

## Electronic properties of superconducting $\text{LuNi}_2\text{B}_2\text{C}$ and related boride carbide phases

L. F. Mattheiss

*AT&T Bell Laboratories, Murray Hill, New Jersey 07974*

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Band properties of the quaternary superconductor  $\text{LuNi}_2\text{B}_2\text{C}$  ( $T_c \approx 17$  K) and related intermetallic phases (including, for example,  $\text{LuNiBC}$  and  $\text{YNi}_2\text{B}_2\text{C}$ ) have been calculated in the local-density approximation via the linear augmented-plane-wave method. The results feature a density-of-states peak near the top of nearly filled  $\text{Ni}(3d)$  bands, with only modest B and C orbital admixture. The calculated  $\text{LuNi}_2\text{B}_2\text{C}$  bands suggest that these interesting boride carbide phases represent a family of conventional (rather than high- $T_c$ ) superconductors.

Recently, Cava *et al.*<sup>1</sup> have discovered superconductivity at 23 K in the quaternary Y-Pd-B-C system, thereby equaling the previous  $T_c$  record<sup>2</sup> that existed prior to the advent in 1986 of the high- $T_c$  era. Unfortunately, the composition of the superconducting phase has not yet been identified. (Multiphase superconducting samples were obtained from melt compositions with an average stoichiometry of  $\sim \text{YPd}_5\text{B}_3\text{C}_{0.3}$ .) However, superconductivity has also been observed<sup>3</sup> in the closely related R-Ni-B-C system ( $R=\text{Y}$ , Tm, Er, Ho, and Lu), though with slightly lower critical temperatures (up to  $\sim 16.6$  K). In this case, the composition<sup>3</sup> (i.e.,  $\text{RNi}_2\text{B}_2\text{C}$ ) and structure<sup>4</sup> ( $\text{ThCr}_2\text{Si}_2$ -like with filled  $2b$  sites) of the superconducting phase have been determined. These findings provide a possible explanation for the earlier report<sup>5</sup> of superconductivity at 12 K in the Y-Ni-B system.

The purpose of the present investigation is to calculate the electronic properties of these superconductors in order to help clarify the issue as to whether these materials should be regarded as conventional intermetallic superconductors or the first members of a possible family of high- $T_c$  boride carbide phases. We consider here only the nonmagnetic compounds for which the rare-earth  $f$  levels are either completely filled (i.e.,  $R=\text{Lu}$ ) or completely empty ( $R=\text{Y}$ ). Based on the results that are presented below, it appears that  $\text{LuNi}_2\text{B}_2\text{C}$  (and related compounds) are conventional intermetallic superconductors whose promising transition temperatures ( $\sim 16.6$  K at present) are derived from a standard electron-phonon mechanism in which the light mass of the B and C constituents leads to high-frequency phonons and enhanced  $T_c$ 's.

The structure of the  $\text{LuNi}_2\text{B}_2\text{C}$  phase<sup>4</sup> is illustrated in Fig. 1(a). It consists basically of the body-centered-tetragonal (bct)  $\text{ThCr}_2\text{Si}_2$  structure (space group  $I4/mmm$ ), but with an extra constituent (namely, carbon) at the  $2b$  site in the basal plane. Thus,  $\text{LuNi}_2\text{B}_2\text{C}$  is related structurally to the well-known heavy-Fermion superconductor  $\text{CeCu}_2\text{Si}_2$  (Ref. 6). The observed  $\text{LuNi}_2\text{B}_2\text{C}$  lattice parameters ( $a \approx 3.46$  Å,  $c \approx 10.63$  Å) and atom-position parameters<sup>4</sup> produce nearly ideal  $\text{NiB}_4$  tetrahedra with B-Ni-B bond angles of  $108.8^\circ$  and  $110.9^\circ$  and a Ni-B bond distance of 2.1 Å. The  $\sim 2.45$  Å Ni-Ni bondlength is slightly smaller than the

fcc Ni value ( $\sim 2.50$  Å), thus suggesting the presence of metal-metal bonds in this phase. The B-C bondlength ( $\sim 1.47$  Å) is slightly smaller than the value for hexagonal  $\text{B}_4\text{C}$  ( $\sim 1.64$  Å). Although this bct structure is analogous to those adopted by many of the cuprate high- $T_c$  superconductors, the calculated electronic properties for  $\text{LuNi}_2\text{B}_2\text{C}$  are shown below to exhibit three-dimensional (3D) rather than 2D characteristics.

