Phase control of La_2CuO_4 **in thin film synthesis**

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The lanthanum copper oxide, La_2CuO_4 , which is an end member of the prototype high- T_c superconductors $(La, Sr)_{2}CuO_{4}$ and $(La, Ba)_{2}CuO_{4}$, crystallizes in the K₂NiF₄ structure in high-temperature bulk synthesis. The crystal chemistry, however, predicts that La_2CuO_4 is at the borderline of the K₂NiF₄ stability and that it can crystallize in the Nd_2CuO_4 structure at low synthesis temperatures. In this article we demonstrate that lowtemperature thin-film synthesis actually crystallizes La_2CuO_4 in the Nd_2CuO_4 structure. We also show that the phase control of K₂NiF₄-type La₂CuO₄ versus Nd₂CuO₄-type La₂CuO₄ can be achieved by varying the synthesis temperature and using different substrates.

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I. INTRODUCTION

The rare earth copper oxides of the general chemical formula R_2 CuO₄ take two different crystal structures: K₂NiF₄ (abbreviated as T) and Nd_2CuO_4 (T'). The structural difference between T and T' can be viewed simply as the difference in the *R*-O arrangements: rocksaltlike versus fluoritelike. With regard to the Cu-O coordination, however, there is a significant difference: T has octahedral $CuO₆$, whereas T' has two-dimensional square-planar CuO₄. Empirically, the former accepts only hole doping, the latter only electron doping. The T structure is formed with large La^{3+} ions, while the T' structure is formed with smaller R^{3+} ions, such as $R = Pr$, Nd, Sm, Eu, and Gd.¹ The T-T' boundary lies between La^{3+} and Pr^{3+} . Namely, La_2CuO_4 is at the borderline of the T-phase stability.

The crystal chemistry of the rare earth copper oxides has been explained in terms of the crystallographic tolerance factor (t) , 2,3 which is defined as

$$
t = \frac{r_i(R^{3+}) + r_i(O^{2-})}{\sqrt{2}\{r_i(Cu^{2+}) + r_i(O^{2-})\}},
$$
\n(1)

where $r_i(R^{3+})$, $r_i(Cu^{2+})$, and $r_i(O^{2-})$ are the ionic radii for R^{3+} , Cu²⁺, and O²⁻ ions. The *t* values for La₂CuO₄ and Pr_2CuO_4 are evaluated as 0.8685 and 0.8562 using the roomtemperature ionic radii by Shannon and Prewitt. 4 From the extensive data collected for a variety of $R_2CuO₄$ -type cuprates, the critical (room-temperature) value for the $T\rightarrow T'$ transition is presumed to be t_c =0.865, below which T is unstable. $2,3$

The different thermal expansion ("thermal-expansion") mismatch'') between the R -O and Cu-O bond lengths plays an important role in the T -versus- T' stability as pointed out initially by Manthiram and Goodenough.⁵ The "ionic" *R*-O bond has a larger thermal expansion than the ''covalent'' Cu-O bond, which leads to the increase in *t* with increasing temperature. Hence the T phase is stable at high temperatures whereas the T' phase is stable at low temperatures. In the case of La_2CuO_4 , the transition from T to T' is predicted to occur at around 700 K (427 °C), where $t(700 \text{ K}) \sim 0.88$. There have been a few attempts to stabilize the $T[']$ phase of $La₂CuO₄$ in the past. However, a conventional solid-state reaction method requires a firing temperature of at least 500 °C even with coprecipitated fine powders, so it could not produce single-phase T' -La₂CuO₄. Bulk synthesis of T' -La₂CuO₄ has been achieved only by a very special recipe as given by Chou *et al.*⁶ Their method consists of the following two steps. The first step is to reduce $T-La_2CuO_4$ with hydrogen around 300 °C and obtain the Sr_2CuO_3 -like phase. The second step is to convert the $Sr₂CuO₃$ -like phase to T'-La₂CuO₄ by reoxygenation below 400 °C. The resultant product was single-phase T' , although x-ray peaks were broadened due to the considerable lattice disorder and defects.

In thin-film synthesis, the reaction temperature can be lowered significantly, since reactants are much smaller in size and also more reactive than in bulk synthesis. The reactants in thin-film synthesis are atoms or molecules or ions or clusters, depending on the technique employed. The limiting case is achieved by reactive coevaporation from metal sources, in which the reactants are atoms and the oxidation reaction is initiated on a substrate. Using this reactive coevaporation technique, we have learned from our ten-year experience that cuprate films crystallize at temperatures as low as 400 °C. This enabled us to synthesize single-phase T' -La₂CuO₄. In this article we describe the phase control of K_2NiF_4 -type La_2CuO_4 versus Nd_2CuO_4 -type La_2CuO_4 by varying the synthesis temperature and using different substrates.

