

**Electronic structure of  $\text{Ca}_3\text{CoXO}_6$  ( $X=\text{Co, Rh, Ir}$ ) studied by x-ray photoemission spectroscopy**K. Takubo,<sup>1</sup> T. Mizokawa,<sup>1,2</sup> S. Hirata,<sup>1</sup> J.-Y. Son,<sup>2</sup> A. Fujimori,<sup>1</sup> D. Topwal,<sup>3</sup> D. D. Sarma,<sup>3</sup> S. Rayaprol,<sup>4</sup> and E.-V. Sampathkumaran<sup>4,5</sup><sup>1</sup>*Department of Physics and Department of Complexity Science and Engineering, University of Tokyo, 5-1-5 Kashiwanoha, Chiba 277-8581, Japan*<sup>2</sup>*PRESTO, Japan Science and Technology Agency, 4-1-8 Honcho Kawaguchi, Saitama, Japan*<sup>3</sup>*Solid State and Structural Chemistry Unit, Indian Institute of Science, Bangalore 560012, India*<sup>4</sup>*Tata Institute of Fundamental Research, Homi Bhabha Road, Mumbai 400005, India*<sup>5</sup>*Institute of Solid State Physics, University of Tokyo, 5-1-5 Kashiwanoha, Chiba 277-8581, Japan*

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We have obtained Co  $2p$ , Rh  $3d$ , Ir  $4f$ , and valence-band x-ray photoemission spectra of  $\text{Ca}_3\text{Co}_2\text{O}_6$ ,  $\text{Ca}_3\text{CoRhO}_6$ ,  $\text{Ca}_3\text{CoIrO}_6$ , and  $\text{Sr}_3\text{ZnIrO}_6$  in order to determine the valence states of Co. The spectral features are consistent with the picture that Co is in the trivalent state at both trigonal prismatic and octahedral sites with high- and low-spin states, respectively, in  $\text{Ca}_3\text{Co}_2\text{O}_6$ , whereas Co is in the high-spin  $\text{Co}^{2+}$  state for  $\text{Ca}_3\text{CoRhO}_6$  and  $\text{Ca}_3\text{CoIrO}_6$ . This conclusion resolves a recent controversy regarding the valence and spin states of Co in  $\text{Ca}_3\text{Co}_2\text{O}_6$  and  $\text{Ca}_3\text{CoRhO}_6$ . This is the first x-ray photoemission spectroscopic study of this family of oxides.

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Recently, Co containing spin-chain compounds, e.g.,  $\text{Ca}_3\text{MXO}_6$  ( $M=\text{Co}$ ,  $X=\text{Co}$ , Rh, and Ir), obtained from the  $\text{K}_4\text{CdCl}_6$ -derived rhombohedral structure, have been attracting a lot of attention due to many exotic features exhibited by them in their properties (see, for instance, Refs. 1–25). These compounds are characterized by the presence of  $M$ - $X$  spin chains running along the  $c$  direction intervened by Ca ions, with a (distorted) trigonal prismatic and octahedral coordination around  $M$  and  $X$ , respectively. The compound  $\text{Ca}_3\text{Co}_2\text{O}_6$  is also found to be characterized by large thermopower and therefore there are discussions in the literature about the applicability of this compound or its derivatives for thermoelectric applications.<sup>12</sup>

A survey of the current literature, which can be found in Ref. 19, suggests that there is a dispute, particularly aroused by a report of Vidya *et al.*,<sup>17</sup> on the nature of valence/spin states of Co in  $\text{Ca}_3\text{Co}_2\text{O}_6$ , broadly speaking, between two possibilities: (i) Co at both the sites is trivalent,<sup>3,18</sup> (ii) Co exists in the high-spin  $2+$  and low-spin  $4+$  states (since the octahedral coordination in this compound is known to favor the low-spin state, according to crystal-field calculations).<sup>17</sup>

In order to enable further understanding of this compound, particularly for any theoretical framework, it is crucial to resolve this issue. Recently, <sup>59</sup>Co NMR measurements were performed<sup>19</sup> keeping this issue in mind, and the results established that Co is in the high- and low-spin  $d^6$  configurations at the  $M$  and  $X$  sites, respectively, at least below 15 K. However, the NMR signal could not be observed at higher temperatures and there is an urgent need to explore whether there is a change in the valence/spin state of Co (favoring the second alternative) as the temperature is increased, as often discussed in perovskite-based cobaltites.

