

Percolation in three-dimensional random field Ising magnets

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The structure of the three-dimensional (3D) random field Ising magnet is studied by ground-state calculations. We investigate the percolation of the minority-spin orientation in the paramagnetic phase above the bulk phase transition, located at $[\Delta/J]_c \approx 2.27$, where Δ is the standard deviation of the Gaussian random fields ($J=1$). With an external field H there is a disorder-strength-dependent critical field $\pm H_c(\Delta)$ for the down (or up) spin spanning. The percolation transition is in the standard percolation universality class. $H_c \sim (\Delta - \Delta_p)^\delta$, where $\Delta_p = 2.43 \pm 0.01$ and $\delta = 1.31 \pm 0.03$, implying a critical line for $\Delta_c < \Delta \leq \Delta_p$. When, with zero external field, Δ is decreased from a large value there is a transition from the simultaneous up- and down-spin spanning, with probability $\Pi_{\uparrow\downarrow} = 1.00$ to $\Pi_{\uparrow\downarrow} = 0$. This is located at $\Delta = 2.32 \pm 0.01$, i.e., above Δ_c . The spanning cluster has the fractal dimension of standard percolation, $D_f = 2.53$ at $H = H_c(\Delta)$. We provide evidence that this is asymptotically true even at $H=0$ for $\Delta_c < \Delta \leq \Delta_p$ beyond a crossover scale that diverges as Δ_c is approached from above. Percolation implies extra finite-size effects in the ground states of the 3D random field Ising model.

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I. INTRODUCTION

The random field Ising model (RFIM) is one of the most basic models for random systems.¹⁻³ Its beauty is that the mixture of random fields and the standard Ising model creates rich physics and leaves many still unanswered problems. By now it is known that three dimensions (3D) is the cornerstone of the model, since it presents a phase transition where the randomness proves to be a relevant perturbation to the pure 3D Ising model. For the last 15 years, since the seminal paper by Ogielski,⁴ the studies of the transition have centered around zero-temperature ground-state computations because the temperature is due to renormalization group arguments believed to be an (perhaps dangerously so) irrelevant variable.

Many such works exist so far, the most recent and comprehensive being due to Middleton and Fisher.⁵ In spite of all the effort many uncertainties remain concerning the nature of the phase transition. The question is if the transition is of the second order, of traditional first-order type, or finally some other kind of discontinuous transition. The order-parameter exponent β may have a finite value or it can be equal to zero.⁵⁻¹² Its very small value makes it unlikely that insight will be obtained in the near future, in spite of the fact that the optimization algorithms used can at best scale almost linearly with the number of spins in the system. Moreover, a controversy exists with regards to the role of disorder: the available simulations are not able to settle the question whether the critical exponents depend on the particular choice of the distribution for the random fields, analogously to the mean-field theory of the RFIM where binary ($\pm h$) disorder results in a first-order transition and Gaussian (see below) in a second-order one.¹³

In this paper we focus on a novel aspect of the three-dimensional RFIM: namely, percolation.¹⁴ The goal is to explore percolation critical phenomena in the 3D RFIM. The

work is an extension to our studies of percolation in the two-dimensional RFIM.¹⁵ In the traditional 3D Ising model, without disorder, the percolation behavior in an applied field and its consequences, as whether the phase transition critical exponents would be affected by the percolation criticality, have been known for a long time as the “trouble with Kertész.”^{14,16} This problem was solved by Wang¹⁷ by studying Fortuin-Kasteleyn or Coniglio-Klein^{18,19} clusters using so-called ghost spins. In the RFIM the situation is different in that at small temperatures one has a *nonzero* spin-spin overlap q with the ground state: thus the existence of a ground-state percolation transition (even without an external field) implies measurable consequences even at finite temperatures. It also complicates the phase diagram by its existence.

There is one fundamental difference between two and three dimensions (besides the fact that there is no phase transition in two dimensions, and hence there systems are always paramagnetic). In two-dimensional square lattices the critical percolation site-occupation probability is 0.592 746, i.e., above one-half, and in three-dimensional cubic lattices well below one-half, 0.3116. Therefore in three dimensions, deep in the paramagnetic phase, both the spin orientations should span the system (this has been noted by Esser *et al.* to be true for the RFIM; see Ref. 20). Thus introducing an external field in paramagnetic systems leads in two dimensions to the percolation of the spin direction parallel with the external field. In three dimensions, on the other hand, the external field destroys the spanning property of the spin orientation opposite to the external field.

Consequences of the percolation type of order at the paramagnetic phase are manifold. There are experimentally accessible random field magnets, so-called diluted antiferromagnets in an external field (DAFF),²¹ in which the percolation order could be seen, should it exist for zero external fields. It is already known that the percolation of the diluted atoms has a strong contribution to the behavior of the structure factor line shapes of the 3D DAFF.²²⁻²⁴ Near the

thermodynamical phase transition point the universality class of the transition is determined by several exponents, among them the correlation length exponent (*if* the transition is continuous). The critical percolation phenomenon near the thermodynamical phase transition point may contribute there and introduce extra corrections, which have to be taken into account when the thermodynamical correlation length exponent is determined.

This paper is organized so that it starts with an introduction of the random field Ising model in the next section. Also the numerical method solving exactly the ground states is introduced. In Sec. III the percolation phenomenon is studied with a nonzero external field. The universality class of the percolation behavior is determined and the dependence of the critical external field on the random field strength is investigated. Section IV concentrates on the percolation phenomenon without an external field and compares it with the cases when an external field is applied. The properties of the spanning cluster are studied in Sec. V. Implications of the percolation to the phase diagram are discussed, together with the conclusions in Sec. VI.

