

Anisotropy of optical absorption in wurtzite-phase $\text{Zn}_{0.85}\text{Mn}_{0.15}\text{S}$

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Optical-absorption results on single-crystal samples of wurtzite-phase $\text{Zn}_{0.85}\text{Mn}_{0.15}\text{S}$ are reported. Five absorption bands associated with the Mn^{2+} intraion transition are seen. Polarization results for the four lowest bands are found to be similar to previous observations in wurtzite-phase $\text{Zn}_{1-x}\text{Mn}_x\text{Se}$, indicating systematic polarization in the Zn-based dilute magnetic semiconductors.

I. INTRODUCTION

In a recent paper,¹ we reported on optical absorption of the dilute magnetic semiconductor (DMS) $\text{Zn}_{1-x}\text{Mn}_x\text{Se}$. The investigation revealed the first evidence for polarization effects in the broad Mn^{2+} intraion transitions for samples in the wurtzite region ($0.3 \lesssim x < 0.57$). These results are of interest since extended x-ray-absorption fine structure (EXAFS) suggests that distortion of the local environment of Mn from T_d occurs in ternary alloys of $\text{Zn}_{1-x}\text{Mn}_x\text{Se}$.² In our work, an explanation for the polarization was sought based on the model of a bimodal distribution of Mn—Se Zn—Se bond lengths and the splitting of $\text{Mn}^{2+} 3d^5$ states due to bond bending. Applications of this model for a random distribution of Mn in the lattice suggested instead that no polarization should be present in any band, even in the C_{3v} wurtzite phase.

In order to clarify this question, we have extended the investigation to include study of the polarization effects in $\text{Zn}_{1-x}\text{Mn}_x\text{S}$, a second Zn-based DMS which can be grown in the wurtzite phase. The width of the band gap in this material³ (~ 3.7 eV) as compared to the selenide⁴ (~ 3.0 eV) gives a special opportunity for determining the behavior of several higher Mn transitions which are inaccessible in the lower gap material at atmospheric pressure. Choice of x values for investigation was made based on phase diagram information for the ZnS—MnS system,⁵ available samples, and interest in comparing the broad band features in the two materials.

II. EXPERIMENT

The monocrystals used in this study were grown by the chemical vapor transport method.⁶ Single crystals of external hexagonal morphology and good optical quality were chosen for study. The wurtzite structure was confirmed by x-ray measurements. The crystals were cut and oriented with the c axis in the plane of the polished surface. Two samples of different thickness were required

to cover the spectral region investigated: thicknesses chosen were 87.5 ± 1 and 370 ± 1 μm . The sample thicknesses were measured interferometrically. Adjacent slices of the same crystal were used to ensure that they had the same manganese concentration. The manganese composition was determined to be $x \approx 0.15 \pm 0.005$ by microprobe analysis. Sample preparation techniques and measurement procedures are identical to those described previously in the $\text{Cd}_{1-x}\text{Mn}_x\text{Se}$ alloy system.⁷

III. RESULTS

Absorption coefficients were calculated using the formula appropriate for multiple reflections at normal incidence for samples with parallel faces,

$$T = \frac{(1-R)^2 e^{-\alpha t}}{1 - R^2 e^{-2\alpha t}}$$

where T is the fraction of transmitted light, t is the sample thickness, α is the absorption coefficient, and R is the sample reflectance.

The above equation yields the absorption coefficient, provided the reflectance is known as a function of wavelength. Published values for the reflectance of $\text{Zn}_{1-x}\text{Mn}_x\text{S}$ are not yet available in the literature. The reflectance of ZnS has been measured⁸ to vary from ~ 0.27 to ~ 0.30 in the range from ~ 2.0 eV to ~ 3.0 eV. Unfortunately, these values for R yield negative absorption coefficients below the Mn threshold ($E < 2.2$ eV) in our samples. Instead the reflectance was derived by using two samples of the same manganese concentration but of different thickness; the value of R was then adjusted until both samples gave the same absorption coefficient at the same wavelength. In our work R was chosen to match the absorption coefficients of the $370\text{-}\mu\text{m}$ sample and the $87.5\text{-}\mu\text{m}$ sample at 515 nm. This wavelength was chosen because it lies in the midrange of the onset of the Mn^{2+} absorption threshold, where both samples showed

significant absorption ($\alpha \sim 23 \text{ cm}^{-1}$). To obtain a good match, it was found that different values of R were necessary for different polarization directions. Optimum agreement was obtained for $\mathbf{E} \parallel \mathbf{c}$ using $R_{\parallel} = 0.10$ and $\mathbf{E} \perp \mathbf{c}$ using $R_{\perp} = 0.153$. These choices are consistent with the earlier work on $\text{Zn}_{1-x}\text{Mn}_x\text{Se}$ (Ref. 1) where the difference in R_{\perp} (R_{\parallel}) with \mathbf{E} -field direction gave the trend $R_{\perp} > R_{\parallel}$.¹ Changes of the index of refraction of ZnS have been shown to be less than 10% in the spectral range of interest;⁹ as a result, no dispersion in the value of R_{\perp} (R_{\parallel}) was used in calculating the absorption coefficient. We note, however, that the R_{\perp} (R_{\parallel}) values chosen above gave a mismatch of the absorption coefficients for the thick and thin samples at the shorter wavelengths, possibly indicating some dispersion in R_{\perp} (R_{\parallel}). As a consequence, the sign of the polarization is uncertain for the highest-energy Mn absorption band found.

