

Giant moment of Fe_{16}N_2 as evidenced by ^{57}Fe NMR studies

Y. D. Zhang, J. I. Budnick, and W. A. Hines

Department of Physics, Connecticut Advanced Technology Center for Precision Manufacturing, and Institute of Materials Science, University of Connecticut, Storrs, Connecticut 06269

M. Q. Huang and W. E. Wallace

Department of Materials Science and Engineering, Carnegie Mellon University, Pittsburgh, Pennsylvania 15213

(Received 8 March 1996)

A nuclear-magnetic-resonance study of Fe_{16}N_2 , Fe_4N , and bcc α -Fe phases, along with magnetization experiments on samples containing these phases, show that the average Fe moment for Fe_{16}N_2 is between $2.7\mu_B$ and $2.9\mu_B$. This is in contrast to previous studies which report values ranging from $2.3\mu_B$ to $3.5\mu_B$. The results reported here indicate an average Fe moment for Fe_{16}N_2 which is significantly enhanced compared with that for α -Fe, and a moment for Fe atoms at the $4d$ sites which is greater than $3.0\mu_B$. [S0163-1829(96)01225-8]

Although the α' - Fe_{16}N_2 phase was discovered more than 40 years ago,¹ its magnetic properties have attracted considerable attention only recently.²⁻¹¹ The issue is whether or not this phase has an average Fe moment 30% greater than that for bcc α -Fe, i.e., a so-called "giant moment." If such a giant Fe moment exists, such as that first reported in 1971,¹² then the fact that an Fe atom can carry a moment of about $3\mu_B$ in Fe metallics should stimulate a major effort to explore additional metallic materials with high saturation magnetization. Also, this will have an impact on the theory of magnetism of $3d$ metals since such a large moment is unexpected from band theory. However, after four decades of study on this phase, pure single-phase Fe_{16}N_2 has not yet been produced in bulk form due to fundamental difficulties.^{1,6} All of the samples studied to date contained a large amount of bcc α -Fe and γ -Fe(N) (a nonferromagnetic fcc phase). The saturation magnetization for Fe_{16}N_2 has to be obtained by subtracting the contributions of these phases from the total saturation magnetization, which involves the difficult task of determining the phase composition quantitatively. The average Fe moment of the Fe_{16}N_2 phase reported by different groups varies over a wide range from $2.3\mu_B$ to $3.5\mu_B$ (see the review by Metzger, Bao, and Carbuicchio⁹). If one relies solely on the magnetization measurements, then the controversy regarding the Fe moment of Fe_{16}N_2 will remain until a single-phase sample is prepared, which is unlikely to be realized anytime soon. In order to resolve the controversy, it is essential to utilize an approach which provides information concerning the Fe moments for each individual phase (or Fe site) without requiring a knowledge of the phase composition. Hyperfine interaction experiments offer such an alternate means as the Fe hyperfine field (HF) for a given site provides a measure of the Fe moment. Within the Fe-N family, there are some similarities in the structure and magnetic properties between the Fe_4N and Fe_{16}N_2 phases. Therefore, the relationship between hyperfine fields and atomic moments for Fe_4N provides important insight in the understanding of the local magnetic properties of Fe_{16}N_2 . In this paper, we present an observation of nuclear magnetic resonance (NMR) in the Fe_{16}N_2 phase, and a comparative HF

study of Fe_{16}N_2 , Fe_4N , and α -Fe, focusing on the question of a giant Fe magnetic moment in the Fe_{16}N_2 phase.

The Fe_4N single phase and Fe_{16}N_2 samples used in this study were prepared following essentially the method originated by Lehrer¹³ and Jack¹ several decades ago. We refer the reader to Refs. 4-6 for the details of the sample preparation. The phase constituents of the samples were identified using x-ray diffraction (XRD) and NMR. We note that NMR is a much more sensitive technique than XRD for detecting a trace amount of α -Fe, which is characterized by a sharp ^{57}Fe resonance peak at 46.7 MHz. The Fe_{16}N_2 sample consisted of three phases: α' - Fe_{16}N_2 , α -Fe, and γ -Fe(N). The α' - Fe_{16}N_2 content was about 50 wt. %. A careful search for α -Fe in the Fe_4N sample was carried out at 1.3 K by NMR with various experimental conditions, and none was observed.