A survey of the literature for  $\text{Ca}_3\text{CoRhO}_6$  suggests that there are similar disputes in this compound as well. While the magnetic susceptibility data<sup>20</sup> indicates a divalent state

for the Co site, the neutron-diffraction experiment<sup>23</sup> and band calculations<sup>18</sup> appear to favor the trivalent (high-spin) state. Similar to  $\text{Ca}_3\text{Co}_2\text{O}_6$ ,  $\text{Ca}_3\text{CoRhO}_6$  has two possibilities: (i) Co and Rh are trivalent, (ii) Co exists in the high-spin  $2+$  state and Rh in the low-spin  $4+$  state.

In order to address these issues in the high-temperature range of these compounds, we have performed photoemission measurements at room temperature. The results for  $\text{Ca}_3\text{Co}_2\text{O}_6$  conclusively establish that Co is indeed in the trivalent state favoring the first possibility, even at 300 K, thereby resolving the major controversy. On the other hand, the photoemission data for  $\text{Ca}_3\text{CoRhO}_6$  indicate that Co is in the divalent state favoring the second possibility. We have in addition carried out similar measurements on  $\text{Ca}_3\text{CoIrO}_6$ ,<sup>10,11</sup> as well as on  $\text{Sr}_3\text{ZnIrO}_6$ , as suitable reference materials, since it is known that Co (Zn) and Ir are in the  $2+$  and  $4+$  states in  $\text{Ca}_3\text{CoIrO}_6$  and  $\text{Sr}_3\text{ZnIrO}_6$ .<sup>24,25</sup>

Polycrystalline samples were synthesized by a solid-state reaction<sup>8</sup> and characterized by x-ray diffraction to be single phase. The photoemission spectroscopy measurements for  $\text{Ca}_3\text{CoXO}_6$  ( $X=\text{Co}$ , Rh and Ir) were performed at room temperature using a JPS9200 spectrometer equipped with a monochromatized Al  $K\alpha$  x-ray source ( $h\nu=1486.6$  eV). The total-energy resolution was  $\sim 0.6$  eV. The pressure of the spectrometer was  $\sim 8 \times 10^{-10}$  mb during the measurements. The polycrystalline samples were fractured *in situ* in order to obtain clean surfaces. All photoemission data were collected within 24 h after fracturing. The photoemission spectroscopy for  $\text{Sr}_3\text{ZnIrO}_6$  has been performed at room temperature using a custom built multitechnique VSW electron spectrometer with a resolution of 0.5 eV and a vacuum of  $\sim 8 \times 10^{-10}$  mb.

The high-spin  $\text{Co}^{2+}$  state is known to result in an intense satellite structure in the Co  $2p$  XPS spectrum.<sup>26</sup> The main peak at  $\sim 780$  eV and the satellite structure at  $\sim 787$  eV are

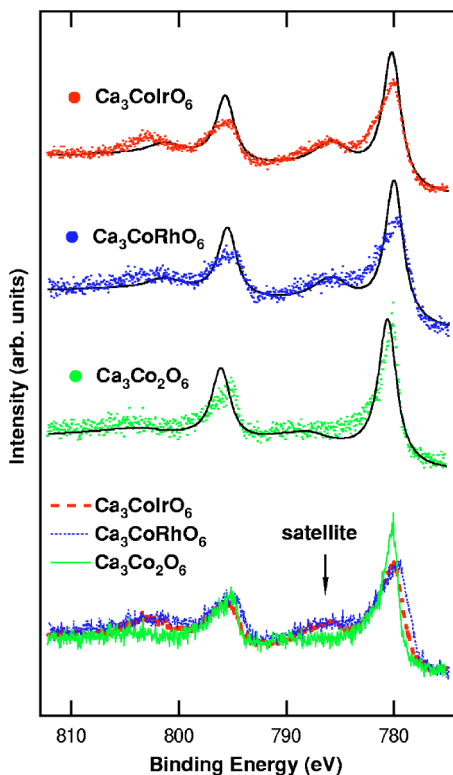


FIG. 1. Co  $2p$  photoemission spectra of  $\text{Ca}_3\text{Co}_2\text{O}_6$ ,  $\text{Ca}_3\text{CoRhO}_6$ , and  $\text{Ca}_3\text{CoIrO}_6$ . The calculated spectra are shown by solid curves. The Co  $2p$  spectra of  $\text{Ca}_3\text{CoRhO}_6$  and  $\text{Ca}_3\text{CoIrO}_6$  are overlaid with that of  $\text{Ca}_3\text{Co}_2\text{O}_6$  to show the difference in the satellite region.