TABLE I. Crystal structure and *a*-axis lattice constant (a_s) for the substrates used in this work. The in-plane lattice constants (a_0) for T'-La₂CuO₄ and T-La₂CuO₄ are also included.

Substrate	Abbreviation	a_s or a_0 (Å)	Crystal structure
MgO (100)	MGO	4.212	NaCl
$KTaO3$ (100)	KTO	3.989	perovskite
$SrTiO3$ (100)	STO	3.905	perovskite
$LaSrGaO4$ (001)	LSGO	3.843	K_2NiF_4
$NdGaO3$ (100)	NGO.	3.838	perovskite
LaAlO ₃ (100)	LAO	3.793	perovskite
$LaSrAlO4$ (001)	LSAO	3.755	K_2NiF_4
$PrSrAlO4$ (001)	PSAO	3.727	K_2NiF_4
$YAlO3$ (100)	YAO	3.715	perovskite
$NdSrAlO4$ (001)	NSAO	3.712	K_2NiF_4
$NdCaAlO4$ (001)	NCAO	3.688	K_2NiF_4
$ZrO_2(Y)$ (100)	YSZ.	3.616	fluorite
T' -La ₂ CuO ₄		4.005	Nd_2CuO_4
$T-La_2CuO_4$		3.803	K_2NiF_4

II. EXPERIMENT

We grew La_2CuO_4 thin films in a customer-designed molecular beam epitaxy chamber from metal sources using multiple electron-gun evaporators.⁷ The stoichiometry adjustment was performed by controlling the evaporation beam flux of each element using electron impact emission spectrometry via feedback loop to electron guns, and was refined with the aid of reflection high-energy electron diffraction $(RHEED)$ monitoring so as to minimize precipitates of the impurity phase on the surface.8 During growth, rf-activated atomic oxygen was used for oxidation. The chamber pressure during growth was 6×10^{-6} Torr. The substrate temperature was varied from $425\,^{\circ}\text{C}$ to $725\,^{\circ}\text{C}$. The growth rate was \sim 1.5 Å/s, and the film thickness was typically 450 Å. After the evaporation, most of the films were cooled to temperatures lower than 200 \degree C at a rate lower than 20 \degree C/min in 1×10^{-5} Torr molecular oxygen to avoid phase decomposition. Some of the films were cooled in vacuum or in ozone to investigate the change of the transport properties by excess oxygen.

In order to examine the substrate influence on the selective phase stabilization, 9 we used various substrates as listed in Table I. The in-plane lattice constant (a_s) ranges from 3.6 Å to 4.2 Å, which should be compared to $a_0 = 3.803$ Å for T-La₂CuO₄ and a_0 =4.000–4.010 Å for T'-La₂CuO₄ $(T-La_2CuO_4$ has an orthorhombic structure with *a'* = 5.3574 Å and $b' = 5.4005$ Å, and a_0 is calculated as $\sqrt{(a'b')/2}$. The crystal structures include perovskite, K_2N i F_4 , NaCl, and Ca F_2 (fluorite). We deposited films simultaneously on all the substrates listed in Table I, which were pasted to one substrate holder by Ag paint. This avoids run-to-run variations.

The lattice parameters and crystal structures of the films were determined using a standard x-ray diffractometer. Resistivity was measured by the standard four-probe method using electrodes formed by Ag evaporation.

FIG. 1. XRD patterns for La_2CuO_4 films grown on NCAO substrates at T_s =725–425 °C. The top two patterns are simulations for the T and T' structure. The broken and dotted lines indicate the peak positions of the (008) line for the T and T' structure, respectively. Peak positions of NCAO are indicated in the lowest figure.

III. RESULTS AND DISCUSSION

A. Effect of synthesis temperature on the selective phase stabilization

Figure 1 shows the x-ray diffraction (XRD) patterns of La_2CuO_4 films grown on NdCaAlO₄ (NCAO) substrates with different synthesis temperatures (T_s) . Since the *c*-axis lattice constant (c_0) is distinct between T and T' $[c_0(T)]$ = 13.15 Å versus $c_0(T)$ = 12.55 Å], the phase identification

FIG. 2. XRD patterns for La_2CuO_4 films grown on various substrates at $T_s = 525$ °C. The top two patterns are simulations for the T and T' structure. Substrate peaks are removed. The broken and dotted lines indicate the peak positions of the (008) line for the T and T' structure, respectively.

is rather straightforward. The calculated patterns for T and T' are also included in Fig. 1. The films grown at T_s >625 °C are single-phase T, while the films grown at T_s =500–550 °C are single-phase T'. The films grown at T_s $=$ 575–600 °C are a two-phase mixture of T and T' with T' more dominant for lower T_s . The films grown below T_s =475 °C show unidentified peaks at $2\theta \sim 31.4^{\circ}$ and 65.5°.