The magnetization curves for the two samples were obtained at 4.5 K using a superconducting quantum interference device magnetometer. The results are shown in Fig. 1. Also shown in Fig. 1 is the magnetization curve for α -Fe as a reference. The saturation magnetization values for these

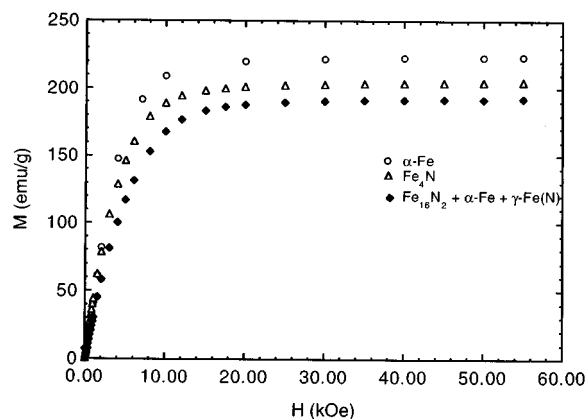


FIG. 1. Magnetization M in emu/g vs magnetic field H in kOe measured at 4.5 K: open circles, bcc α -Fe; open triangles, Fe_4N ; closed diamonds, Fe_{16}N_2 .

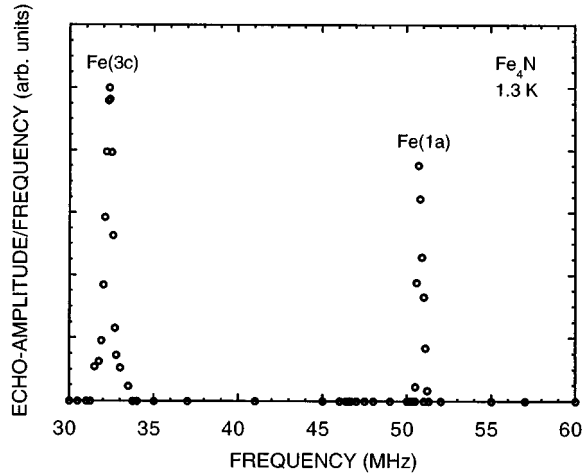


FIG. 2. ^{57}Fe spin-echo NMR spectrum obtained from the Fe_4N sample at 1.3 K.

samples are 226 emu/g (or an Fe moment of $2.26\mu_B$) for the α -Fe phase, 205 emu/g (or an average Fe moment $\bar{\mu}_{\text{Fe}}=2.18\mu_B$) for the Fe_4N phase, and 196 emu/g for the Fe_{16}N_2 sample from which the saturation magnetization for the Fe_{16}N_2 phase was deduced to be 280 emu/g ($\bar{\mu}_{\text{Fe}}=2.9\mu_B$).⁶

The ^{57}Fe hyperfine field measurements were carried out at 1.3 K using a coherent spin-echo NMR spectrometer for frequencies ranging from 10 to 300 MHz. Figure 2 shows the NMR spectrum obtained from the Fe_4N sample. Two ^{57}Fe peaks were observed at 32.3 and 50.7 MHz with a relative intensity (area) ratio of about 3:1, respectively. These correspond to ^{57}Fe HF's of 235 and 369 kOe, respectively. Based on the intensities, the 32.3- and 50.7-MHz peaks are assigned to the face center [$\text{Fe}(3c)$] and the corner [$\text{Fe}(1a)$] Fe sites, respectively, consistent with the previous assignment for this compound.¹⁴ Figure 3 shows the ^{57}Fe spectrum obtained from the Fe_{16}N_2 sample. In addition to the 46.7-MHz α -Fe peak, there are three peaks centered at 42.1, 45.2, and 57.3 MHz, which correspond to ^{57}Fe HF's of 307, 329, and 418 kOe, respectively. The relative intensities (areas) for the 42.1- and 57.3-MHz peaks are about the same, and are

