

Experimental determination of the unoccupied bands of W(110)

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Inverse photoemission from W(110) has been used to determine the dispersion of the unoccupied energy bands along the tungsten $\Gamma\Sigma N$ line. The critical points are determined to be Γ_{12} at 2.6 eV, N_1 at 2.2 eV, and the lower limit of N_4 at 3.2 eV. The upper bound of the lifetime broadening is estimated as 0.3–0.4 eV at 2–3 eV above E_F . The image state is observed at 4.5 eV above E_F , or 0.7 eV in binding energy, relative to the vacuum level.

I. INTRODUCTION

The band structure of tungsten has been studied extensively^{1–17} for a variety of reasons. Tungsten is a useful high-temperature material, and is used in field-emission- and scanning-tunneling-microscope tips. In addition, tungsten serves as a convenient substrate for the growth of epitaxial ferromagnetic-rare-earth films. The W(100) surface exhibits an unusual reconstruction,^{8,18} but the W(110) surface is more stable and therefore suitable for bulk-band-structure studies. While the occupied part of the band structure has been extensively investigated with photoemission for different surfaces,^{9–12} the unoccupied bands have not been fully characterized experimentally.^{14–17} The previous measurements of empty bands with inverse-photoemission were limited and did not map out the band dispersion experimentally.

We used a high-resolution, tunable inverse-photoemission system to map the unoccupied bands at the close-packed W(110) surface. In the following, we will show the experimental dispersion of the bulk bands along the $\Gamma\Sigma N$ line, and compare our results with the previously published experimental data and the band calculations. The experimental results for the unoccupied W $5d$ Σ_1 and Σ_4 bands reveal that the dispersion is qualitatively similar to that obtained by theoretical calculations but are shifted to the higher energies by some 0.1–0.7 eV for different parts of the bands. The real and imaginary parts of the self-energy will be discussed. We find an image potential surface state at 4.5 eV that has not been previously reported.

II. EXPERIMENT

The experiment was carried out on a high-resolution inverse-photoemission spectrometer with a photon energy range of 8–30 eV and up to 44 eV in second order. The energy resolution was 0.26 eV [half-width at full maximum (HWFHM)] at 9 eV and was larger at higher energies (0.6 eV at 18 eV). The accuracy of the energy calibration is 0.1 eV. All the spectra in this work were taken

at room temperature. The pressure of the spectrometer chamber was in the 10^{-11} -torr range. The W(110) single crystal has been cleaned with the normal procedure (oxygen treatment followed by flashing)¹⁹ in the attached preparation chamber. The surface cleanliness and order of the crystal after such a procedure were checked with photoemission and low-energy electron diffraction (LEED) in another separate chamber, and with LEED in the preparation chamber.

III. RESULTS

Figure 1 shows a typical inverse-photoemission spectrum with an initial electron energy of 14.5 eV. The feature at 4.54 eV is very sensitive to contamination and shows no k_{\perp} dependence. We interpret this feature as the image state of W(110). Taking the work function of W(110) as 5.25 eV, the binding energy of this state (energy below the vacuum level) is 0.7 eV, the same as that observed for W(001) (Ref. 8) within the experimental error. The appearance of this image state is an excellent indication that the surface is contamination-free. The three features at lower energies exhibit energy dispersion (relative to E_F) with the incident electron energy (dependence upon k_{\perp}). From a comparison with existing band-structure calculations, we can identify these features as the bulk bands Σ_1 , Σ_1 , and Σ_4 , respectively (see Fig. 1).

In order to determine the value of k_{\perp} experimentally, we measured the intensity at E_F for the different initial electron energies (Fig. 2). It is clear that there are two Fermi-level crossing points at 11.5 eV and 20.5 eV. According to the existing de Haas–van Alphen (dHvA) data,²⁰ Fermi-level crossings should occur at wave vectors of $0.09(N\Gamma)$ and $0.67(N\Gamma)$, respectively. Based on these two experimentally determined points of the initial band, we can assign k_{\perp} values as noted below. We take a free-electron parabola with an inner potential of -0.63 eV relative to the Fermi level to account for the higher energy part of the initial band. The Γ point is reached

