## InP(110) oxidation with  $O_2$ , NO, and N<sub>2</sub>O at 20 K: Temperature and photon-energy dependencies

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Interfaces formed at 20 K by condensing  $O_2$ , NO, and N<sub>2</sub>O on InP(110) reveal striking differences in reactivity when they are heated, illuminated with visible light, or irradiated with soft x rays  $[hv=170 \text{ eV}, \text{ photon exposures of } (1-1000)\times 10^{14} \text{ cm}^{-2}]$ . Synchrotron-radiation photoemission studies show that  $O_2$  and  $N_2O$  desorb from the surface without reacting when annealed above  $\sim$  100 K but that oxidelike In and P bonding configurations are produced by NO reactions when heated to  $\sim$  150 K. Flood-lamp illumination induces reaction for O<sub>2</sub>/InP and NO/InP at 20 K but not for  $N_2O/InP$  (hv=0.5–4.5 eV, photon exposure  $\sim 10^{19}$  cm<sup>-2</sup>). Irradiation with 170-eV photons yields surface oxides for all three oxygen-bearing condensates. Investigations of the reaction cross sections for these systems indicate that surface chemistry is mediated by the capture of photogenerated lowenergy electrons. Electron-attachment processes that produce negative ions are less likely for  $N_2O$ than for NO or  $O<sub>2</sub>$ .

Photon-induced oxidation of III-V semiconductor surfaces is scientifically challenging and technologically important. Many excellent studies of GaAs(110) oxidation have been conducted as a function of temperature,  $1-3$  oxidizing gas,  $2,4-6$  and photon energy.<sup>7</sup> The most recent work has focused on the dynamics of photon-induced reactions at temperatures where the reactant is condensed on the surface. $^{\widehat{6},8-10}$  These condensed-layer systems have exhibited temperature and photon-energy dependencies, and it has been shown that the experimental probe (photoemission) induces both reaction and desorption.<sup>8,9</sup> In general, it has been concluded that substrate electronic excitations are responsible for photon-induced oxidation, but a consensus on the role of these excitations has yet to be reached.

In this paper, we report the first studies of the interactions of the condensed oxidizers  $O_2$ , NO, and N<sub>2</sub>O on InP(110) at 20 K. These gases were chosen because they In $P(110)$  at 20 K. These gases were chosen because they oxidize In $P$ ,  $1,11-13$  they are sensitive to light exposure and they exhibit different thresholds for photon- and electron-induced decomposition. Core-level and valence-band studies indicate that each is adsorbed molecularly on InP(110) at 20 K, but that striking differences are revealed upon heating above  $\sim$  100 K or illuminating with visible light and soft x rays. In particular, condensed NO reacts with InP when annealed to 150 K but condensed  $O_2$  and  $N_2O$  simply desorb. Illumination of the condensed-gas/InP(110) surface with a highintensity flood lamp induces InP oxidation with  $O<sub>2</sub>$  and NO but not with N<sub>2</sub>O (total exposure  $\sim 10^{19}$  photons  $cm^{-2}$ ). Irradiation with 170-eV photons promotes reaction with all three gases,  $O_2$ , NO, and N<sub>2</sub>O. The differences can be understood by considering the effect of temperature on the adsorbate-semiconductor system and the role of negative ions formed via capture of photogenerated electrons.

High-resolution photoemission experiments were conducted at the Wisconsin Synchrotron Radiation Center. Energy distribution curves (EDC's) were acquired with a double-pass cylindrical mirror analyzer that, when combined with monochromator contributions, yielded an overall resolution of  $\sim$ 400 MeV for the P 2p core levels that will be the focus of this paper. Valence-band and In 4d core-level results supporting our conclusions will be reported elsewhere. Posts of InP (S doped, reported elsewhere. Posts of InP (S doped,  $N_D = 2.7 \times 10^{18}$  cm<sup>-3</sup>) were cleaved at pressures of<br>  $\sim 5 \times 10^{-11}$  Torr to yield pristine (110) surfaces. Nearly identical chemical trends were measured for p-type InP. All of the results were obtained with samples for which the clean surface Fermi level  $E_F$  was less than 60 meV from the band edge at 300 K.