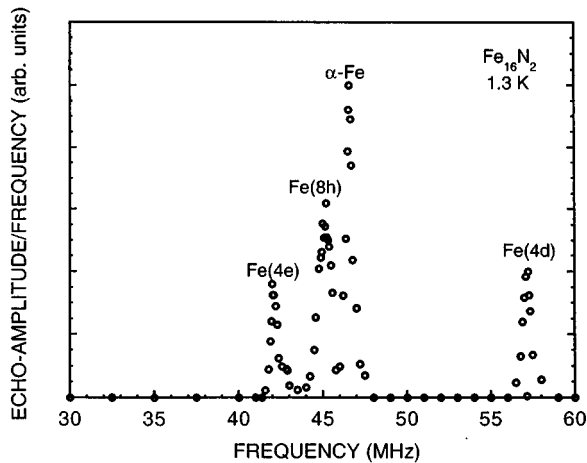


FIG. 3. ^{57}Fe spin-echo NMR spectrum obtained from the Fe_{16}N_2 sample at 1.3 K.

about half of that for the 45.2-MHz peak. The ^{57}Fe NMR frequency in Fe_3N is 35.6 MHz.¹⁴ Therefore, it is concluded that the Fe_{16}N_2 sample studied contains no Fe_3N or Fe_4N phases. Based on the relative intensities (areas) and the corresponding HF values, the 42.1-, 45.2-, and 57.3-MHz peaks are assigned to the $\text{Fe}(4e)$, $\text{Fe}(8h)$, and $\text{Fe}(4d)$ sites, respectively. The HF values obtained from the present NMR experiments are close to most Mössbauer effect (ME) experimental results,^{10,15,16} with one exception in that the HF value for the $\text{Fe}(4d)$ site obtained by NMR is slightly greater than that obtained from ME. The very sharp NMR spectra shown in Figs. 2 and 3 indicates the small orbital and dipolar fields in these two systems.

The total HF, H_{hf} , produced by s electron polarization at a given $3d$ atom, in this case Fe, is often thought of as arising from three contributions,¹⁷ and can be written

$$H_{\text{hf}}(i) = H_{\text{cp}}(i) + H_s(i) + H_{\text{sp}}(i). \quad (1)$$

In Eq. (1), H_{cp} and H_s denote the contributions from the spin polarizations of the inner-core s electrons and the outer $4s$ electrons, respectively, by the $3d$ electrons of the Fe atom itself and, consequently, are approximately proportional to the on-site moment $\mu_{\text{Fe}}(i)$. H_{sp} denotes the contribution from the $4s$ electron polarization due to the nearest-neighbor Fe atoms, and is approximately proportional to the number of the nearest Fe neighbors, $n(i)$, and their average atomic moments $\bar{\mu}_{\text{Fe}}$. Hence,

$$H_{\text{cp}}(i) + H_s(i) = a\mu_{\text{Fe}}(i) \quad \text{and} \quad H_{\text{sp}}(i) = cn(i)\bar{\mu}_{\text{Fe}}, \quad (2)$$

where a and c are the hyperfine coupling constants. The fact that $H_{\text{cp}} + H_s$ is dominant for Fe-based systems, and that the average moment of Fe atoms in the nearest neighbor shell is not significantly different from that of the on-site moment, lead to the following approximate expression which is generally used within a specific system:

$$H_{\text{hf}}(i) \approx A\mu_{\text{Fe}}(i), \quad (3)$$

where A is the effective hyperfine coupling constant due to spin polarization. For a given Fe atom, A is dependent on its electronic state and the configuration of the neighboring Fe atoms, including the symmetry of the distribution of these Fe atoms and distances of the Fe atom to these neighboring Fe atoms. For Fe ionic compounds, there are no conduction electrons and, hence, $H_s = 0$ and $H_{\text{sp}} \approx 0$. The experiments on ferrites have shown that the contribution of core polarization to the Fe HF leads to $a \approx 110 \text{ kOe}/\mu_B$.¹⁸ For Fe metallic systems, H_{sp} has a significant contribution to the total HF, leading to $A > a$. α -Fe, for which $H_{\text{hf}} = 340 \text{ kOe}$ (Fig. 3) and $\mu_{\text{Fe}} = 2.26\mu_B$ (Fig. 1), is a typical example of such Fe metallics, and has $A = 150 \text{ kOe}/\mu_B$. For a majority of the Fe-based materials, A varies over a range of values from approximately 110 to 150 kOe/μ_B (from ionic to metallic).

Fe_4N has the fcc structure. The number of the nearest Fe neighbors for the $1a$ and $3c$ Fe sites in Fe_4N is the same (12). The difference between these two sites is that there are two N atoms in the nearest-neighbor shell surrounding the $3c$ site, while no N atoms surround the $1a$ site. The Fe moments at the $1a$ and $3c$ sites have been determined by neutron diffraction (ND) to be $2.98\mu_B$ and $2.01\mu_B$, respectively.¹⁹ Substituting these values into Eq. (3), the hy-

perfine coupling constant values are found to be 124 and 117 kOe/ μ_B for the Fe(1*a*) and Fe(3*c*), respectively. Combining the total Fe moments obtained from magnetization measurements (Fig. 1), and taking the NMR frequencies (Fig. 2) as relative measures of the Fe moments at the 1*a* and 3*c* sites, we have

$$\mu_{\text{Fe}}(1a) + 3\mu_{\text{Fe}}(3c) = 8.71\mu_B$$

and

$$\mu_{\text{Fe}}(1a) = 1.57\mu_{\text{Fe}}(3c). \quad (4)$$

Thus it is found that $\mu_{\text{Fe}}(1a) = 2.99\mu_B$, $\mu_{\text{Fe}}(3c) = 1.91\mu_B$, and $A(\text{Fe}_4\text{N}) = 123$ kOe/ μ_B . The excellent agreement between the NMR and ND results for Fe_4N indicates that HF experiments might be an effective means of obtaining information concerning the Fe atomic moment in Fe-N systems. Also, these results indicate that the N-Fe interaction results in a decrease in the Fe moment and, consequently, the Fe HF, whereas it does not affect the hyperfine coupling constant.

Fe_{16}N_2 has a definite structural relationship with α -Fe and Fe_4N . If one takes the number of Fe atoms within 2.866 Å (the lattice parameter of bcc α -Fe) as the number of Fe nearest neighbors, then the Fe_{16}N_2 lattice can be regarded as a distorted Fe_4N structure. The 4*d*, 8*h*, and 4*e* sites have 12, 12, and 13 Fe atoms, respectively, with the same average distance of 2.65 Å. Based on the similarity between the two phases, it is expected that $A(\text{Fe}_{16}\text{N}_2)$ and $A(\text{Fe}_4\text{N})$ are close to each other. Furthermore, the local atomic environments for Fe sites in Fe_{16}N_2 are similar to that in Fe_4N , for which each Fe has 12 Fe atoms as the nearest neighbors at 2.68 Å. In particular, the 4*d* site for Fe_{16}N_2 is similar to the 1*a* site for Fe_4N in that both sites have no N as nearest neighbors. The 8*h* and 4*e* sites for Fe_{16}N_2 are similar to the 3*c* site for Fe_4N , with the principal difference being that the 8*h* and 4*e* sites in Fe_{16}N_2 have one N nearest neighbor, while the 3*c* site in Fe_4N has two N nearest neighbors. From these considerations, one expects that $\mu_{\text{Fe}}(4d)$ is much greater than $\mu_{\text{Fe}}(8h)$ and $\mu_{\text{Fe}}(4e)$, and that $\mu_{\text{Fe}}(8h)$ and $\mu_{\text{Fe}}(4e)$ are nearly equal to each other with $\mu_{\text{Fe}}(8h)$ being slightly greater. As shown in Fig. 3, the high ^{57}Fe HF value for the 4*d* site and the much lower and very close HF values for the 4*e* and 8*h* sites are consistent with the above-mentioned picture. Therefore, it is expected that the effective hyperfine coupling constant values for the three Fe sites are, just as found in Fe_4N , nearly equal to each other. Hence, the ^{57}Fe HF value for a given Fe site can be taken as a measure of the relative magnitude of the Fe moment at that site, namely,