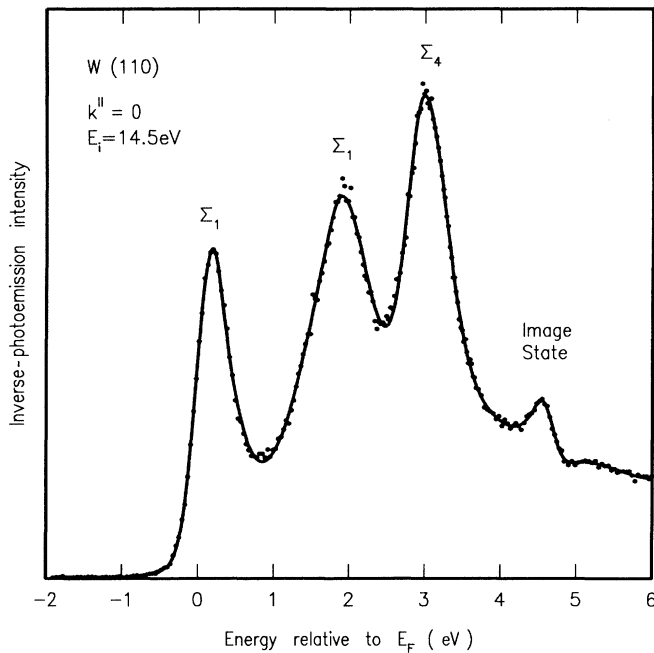


FIG. 1. A typical inverse-photoemission spectrum at normal incidence on the W(110) surface. The image state and the band assignments are indicated in the figure.

with an electron impact energy of 29.7 eV. For the lower energy part of the initial band (near the N point), a relativistic band calculation² was used to reflect the gap near the Brillouin-zone boundary. This assignment of incidence electron energies to k_{\perp} values is shown in the comparison of the upper and lower abscissas of Fig. 3.

This assignment of k_{\perp} to electron kinetic energies is also supported by other evidence. Similar plots of intensity (at the fixed energies above E_F) versus incident electron energies also indicate when the Σ_1 and Σ_3 bands disperse through at a particular k_{\perp} point. Also, a minimum width of the principal inverse-photoemission (IPE) feature (the combination of the flat Σ_1 and Σ_4 bands, unresolved at high energies) is reached around 30 eV, which confirms our assignment of the critical point Γ_{12} (at 29.7 eV). Our k_{\perp} assignment is therefore accurate within $\pm 5\%$.

Figure 3 shows the experimentally determined band dispersion of the unoccupied bands of W(110) using the initial band as described above. For the $\langle 110 \rangle$ direction of a bcc crystal, the allowed final states in dipole transitions are Σ_1 , Σ_3 , and Σ_4 , owing to the fact that the initial state of normal incident electrons should be Σ_1 and no light polarization detection is made in the current setup. In Fig. 3, we mapped out the relatively flat bands of Σ_1 and Σ_4 with an energy accuracy of 0.1–0.2 eV and momentum accuracy of 5%. By averaging the split Σ_1 and Σ_4 bands near Γ , the critical point Γ_{12} is determined as 2.6 ± 0.2 eV. By extrapolating the experimental points near N , the N_1 point is determined as 2.2 ± 0.2 eV, and the lower limit of N_4 as 3.2 eV. Compared with a relativistic pseudopotential band calculation,³ the experimental bands