derived from charge-transfer-type screening on the Co  $2p$  core hole. The main and satellite peaks correspond to the well screened and poorly screened final states, respectively. In contrast to the  $\text{Co}^{2+}$  state, the satellite structure of the Co  $2p$  XPS spectrum for the  $\text{Co}^{3+}$  state has only a small intensity.<sup>26–28</sup> The difference arises from an enhanced screening in the case of  $\text{Co}^{3+}$  due to a larger Co  $3d$ –O  $2p$  hopping strength and a smaller charge-transfer energy compared to those for the  $\text{Co}^{2+}$  state.<sup>26,29,30</sup> Therefore the Co  $2p$  XPS spectrum is very useful to distinguish the two possibilities:  $M$  and  $X$  are trivalent or  $M$  exists in the high-spin  $2+$  state and  $X$  in the low-spin  $4+$  state.

Figure 1 shows the Co  $2p$  XPS spectra of  $\text{Ca}_3\text{Co}_2\text{O}_6$ ,  $\text{Ca}_3\text{CoRhO}_6$ , and  $\text{Ca}_3\text{CoIrO}_6$ . The Co  $2p$  XPS spectra of  $\text{Ca}_3\text{CoRhO}_6$  and  $\text{Ca}_3\text{CoIrO}_6$  show intense satellite structures that are very similar to that of  $\text{CoO}$  with the high-spin  $\text{Co}^{2+}$  configuration and are very different from that of  $\text{LiCoO}_2$  with the low-spin  $\text{Co}^{3+}$  configuration.<sup>26,27</sup> This result indicates that  $\text{Ca}_3\text{CoRhO}_6$  and  $\text{Ca}_3\text{CoIrO}_6$  have the high-spin  $\text{Co}^{2+}$  species. On the other hand, the satellite structure is not visible at all in  $\text{Ca}_3\text{Co}_2\text{O}_6$  as shown in Fig. 1. The apparent absence of the satellite shows that  $\text{Ca}_3\text{Co}_2\text{O}_6$  has no high-spin  $\text{Co}^{2+}$  site and consists solely of  $\text{Co}^{3+}$ . These observations for the Co  $2p$  spectra indicate that the prismatic Co site is in the  $\text{Co}^{2+}$  high-spin configuration in both  $\text{Ca}_3\text{CoRhO}_6$  and  $\text{Ca}_3\text{CoIrO}_6$ . In contrast,  $\text{Ca}_3\text{Co}_2\text{O}_6$  consists of  $\text{Co}^{3+}$  ions only.

In order to confirm this conclusion, we have performed

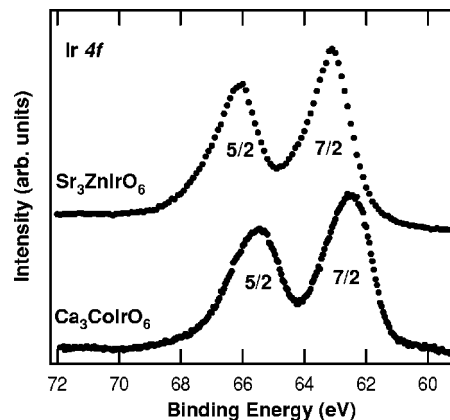


FIG. 2. Ir  $4f$  photoemission spectra of  $\text{Ca}_3\text{CoIrO}_6$  and  $\text{Sr}_3\text{ZnIrO}_6$ .