II. RANDOM FIELD ISING MODEL AND NUMERICAL METHOD

The random field Ising model is defined by its energy Hamiltonian

$$\mathcal{H} = -J \sum_{\langle ij \rangle} S_i S_j - \sum_i (h_i + H) S_i, \quad (1)$$

where $J > 0$ (throughout this paper we set $J = 1$, since the relevant value is its ratio with the random field strength) is the coupling constant between nearest-neighbor spins S_i and S_j . We use here cubic lattices. H is a constant external field, which if nonzero is assigned to all of the spins, and h_i is the random field, acting on each spin S_i . We concentrate only on a Gaussian distribution for the random field values,

$$P(h_i) = \frac{1}{\sqrt{2\pi}\Delta} \exp\left[-\frac{1}{2}\left(\frac{h_i}{\Delta}\right)^2\right], \quad (2)$$

with the disorder strength given by Δ (in this paper Δ actually denotes the ratio between disorder strength and the coupling constant), the standard deviation of the distribution. The arguments presented in this paper could be extended to other lattices and other distributions, e.g., uniform and bimodal, too. However, discrete distributions, such as the bimodal one, suffer from degeneracies, and when calculating thermodynamical quantities extra averaging, over the degeneracies has to be done when using discrete distributions.^{25,26}

To find the ground-state structure of the RFIM means that the Hamiltonian (1) is minimized, in which case the positive ferromagnetic coupling constants prefer to have all the spins aligned in the same direction. On the other hand, the random field contribution is to have the spins to be parallel with the local field and thus has a paramagnetic effect. This competition of ferromagnetic and paramagnetic effects leads to a complicated energy landscape and the finding the ground state becomes a global optimization problem. An interesting

detail of the RFIM is that for $H = 0$ it has an experimental realization as a diluted antiferromagnet in a field. By gauge-transforming the Hamiltonian of DAFF,

$$\mathcal{H} = -J \sum_{\langle ij \rangle} S_i S_j \epsilon_i \epsilon_j - B \sum_i \epsilon_i S_i, \quad (3)$$

where the coupling constants $J < 0$, ϵ_i is the occupation probability of a spin S_i , and B is now a uniform external field, one gets the Hamiltonian of RFIM (1) with $H = 0$.^{27,28,21} The ferromagnetic order in the RFIM corresponds to antiferromagnetic order in the DAFF, naturally.

For the numerical calculations a graph-theoretical combinatorial optimization algorithm has been used. The Hamiltonian (1) is transformed into a random flow graph widely used in computer science with two extra sites: the source and the sink. The positive field values h_i correspond flow capacities c_{it} connected to the sink (t) from a spin S_i , similarly the negative fields with c_{is} are connected to the source (s), and the coupling constants $2J_{ij} \equiv c_{ij}$ between the spins correspond to flow capacities $c_{ij} \equiv c_{ji}$ from a site S_i to its neighboring one S_j .²⁹ In the case the external field is applied, only the local sum of fields, $H + h_i$, is added to a spin toward the direction it is positive. The algorithms—namely, maximum-flow minimum-cut algorithms—enable us to find the bottleneck, which restricts the amount of the flow which is possible to get from the source to the sink through the capacities, of such a random graph. This bottleneck, path P , which divides the system in two parts—sites connected to the sink and sites connected to the source—is the global minimum cut of the graph and the sum of the capacities belonging to the cut $\sum_P c_{ij}$ equals the maximum flow and is smaller than of any other path cutting the system. The value of the maximum flow gives the total minimum energy of the system and the minimum cut defines the ground-state structure of the system, so that all spins in the source side of the cut are the spins pointing down, and the spins in the sink side of the cut point up. The maximum-flow algorithms can be proved to give the exact minimum cut of all the random graphs, in which the capacities are positive and with a single source and sink.³⁰ We have used a sophisticated method for solving the maximum-flow–minimum-cut problem called push-and-relabel by Goldberg and Tarjan,³¹ which we have optimized for our purposes. It scales almost linearly, $\mathcal{O}(n^{1.2})$, with the number of spins and gives the ground state in about minute for 10^6 spins in a workstation.

We have used periodic boundary conditions in all of the cases. Also the percolation is tested in the periodical or cylindrical way; i.e., a cluster has to meet itself when crossing a boundary in order to span a system. Finding the spanning cluster has been done using the usual Hoshen-Kopelman algorithm.³²

III. PERCOLATION WITH AN EXTERNAL FIELD

As a start of the percolation studies of the 3D RFIM we draw in Fig. 1(a) the spanning probabilities of down spins Π_{\downarrow} with respect to the uniform external field H pointing up for several system sizes L and for a fixed random field