$$\mu_{\text{Fe}}(4e) : \mu_{\text{Fe}}(8h) : \mu_{\text{Fe}}(4d) = 1.00 : 1.07 : 1.36. \quad (5)$$

The value of the *average* Fe moment for Fe_{16}N_2 reported by the researchers in favor of a giant Fe moment ranges from $2.8\mu_B$ to $3.5\mu_B$, while the value reported by those who argue against such a giant moment ranges between $2.3\mu_B$ and $2.7\mu_B$.^{9,10} This situation clearly shows the difficulty involved in a determination of the phase composition. The NMR results reported here offer two critical points: (1) the ratio of Fe moments follows Eq. (5), and (2) the hyperfine

coupling constant for the three sites should lie between 120 and 150 kOe/ μ_B , and is closer to 120 kOe/ μ_B . Hyperfine experiments cannot provide exact values for the Fe moments; however, these two points are valuable in the consideration of the reasonableness of the reported magnetization values for the Fe_{16}N_2 phase.

Taking $A(\text{Fe}_{16}\text{N}_2) = 120$ kOe/ μ_B and using Eq. (5), the present NMR measurements yield the maximum possible Fe moment values: $\mu_{\text{Fe}}(4e) = 2.6\mu_B$, $\mu_{\text{Fe}}(8h) = 2.7\mu_B$, and $\mu_{\text{Fe}}(4d) = 3.5\mu_B$, and an average moment $\bar{\mu}_{\text{Fe}} = 2.9\mu_B$. This is in agreement with the result obtained by Huang and Wallace from magnetization and XRD measurements.⁶ The NMR results reported here are consistent with an average Fe moment in the range of 2.7 to $2.9\mu_B$, which is the common ground for the two points of view. The present NMR study does not support the idea of either a very low or very high average Fe moment for the Fe_{16}N_2 phase. If we follow the claim that $\bar{\mu}_{\text{Fe}} = 2.3\mu_B$ in Fe_{16}N_2 ,^{10,11} then, from Eqs. (3) and (5), $\mu_{\text{Fe}}(4e) = 2.0\mu_B$, $\mu_{\text{Fe}}(8h) = 2.2\mu_B$, and $\mu_{\text{Fe}}(4d) = 2.8\mu_B$. In order to fit this low value of average Fe moment to the NMR data, the value of A has to be 150 kOe/ μ_B , which, as discussed above, is unreasonably large. On the other hand, if we follow the claim that $\bar{\mu}_{\text{Fe}} = 3.5\mu_B$,³ this corresponds to $\mu_{\text{Fe}}(4e) = 3.1\mu_B$, $\mu_{\text{Fe}}(8h) = 3.3\mu_B$, and $\mu_{\text{Fe}}(4d) = 4.3\mu_B$, and A has to be 97 kOe/ μ_B , which is also unreasonable since it is significantly less than the hyperfine coupling constant for the Fe^{3+} ion (110 kOe/ μ_B).

From the point of view of a fundamental study, the determination of the saturation magnetization for Fe_{16}N_2 relates to the issue of whether or not an Fe atom in a metallic system can carry a moment greater than $3.0\mu_B$. In this regard, the Fe moment at the 4*d* site is of most importance. According to Eq. (5), an average Fe moment greater than $2.5\mu_B$ corresponds to a Fe(4*d*) moment greater than $3.0\mu_B$ (and a value of A less than 140 kOe/ μ_B). Despite the large controversy over the average Fe moment among various research groups, most of the reported data implies $\mu_{\text{Fe}}(4d) > 3.0\mu_B$.