configuration-interaction (CI) calculations on the prismatic and octahedral  $\text{CoO}_6$  cluster models. We have followed the conventional CI method that has been widely used to analyze photoemission spectra of transition-metal oxides.<sup>26,29,30</sup> In the CI approach, the ground state is expressed by a linear combination of  $d^n$ ,  $d^{n+1}L$ ,  $d^{n+2}L^2$ , etc., where  $L$  denotes a hole at the O  $2p$  level. The values of  $n$  are 6 and 7 for the  $\text{Co}^{3+}$  and  $\text{Co}^{2+}$  cases, respectively. The energy difference between  $d^n$  and  $d^{n+1}L$  is given by charge-transfer energy  $\Delta$  and that between  $d^{n+1}L$  and  $d^{n+2}L^2$  by  $\Delta+U$ , where  $U$  is the Coulomb repulsion energy between the Co  $3d$  electrons. The final states are expressed by linear combinations of  $cd^n$ ,  $cd^{n+1}L$ ,  $cd^{n+2}L^2$  etc., where  $c$  denotes a hole at the Co  $2p$  core level. The energy difference between  $cd^n$  and  $cd^{n+1}L$  is given by  $\Delta-Q$  and that between  $cd^{n+1}L$  and  $cd^{n+2}L^2$  by  $\Delta-Q+U$  where  $Q$  is the Coulomb energy between the Co  $2p$  core hole and the Co  $3d$  electron. Here, the ratio  $U/Q$  is fixed at  $\sim 0.7$ . The transfer integrals between  $cd^n$  and  $cd^{n+1}L$ , etc., are expressed by  $(pd\sigma)$  and  $(pd\pi)$  with the ratio  $(pd\sigma)/(pd\pi)$  of  $-2.16$ . As a consequence, the Co  $2p$  spectrum can be calculated with the three adjustable parameters  $\Delta$ ,  $U$ , and  $(pd\sigma)$ . As shown in Fig. 1, the calculated spectra can explain the satellite structure observed in  $\text{Ca}_3\text{CoRhO}_6$  and  $\text{Ca}_3\text{CoIrO}_6$ . The obtained parameters are  $\Delta=5.0$  eV,  $U=6.5$  eV,  $(pd\sigma)=-1.2$  eV for  $\text{Ca}_3\text{CoRhO}_6$  and  $\Delta=5.0$  eV,  $U=6.5$  eV,  $(pd\sigma)=-1.1$  eV for  $\text{Ca}_3\text{CoIrO}_6$ . These values are reasonable for  $\text{Co}^{2+}$  oxides. On the other hand, it is almost impossible to fit the Co  $2p$  spectrum of  $\text{Ca}_3\text{Co}_2\text{O}_6$  by adjusting the parameters since it has the contributions from the two different Co sites. Here, we estimate  $(pd\sigma)$  for the two Co sites using Harrison's rule. Also we employ typical  $\Delta$  value for  $\text{Co}^{3+}$  oxides.<sup>26,29,30</sup> Using the parameter sets of  $\Delta=2.0$  eV,  $U=6.5$  eV,  $(pd\sigma)=-1.7$  eV for the octahedral site (low-spin  $\text{Co}^{3+}$ ) and  $\Delta=2.0$  eV,  $U=6.5$  eV,  $(pd\sigma)=-1.3$  eV for the prismatic site (high-spin  $\text{Co}^{3+}$ ), the calculated Co  $2p$  spectrum agrees well with the experimental result as shown in Fig. 1.

Figure 2 shows the Ir  $4f$  XPS spectra of  $\text{Ca}_3\text{CoIrO}_6$  and  $\text{Sr}_3\text{ZnIrO}_6$ . Since Sr and Zn are definitely divalent, Ir in  $\text{Sr}_3\text{ZnIrO}_6$  must be tetravalent or  $\text{Ir}^{4+}$ . Therefore  $\text{Sr}_3\text{ZnIrO}_6$  is a good reference sample to judge whether Ir is trivalent or tetravalent in  $\text{Ca}_3\text{CoIrO}_6$ . As shown in Fig. 2, the binding

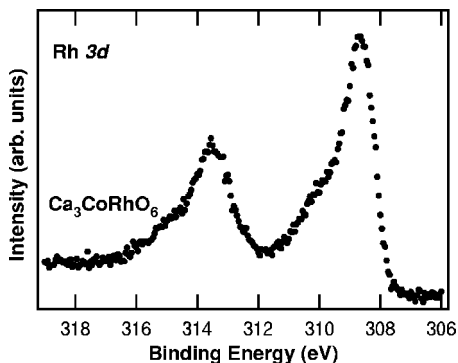


FIG. 3. Rh 3d photoemission spectra of  $\text{Ca}_3\text{CoRhO}_6$ .