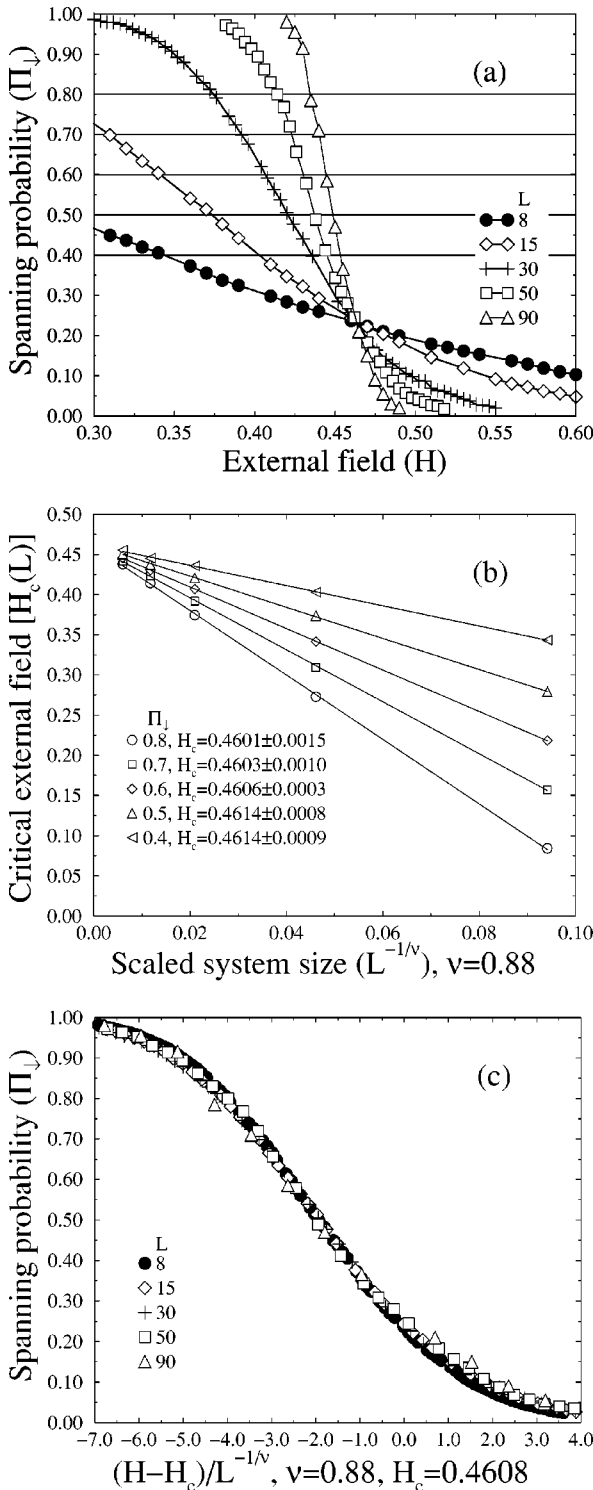


FIG. 1. (a) The spanning probabilities of minority down spins Π_{\downarrow} as a function of upward external field H for $\Delta = 3.5$ with $L^3 \in [8^3 - 90^3]$. The number of realizations varies between 5000 realizations for $L = 8$ and 200 for $L = 90$. (b) The finite-size scaling of the fields $H_c(L)$, which are from the crossing points of the spanning probability curves with the horizontal lines in (a), leading to the estimate of the critical $H_c = 0.461 \pm 0.001$ using $L^{-1/\nu}$, $\nu = 0.88$. The error bars in the labels of the figure for different H_c are the errors of the least-squares fits. (c) The data collapse of different system sizes with the corresponding critical $H_c = 0.4608$.

strength $\Delta = 3.5$. The curves look rather similar to standard percolation, except that in site percolation the systems span at the high occupation probability limit, and here the down spins do not span, when the external positive field has a large value, and thus the step in the spanning probability is inverse compared to the one in the occupation percolation. It is interesting to note, also, that since we are using periodic boundary conditions in all of the directions, also for spanning, the $\Pi_{\downarrow}(L)$ lines for various system sizes cross at rather low Π_{\downarrow} values. This is the case for the other Δ , too. Similar boundary-condition-dependent behavior has been seen in the standard percolation, too.³³⁻³⁵ When we take the crossing points $H_c(L)$ of the spanning probability curves with fixed spanning probability values $\Pi_{\downarrow} = 0.4, 0.5, 0.6, 0.7,$ and 0.8 , for each systems size L , we get an estimate for the critical external field H_c using finite-size scaling; see Fig. 1(b). There we have attempted with success to find the value for H_c using the standard short-range-correlated 3D percolation correlation length exponent $\nu = 0.88$.¹⁴ Using the estimated $H_c = 0.461 \pm 0.001$ for $\Delta = 3.5$ we show a data collapse of Π_{\downarrow} versus $(H - H_c)/L^{-1/\nu}$ in Fig. 1(c), which confirms the estimates of H_c and $\nu = 0.88$. We get similar data collapses for various other random field strength values Δ as well.

Considering the percolation and critical external field with respect to the random field strength, there is an obvious constraint on the phase diagram H vs. Δ . Below the phase transition critical point, $\Delta_c \approx 2.27$,^{5,7,8} only one of the spin orientations may span a system, since in a ferromagnetic system the magnetization has a finite, positive or negative, value and thus there cannot be a massive percolation cluster of the opposite spin direction. Since the earlier studies of the phase transition at the 3D (Refs. 6-9,11, and 12) have shown that the order parameter exponent β has a value close to zero, if not zero, the transition is sharp and therefore the simultaneous percolation of both (up and down) spin directions should vanish or have vanished at Δ_c when approaching from above. The question now remains whether this takes place exactly at the phase transition point, so that the critical points would coincide, or for a $\Delta_p > \Delta_c$. In the latter case it is also of interest what happens for $H = 0$ between the critical points, on the line $\Delta_c < \Delta < \Delta_p$. We now propose a phase diagram, Fig. 2, for the percolation phenomenon, and ask at which value do the dashed lines in the diagram meet. Above we showed that in the direction of the vertical arrow at $H > 0$ the universal standard percolation correlation length exponent is valid. What about at the vertical arrow, what are the critical exponents there?

To answer the question how the percolation critical external field H_c behaves with respect to the random field strength Δ , we have attempted a critical type of scaling using the calculated $H_c(\Delta)$ for various $\Delta = 2.5, 2.6, 2.75, 3.0, 3.25, 3.5, 4.0,$ and 4.5 . We have been able to use the ansatz

$$H_c \sim (\Delta - \Delta_p)^\delta, \tag{4}$$

where $\delta = 1.31 \pm 0.03$ by assuming $\Delta_p = 2.43$; see Fig. 3(a). In Fig. 3(b), on the other hand, we have plotted the calculated Δ values versus the scaled critical external field $[H_c(\Delta)]^{1/1.31}$ and it gives an estimate for $\Delta_p = 2.43 \pm 0.01$.