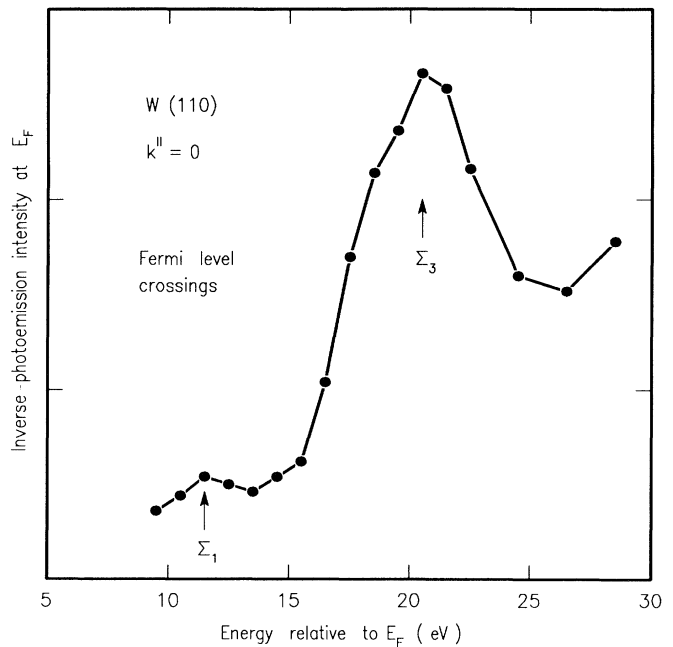


FIG. 2. Normalized intensities at E_F of a set of inverse-photoemission spectra at normal incidence, taken with different initial electron energies. The peaks at 11.5 and 20.5 eV indicate the Σ_1 and Σ_3 bands, respectively, dispersing through E_F .

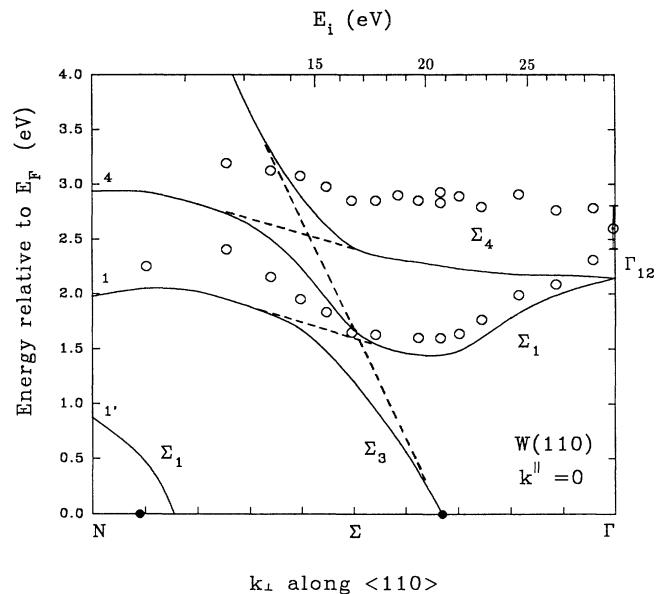


FIG. 3. Unoccupied band dispersion (open circles) along $\Gamma\Sigma N$ of W(110), derived from inverse-photoemission spectra. The solid circles are from dHvA data (Ref. 19), which were used to determine the k_{\perp} as described in the text. The corresponding initial electron energies for different k_{\perp} are shown on the upper horizontal axis. A relativistic band calculation (Ref. 3) is given for the solid lines for comparison. The band assignment is given for the nonrelativistic limit, which is shown as the dashed lines.

are shifted to higher energies away from E_F . The significant broadening makes it difficult to map out steep bands, such as the Σ_1 and Σ_3 bands that cross E_F . We attribute the feature near E_F to the rapidly dispersing Σ_1 and Σ_3 bands.

IV. DISCUSSION

The unoccupied states of W(110) were previously studied with inverse photoemission by Collins *et al.*¹⁴ and Funnemann and Merz.¹⁵ The latter study presented very similar spectra to those reported here, although all their energy positions were shifted to higher energies by about 1 eV, possibly due to an uncertainty in the position of the Fermi level. Collins *et al.*¹⁴ showed a sequence of IPE spectra with photon energies of 15–30 eV, but observed no image state and their features were less pronounced. They compared their data with photocurrent calculations, yet no attempt of a direct band mapping was made because of the difficulty of determining k_{\perp} . In the current work, with the experimentally determined k_{\perp} and better spectral resolution, we have demonstrated that the standard three-step model of the photoemission process can still be used to understand the band structure of tungsten.

As seen in Fig. 3, there is clearly a discrepancy between the theoretical band structure and our experimental results, especially for the Σ_4 band. The experimentally determined Σ_4 band is shifted upwards for about 0.5–0.6 eV. At the critical point Γ_{12} , this shift is 0.5 ± 0.2 eV. Several different types of band calculations, such as the relativistic augmented-plane-wave (RAPW) calculation done by Christensen and Feuerbacher,¹ and the semirelativistic linearized APW calculations done by Jansen and Freeman⁴ and Wei, Krakauer, and Weinert,⁵ agree with each other within 0.1–0.2 eV. In the data of Collins *et al.*,¹⁴ a similar shift away from E_F also appears compared to their own photocurrent calculation, though the differences are less pronounced than is the case in our results. This consistent discrepancy between all these theoretical band structures and our experimental results cannot be explained by finite angular acceptance or crystal misorientation. If a nonzero k_{\parallel} is the cause of this large shift, then we should expect a large band dispersion with changing incident angle and a shift to lower energies at a certain angle. Instead, as shown in Fig. 4, tilting the crystal along either the $\langle 001 \rangle$ direction or the $\langle 1\bar{1}0 \rangle$ direction for several degrees only results in the Σ_4 band dispersing upward slightly. The pronounced changes in the image state also indicate that the angular resolution of the system is much better than 5° .