energy of Ir  $4f_{7/2}$  is 62.5 eV for  $\text{Ca}_3\text{CoIrO}_6$  and is 63.1 eV for  $\text{Sr}_3\text{ZnIrO}_6$ . The valence-band spectrum of  $\text{Sr}_3\text{ZnIrO}_6$  is shifted by  $\sim 0.5$  eV to the higher binding energy side compared to that of  $\text{Ca}_3\text{CoIrO}_6$  (see below). The binding-energy shift in  $\text{Sr}_3\text{ZnIrO}_6$  is probably caused by charging effect. After correction for the binding-energy shift, the Ir  $4f_{7/2}$  binding energy of  $\text{Sr}_3\text{ZnIrO}_6$  becomes  $\sim 62.6$  eV, a value which is very close to the Ir  $4f_{7/2}$  binding energy of  $\text{Ca}_3\text{CoIrO}_6$ . This result suggests that Ir in  $\text{Ca}_3\text{CoIrO}_6$  also has  $\text{Ir}^{4+}$ . This argument with the Ir  $4f$  spectra naturally leads us to the conclusion that the octahedral Ir site has  $\text{Ir}^{4+}$  ions and, consequently, the prismatic Co site should be in  $\text{Co}^{2+}$  configuration. This result is consistent with the conclusion from magnetic susceptibility experiments.<sup>10</sup> Therefore comparison of the  $\text{Ir}^{4+}$  spectra between  $\text{Ca}_3\text{CoIrO}_6$  and  $\text{Sr}_3\text{ZnIrO}_6$  also supports the valence states obtained from the Co  $2p$  spectra. Figure 3 shows the Rh  $3d$  XPS spectra of  $\text{Ca}_3\text{CoRhO}_6$ . The Rh  $3d_{5/2}$  binding energy of  $\text{Ca}_3\text{CoRhO}_6$  (308.7 eV) is much higher than that of  $\text{Rh}_2\text{O}_3$  (308.2 eV) and is rather close to that of  $\text{RhO}_2$  (308.5 eV).<sup>31</sup> This result also suggests that the valence state of the octahedral Rh site should be far away from 3+, which means that Rh is 4+. A higher binding energy compared to  $\text{RhO}_2$  also should not be taken as an indication for a still higher valence, as this yields an unphysical valence for Co (close to 1+). This conclusion that is consistent with that from the Co  $2p$  spectra.

Figure 4 shows the valence-band XPS spectra of  $\text{Ca}_3\text{Co}_2\text{O}_6$ ,  $\text{Ca}_3\text{CoRhO}_6$ ,  $\text{Ca}_3\text{CoIrO}_6$ , and  $\text{Sr}_3\text{ZnIrO}_6$ . Since the photoionization cross sections of the Rh  $4d$  and Ir  $5d$  orbitals are much larger than that of the Co  $3d$  orbital at photon energy of 1486.6 eV, the valence-band spectra of  $\text{Ca}_3\text{CoRhO}_6$ ,  $\text{Ca}_3\text{CoIrO}_6$ , and  $\text{Sr}_3\text{ZnIrO}_6$  are dominated by contributions from the octahedral Rh or Ir sites. On the other hand, in the valence-band spectrum of  $\text{Ca}_3\text{Co}_2\text{O}_6$ , contributions from the octahedral Co site are comparable to those from the prismatic Co site. The energy splitting between structures A and B in  $\text{Ca}_3\text{CoIrO}_6$  and  $\text{Sr}_3\text{ZnIrO}_6$  is larger than that between structures A'' and B'' in  $\text{Ca}_3\text{Co}_2\text{O}_6$ . Structures A (A'') and B (B'') can be assigned to the antibonding and bonding bands of the Ir  $5d$  (Co  $3d$ ) and O  $2p$  orbitals. The energy splitting becomes smaller in going from  $\text{Ca}_3\text{CoIrO}_6$  to  $\text{Ca}_3\text{CoRhO}_6$  to  $\text{Ca}_3\text{Co}_2\text{O}_6$  since the hybridization strength with the O  $2p$  orbitals becomes smaller in going from Ir  $5d$  to Rh  $4d$  to Co  $3d$ . Another important point is that sharp structures A, A', and A'' are rather intense and located just

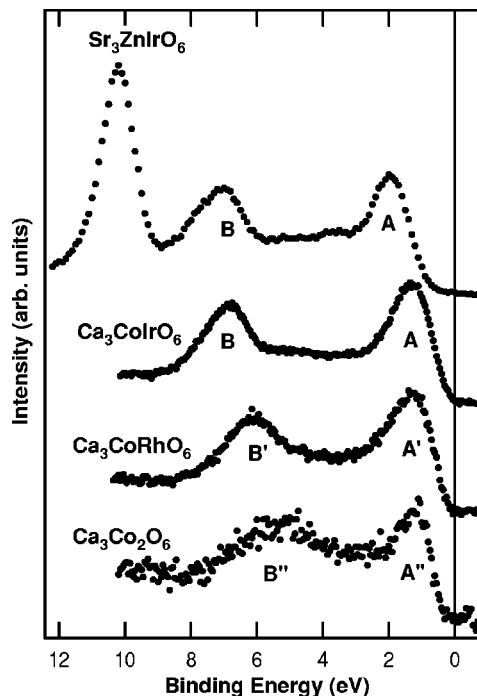


FIG. 4. Valence-band photoemission spectra of  $\text{Ca}_3\text{Co}_2\text{O}_6$ ,  $\text{Ca}_3\text{CoRhO}_6$ ,  $\text{Ca}_3\text{CoIrO}_6$ , and  $\text{Sr}_3\text{ZnIrO}_6$ .