The structure of a closely related simple-tetragonal (st)  $\text{LuNiBC}$  phase<sup>4</sup> (which has not yet been found to superconduct) is illustrated in Fig. 1(b). This compound contains a rocksalt-type  $\text{Lu}_2\text{C}_2$  double layer in place of the  $\text{LuNi}_2\text{B}_2\text{C}$  monolayer, and this changes the Bravais lattice from bct to st ( $a \approx 3.50$  Å,  $c \approx 7.75$  Å). The resulting nonsymmorphic space group exhibits  $P4/nmm$  symmetry. As Siegrist *et al.*<sup>4</sup> have noted, both structures possess remarkably similar bond distances and bond angles. This strongly suggests that the electronic and bonding characteristics of both materials are comparable.

In the present study, the self-consistent band structure for bct  $\text{LuNi}_2\text{B}_2\text{C}$  and st  $\text{LuNiBC}$  have been calculated in the local-density approximation (LDA) with the use

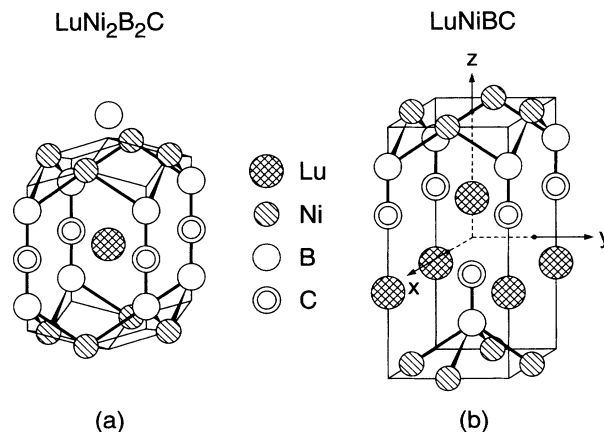


FIG. 1. Comparison of the primitive unit cells for (a) body-centered-tetragonal (bct)  $\text{LuNi}_2\text{B}_2\text{C}$ , and (b) simple-tetragonal (st)  $\text{LuNiBC}$ .

of a scalar-relativistic version<sup>7</sup> of the linear augmented-plane-wave (LAPW) method.<sup>8</sup> The present LAPW implementation imposes no shape approximations on either the charge density or potential. The LAPW basis has included plane waves with a 1.36 Ry cutoff ( $\sim 60$ –70 LAPW's per atom) and spherical-harmonic terms through  $l = 8$  (Lu, Ni) or  $l = 6$  (B, C) within the corresponding muffin-tin spheres. The crystalline charge density and potential have been expanded using  $\sim 4500$ –5500 plane waves (64 Ry) in the interstitial region and lattice-harmonic expansions ( $l_{\max} = 6$ ) within the muffin tins. A twelve-point  $\mathbf{k}$  sample in the  $\frac{1}{8}$  irreducible wedge has been used to carry out Brillouin-zone (BZ) integrations.

Exchange and correlation effects have been included within the LDA framework with the use of the Wigner interpolation formula.<sup>9</sup> The atomic Lu( $5d^{16}s^2$ ), Ni( $3d^9 4s^1$ ), B( $2s^2 2p^1$ ), and C( $2s^2 2p^2$ ) states are treated as valence electrons in these calculations, while a frozen-core approximation has been applied to the more tightly bound core-type levels, including the Lu( $4f$ ) electrons. This rigid-core Lu( $4f$ ) treatment, which has been validated by a separate calculation where these electrons are treated as valence states, helps to simplify an already complicated valence-band manifold. The assumed LAPW muffin-tin radii are 1.32 (Lu), 1.20 (Ni), 0.75 (B), and 0.72 (C) Å.

The results of the present LAPW calculations for LuNi<sub>2</sub>B<sub>2</sub>C are shown in the left-hand panel of Fig. 2.