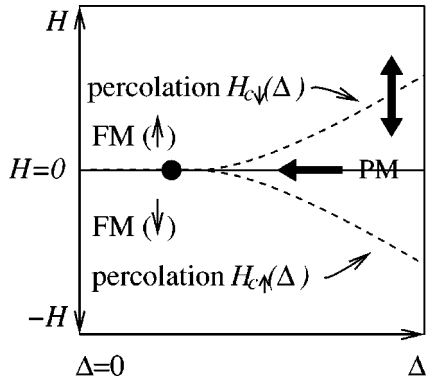


FIG. 2. The phase diagram for the minority-spin percolation of the 3D RFIM with disorder strength Δ and an applied external field H . The dashed lines define the percolation thresholds $H_c(\Delta)$ for up and down spins to lose their spanning property, below and above which the minority spins do not percolate anymore. The phase transition point for the ferro- and paramagnetic phases at $\Delta=2.27$, $H=0$ is shown as a circle.

This indicates that the percolation probability lines for up and down spins to lose their spanning property meet at $\Delta_p = 2.43 \pm 0.01$. Note that our studies in the two-dimensional RFIM gave the values $\Delta_p = 1.65 \pm 0.05$ and $\delta = 2.05 \pm 0.10$ for systems *to span*, not *to lose*, the spanning property as here.¹⁵ We also tested various exponential scaling assumptions for the $H_c(\Delta)$ scaling, but none of them worked. However, here we know that H_c has to vanish at some finite Δ_p value, which is greater than or equal to Δ_c .

We have also calculated the order parameter of the percolation, the probability that a down spin belongs to the down-spin spanning cluster P_∞ . Using the scaling for the correlation length

$$\xi_{perc} \sim |H - H_c|^{-\nu} \quad (5)$$

and for the order parameter, when $L < \xi_{perc}$,

$$P_\infty(H) \sim (H_c - H)^\beta, \quad (6)$$

we get the limiting behaviors

$$P_\infty(H, L) \sim \begin{cases} (H_c - H)^\beta, & L < \xi_{perc}, \\ L^{-\beta/\nu}, & L > \xi_{perc}, \end{cases} \quad (7)$$

and thus the scaling behavior for the order parameter becomes

$$P_\infty(H, L) \sim L^{-\beta/\nu} F\left[\frac{(H_c - H)^{-\nu}}{L}\right] \sim L^{-\beta/\nu} f\left(\frac{H_c - H}{L^{-1/\nu}}\right). \quad (8)$$

Note that here and later in this article β denotes the percolation order parameter exponent as opposed to the bulk phase transition order parameter exponent discussed earlier in this paper. We have done a successful data collapse, i.e., plotted the scaling function f , for various Δ using the standard 3D short-range-correlated percolation exponents $\beta=0.41$ and $\nu=0.88$, of which the case $\Delta=4.5$ with $H_c=1.0441$ is shown in Fig. 4. Note that only the left part (below zero) of the

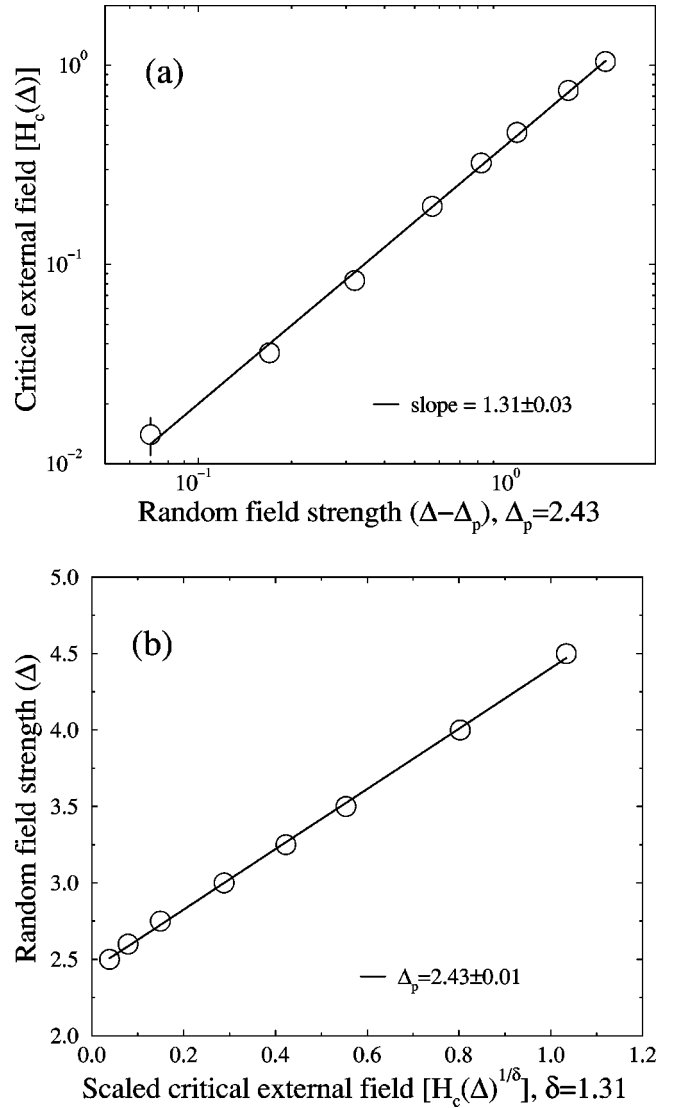


FIG. 3. (a) For each calculated Δ the critical positive $H_c(\Delta)$ for down-spin spanning versus $\Delta - \Delta_p$, where Δ_p is estimated to be 2.43. The power-law behavior suggests a scaling $H_c \sim (\Delta - \Delta_p)^\delta$, where $\delta = 1.31 \pm 0.03$. The error bar for δ is the error of the least-squares fit. (b) The same data but plotted as each Δ vs $[H_c(\Delta)]^{1/\delta}$, where $\delta = 1.31$, which estimates that, at $\Delta_p = 2.43 \pm 0.01$, $H_c = 0$. Again the error bar for Δ_p is the error of the least-squares fit. The other details are as in Fig. 1.