The gases  $O_2$ , NO, and N<sub>2</sub>O were introduced into the experimental chamber through a calibrated leak valve. The gas line was flushed during and after bakeout because NO and  $N_2O$  decompose by reaction with the chamber walls.<sup>4,8</sup> Inspection of the valence-band emission immediately after gas exposure gave no evidence for contaminants. Gas exposures (in langmuirs where<br> $1 \text{ L} = 1 \times 10^{-6}$  Torrsec) were measured with the ion gauge without correction for gauge sensitivity. Results acquired for dosing the sample with the gauge on and off were nearly identical, indicating a negligible effect on molecules condensed on the surface. The ion gauge was off except during gas dosing or sample heating, and the sample remained in the dark except during irradiation with synchrotron radiation or the flood lamp.

Figure 1(a) presents P 2p core-level results obtained after the exposure of InP(110) to 4-L NO, 2-L  $O_2$ , and 4-L N<sub>2</sub>O at 20 K followed by warming above 100 K. The purpose was to determine whether adsorption and desorption were reversible in the absence of an electron source or photon probe, i.e., whether there would be thermal conversion from physisorbed to chemisorbed states. Gas exposures were chosen to condense approximately equal numbers of oxygen atoms on the surface. Each spectrum has been normalized to constant height to emphasize line-shape changes, and the bulk features have been aligned. EDC's obtained after  $N_2O$  and  $O_2$  conden-



FIG. 1. (a) P 2p core-level EDC's for 4-L NO, 2-L O<sub>2</sub>, and 4-L N<sub>2</sub>O condensed on InP(110) in the dark at 20 K and then annealed to desorb the gas. The N<sub>2</sub>O and O<sub>2</sub> spectra show no substrate disruption. In contrast, annealing NO-covered surfaces to 150 K results in the formation of bonding configurations characteristic of an intermediate oxide. (b) P  $2p$  EDC's after illumination by a floodlamp ( $h v=0.5-4.4$  eV). Surfaces exposed to  $O_2$  and  $N_2O$  were annealed above 100 K before data acquisition. The spectra for the NO-exposed surface were acquired with a minimum of photon irradiation to avoid x-ray-induced oxidation. No reaction was observed for N<sub>2</sub>O, but two reaction-induced P peaks appeared for O<sub>2</sub> and NO. (c) P 2p EDC's after irradiation with soft x rays  $(hv=170 \text{ eV})$  showing reaction for all three molecules.

sation at 20 K and heating were indistinguishable from those of the clean surface, exhibiting only surface-shifted and bulk P peaks (although the Gaussian linewidths did increase 20—50 meV due to thermal broadening). This indicates that  $N_2O$  and  $O_2$  desorb from the surface without inducing measureable oxidation or surface unrelaxation. (The spectral linewidths are broader for  $O<sub>2</sub>/InP$  because of the monochromator resolution in that experiment.) In contrast, warming from 20 to 150 K for 4-L NO/InP(110) resulted in strong adatom-substrate interactions as a component formed  $\sim$  1.2 eV from the bulk P feature and as emission from the surface-shifted peak decreased. This new feature corresponds to an intermedidecreased. This new feature corresponds to an intermedi-<br>ate P oxide, as reported by Hollinger *et al*.<sup>11</sup> in their high-resolution study of  $O_2/InP(110)$  reaction at 300 K.

In Fig. 2 we present schematic Lennard-Jones —type potential-energy curves for physisorption and chemisorption to the observed low-coverage interactions of  $O_2$ , NO, and  $N_2O$  with InP(110) as a function of temperature. Physisorption is represented by a van der Waals potential that has a minimum at  $E_p$ . Splitting the molecule into atoms far from the surface requires a dissociation energy  $E_d$  while atom-substrate interactions lower the energy of the system to  $E_c$ , the chemisorption energy. The activation energy for dissociative chemisorption  $E_a$  is the energy difference between  $E_p$  and the intersection of the chemisorption and physisorption potential-energy curves, chemisorphion and physisorphion potential-energy curves,<br>as sketched. For NO,  $E_a \ll E_p$ , and the incoming molecule can be trapped in a physisorption well without reacting with the substrate.<sup>14</sup> Heating the NO/InP(110) system excites vibrational modes of the molecule and, since the barrier for reaction is smaller than for molecular desorption, NO-InP chemical interactions are observed, as in Fig.  $1(a)$ .