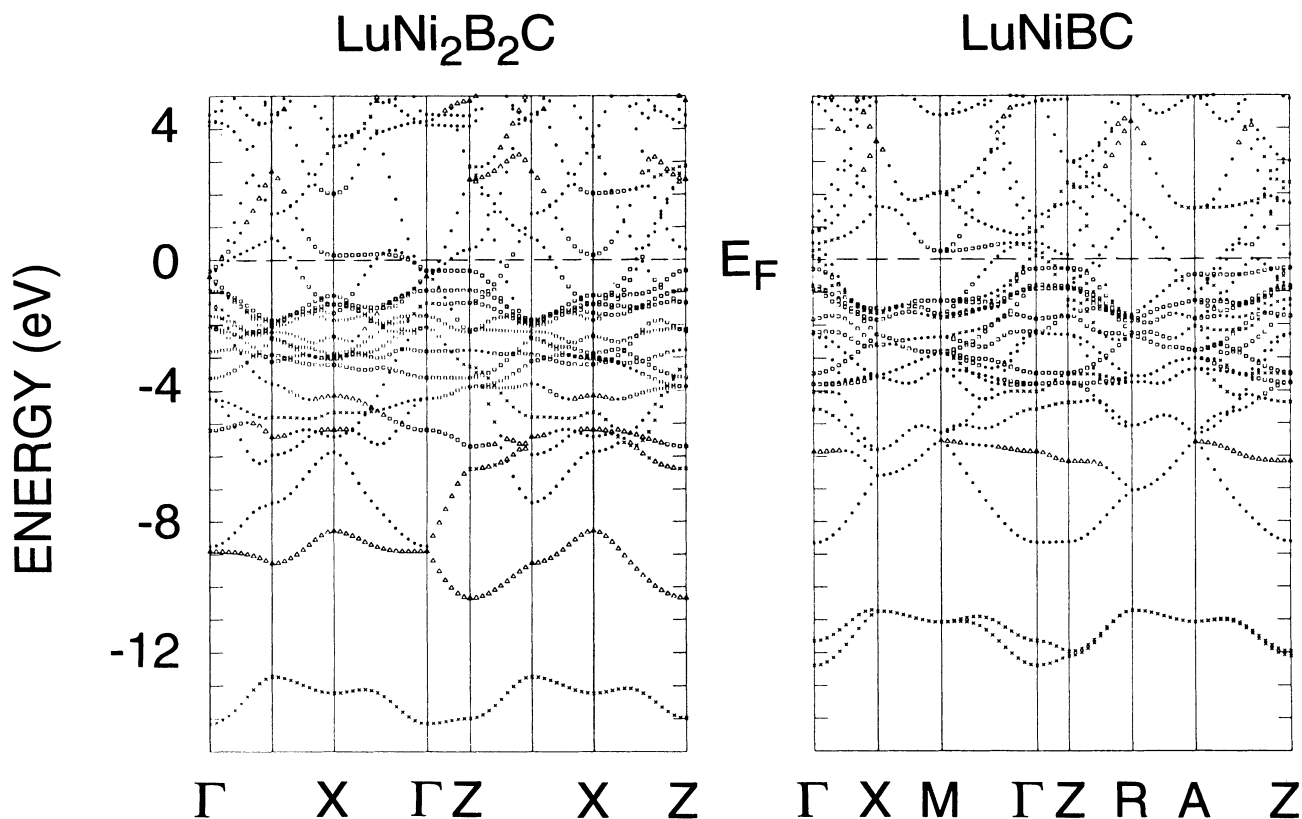


FIG. 2. LAPW energy-band results along comparable Brillouin-zone directions for bct LuNi<sub>2</sub>B<sub>2</sub>C and st LuNiBC. Bands with significant Ni, B, and C orbital character are labeled with squares ( $\omega_{\text{Ni}} > 0.5$ ), triangles ( $\omega_{\text{B}} > 0.12$ ), and crosses ( $\omega_{\text{C}} > 0.12$ ), respectively. The rigid-core approximation places the filled Lu( $4f$ ) levels (not shown) at binding energies of  $\sim 3.8$  and 2.9 eV, respectively.