scaling function is shown, since $P_\infty(H, L)$ is limited between $[0, 1]$. When one divides it by $L^{-\beta/\nu}$ the part where nonscaled $P_\infty(H, L)$ had a value of unity the scaled $P_\infty(H, L)/L^{-\beta/\nu}$ saturates at different values depending on L . One can easily see that the smallest system size $L^3 = 8^3$ does not scale (the rest are scattered around each other and do not have any trend). We believe that this is due to an intrinsic length scale over which the spins are correlated and which depends on the random field strength value. This will be discussed in more detail in Sec. V, when the scaling of the spanning cluster is studied.

Hence, we conclude that the percolation transition for a fixed Δ versus the external field H is in the standard 3D short-range-correlated percolation universality class.¹⁴ This

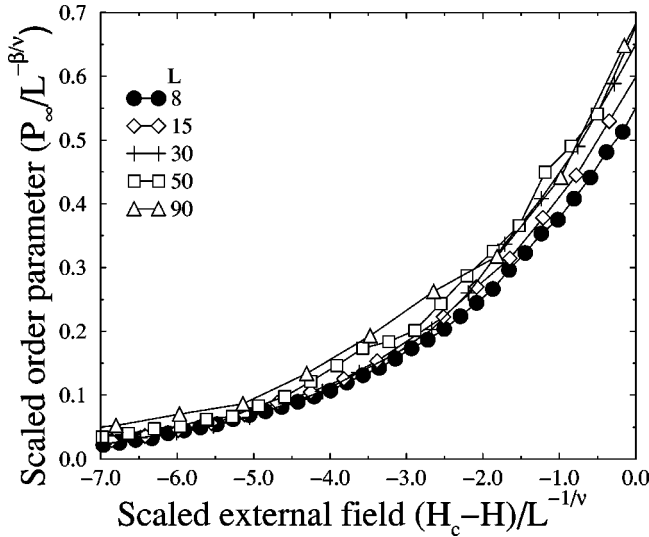


FIG. 4. The scaled order parameter probability that a down spin belongs to the down-spin spanning cluster, $P_\infty/L^{-\beta\nu}$, $\beta=0.41$, $\nu=0.88$ vs the scaled external field $(H_c - H)/L^{-1/\nu}$, for $\Delta=4.5$ with $L^3 \in [20^3 - 90^3]$. The data points are disorder averages over 200–5000 realizations. The corresponding critical $H_c(\Delta=4.5) = 1.0441$.

is confirmed by the fractal dimension of the spanning cluster, too, as discussed below. The fact that the critical behavior of the percolation with respect to the external field belongs to the standard short-range-correlated percolation universality class is not surprising, since the strong-disorder limit can be seen to be related with the site percolation problem and that, e.g., the positive external field decreases the number of the occupied down spins. Also other exponents could be measured, such as γ for the average size $\langle s \rangle$ of the clusters and σ and τ for the cluster size distribution as well as the fractal dimension of the backbone of the spanning cluster, the fractal dimension of the chemical distance, the hull exponent, etc.

IV. PERCOLATION AT $H=0$

In the previous section we learned that the dashed lines in the phase diagram, Fig. 2, meet at the value $\Delta_p = 2.43 \pm 0.01$, which is well above the phase transition critical point $\Delta_c = 2.27$. This raises the question how this is seen when the external field $H=0$ and what happens between Δ_c and Δ_p . Thus we study the phase diagram in the direction of the horizontal arrow in Fig. 2. There are two strategies for this that we employ separately to evaluate their advantages and disadvantages. That is, one can take the Δ_p to be *a priori* the same for all $\Pi_{\uparrow\downarrow}$, the probability for simultaneous spanning of up *and* down spins. Or then this can be let to vary with $\Pi_{\uparrow\downarrow}$, as in two dimensions.¹⁵

In Fig. 5(a) we have plotted the probability for simultaneous spanning of up *and* down spins $\Pi_{\uparrow\downarrow}$ as a function of the Gaussian random field strength Δ for various system sizes $L^3 = 8^3, 15^3, 30^3, 50^3, 90^3$, and 120^3 . This case now resembles the standard occupation percolation in the sense that the step in the percolation probability is from a low

value to a large value when Δ is increased. By estimating that the $\Delta_{p,H=0}$ at the thermodynamic limit has a value of 2.32 using fixed $\Pi_{\uparrow\downarrow} = 0.2, 0.4, 0.6$, and 0.8 for the $\Delta_{p,H=0}(L)$ we find that the effective ν gets a value of 0.97 ± 0.05 when approaching the critical point in this direction; see Fig. 5(b). On the other hand, assuming that $\nu=1.0$, the $\Delta_{p,H=0}$ becomes 2.32 ± 0.01 ; see Fig. 5(c). These plots show that the estimates should be correct. However, the data collapse, Fig. 5(d), using the estimates above could be better. Obviously the smallest system size $L^3 = 8^3$ does not scale.