At first sight, it might be difficult to accept a 25% increase in the average Fe moment as a consequence of lattice enlargement. In this regard, the enhancement of the Fe moment in $R_2\text{Fe}_{17}$ systems due to lattice expansion upon nitrogenation is instructive.²⁰ Studies of various $R_2\text{Fe}_{17}$ compounds and their nitrides show that nitrogen insertion into the interstitial sites of the 2:17 phase increases the volume of the lattice by about 6%, which leads to an increase in the average Fe moment of 12–15%.²⁰ In comparison with the $R_2\text{Fe}_{17}$ systems, the transformation from α -Fe to Fe_{16}N_2 leads to a volume increase of 11%. Thus one might expect a larger increase in the average Fe moment.

Finally, it should be pointed out that the giant Fe moment found in Fe_{16}N_2 mainly originates from electron spins, and does not have a significant orbital contribution. Otherwise, as the hyperfine coupling constant for an orbital moment is much higher than that for a spin moment, one would have observed much higher Fe NMR peak frequencies.

In conclusion, the ^{57}Fe HF values for Fe_{16}N_2 and Fe_4N samples have been obtained by NMR experiments. Our analysis makes use of information concerning the structural, hyperfine interaction, and magnetization data. It is shown that, in order to reconcile the magnetic data with the NMR

results, the Fe atoms, especially for the 4*d* site Fe atoms, should have large atomic moments. The enhancement appears to be a volume effect. The fact that an Fe atom in Fe metallics can possess a spin moment of greater than $3.0\mu_B$ is a significant result in the field of magnetism.

This work was supported in part by NSF Grant No. DMR9319367, by the Army Research Office, and by the Connecticut Advanced Center for Manufacturing. The authors would like to thank G. Fernando for many useful discussions.

-
- ¹K. H. Jack, Proc. R. Soc. London Ser. A **195**, 41 (1949); **208**, 200 (1951); **208**, 216 (1951); Acta Crystallogr. **5**, 404 (1952).
²K. H. Jack, J. Appl. Phys. **76**, 6620 (1994).
³Y. Sugita *et al.*, J. Appl. Phys. **70**, 5977 (1991).
⁴M. Q. Huang *et al.*, J. Magn. Magn. Mater. **135**, 226 (1994).
⁵M. Q. Huang *et al.*, J. Appl. Phys. **75**, 6574 (1994).
⁶W. E. Wallace and M. Q. Huang, J. Appl. Phys. **76**, 6648 (1994).
⁷M. Komuro *et al.*, J. Magn. Soc. Jpn. **13**, 301 (1989).
⁸M. Komuro *et al.*, J. Appl. Phys. **69**, 5126 (1990).
⁹R. M. Metzger, X. H. Bao, and M. Carbucicchio, J. Appl. Phys. **76**, 6626 (1994).
¹⁰J. M. D. Coey, J. Appl. Phys. **76**, 6632 (1994).
¹¹M. Takahashi *et al.*, J. Appl. Phys. **76**, 6642 (1994).
¹²T. Kim and M. Takahashi, Appl. Phys. Lett. **20**, 492 (1972).
¹³E. Lehrer, Z. Elektrochem. **36**, 383 (1930).
¹⁴T. D. Xiao *et al.*, Nanostruct. Mater. **2**, 285 (1993).
¹⁵T. Moriya *et al.*, J. Phys. Soc. Jpn. **35**, 1378 (1973).
¹⁶X. Bao, R. M. Metzger, and M. Carbucicchio, J. Appl. Phys. **73**, 6734 (1993).
¹⁷Y. D. Zhang *et al.*, J. Magn. Magn. Mater. **100**, 13 (1991).
¹⁸J. J. Van Loef, Physica **32**, 2102 (1966).
¹⁹B. C. Frazer, Phys. Rev. **112**, 751 (1958).
²⁰J. M. D. Coey and H. Sun, J. Magn. Magn. Mater. **87**, L251 (1990).