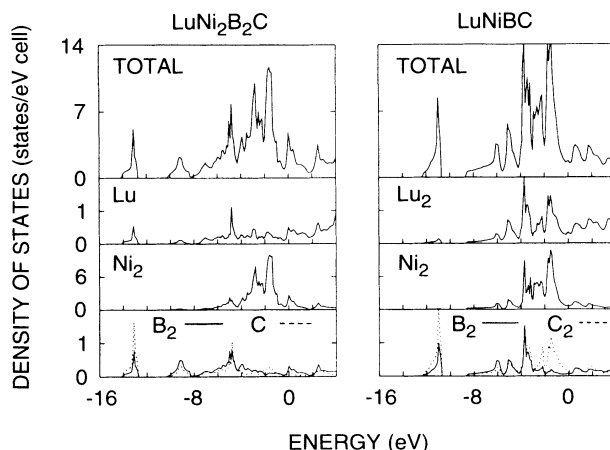


FIG. 3. Total and muffin-tin-projected density-of-states results for LuNi<sub>2</sub>B<sub>2</sub>C and LuNiBC.

The lowest band originates from the C( $2s$ ) levels while the upper valence-band complex that begins at  $\sim -10$  eV evolves gradually from B( $2s$ ) to B( $2p$ )-C( $2p$ ), and finally to Ni( $3d$ ) character near  $E_F$ . The LuNi<sub>2</sub>B<sub>2</sub>C Fermi level occurs near the top of the Ni( $3d$ ) bands. However, the lowest subband of the Lu( $5d$ ) manifold (which exhibits mainly  $d_{x^2-y^2}$  character) dips slightly below  $E_F$  at  $\Gamma$ , thereby producing a small pocket of Lu( $5d$ ) carriers in this material.

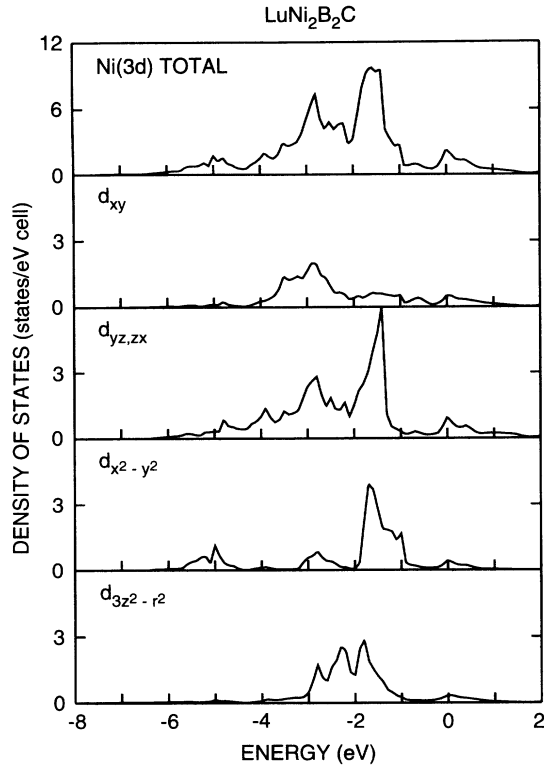


FIG. 4. Projected Ni(3d) density-of-states results for  $\text{LuNi}_2\text{B}_2\text{C}$ , including total and individual-orbital components.

The LAPW band results for st  $\text{LuNiBC}$  that are shown to the right in Fig. 2 exhibit qualitative similarities with the bct  $\text{LuNi}_2\text{B}_2\text{C}$  results, although structural differences preclude a direct band-by-band comparison. To minimize these structural differences, the bct  $\text{LuNi}_2\text{B}_2\text{C}$  results in Fig. 2 are plotted along symmetry lines of the st BZ, thus ignoring the actual shape of the bct BZ. In both compounds, the Lu(4f) levels (not shown) lie near the bottom of the Ni(3d) bands, well below  $E_F$ . This is confirmed below by a separate  $\text{LuNi}_2\text{B}_2\text{C}$  calculation with Lu(4f)-type valence states, where relaxation increases the Lu(4f) rigid-core binding energy by  $\sim 0.7$  to  $\sim 4.5$  eV.