There are a couple of points one should note from the scaling. First, the estimate for the $\Delta_{p,H=0} = 2.32 \pm 0.01$ is still above the phase transition point $\Delta_c = 2.27$. Another point is that $\Delta_{p,H=0}$ is reasonably far away from $\Delta_p = 2.43 \pm 0.01$ [note that the error bar in the finite-field case is the error bar of the least-squares fit in Fig. 3(b) and does not take into account other sources for the error, e.g., the error of δ , statistics, etc., and thus is a lower limit]. The third point is that $\Pi_{\uparrow\downarrow} = 0.0$ at $\Delta_{p,H=0}$ and $\Pi_{\downarrow} = 0.25$ at $H_c(\Delta)$ [for $\Delta = 3.5$; see Fig. 1(c)]. Our take on the two different estimates is that they are compatible with the following scenario. For Δ values that are slightly below 2.43 one can have *only one critical spanning cluster*, and the probability for this is then Π_{\downarrow} , about 0.25. Both orientations do span simultaneously, as they can do for all Δ values above Δ_c , but they should not be both critical, unless one decreases the disorder strength further.

For the estimate of the correlation length exponent, deviations from normal percolation are seen since $\nu = 1.00 \pm 0.05$ instead of $\nu = 0.88$. In our opinion this reflects the fact that for $H \neq 0$ the correlations from the proximity of Δ_c are negligible, whereas here the spin-spin correlations change with system size. The correlation length exponent is higher than that for percolation, so clear-cut percolation scaling cannot be expected. Differences between the $H=0$ and $H \neq 0$ cases were found also in the two-dimensional case.¹⁵ Note that in two dimensions, the exponent was *dependent on the spanning probability* and the standard correlation length exponent was found where the spanning probability for either of the spin directions to span $\Pi_{\uparrow/\downarrow}$ had a nonzero value (remember that in two-dimensional square lattices without an external field at large Δ neither of the spin directions span and with small Δ either of them start to span).

Here we tried, as in two dimensions, to do fits using several criteria for $\Pi_{\uparrow\downarrow} = [0.05, 0.15, 0.20, \dots, 0.95]$ and letting both Δ_p and ν vary depending on $\Pi_{\uparrow\downarrow}$. Indeed, we obtained monotonous behaviors depending on $\Pi_{\uparrow\downarrow}$ for both ν and Δ_p . However, this may just reflect how finite-size effects depend on the criterion. It is anyhow worth noting that for $\Pi_{\uparrow\downarrow}$ approaching zero, Δ_p gets also closer and closer to 2.27, i.e., the accepted value for the phase transition point Δ_c . Moreover the correlation length exponent moves towards $\nu = 1.3 \pm 0.1$, in the neighborhood of the phase transition correlation length exponents reported in the literature.^{5,8,11} Similarly, if $\Pi_{\uparrow\downarrow}$ is let to approach unity, Δ_p closes on the value $\Delta_p = 2.43$ obtained above, in the finite-field case. This behavior may be just coincidence or related to the (Δ -dependent) correlations in the system to how they

