

Lattice contraction in carbon-doped GaAs epilayers

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Epitaxial GaAs layers, heavily doped with carbon, have been grown by gas-source molecular-beam epitaxy using an electron-beam-evaporating graphite source for doping. Lattice strain and electrical compensation in these samples have been studied theoretically and experimentally. Most of the carbon acceptors can be activated by post-growth annealing at 490 °C in the N₂ ambient, probably due to the reduction of C-H bonds formed during growth. Annealing at higher temperatures reduces the electrical activity and causes an abrupt change in lattice strain, likely due to the formation of C-C interstitial couples. [S0163-1829(98)04920-0]

Carbon is an attractive alternative to beryllium as a *p*-type dopant for GaAs and related compounds because of its lower diffusion coefficient and possibility of obtaining a higher free-carrier concentration.¹ Very high doping levels in GaAs have, indeed, been achieved.²⁻⁵ On the other hand, self-compensation effects begin to play an important role at carbon concentrations greater than about $5 \times 10^{19} \text{ cm}^{-3}$.⁶ The compensation effects are associated with the generation of donors, mainly due to the formation of interstitial carbon, substitutional carbon in gallium sites, and the precipitation of carbon.^{7,8} If hydrogen is present in the growth ambient, as in the case of gas-source molecular-beam epitaxy, an additional compensation effect appears, namely, hydrogen passivation that is caused by the formation of C-H bonds.⁹

Besides electrical compensation, C doping in excess of 10^{19} cm^{-3} causes lattice contraction. This is because C has a small covalent radius (0.77 Å) compared to As (1.20 Å) or Ga (1.26 Å).¹⁰ Lattice contraction can be determined by a high-precision measurement, such as double-crystal x-ray diffractometry (DCXRD), and be related to carbon concentration in a rather straightforward way.¹¹ The situation is somewhat more complicated if the C-H bonds are present.

In this work we shall study lattice strain and electrical properties of heavily doped GaAs:C grown by gas-source molecular-beam epitaxy (MBE) (for gas-source MBE, see Ref. 12). We first describe the experiments and then discuss the results obtained.

The GaAs:C samples were grown homoepitaxially at 580 °C. The background pressure was 10^{-5} mbar during growth, due to hydrogen from cracked AsH₃. Carbon was evaporated from a graphite rod of an electron-beam-heated source. The C flux was controlled by the filament voltage and emission current. Upon growth, the samples were thermally annealed (RTA) *ex situ* in 100% N₂. The annealing procedure consisted of an initial N₂ purge, a temperature ramp at 50 °C/s to a final temperature, and a 5 min anneal at this temperature.

The carrier density and van der Pauw Hall mobility were measured at room temperature. The DCXRD rocking curves were measured in a nondispersive Bragg geometry to determine the lattice strain and carbon concentration.

Carbon incorporates substitutionally into the arsenic sublattice where it becomes an acceptor (C_{As}⁻). At high doping levels, greater than 10^{19} cm^{-3} , it may also occupy gallium

sites (C_{Ga}), acting as a donor. The density of C_{Ga}⁺ is, however, four orders of magnitude smaller than the density of C_{As}⁻ because of a higher formation energy of C_{Ga}⁺.¹³ Carbon may also occupy an interstitial site and reduce lattice strain. Interstitial carbon is unstable. It readily captures a neighboring C_{As}⁻ acceptor to form a C-C interstitial couple that acts as a donor. The formation of C-C bonds can become a dominant compensation effect in GaAs:C at a very high doping concentration.