A comprehensive overview of the  $\text{LuNi}_2\text{B}_2\text{C}$  and  $\text{LuNiBC}$  electronic properties is provided by the density-of-states results  $N(E)$  that are shown in Fig. 3. The  $\text{LuNi}_2\text{B}_2\text{C}$  Fermi level coincides with a  $N(E)$  peak where  $N(E_F) \approx 2.4$  states/eV Ni. This is about 1.8 times larger than that calculated<sup>10</sup> for a typical cuprate superconductor such as  $\text{La}_2\text{CuO}_4$ , where  $N(E_F) \approx 1.3$  states/eV Cu. In fact, the  $\text{LuNi}_2\text{B}_2\text{C}$  value for  $N(E_F)$  is comparable (on a per-atom basis) to that calculated<sup>11</sup> for  $\text{V}_3\text{Si}$ , where  $N(E_F) \approx 2.5$  states/eV V. Thus, even ignoring enhancement effects, this calculated band-structure value for  $N(E_F)$  places  $\text{LuNi}_2\text{B}_2\text{C}$  among the “conventional” superconductors when this material is added to the plot<sup>12</sup> of  $T_c$  vs  $\gamma$ .

Comparing the  $\text{LuNi}_2\text{B}_2\text{C}$  and  $\text{LuNiBC}$   $N(E)$  results, one concludes that an important difference between the two phases is the fact that the  $\text{LuNiBC}$  Fermi level oc-

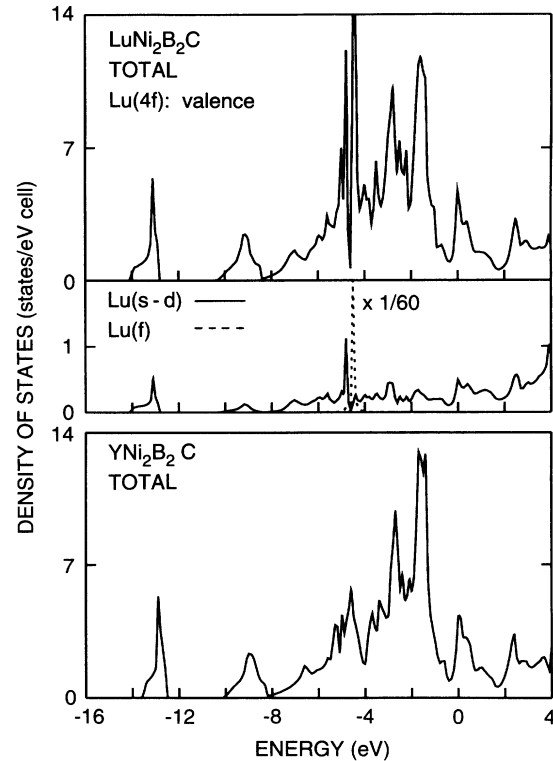


FIG. 5. Total and Lu-projected density-of-states results for  $\text{LuNi}_2\text{B}_2\text{C}$  in which the Lu(4f) states are treated as valence electrons. Total  $\text{YNi}_2\text{B}_2\text{C}$  density-of-states results are shown below.

curs at an  $N(E)$  minimum, just below the corresponding  $\text{LuNi}_2\text{B}_2\text{C}$  peak. This suggests that the superconducting prospects for  $\text{LuNiBC}$  will be improved if the Fermi level can be raised (by about 1 electron/cell) by substitutional doping until it coincides with the unoccupied  $\text{LuNiBC}$  peak. In this regard, it is noteworthy that the Lu contribution to  $N(E_F)$  is not totally negligible in these materials. Thus, substitutional doping at this site is less favorable than in typical cuprate superconductors such as  $\text{La}_2\text{CuO}_4$  where the La component is small.<sup>10</sup>

The present boride carbide phases differ from the cuprates in that they suggest no obvious “few-band” model that summarizes the essential features of their electronic structure. This is evident from the projected Ni(3d) density-of-states results in Fig. 4. All five Ni(3d) subbands contribute in roughly equal proportions to the  $N(E)$  peak near  $E_F$ . In fact, the  $d_{xy}$  and  $d_{yz,zx}$   $N(E_F)$  components exhibit a slight enhancement (by  $\sim 16\%$  and  $6\%$ , respectively) relative to the corresponding  $d_{x^2-y^2}$  and  $d_{3z^2-r^2}$  terms.

