Erratum: Correlation length and free energy of the $S=\frac{1}{2}$ XYZ chain [Phys. Rev. B 43, 5788 (1991)]

Minoru Takahashi

In the left-hand column of page 5790, the second line from the bottom, "$v_l \rightarrow v_l + v$ in . . ." should read "$v_l \rightarrow v_l - v$ in . . . ." In Eq. (12),

$$\sum_{l=1}^{M} \cdots \text{ should read } \prod_{l=1}^{M} \cdots .$$

On page 5791, the line below Eq. (19), "If we put $\alpha \equiv -\frac{1}{2}(\eta + v_1)$, . . ., should read "If we put $\eta + v_1 = 2K - 2\alpha$, $2\alpha = \mathrm{sn}^{-1}[\mathrm{sn}(2\eta)a/c]$, ..." In Eq. (37),

$$\left| \left[\frac{J_z}{4\pi T(l-1)} \right] \right|^2 \cdots \text{ should read } \left[\frac{J_z}{4\pi T(l-1)} \right]^2 \cdots .$$

Erratum: Mass enhancement close to a Mott transition [Phys. Rev. B 43, 6292 (1991)]

Mucio A. Continentino

On the top of p. 6293, left-hand column, line 8, should read: "For this purpose we first notice that the ground-state-energy density can be expanded as $E_T = E_0 + a_1 Z(k_F^2/2m) + O(Z^2)$, where k_F is the Fermi momentum of the noninteracting system according to Luttinger's theorem[8,9] and a_1 is a constant. We now consider a finite hypercubic system of finite size L and calculate the difference in ground-state-energy density ΔE_T between the situations where we impose periodic and antiperiodic boundary conditions.[10]" The scaling properties of the effective mass signaling the metal-insulator transition can be obtained more rigorously if we consider the charge stiffness D_c as defined by Shastry and Sutherland [Phys. Rev. Lett. **65**, 243 (1990)]. In fact, D_c is proportional to the inverse of the conductivity effective mass and scales, as we obtained for Z, as $D_c \propto \xi^{-(d+z-2)}$ close to the transition.