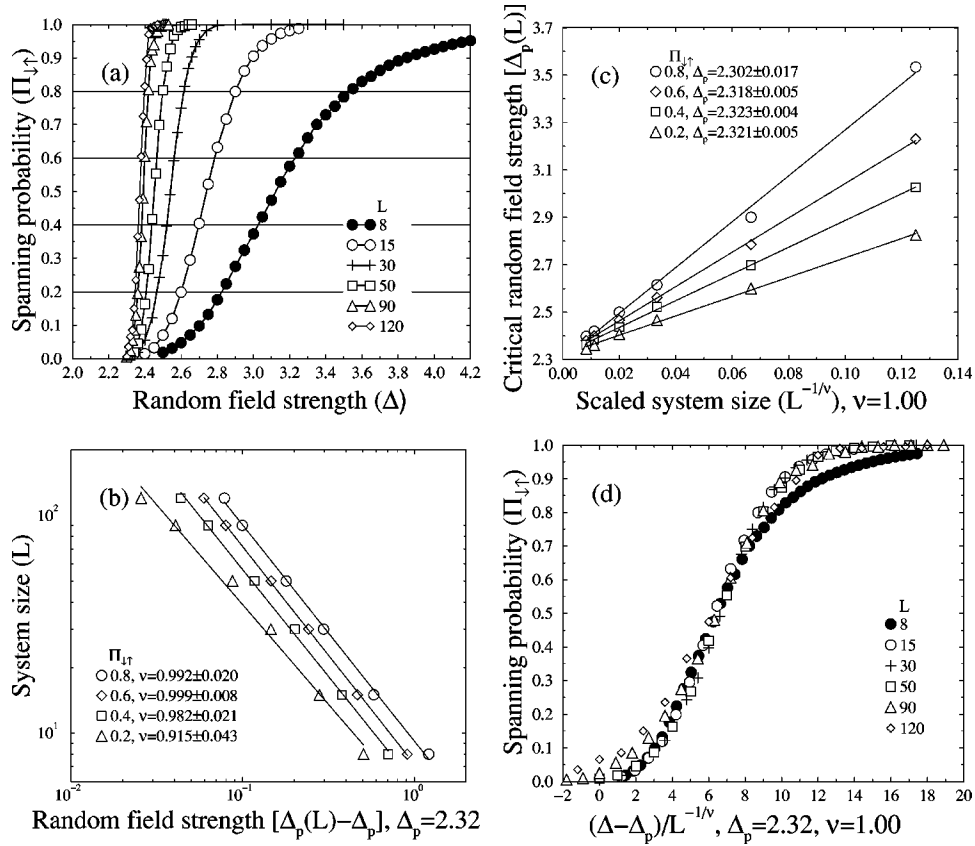


FIG. 5. (a) The spanning probabilities for system sizes $L^3 \in [8^3 - 120^3]$ of simultaneous up- and down-spin spanning $\Pi_{\uparrow\downarrow}$ as a function of Δ for $H=0$. The data points are disorder averages over 200–5000 realizations. (b) Each system size L vs $\Delta_p(L) - \Delta_p$, where $\Delta_p(L)$'s are the corresponding crossing points of the spanning probability curves with the horizontal lines of $\Pi_{\uparrow\downarrow} = 0.2, 0.4, 0.6,$ and 0.8 in (a) and Δ_p is estimated to be 2.32. The power-law behavior suggests a scaling $L \sim (\Delta - \Delta_p)^{-\nu}$, where $\nu = 0.97 \pm 0.05$. The error bars in the labels of the figure for different ν 's are the errors of the least-squares fits. (c) The same data as in (b), but now plotted as random field strength values $\Delta_p(L)$ vs the scaled system size $L^{-1/\nu}$, where $\nu = 1.0$, leading to a same estimate of $\Delta_p = 2.32 \pm 0.01$. The error bars in the labels of the figure for different Δ_p are the errors of the least-squares fits. (d) The data collapse of (a) with the corresponding critical $\Delta_p = 2.32$ and $\nu = 1.0$.

change the universality class of percolation in the vicinity of the phase transition. We return to this in the conclusions in Sec. VI.

Hence we have shown that at large Δ both spin directions span simultaneously, and by decreasing random field strength we find a critical $\Delta_{p,H=0}$, which is above the phase transition point Δ_c and below which there is no simultaneous spanning. Therefore we conclude that in the whole regime $\Delta_c < \Delta \leq \Delta_{p,H=0}$ there is geometrical criticality in 3D RFIM, since always only either of the spin directions spans the system. However, the spanning cluster cannot be massive there, i.e., scale with Euclidean dimension ($d=3$), the system still being paramagnetic, but has to be a fractal. The scaling of the spanning cluster is studied in the next section and the implications of the critical region in Sec. VI.

V. SPANNING CLUSTER

In Fig. 6(a) we have plotted the mass of the spanning cluster of down spins with respect to the system size at $H_c(\Delta) > 0$ for four random field strength values $\Delta = 2.75, 3.0, 3.5,$ and 4.0 up to system size $L^3 = 120^3$. As a guide to

the eye the fractal dimension $D_f = 2.53$ of the standard percolation is drawn in the figure and the systems can be seen asymptotically approaching the same scaling. However, there are obvious finite-size effects, which depend on Δ . We have estimated roughly the crossover system sizes for the systems to reach the correct scaling, $L_x \approx 30, 20, 10,$ and 5 for $\Delta = 2.75, 3.0, 3.5,$ and 4.0 , respectively. This hints at an exponential scaling with a slope of -1.42 ± 0.03 for the crossover length scale; see solid diamonds in the inset of Fig. 6(a). The above scaling predicts, for $\Delta = 4.5$, $L_x \approx 3$, smaller than $L = 8$ (in Fig. 4, this size does not scale), but note that the prefactors of the scaling behaviors need not to be the same. In Fig. 6(b) we have drawn for three $\Delta \leq \Delta_p$, i.e., $\Delta = 2.35, 2.38,$ and 2.45 (which is so close to Δ_p that its H_c is practically zero with respect to the numerical precision, 10^{-3}), at $H=0$, the scaling of the mass of the spanning cluster of either of the spin orientations up to system size $L^3 = 120^3$. There one can see that the fractal dimension $D_f = 2.53$ of the standard percolation is asymptotically met, too, but at much larger system sizes. Here we have estimated the crossover system sizes $L_x \approx 80, 60,$ and 50 , for $\Delta = 2.35, 2.38,$ and 2.45 , respectively. They are plotted as open circles