The thermally induced reactions for NO/InP(110) contrast with what is seen for  $O_2$  and  $N_2O$  on InP(110).<sup>15</sup> Since  $O_2$  and N<sub>2</sub>O desorb upon heating above ~100 K,



FIG. 2. Schematic potential-energy functions for dissociative chemisorption that describe the temperature-dependent results of Fig. 1(a). The curves for NO/InP(110) indicate that the physisorbed molecule can overcome the chemisorption barrier upon vibrational excitation (annealing). For  $O_2$  and  $N_2O$ , physisorption and chemisorption potentials intersect above zero, favoring desorption upon heating.

we conclude that the intersection of their potentialenergy curves for physisorption and chemisorption is such that  $E_a > E_p$ , as drawn in Fig. 2. The observed behavior for  $O_2/\text{InP}(110)$  was not surprising, based on the results of Bertness et al.<sup>4</sup> for  $O_2/GaAs(110)$ . They found that  $E_a$  was ~0.45 eV and  $E_p$  was ~0.15 eV for  $O_2/GaAs(110)$  at 300 K, in reasonable agreement with the barrier of  $\sim 0.3$  eV reported by Mönch.<sup>1</sup> For  $N<sub>2</sub>O/InP(110)$ , however, the results of Fig. 1(a) demonstrate that there was desorption with little reaction, and this is opposite the conclusion of Bertness et  $al<sup>4</sup>$  for  $N_2O/GaAs(110)$ . Indeed, they reported values of  $\overline{E}_a$  ~0.15 eV and  $E_n$  ~0.25 eV, indicating that chemisorption should be favored over desorption.<sup>4</sup> For  $N_2O$ , the apparent contradiction probably reflects the fact that the experiments were done at different temperatures. At low temperature, physisorption is an important step, as is the case for our results. At high temperature, it is not necessarily the precursor state of chemisorption, and direct impact of  $N_2O$  with the surface may yield the chemisorbed state directly. This reaction path was not considered by Ref. 4. We also note that studies of  $N_2O$ -GaAs(110) interactions at 20 K have shown that  $E_a > E_p$ and the  $N_2O$  condensed at low temperature does not react with GaAs(110) to any appreciable extent when warmed to 300 K.<sup>16</sup> Thus, thermally activated interactions of  $O_2$  and  $N_2O$  with GaAs(110) and InP(110) are analogous and they can be understood from Fig. 2.

To further examine the stability of  $O_2$ , NO, and N<sub>2</sub>O condensates on  $InP(110)$ , we illuminated these systems with a high-intensity flood lamp. The bottom EDC of Fig. 1(b) shows the result of illumination of 4 L  $N_2O/InP(110)$  for 15 min (exposure  $\sim 10^{19}$  photons  $cm^{-2}$ , spectral range 0.5–4.5 eV) followed by warming to 140 K to remove any unreacted physisorbed molecules. Since no significant line-shape changes relative to the clean surface were observed, we conclude that lowenergy photons induced minimal reaction between  $N_2O$ and InP(110). In contrast, illumination of 2 L  $O_2/InP(110)$  followed by heating to 100 K produced considerable surface oxidation, as indicated by the features shifted  $\sim$  1.2 and  $\sim$  4 eV from the bulk peak of Fig. 1. Analogous features have been observed for  $O_2/InP(110)$ at 300 K, although at  $10^6 - 10^{12}$  higher O<sub>2</sub> exposures.<sup>3,11</sup> Hence, low-energy photon irradiation can effectively promote oxide growth at 20 K for  $O_2$  condensed on InP(110). For  $4-L$  NO/InP(110), flood-lamp illumination induced surface modifications, as is evident by the growth of structure  $\sim$  1.2 and  $\sim$  3.5 eV from the bulk peak, as shown in Fig. 1. Note that this EDC was acquired immediately after flood-lamp illumination at 20 K because heating by itself alters the surface chemistry [Fig. 1(a)]. Moreover, the x-ray exposure needed to acquire the data was much less than that used to induce the reaction depicted in Fig. 1(c). We also note that annealing the NOcovered InP(110) surface above 150 K produced the  $\sim$ 3.5-eV peak (not shown). Warming to 300 K or exposure to additional NO and flood-lamp illumination shifted the  $3.5-eV$  peak toward  $4 eV$ . Hence, flood-lamp illumination produced similar oxides for NO/InP and  $O_2/$ InP but there was almost no effect for  $N_2O/InP$ .