FIG. 3. Film's c_0 versus substrates a_s for La₂CuO₄ films grown at $T_s = 525$ °C on different substrates. The lattice constants of bulk T- and T'-La₂CuO₄ (a_0 =3.803 Å, c_0 =13.15 Å for T and a_0 =4.005 Å, c_0 =12.55 Å for T') are indicated by arrows together with a_s of the substrates. The circles connected by the vertical dotted lines indicate multiphase formation. The c_0 values of the T structure is noticeably substrate dependent because of epitaxial strain: the longest (c_0 =13.25 Å) for LSAO and the shortest (c_0 = 13.05 Å) for LSGO. The c_0 value of 12.8 Å on STO seems to correspond to the T*-like phase.

From this result, we can see the following trend for synthesis temperature on the selective phase stabilization. High T_s stabilizes T and low T_s stabilizes T'.

B. Effect of substrates on the selective phase stabilization

Figure 2 shows the XRD patterns of La_2CuO_4 films grown at $T_s = 525$ °C on different substrates. Of these films in this figure, the films on $KTaO₃$ (KTO), NCAO, and $ZrO_2(Y)$ (YSZ) are single-phase $T^{'},^{10}$ while the films on $LaSrGaO₄$ (LSGO), LaAlO₃ (LAO), LaSrAlO₄ (LSAO), $PrSrAIO₄$ (PSAO), and NdSrAlO₄ (NSAO) are single-phase T. On YAlO₃, the film is dominantly T' with a trace amount of T. On SrTiO₃ (STO) and NdGaO₃ (NGO), the films are clearly a mixture of T and T' . The film on STO contains some amount of the T^* -like phase.¹¹ On MgO (MGO), no clear peak is observed. The c_0 values of these films together with films on other substrates are summarized in Fig. 3. Because of epitaxial strain,¹² $c₀$ of the T structure is noticeably substrate dependent: the longest $(c_0=13.25 \text{ Å})$ for LSAO and the shortest $(c_0=13.05 \text{ Å})$ for LSGO. From these results, we can see the following trend for a substrate lattice parameter on the selective phase stabilization. Substrates with a_0 of 3.70–3.85 Å stabilize T, and substrates with a_0 of $>$ 3.90 Å or $<$ 3.70 Å stabilize T' (or destabilize T).

Next we mention the effect of substrate crystal structure on the selective phase stabilization. If, in Fig. 2, one compares the films grown on perovskite and K_2NiF_4 -type substrates with almost the same a_0 , for example, NGO (a_0) $=$ 3.838 Å) versus LSGO (a_0 = 3.843 Å) or YAO (a_0 $=$ 3.715 Å) versus NSAO (a_0 = 3.712 Å), one can notice the

FIG. 4. Phase diagram on the selective stabilization of T versus T' in the T_s - a_s plane. The crosses indicate no phase formation. The open circles represent single-phase T while the filled circles represent single-phase T'. The partially filled circles represent a twophase mixture. The size (area) of the circles is proportional to the XRD peak intensities of the (006) lines. For two-phase mixed films, the ratio of the unshaded and the shaded areas represent the ratio of the T and T' peak intensity of the (006) line. The results on LSGO and NSAO substrates are not included to avoid overlapping with the results on NGO and YAO. On LSGO and NSAO, the T structure is formed for 725 °C $> T_s$ $>$ 475 °C, and the T' structure is not formed for any T_s . The gray area at $a_s = 3.905$ Å (STO) indicate the formation of the *T**-like phase.

trend that K_2NiF_4 -type substrates have a tendency to stabilize the T structure rather than the $T[']$ structure.

C. Phase diagram in the T_s **-** a_s plane

Our survey was performed at T_s from 425 °C to 725 °C on all substrates in Table I. Figure 4 summarizes the results, which show the phase diagram on the selective stabilization of T versus T' in the T_s - a_s plane.

*1. High Ts (625***–***725 °C)*

The films on most of the substrates are single-phase T. There are three exceptional substrates: KTO, YAO, and YSZ. The films on KTO and YSZ do not show any definite x-ray peak. The film on YAO is a mixture of T and T' even at the highest temperature investigated. This can be explained by interdiffusion of Y from YAO substrates into La_2CuO_4 since Y substitution for La is known to stabilize the T' structure.