The origin of the discrepancy between the IPE experimental bands and the calculations is unclear. Several effects need to be considered, as discussed below.

(1) The one-dimensional density of states (1D DOS) could dominate the experimental spectra under certain circumstances. Most of the existing photoemission data (with a photon energy of about 25–50 eV) from clean W(110) do not show clear band dispersion with k_{\perp} and are therefore considered to be more representative of the 1D DOS along the $\Gamma\Sigma N$ line, instead of the direct transition.^{10,12,13} This could be caused by evanescent up-

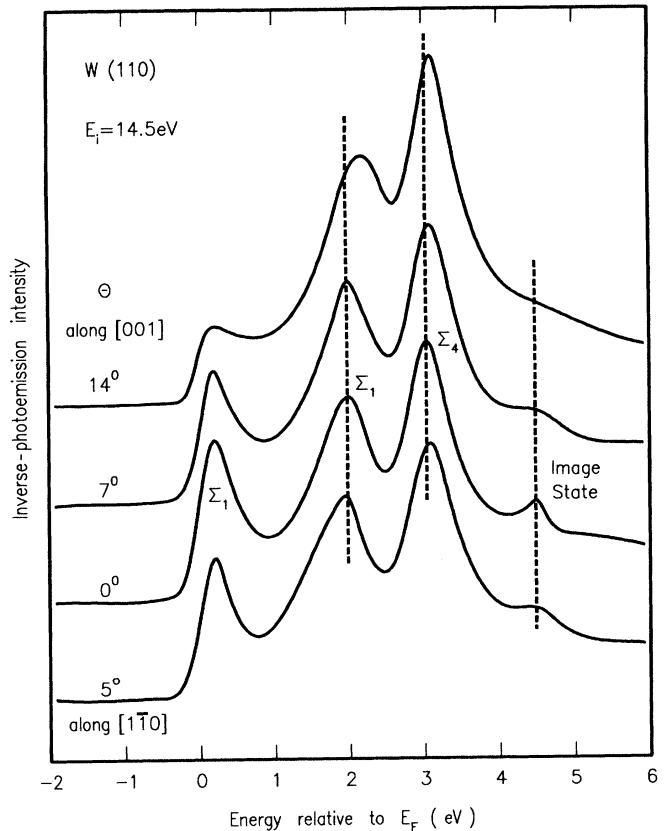


FIG. 4. Inverse-photoemission spectra of W(110) with different incident angles for $E_i = 14.5$ eV. The dispersion of the Σ_4 band is discussed in the text. The image state disperses to higher energy and becomes broader for off-normal incidence angles, as expected.

per states or by multiple upper bands. In general, we cannot assign our data to a 1D DOS since we observe dispersion for changing k_{\perp} quite clearly, especially for low incidence energies. With higher incidence electron energies, there are more states along $\Gamma\Sigma N$ available (seen in the band calculation²) and more broadening in the initial band of IPE. It is therefore possible that the spectra become more representative of a 1D DOS at higher incidence energies. This effect can result in the IPE features shifted to higher energies from Γ_{12} and a finite apparent splitting even at the Γ_{12} point, where the Σ_1 and Σ_4 bands should become degenerate, even in the relativistic case. Therefore we take the average between the Σ_1 and Σ_4 peaks as the Γ_{12} point. The splitting of the IPE features at large electron energies accounts for some of the differences between experiment and theory but not all. This 1D DOS-like effect does not explain the differences that extend over a large range of the Brillouin zone, in particular for our highest Σ_4 point at 3.2 eV. It lies 0.4–0.5 eV above the calculated band maximum at N_4 . It would be desirable to compare photocurrent calculations, which use the real upper bands, with high-quality experimental data. The photocurrent calculation of Collins *et al.*,¹⁴ however, still produces a spectrum that lies lower

in energy than the data at higher energies (near Γ).