below the Fermi level in  $\text{Ca}_3\text{CoXO}_6$  ( $X=\text{Co}, \text{Rh}, \text{Ir}$ ). For typical low-spin compounds such as  $\text{LiCoO}_2$ , the valence-band XPS spectra commonly have a sharp and intense peak (with a width of  $\sim 1.5$  eV) just below the Fermi level that is assigned to the  $t_{2g}$  band fully occupied by electrons.<sup>26</sup> On the other hand, since the  $t_{2g}$  and  $e_g$  bands are partially occupied in the high-spin configurations, an antibonding band for a high-spin state consists of  $t_{2g}$  and  $e_g$  components and is not so sharp as that for the low-spin state. In addition, for the  $\text{Co}^{3+}$  case, while the low-spin ground state ( $S=0$ ) has only the  $S=1/2$  final state, the high-spin ground state ( $S=2$ ) gives the  $S=5/2$  and  $S=3/2$  final states that are split by the exchange interaction. As a result, for the high-spin  $\text{Co}^{3+}$  state, one should have seen a much broader peak near the Fermi level, which is not observed. Therefore, for  $\text{Ca}_3\text{CoXO}_6$  ( $X=\text{Co}, \text{Rh}, \text{Ir}$ ), the existence of structures A, A', and A'' indicates that one of the Co and X sites should have a low-spin configuration in which the  $t_{2g}$  band is fully (or almost fully) occupied and the  $e_g$  band is fully unoccupied by electrons. Since it is well established that the prismatic  $M$  site is in the high-spin state, the above argument leads us to the conclusions that the  $\text{Ir}^{4+}$  and  $\text{Rh}^{4+}$  sites in  $\text{Ca}_3\text{CoXO}_6$  ( $X=\text{Rh}, \text{Ir}$ ) take the low-spin  $d^5(t_{2g}^5)$  state and that the octahedral  $\text{Co}^{3+}$  site in  $\text{Ca}_3\text{Co}_2\text{O}_6$  takes the low-spin  $d^6(t_{2g}^6)$  state.

In conclusion, we have studied the electronic structure of quasi-one-dimensional Co oxides  $\text{Ca}_3\text{CoXO}_6$  ( $X=\text{Co}, \text{Rh}, \text{Ir}$ ) using x-ray photoemission spectroscopy. The Co  $2p$  spectra of  $\text{Ca}_3\text{CoRhO}_6$  and  $\text{Ca}_3\text{CoIrO}_6$  show intense satellite structure indicating that the prismatic Co site is in the high-spin  $\text{Co}^{2+}$  configuration. The Ir  $4f$  and valence-band spectra of  $\text{Ca}_3\text{CoIrO}_6$  and  $\text{Sr}_3\text{ZnIrO}_6$  indicate that the octahedral Ir sites take the low-spin  $\text{Ir}^{4+}$  configuration in

$\text{Ca}_3\text{CoIrO}_6$ . In contrast, the Co  $2p$  and valence-band spectra of  $\text{Ca}_3\text{Co}_2\text{O}_6$  show that  $\text{Ca}_3\text{Co}_2\text{O}_6$  consists of  $\text{Co}^{3+}$  species and that one of the prismatic and octahedral sites should take the low-spin configuration. Considering the fact that some Co sites have considerable magnetic moment in  $\text{Ca}_3\text{Co}_2\text{O}_6$ , the prismatic Co and octahedral Co sites take the  $\text{Co}^{3+}$  high-spin and  $\text{Co}^{3+}$  low-spin configurations, respectively, in

$\text{Ca}_3\text{Co}_2\text{O}_6$ . The present results are able to resolve the issues on Co valence for  $\text{Ca}_3\text{Co}_2\text{O}_6$  and  $\text{Ca}_3\text{CoRhO}_6$  at least around 300 K.

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