The  $\text{LuNi}_2\text{B}_2\text{C}$  results in Figs. 2–4 are essentially unchanged when the Lu(4f) levels are treated as valence bands rather than rigid-core states. As shown in Fig. 5, the  $\text{LuNi}_2\text{B}_2\text{C}$  total and Lu-projected  $N(E)$  results are visually indistinguishable from the previous frozen-core results in Fig. 3. Shown below are the corresponding  $N(E)$  results that have been derived from an analogous LAPW calculation<sup>13</sup> for isostructural  $\text{YNi}_2\text{B}_2\text{C}$ . In view of the obvious  $\text{YNi}_2\text{B}_2\text{C}$ - $\text{LuNi}_2\text{B}_2\text{C}$  band similarities, it is

not surprising that the  $T_c$ 's for these two phases ( $\sim 15.6$  K for  $\text{YNi}_2\text{B}_2\text{C}$  and  $\sim 16.6$  K for  $\text{LuNi}_2\text{B}_2\text{C}$ ) are comparable.

Finally, we consider the question of whether these Ln-Ni-B-C phases should be regarded as conventional rather than high- $T_c$  superconductors. The present boride carbide phases do not satisfy the *ad hoc* band criteria that have been proposed<sup>14</sup> for identifying possible high- $T_c$  candidates. The key element of these criteria is the presence of half-filled  $\sigma^*$  bands that are due to strong nearest-neighbor bonds between nearly degenerate cation-anion levels. Compared to the cuprates, the reduced B and C binding energies (relative to oxygen) yield nearly degenerate Ni(3d)-B(2s) or Ni(3d)-C(2p) orbital combinations. However, there are no readily identifiable  $\sigma^*$  bands near  $E_F$  in these Ln-Ni-B-C phases that satisfy the "half-filled" requirement.

The tetrahedrally coordinated Ni produces Ni-B bonds that are relatively weak, and this leads to Ni-B hybridization effects that are moderate in strength. As a result, the bands near the  $\text{LuNi}_2\text{B}_2\text{C}$  Fermi level exhibit predominant Ni(3d) orbital character. A simple model calculation illustrates the basic differences between the

planar  $p_\sigma$ - $d_{x^2-y^2}$  geometry of typical cuprate high- $T_c$  superconductors<sup>10</sup> and the tetrahedral  $\text{NiB}_4$  coordination of  $\text{LuNi}_2\text{B}_2\text{C}$  in their ability to produce strong  $\sigma$  bonds and a large  $\sigma$ - $\sigma^*$  bandwidth. For example, a tight-binding calculation<sup>14</sup> for a hypothetical  $\text{CaNiN}_2$  phase that contains  $\text{NiN}_2$  planes and degenerate Ni(3d)-N(2p) orbitals yields a total valence bandwidth of 11.6 eV. However, an analogous calculation for  $\text{LuNi}_2\text{B}_2\text{C}$  that includes the same tight-binding parameters produces an overall bandwidth of only 6.1 eV for this geometry.

In summary, the results of LAPW band calculations for the recently discovered intermetallic superconductors  $\text{LuNi}_2\text{B}_2\text{C}$  and  $\text{YNi}_2\text{B}_2\text{C}$  predict a density-of-states peak at  $E_F$  that exhibits predominant Ni(3d) orbital character. Analysis of these LDA results suggests that the observed superconducting properties seem to place these materials in the category of conventional rather than high- $T_c$  superconductors.

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<sup>13</sup> This calculation has utilized the  $\text{YNi}_2\text{B}_2\text{C}$  lattice parameters ( $a = 3.53$  Å,  $c = 10.57$  Å) in combination with the  $\text{LuNi}_2\text{B}_2\text{C}$  atom-position parameters (Refs. 3 and 4).

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