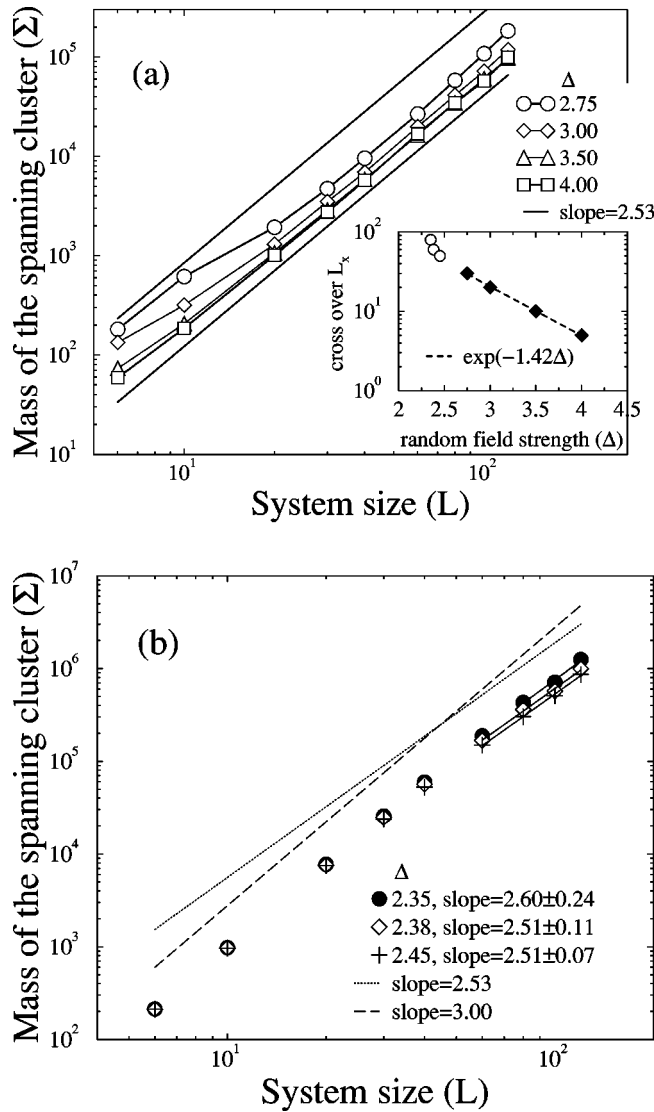


FIG. 6. (a) The average mass of spanning cluster of down spins for random field strength values $\Delta = 2.75, 3.0, 3.5,$ and 4.0 at the critical positive external field value $H_c(\Delta)$ [see the values from Fig. 3(b)]. The 3D percolation fractal dimension $D_f = 2.53$ is indicated with solid lines. In the inset a crossover length scale at which system size the asymptotic behavior is met for each random field strength is plotted as solid diamonds. The least-squares fit estimates an exponential behavior with a slope of -1.42 ± 0.03 . (b) The average mass of spanning clusters of either spin orientation for random field strength values $\Delta = 2.35, 2.38,$ and 2.45 when $H = 0$. The solid lines are the least-squares fits to the data with the slopes indicated in the labels. The dotted line with a slope of $D_f = 2.53$ and the dashed line with a slope of $d = 3$ are guides to the eye. The estimated crossover length scales are plotted in the inset of Fig. 6(a) as open circles.

in the inset of Fig. 6(a) and are obviously diverging from the exponential behavior mentioned above when approaching phase transition Δ_c . These large values for L_x do not leave much room for the asymptotic scaling, since it is difficult to go above $L^3 = 120^3$. However, the crossover is visible. There is one other thing one notices from Figs. 6(a) and 6(b). In the case we plot the mass of the spanning cluster of the down

spins in $\Delta > \Delta_p$ and $H_c(\Delta) > 0$ the crossover is from a smaller slope to the asymptotic $D_f = 2.53$ one. In the case $\Delta \leq \Delta_p$ the crossover is from the Euclidian dimension (slope of 3, i.e., effective ferromagnetism) to the asymptotic $D_f = 2.53$. There it is obviously affected by the vicinity of the phase transition point.

VI. DISCUSSION AND CONCLUSIONS

In this paper we have studied the character of the ground state of the three-dimensional random field Ising magnet in, mostly, the paramagnetic phase. A geometrical critical phenomenon exists in these systems: for cubic lattices in ordinary percolation both occupied and unoccupied sites span the systems when the occupation probability is one-half. In the RFIM this corresponds to the case with a high random field strength value, without an external field. When an external field is applied and the random field strength decreased, a percolation transition, for the other spin orientation to lose the spanning property, can be seen. The transition is shown to be in the standard 3D short-range-correlated percolation universality class when studied as a function of the external field. Hence, the correlations in the three-dimensional random field Ising magnets are only of finite extent as could be expected in this region of the bulk phase diagram. Based on our numerical results both critical points $\pm H_c(\Delta)$ approach when Δ is decreased and finally meet at a $\Delta_p \approx 2.43 > \Delta_c$, at which $H_c = 0$. When the percolation transition is studied without an external field and tuning the random field strength similar behavior is found, i.e., signatures of a percolation line (a $\Delta_p > \Delta_c$). This might cause puzzling consequences when studying the character of the ground states, because the percolation correlations may influence the magnetization correlation length.

The major theoretical implications have to do with the phase transition. Note that earlier ground-state studies of the domain structure implied that there is only a “one-domain state” below the critical field and a “two-domain state” in the paramagnetic phase (extending down from high disorder values).²⁰ If the transition is first order, then one expects the percolation properties of the paramagnetic phase to be discontinuous in the thermodynamic limit. If the transition is second order, then one may ask what is the correct way to link the presence of the percolation transition to the critical phase. At Δ_c , one expects that the spin-spin correlations show power-law correlations. For a normal percolation transition, these are (as in the disordered phase in general) of short-range character. There is a divergent length scale as the transition is approached from the paramagnetic phase, below which the spin-spin correlations matter and the scaling of the spanning clusters is volume like.

Assume that the properties of the largest cluster are governed by the power-law correlations. An old result by Weinrib gives a Harris criterion for this approach to check how this would change its structure from ordinary percolation.³⁶ If the site occupation probability correlations decay as r^{-a} , one has that the decay is relevant if $a\nu_{old} - 2 < 0 \rightarrow \nu_{new} = 2/a$, where now $\nu_{old} = 0.88$ for 3D site percolation. One gets a critical decay exponent $a_c = 2.27$, much larger than

that found by Middleton and Fisher,⁵ which is very close to zero. An application of the theory of correlated percolation would thus imply that the spin-spin correlations at Δ_c are relevant for percolation. They would change the universality class, of percolation, in a way that would reflect such correlations. This reasoning would need further consideration.

One should note also that although this study was done using cubic lattices, it can be extended to other lattices, too, since all common three-dimensional lattices have $p_c < 0.5$. Thus the transition from both spin orientations spanning phase to only one spin orientation spanning phase should exist. In the case of diluted antiferromagnets the percolation is already seen as percolation of diluted spins. The implica-

tion of this paper is that the influence of percolation is even more rich. Lately there has been interest in studying domain walls and excitations in RFIM. In both cases the underlying percolation criticality should affect the structure of the clusters that result from varying the boundary conditions.

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