(2) Surface effects, which are not included in the bulk band calculations,¹⁻⁵ can also show up in the experimental data. Feuerbacher and Christensen¹ argued that the apparent shift (about 1 eV) of the projected gap at $\bar{\Gamma}$ of the occupied bands in their photoemission data is caused by the narrowing of the bandwidth at the surface. For the unoccupied part of the band structure, the narrowing of the surface DOS can only result in a shift toward E_F instead of away from E_F . The existing surface density of state calculations^{6,7} do not reveal the presence of a surface state with an energy slightly higher than the bulk Σ_4 band necessary to explain the band shift to greater energies.

(3) The band-structure calculations are for the ground state of N electrons while the inverse photoemission involves the excited state of $(N + 1)$ electrons. This means there is always an energy difference between the experimental and theoretical band structures. This can be calculated as the real part of the self-energy term and should explain part of the discrepancy between the band calculations and experiment. Such a self-energy calculation for tungsten, however, does not exist, at present, to the authors' knowledge.

The imaginary part of the self-energy is the lifetime broadening of the excited state. For photoemission in the EDC (energy distribution curve) mode with normal emission electrons, the linewidth of the observed transition is often written as^{21,22}

$$\Gamma = \left| \frac{\Gamma_i + \left| \frac{V_i}{V_f} \right| \Gamma_f}{1 - \frac{V_i}{V_f}} \right|,$$

where Γ_f and Γ_i are the inverse lifetimes of the electrons in final and initial states, respectively. V_f and V_i are the group velocities along the normal direction in the final and initial states. As Smith, Thiry, and Petroff²³ recently pointed out, however, the expression of the measured linewidth is different for the different modes, i.e., EDC, CIS (constant initial state spectrum), and CFS (constant final state spectrum). For our inverse-photoemission setup (using fixed incidence energy and measuring the emitted photon spectra), the expression for CFS in Ref. 22 is the suitable form:

$$\Gamma = \Gamma_f + \left| \frac{V_f}{V_i} \right| \Gamma_i,$$

where the initial and final states are reversed as compared to the case in photoemission. While V_i and V_f can be measured from the initial and final band dispersion, Γ_i can be estimated from the Fermi-level crossing at 11.5 eV in Fig. 2 as about 2 eV. This corresponds to a mean free path of about 8 Å, which is reasonable for this energy

range. By measuring the widths of the 2.0 eV and 3.1 eV features in the $E_i=14.5$ eV spectrum and correcting for the instrumental resolution (about 0.3 eV at these energies), we can therefore estimate the inverse lifetime of the electrons in the lower final state band as

$$\Gamma_f(3.1 \text{ eV}) = 0.5 \pm 0.1 \text{ eV}$$

and

$$\Gamma_f(2.0 \text{ eV}) = 0.4 \pm 0.1 \text{ eV}.$$

Here we have ignored the possible k -dependence effects. Suppose this lifetime broadening also follows a linear relationship with the energy away from E_F [$\Gamma_f(E) = C(E - E_F)$], as is the case for the $3d$ bands of some transition metals, then the slope C for the tungsten $5d$ bands is 0.1–0.2. This value is similar to that for single-crystal Cu(0.1),²² but much lower than that for Fe (0.6).²⁴ Given that there exist other sources of peak broadening, such as phonon contribution,¹¹ surface imperfections, etc., the values given here should be considered as the upper bounds of the intrinsic lifetime broadening.

By taking luminescence data at different high incident electron energies (for example, 40 eV), several constant photon energy features are revealed at 10.3 eV, 14.5 eV, and 20.8 eV. These features are consistent with the known surface-plasmon modes of 9.7 eV, 14.8 eV, and 20.8 eV, as obtained by optical measurements.²⁵ Although some of these peaks are in the neighborhood of the Fermi-level crossing points, it is clear that the bands indeed cross E_F around 11.5 and 20.5 eV by taking the corresponding plot of intensity versus photon energy for final states above E_F .

