the matrix to produce deep-level defects. When these films are exposed to uv light, the efficiency is strongly affected by C_N . At low dopant concentrations, the photocatalytic decomposition rate was improved, and it reached the highest value near $C_N = 1.7$ at. %. Here, the existence of the optimum value is explained by the slight redshift of the band gap and visible light absorptive N*ⁱ* defects which degrade the efficiency when there are too many. On the contrary, in the case of VIS light, the efficiency is less affected by C_N , especially when it is larger than 1 at. %. In addition, the study of x-ray and Raman spectroscopy revealed that the growth of anatase became rather difficult, and the stable phase was rutile at higher N concentrations. These observations lead us to conclude that optimum C_N is 1–2 at. % provided that TiON is grown under typical conditions.

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