In Fig. 1(c) we show the effects of soft-x-ray irradiation of these condensed layers on InP(110) at 20 K. Surface oxides are produced in each case, and the spectra for soft-x-ray irradiation of  $O_2$  and NO are similar to those for flood-lamp illumination [Fig. 1(b)]. Exposure of 4-L  $N_2O/InP(110)$  to 170-eV photons results in two new P components at  $\sim$  1.2 and  $\sim$  3.5 eV. Comparison of the photon-induced reactions for  $N_2O/InP$  to those for  $O_2$ and NO illustrates the similarity between the oxides formed at 20 K when the amount of condensate is small. Differences can be attributed to differences in the total number of 0 atoms available for reaction, the kinetics of molecular decomposition and reaction for each molecule, and photon-stimulated desorption from the surface. The results of Figs. 1(b) and 1(c) are particularly intriguing because they show that soft-x-ray irradiation led to surface oxidation in all three cases but that low-energy irradiation produced a reacted surface only for NO and  $O<sub>2</sub>$ .

In general, photoinduced oxidation reflects direct adsorbate excitation by the photon and indirect adsorbate excitation via capture of electrons produced by photon absorption in the substrate. The relative roles of these processes can depend on photon energy.<sup>7</sup> For the systems under investigation here, we can show that direct excitation plays a minor role. In particular, in the floodlamp illumination experiments only that portion of the blackbody distribution that was transmitted through the glass viewport was available for excitation  $(0.5-4.5 \text{ eV})$ . These low-energy photons cannot dissociate  $O_2$  ( $E_d$  = 5.1) eV) or NO ( $E_d$  = 6.5 eV) to produce neutral, ground-state atoms.<sup>17</sup> For  $N_2O$ , the lowest-energy mode for optical dissociation (1.67 eV) is spin forbidden, but photolysis into ground-state  $N_2$  and excited O is possible. However, llumination with  $\sim 10^{19}$  photons cm<sup>-2</sup> would cause photolysis of less than 0.1% of the N<sub>2</sub>O molecules,<sup>18</sup> and such a yield is below our detection limit. Thus, the reactions summarized in Fig. 1(b) for  $O_2$ , NO, and N<sub>2</sub>O, are not related to direct photon absorption by the molecule. Indeed, the effect of flood-lamp irradiation is smallest for the molecule for which direct absorption might play a role, namely  $N_2O$ .

Processes that involve direct photoexcitation of overlying molecules and/or substrate atoms in contact with the molecules can also be excluded as major reaction channels by comparing photoionization cross sections to the effective cross sections for the photoinduced reactions. From the photoemission results, it is possible to determine a lower bound for the reaction cross section by measuring the rate of conversion of the surface-shifted P 2p core-level peak as a function of photon exposure. If we assume that reaction for each molecule proceeds by a first-order process and that only the topmost layer of InP atoms is affected initially,<sup>19</sup> then kinetic theory predicts that the relative emission of the surface-shifted component varies with photon exposure as  $\ln[I(E)/I(0)] = -\sigma E$ , where  $I(0)$  is the intensity of the surface component after condensation,  $\sigma$  is the cross section, and the units of exposure E are photons  $cm^{-2}$ . Our analysis of  $O_2$ , NO, and N<sub>2</sub>O interactions with InP(110)

at 20 K reveals cross sections of  $\sim 6 \times 10^{-15}$ ,  $1 \times 10^{-15}$ . and  $4 \times 10^{-16}$  cm<sup>2</sup>/photon, respectively, for 170-eV photons.<sup>20</sup> These reaction cross sections are about two orders of magnitude greater than the photoionization cross sections for either the molecule or the substrate at  $170 \text{ eV}$ .<sup>21</sup> Thus, surface reactions proceed more rapidly than expected for direct photon-absorption processes.

Indirect photon-induced oxidation mechanisms can depend on positive-ion formation, negative-ion forma-' $\text{tion},^{1,3,5}$  and electron-hole recombination events (nonthermal phonons).<sup>2,4</sup> Positive-ion formation can be discounted here because the energies of formation for the positive ions are greater than the dissociation energies of the neutral molecules. $22$  Negative-ion formation can be examined in terms of the potential-energy curves of Fig. 3, where we show results for gas-phase NO,  $O_2$ , N<sub>2</sub>O (solid lines), and for the lowest negative-ion states (dashed lines).  $17,22-24$  Figure 3 focuses on the lowest-energy excitation processes since the experimental results show that low-energy photons can induce reaction, at least for  $O_2$ and NO. We assume that the negative ion is more reactive than the neutral molecule and that reaction with the substrate proceeds once the ion is formed. Dissociation energies are given by the difference between the energy at the equilibrium bond distance and at infinite interatomic spacing. These dissociation energies are 6.5 eV for NO