V. CONCLUSION

The unoccupied tungsten $5d$ band structure has been mapped out with a three-step direct transition model. The discrepancy between the experimental results and the band calculations has been discussed, particularly in terms of the real part of the self-energy, and the one-dimensional density of state effects. The upper bound of the energy-dependent lifetime broadening is estimated as 0.4 eV at 2.0 eV above E_F and 0.5 eV at 3.1 eV above E_F .

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¹N. E. Christensen and B. Feuerbacher, Phys. Rev. B **10**, 2349 (1974); B. Feuerbacher and N. E. Christensen, *ibid.* **10**, 2373 (1974).

²R. F. Willis and N. E. Christensen, Phys. Rev. B **18**, 5140 (1978); N. E. Christensen and R. F. Willis, J. Phys. C **12**,

167 (1979).

³D. M. Bylander and L. Kleinman, Phys. Rev. B **29**, 1534 (1984).

⁴H. J. F. Jansen and A. J. Freeman, Phys. Rev. B **30**, 561 (1984).

- ⁵S.-H. Wei, H. Krakauer, and M. Weinert, *Phys. Rev. B* **32**, 7792 (1985).
- ⁶O. Jepsen and R. O. Jones, *Phys. Rev. B* **34**, 6695 (1986).
- ⁷M. P. López Sancho, J. M. López Sancho, and J. Rubio, *J. Phys. C* **18**, 1803 (1985).
- ⁸W. Drube, D. Straub, F. J. Himpsel, P. Soukiassian, C. L. Fu, and A. J. Freeman, *Phys. Rev. B* **34**, 8989 (1986).
- ⁹M. W. Holmes, D. A. King, and J. E. Inglesfield, *Phys. Rev. Lett.* **42**, 394 (1979).
- ¹⁰R. H. Gaylord, K. H. Jeong, and S. D. Kevan, *Phys. Rev. Lett.* **62**, 2036 (1989); *Phys. Rev. B* **36**, 9337 (1987); **37**, 8491 (1988).
- ¹¹S. Dhar and S. D. Kevan, *Phys. Rev. B* **41**, 8516 (1989).
- ¹²R. C. White, C. S. Fadley, M. Sagurton, and Z. Hussain, *Phys. Rev. B* **34**, 5226 (1986).
- ¹³G. B. Blanchet, N. J. DiNardo, and E. W. Plummer, *Surf. Sci.* **118**, 496 (1982).
- ¹⁴I. R. Collins, A. D. Laine, P. T. Andrews, and P. J. Durham, *J. Phys. Condens. Matter* **3**, 5307 (1991); **4**, 2891 (1992).
- ¹⁵D. Funnemann and H. Merz, *J. Vac. Sci. Technol. A* **5**, 657 (1987).
- ¹⁶I. L. Krainsky, *J. Vac. Sci. Technol. A* **5**, 735 (1987).
- ¹⁷P. O. Nilsson and C. G. Larsson, *Jpn. J. Appl. Phys. Suppl.* **17**, 144 (1978).
- ¹⁸A. Horlacher-Smith, R. A. Barker, and P. J. Estrup, *Surf. Sci.* **136**, 329 (1984); L. D. Roelofs, J. W. Chung, S. C. Ying, and P. J. Estrup, *Phys. Rev. B* **33**, 6537 (1986).
- ¹⁹For example, R. G. Musket, W. McLean, C. A. Colmenowes, D. M. Mako Wieszcki, and W. J. Siekhaus, *Appl. Surf. Sci.* **10**, 143 (1982); Dongqi Li, Jiandi Zhang, P. A. Dowben, and M. Onellion, *Phys. Rev. B* **45**, 7272 (1992).
- ²⁰R. F. Girvan, A. V. Gold, and R. A. Phillips, *J. Phys. Chem. Solids* **29**, 1485 (1968).
- ²¹T. -C. Chiang, J. A. Knapp, M. Aono, and D. E. Eastman, *Phys. Rev. B* **21**, 3513 (1980).
- ²²J. A. Knapp, F. J. Himpsel, and D. E. Eastman, *Phys. Rev. B* **19**, 4952 (1979).
- ²³N. V. Smith, P. Thiry, and Y. Petroff (unpublished).
- ²⁴A. Santoni and F. J. Himpsel, *Phys. Rev. B* **43**, 1305 (1991).
- ²⁵J. H. Weaver, C. G. Olson, and D. W. Lynch, *Phys. Rev. B* **12**, 1293 (1975).