One may calculate the lattice mismatch $\Delta a/a_0$ caused by the carbon-related defects. If the calculation is based on Pauling's covalent bonding radii of atoms and Vegard's law,¹⁴ we obtain the amount of lattice contraction due to the presence of C_{As},

$$\begin{aligned} \Delta a/a_0 &= 4(r_C - r_{As})N_{C,As}/a_0N_0\sqrt{3} \\ &= -7.9478 \times 10^{-18} N_{C,As} \text{ (ppm)}, \end{aligned}$$

where r_C and r_{As} are the covalent radii of C and As, respectively, $N_0 = 2.21 \times 10^{22} \text{ cm}^{-3}$ is the density of Ga or As atoms in pure GaAs, and $N_{C,As}$ is the concentration of C_{As}⁻. Similarly, lattice strain caused by the C-C interstitial couple is¹⁵

$$\begin{aligned} \Delta a/a_0 &= (d_b - r_{Ga} - r_{As})(c_{11} + 2c_{12})N_{C-C}/2N_0c_{11}(r_{Ga} + r_{As}) \\ &= -1.4956 \times 10^{-18} N_{C-C} \text{ (ppm)}, \end{aligned}$$

where $d_b = (\sqrt{3}/3)r_C + \sqrt{[(r_C + r_{Ga})^2 - 2r_C^2/3]}$ is the distance of the C-C couple from its nearest neighbors; c_{11} and c_{12} are the stiffness coefficients of GaAs, and N_{C-C} is the concentration of C-C couples. Our calculations show that interstitial carbon has no contribution to the lattice strain.

The hydrogen passivation effect should be taken into account. Hydrogen is incorporated into interstitial sites forming C-H bonds with C_{As}⁻. Lattice mismatch generated by the C-H bond is

$$\Delta a/a_0 = -6.9432 \times 10^{-18} N_{H-C} \text{ (ppm)},$$

where the effective bond radius is taken to be $(r_C + r_H)/2 = 0.545 \text{ Å}$, and N_{H-C} is the concentration of C-H bonds.

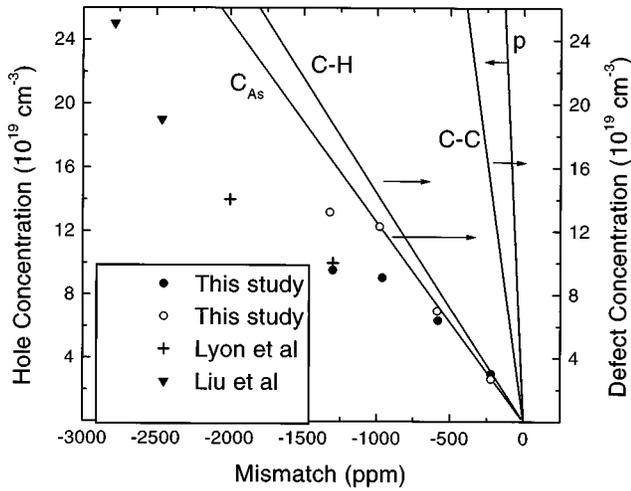


FIG. 1. Plot of the calculated lattice mismatch as a function of C-related defects. The relationship between hole concentration and lattice mismatch is determined from Hall effect and DCXRD measurements for as-grown (●) and RTA treated (○) GaAs:C. Data of Liu *et al.* (Ref. 19) and DeLyon *et al.* (Ref. 7) are included for comparison.

In consideration of a hole in the valence band decreasing the energy of covalent bond,¹⁶ high concentration of holes could also lead to the lattice contraction of GaAs:C, which can be described by

$$\Delta a/a_0 = (D_p/3B)p = -0.4978 \times 10^{-18} p \text{ (ppm)},$$

where p is the hole concentration in the top of the valence band, $D_p = -0.7$ eV is the deformation potential of this maximum,¹⁷ and $B = 7.5 \times 10^6$ N/cm² is the bulk modulus.¹⁸ When all the factors are taken together, one may write

$$\Delta a/a_0 = -(7.9478N_{C,As} + 6.9432N_{C,H} + 1.4956N_{C-C} + 0.4978 p) \times 10^{-18} \text{ (ppm)}.$$

The net hole concentration is $p = N_{C,As} - N_{H-C} - N_{C-C}$ and the total carbon concentration $N_{C,tot} = N_{C,As} + N_{H-C} + 2N_{C-C}$.