Figure 1 gives absorption spectra ranging from the Mn threshold to within ~ 0.2 – 0.3 eV of the band edge for $\mathbf{E} \parallel \mathbf{c}$ at 19, 83, and 293 K. The sharp rise in α beyond $E \sim 3.4 \text{ eV}$ at 293 K suggests the onset of band to band absorption, found at $\sim 3.6 \text{ eV}$ for $x \sim 0.15$ (Ref. 3). Expected thermal shifts of this edge to higher energy could not be studied because of low source strength in the near-uv region. As seen in the figure, five distinct bands emerge with decreasing temperature, with the magnitude and position independent of temperature. Figure 2 gives the absorption in the Mn region for both directions of \mathbf{E} field at 19 K. The assignments of the bands, as given in Fig. 2, are according to the theoretical results and discussion of Richardson and Janssen.¹⁰ We note that some disagreement has existed about the assignment of the fourth and fifth bands to b^4T_2 and b^4T_1 , respectively. Intensities of excitation to the nearly degenerate $a^4A_1(^4G)$ and $a^4E(^4G)$ levels estimated by Parrot and Blanchard¹¹ are such that only a^4E should contribute significantly to the third band.

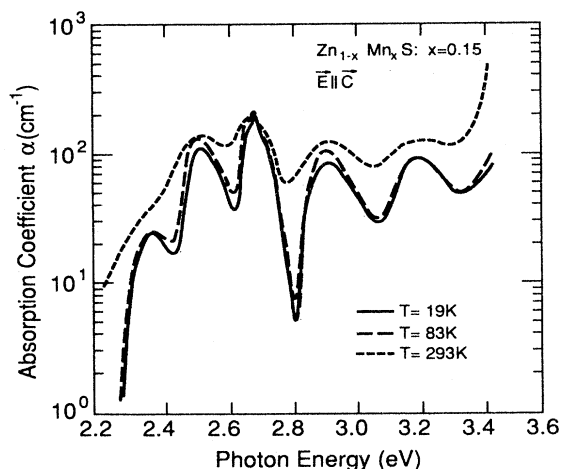


FIG. 1. Absorption spectra for $\mathbf{E} \parallel \mathbf{c}$ at 19, 83, and 293 K.

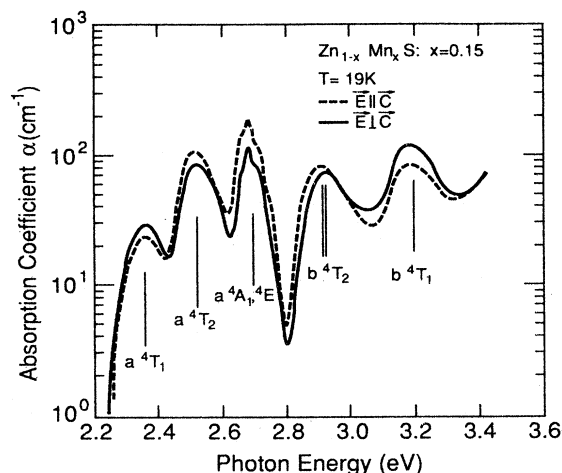


FIG. 2. Absorption spectra for $\mathbf{E} \parallel \mathbf{c}$ and $\mathbf{E} \perp \mathbf{c}$ at 19 K.

IV. DISCUSSION

Polarization features in the manganese absorption spectra reported here are similar to those observed in the wurtzite phase of $\text{Zn}_{1-x}\text{Mn}_x\text{Se}$.¹ For the a^4T_1 transition (at $\sim 2.35 \text{ eV}$), $\alpha(\mathbf{E} \perp \mathbf{c}) > \alpha(\mathbf{E} \parallel \mathbf{c})$, while in the a^4T_2 and $a^4A_1,^4E$ transitions (at ~ 2.5 and $\sim 2.68 \text{ eV}$, respectively) $\alpha(\mathbf{E} \parallel \mathbf{c}) > \alpha(\mathbf{E} \perp \mathbf{c})$. The b^4T_2 transition (at $\sim 2.9 \text{ eV}$) shows a small splitting with $\alpha_{\text{max}}(\mathbf{E} \parallel \mathbf{c})$ at lower energy ($\Delta E \sim 0.02 \text{ eV}$) than $\alpha_{\text{max}}(\mathbf{E} \perp \mathbf{c})$, similar to the behavior of the same transition in the selenide. In turn, the present result strengthens the earlier observation since band-edge overlap effects which may contribute to the b^4T_1 band shape in the selenide are not a problem in the sulfide.

Work by McClure and others^{12,13} shows that in DMS's, narrow features, i.e., zero phonon lines, phonon sidebands, pair spectra, etc., in optical absorption or excitation of Mn disappear for $x \geq 0.1$. We note residual structure in the $a^4A_1,^4E$ transition at 19 K, including a peak at $\sim 2.68 \text{ eV}$, and shoulders at ~ 2.56 and $\sim 2.71 \text{ eV}$; these are similar to features seen by McClure in $\text{Zn}_{0.95}\text{Mn}_{0.05}\text{S}$ (cubic) at 77 K.

The current results, together with the earlier reported work on $\text{Zn}_{1-x}\text{Mn}_x\text{Se}$, suggest that a given pattern for the polarization of absorption holds for the broad Mn d - d intraion transition in Zn-based wurtzite-phase DMS's. It is unlikely that still higher bands may be observed in the sulfide because of the known band-gap closure with increasing x .³ We note that of the DMS's with wurtzite structure, only $\text{Cd}_{1-x}\text{Mn}_x\text{S}$ has a sufficiently wide band gap to study the three lowest transitions in Cd-based systems.¹⁴

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