FIG. 3. Gas-phase potential-energy curves for NO,  $O_2$ , N<sub>2</sub>O, and their ions. The zero-energy point is defined as the groundstate molecule at equilibrium separation. The  $N<sub>2</sub>O$  potential plotted as a function of N-0 separation exhibits no stable minima for the ion. Dissociation of  $N_2O$  requires the capture of a  $\sim$ 2.2-eV electron. NO and O<sub>2</sub> form metastable reactive ions upon capture of  $\sim$  0.2-eV electrons. Differences in reactivity for the three gases are attributed to the relative ease of reactive-ion formation.

and 5.1 eV for  $O_2$ .<sup>22</sup> The potential for N<sub>2</sub>O is plotted as a function of the N-0 separation, and dissociation into ground-state N<sub>2</sub> and O requires  $\sim$  1.7 eV.<sup>24</sup> Again, the experimental results show that the dissociation energy of the molecule is not the critical parameter for predicting photoinduced InP(110) surface oxidation at low temperature. Indeed, they stand in contrast to the expectation that photoinduced surface oxidation should increase for molecules with low dissociation energies. $1 -$ 

The potential-energy curves for  $N_2O - N_2O^-$  are quite different from those for NO–NO and  $O_2-O_2^-$ , as can be seen from Fig. 3. (The vibrational levels of the neutral molecule and the ion that determine the allowed Franck-Condon transitions are not shown.) For gas-phase  $N_2O$ , an electron with  $\sim$  2.2 eV is required to excite the molecule from the neutral ground state to an unstable  $N_2O^$ ion state that can decompose into  $N_2$  and reactive O<sup>-</sup>. In contrast, electron capture to form the corresponding  $NO<sup>-</sup>$  or  $O<sub>2</sub><sup>-</sup>$  ions requires less than 0.2 eV. Comparison of these potential-energy diagrams for gas-phase molecules to the experimental results of Fig. <sup>1</sup> strongly suggests that negative-ion formation controls photoninduced oxidation.

In order to understand the details of electron capture, it is necessary to know the distribution of carriers in the conduction band of InP and the molecular levels of the adsorbate. However, when the molecule approaches a surface, the molecular energy levels are shifted by image potentials that lower (raise) the electron affinity (ionization energy), $25$  and when charge transfer from a conduction-band state to the molecule occurs, the molecule-substrate potential is replaced by an ionsubstrate potential.<sup>26</sup> For ion formation to occur, the energy of a conduction-band electron must align with the molecular (ion) level. Electrons in the conduction band exist because of thermal excitation or photon absorption. The former produces carriers near the conduction-band minimum while the latter provides a large number of secondary electrons with a range of energies. Inspection of Fig. 3 shows that alignment of the electron affinity levels of NO,  $O_2$ , and N<sub>2</sub>O (or equivalently, the negative-ion curves) to a given conduction-band level of InP(110) would require that the electron affinity of  $N_2O$  be reduced by  $\sim$  2 eV more than that of either O<sub>2</sub> or NO when the molecule approaches the surface. This is evident from the fact that the electron affinities of the molecules are  $-2.2$  eV for N<sub>2</sub>O, 0.1 eV for NO, and 0.45 eV for O<sub>2</sub>. From the experimental results, we see that the reaction cross sections for  $N_2O$ , NO, and  $O_2$  on InP(110) correlate with these electron affinities, and this probably explains the very different behavior of these gases on  $InP(110)$ . The experimental results therefore support the idea that negative-ion formation is a limiting factor in photoinduced oxidation. Indeed, the importance of electronattachment processes has been discussed extensive $y, \frac{7, 20, 26 - 28}{7}$  although incorporation of these processes in a quantitative model of photon-induced surface reactions remains to be done.

In summary, photoemission studies of low-temperature oxidation of InP(110) with  $O_2$ , NO, and N<sub>2</sub>O reveal strong temperature and photon-energy dependencies that can be understood by considering the competition between pathways for molecular decomposition and reaction. Photon-energy-dependent results indicate that the probability of negative-ion formation must be considered in low-temperature oxidation. Additional clarification of the role played by excited carriers in photoinduced oxidation should result from investigations conducted with monochromatic low-energy photons<sup>7</sup> and low-energy electron beams, and such studies are underway.

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of  $N_2O$ . Examination of the valence bands for 4 L NO/InP(110) after heating to 150 K reveals the formation of  $N<sub>2</sub>O$  for InP(110) as well (Ref. 16). We note that valenceband spectra obtained immediately after NO dosing and after a 2-h residence at 20 K in the dark were indistinguishable, with features characteristic of only NO and InP. Hence NO reaction with InP $(110)$  is negligible at 20 K but it is initiated by heating or irradiation.

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