2. Low T_s (450–600 °C)

The films on the T-lattice-matched substrates (LSGO, LAO, LSAO, PSAO, and NSAO) are single-phase T. The films on T'-lattice-matched KTO and on fluorite YSZ are single-phase T' . The films on other substrates $(STO, NGO,$ YAO, and NCAO) are a two-phase mixture of T and T' with T' more dominant for lower T_s .

FIG. 5. Comparison of resistivity (ρ) -temperature (T) curves between T-La₂CuO_{4+ δ} and T'-La₂CuO_{4+ δ} films. The solid lines are for films cooled in vacuum (δ ~0) while the broken lines are for films cooled in ozone (δ >0). With vacuum cooling, the T film has much higher resistivity than the T' film. Ozone cooling causes a totally opposite effect on T and T' : the T film gets metallic and superconducting whereas the T' film gets more insulating.

D. Comparison of T -La₂CuO₄ and T' -La₂CuO₄

Next, we make a brief comparison of the physical properties of T-La₂CuO₄ and T'-La₂CuO₄, which have the same chemical formula but different crystal structures. Figure 5 shows the temperature dependences of resistivity for both the phases. The solid lines represent the ρ -*T* curves for the films cooled in vacuum to ambient temperature, which do not have excess oxygen but might have slight oxygen deficiencies $(La_2CuO_{4+\delta}$ with δ ~0). The broken lines represent those for the films cooled in ozone, which have interstitial excess oxygen (δ >0).¹³ The excess oxygen occupies the tetrahedral site in T, and the apical site in T' . The vacuum-cooled T film has much higher resistivity (by several orders of magnitudes at low temperatures) than the vacuum-cooled T' film. In fact, T'-La₂CuO₄ is metallic down to 180 K.¹⁴ The ozone cooling causes a totally opposite effect on T and T' . The resistivity of the T film gets lowered by five orders of magnitudes at room temperature from \sim 50 Ω cm to \sim 5 $\times 10^{-4} \Omega$ cm, indicating that holes doped by excess oxygen are itinerant. Furthermore the film becomes superconducting. In contrast, the resistivity of the T' film increases, indicating that holes doped by excess oxygen are localized.¹⁴

IV. SUMMARY

In summary, we have demonstrated that La_2CuO_4 can crystallize in the Nd_2CuO_4 structure using low-temperature thin-film synthesis. Furthermore the phase control of K_2NiF_4 -type La_2CuO_4 versus Nd_2CuO_4 -type La_2CuO_4 can be achieved by varying the synthesis temperature and also the substrate. The general trends are as follows: (i) high T_s stabilizes T and low T_s stabilizes T', (ii) substrates with a_s \sim 3.70-3.85 Å stabilize T and substrates with *a_s* $>$ 3.90 Å or a_s <3.70 Å stabilize T' (or destabilize T), and (iii) K_2NiF_4 -type substrates stabilize T.

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- ¹H. Müller-Buschbaum and W. Wollsglager, Z. Anorg. Allg. Chem. 414, 76 (1975).
- ² J.F. Bringley, S.S. Trail, and B.A. Scott, J. Solid State Chem. **86**, 310 (1990).
- 3A. Manthiram and J.B. Goodenough, J. Solid State Chem. **87**, 402 $(1990).$
- 4R.D. Shannon and C.T. Prewitt, Acta Crystallogr., Sect. B: Struct. Crystallogr. Cryst. Chem. **25**, 925 (1969).
- 5A. Manthiram and J.B. Goodenough, J. Solid State Chem. **92**, 231 $(1991).$
- 6F.C. Chou, J.H. Cho, L.L. Miller, and D.C. Johnston, Phys. Rev. B 42, 6172 (1990).

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- 7M. Naito and M. Hepp, Jpn. J. Appl. Phys., Part 2 **39**, L485 $(2000).$
- ⁸M. Naito and H. Sato, Appl. Phys. Lett. **67**, 2557 (1995).
- ⁹ I. Tsukada, Phys. Rev. B **64**, 224501 (2001).
- 10 On YSZ, judging from the peak positions, the film seems to be single-phase T' , while the peak intensity ratios do not agree with the calculated ratios.
- 11 H. Sawa, S. Suzuki, M. Watanabe, J. Akimitsu, H. Matsubara, H. Watabe, S. Uchida, K. Kokusho, H. Asano, F. Izumi, and E. Takayama-Muromachi, Nature (London) 337, 347 (1989).
- ¹²H. Sato and M. Naito, Physica C **274**, 221 (1997).
- ¹³ H. Sato, H. Yamamoto, and M. Naito, Physica C **280**, 178 (1997).
- 14 A. Tsukada, H. Yamamoto, H. Shibata, and M. Naito (unpublished).