$N_{C,H}$ can be estimated from the postgrowth annealing experiments. When $N_{C,H}$, $N_{C,As}$, and N_{C-C} are known, one can provide the relationships between $\Delta a/a_0$ of GaAs:C, hole concentration, and carbon-related defect concentration, based on Hall and DCXRD measurements.

Figure 1 shows our calculations together with experimental results. It can be seen that at doping levels lower than 1.0×10^{20} cm⁻³ the carrier compensation is mainly due to hydrogen passivation of C acceptors. At a higher doping level C-C interstitial couples could be formed to further reduce the doping efficiency.

Effects of postgrowth annealing on the electrical properties and $\Delta a/a_0$ of GaAs:C are studied in Figs. 2 and 3. The samples studied in Fig. 2 exhibit the Hall hole densities around 7.0×10^{19} cm⁻³, as measured after growth. Relative changes in hole concentration and Hall mobility are determined as functions of postgrowth annealing temperature (T_{anneal}). The figure shows that there is a significant ($\sim 10\%$) improvement in electrical activation when the sample is annealed at 450–530 °C for 5 min. Figure 3 shows that annealing only causes a slight change in strain level when

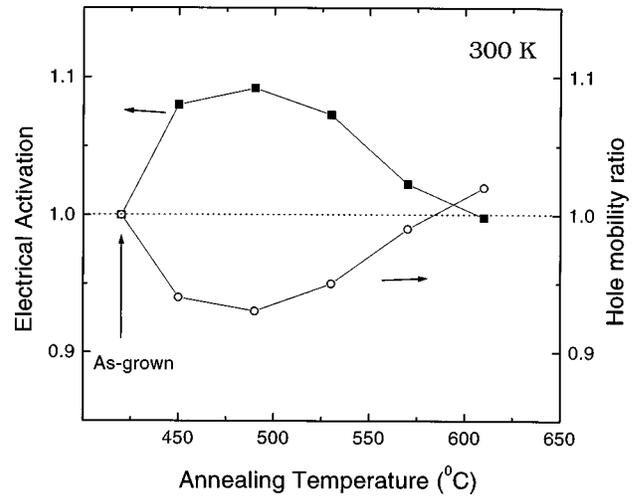


FIG. 2. Electrical activation and hole mobility for postannealed GaAs:C as functions of annealing temperature. The dotted line is a guide to the eye.

$T_{anneal} \leq 530$ °C, but yields a more remarkable, abrupt change of about 10% with the opposite sign of strain if $T_{anneal} > 560$ °C. Similar results were also reported in Ref. 20.

We relate the observed improvement in electrical activity to a postgrowth annealing effect that is thought to remove hydrogen passivated carbon acceptors, the C-H bonds. The optimal annealing temperature is $T_{anneal} \sim 490$ °C. Higher temperatures are not desired because they cause carrier compensation, likely due to the formation of C-C interstitial couples that result in a decrease in the lattice contraction with an increase in mobility. According to the calculations, the relative change of 7.4% in strain observed for $T_{anneal} > 560$ °C corresponds to an increase of 2.3×10^{18} cm⁻³ in the density of C-C couples.

In summary, heavily doped GaAs:C layers have been grown by gas-source molecular-beam epitaxy using an electron-beam-evaporating graphite source for doping. Lattice mismatch, due to C at arsenic sites, C-H bonds, C-C interstitial couples, and free charge, has been calculated

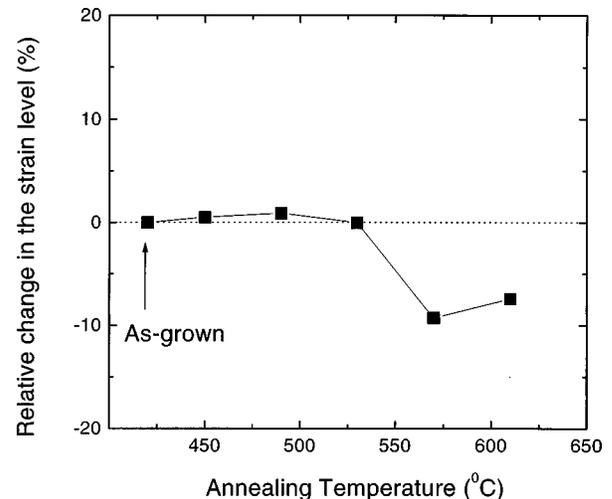


FIG. 3. Relative change in strain for post-annealed GaAs:C as a function of annealing temperature. The dotted line is a guide to the eye.

and compared with that experimentally obtained by x-ray diffraction. The largest mismatch is caused by C_{As} . The C-H bonds are believed to passivate carbon acceptors and cause carrier compensation in as-grown samples. They can be removed to a large extent by postgrowth annealing at 490 °C resulting in an increase of several percent in hole concentra-

tion. Annealing at higher temperatures decreases the hole concentration, likely due to an increase in C-C interstitial couples.

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¹T. H. Chiu, J. E. Cunningham, J. A. Ditzenberger, W. Y. Jan, and S. N. G. Chu, *J. Cryst. Growth* **111**, 274 (1991).

²L. W. Yang, P. D. Wright, V. Eu, Z. H. Lu, and A. Majerfeld, *J. Appl. Phys.* **72**, 2063 (1992).

³R. Pena-Sierra, A. Escobasa, and V. M. Sanchez-R., *Appl. Phys. Lett.* **62**, 2359 (1993).

⁴J. Shirakashi, T. Yamada, M. Qi, S. Nozaki, K. Takahashi, and E. Tokumitsu, *Jpn. J. Appl. Phys.* **30**, L1609 (1991).

⁵T. H. Chiu, J. E. Cunningham, J. A. Ditzenberger, and W. Y. Jan, *Appl. Phys. Lett.* **57**, 171 (1990).

⁶K. Watanabe and H. Yamazaki, *Appl. Phys. Lett.* **59**, 434 (1991).

⁷T. J. DeLyon, J. M. Woodall, M. S. Goorsky, and P. D. Kirchner, *Appl. Phys. Lett.* **56**, 1040 (1990).

⁸G. E. Höfler and K. C. Hsieh, *Appl. Phys. Lett.* **61**, 327 (1992).

⁹R. Iga, H. Sugiura, T. Yamada, and K. Wada, *Appl. Phys. Lett.* **55**, 451 (1989).

¹⁰C. Giannini, A. Fischer, C. Lang, K. Ploog, and L. Tapfer, *Appl. Phys. Lett.* **61**, 183 (1992).

¹¹P. Enquist, *J. Appl. Phys.* **71**, 704 (1992).

¹²M. Pessa, K. Tappura, and A. Ovtchinnikov, *Thin Solid Films* **99**, 265 (1995).

¹³B. H. Cheong and K. J. Chang, *Phys. Rev. B* **49**, 17 436 (1994).

¹⁴L. Pauling, *The Nature of the Chemical Bond and the Structure of Molecules and Crystals*, 3rd ed. (Cornell University Press, Ithaca, NY, 1960), Chap. 7.

¹⁵N. F. Chen, Y. T. Wang, H. J. He, and L. Y. Lin, *Phys. Rev. B* **54**, 8516 (1996).

¹⁶T. Figielski, *Phys. Status Solidi* **1**, 306 (1961).

¹⁷D. D. Nolte, W. Walukiewicz, and E. E. Haller, *Phys. Rev. Lett.* **59**, 501 (1987).

¹⁸M. Leszczynski, S. Podlasin, and T. Suski, *J. Appl. Crystallogr.* **26**, 1 (1993).

¹⁹Q. Liu, A. Brennemann, H. Hardtdegen, A. Lindner, W. Prost, and F.-J. Tegude, *J. Appl. Phys.* **79**, 710 (1996).

²⁰W. E. Hoke, P. J. Lemonias, D. G. Weir, H. T. Hendriks, and G. S. Jackson, *J. Appl. Phys.* **69**, 511 (1991).