973 (1969).

<sup>1</sup>M. J. G. Lee, Phys. Rev. <u>178</u>, 953 (1969). <sup>2</sup>S. T. Inoue, S. Asano, and J. Yamashita, J. Phys. Soc. Japan (to be published).

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## ERRATA

Velocity Acquired by an Electron in a Finite Electric Field in a Polar Crystal, K. K. Thornber and Richard P. Feynman [Phys. Rev. B 1, 4099 (1970)]. In Eq. (18),  $|C_k|$  should be  $|C_k|^2$ ; in Eq. (22b),  $\cos(\frac{1}{2}s_0v\beta)$  should be  $\cos(\frac{1}{2}sv_0\beta)$ ; in Eq. (24),  $C = 4B/A\beta^2$  rather than as given. Equation (27) should read

$$R_{\vec{k}}^{\prime\prime} = -i \{ C_{\vec{k}}^{*} a_{\vec{k}}^{*} \vec{k} \cdot [e^{-i\vec{k}\cdot\vec{z}}, \vec{p}]_{*} - C_{\vec{k}} a_{\vec{k}}^{*} \vec{k} \cdot [e^{+i\vec{k}\cdot\vec{z}}, \vec{p}]_{*} \} / 2m .$$

The third from the last equality in Appendix B should read  $-4\pi i [G'_x(\nu) - G'_x(-\nu)]$ .

Ultrasonic Attenuation Due to Electron-Phonon Interaction in Potassium, T. M. Rice and L. J. Sham [Phys. Rev. B <u>1</u>, 4546 (1970)]. We have discovered a numerical error in the evaluation of the resistivities quoted in Table I and shown in Fig. 3. The correct values of the resistivity  $\rho$  are obtained by multiplying the quoted values by a factor of  $\frac{2}{3}$ . The values quoted for  $\tau_2/\tau_1$  are unchanged. We are grateful to Dr. J. Ekin of Cornell University for a correspondence which led to the discovery of the error.

Anomalous Electron-Phonon Transport Properties of Impure Metals. I. The Electrical Resistivity, M. J. Rice and O. Bunce [Phys. Rev. B 2, 3833 (1970)]. Dr. H. Smith has brought to our attention the following typographical errors in the published paper. The factors of  $(1 + \rho_{ep}^0 / \rho_0)$  occurring explicitly in Eqs. (3.7) and (3.8) should be absent.

On a more serious note we should like to call the attention of the reader of our paper to the recent paper by Dr. H. Smith, [Solid State Commun. 8, 1991 (1970)]. In this paper Dr. Smith has demonstrated that the use of our ansatz (2.15) considerably overestimates the deviation from additivity of the impurity and electron-phonon resistivities in the intermediate temperature range  $T \sim T_0$ . In fact, Dr. Smith argues that the peaks obtained by us in the deviations from additivity in the region  $T \sim T_0$  would be absent in a more rigorous calculation.

Interaction of Dislocations with Electrons and with Phonons, A. Hikata, R. A. Johnson, and C. Elbaum [Phys. Rev. B 2, 4856 (1970)]. The following corrections should be noted:

<sup>3</sup>A. Meyer and W. H. Young, Phys. Rev. Letters 23,

(1) On p. 4861, Eq. (8) should read  $K_2 = \Gamma 3\hbar / 2\pi^2 C^4$ .

(2) On p. 4861 first column, starting with line 37 note that  $\sigma = 1.2 \times 10^{-8}$  cm,  $K_2 = 7.2 \times 10^{-73}$  cgs,  $\Gamma = 2.7 \times 10^{-23}$  cm sec.

(3) On p. 4861, first column, the last line should read  $g(\beta_1)^{1/2} \sim 2.4 \times 10^{10}$  (cgs) and  $\gamma(\beta_2)^{1/2} \sim 0.1$ .

(4) On p. 4861, second column, first paragraph (top) should read:

In order for  $\gamma$  and g to have physically plausible values, these results imply that  $\beta_1$  and  $\beta_2$  have to be of the order of  $10^{-2}$  to  $10^{-3}$ ; we have not been able to evaluate  $\beta_1$  and  $\beta_2$  theoretically. Thus, in conclusion, a comparison of the present results with Leibfried's theory yields a physically plausible value of the scattering width  $\sigma$ . The scattering widths obtained from a comparison of the present results with the theories based on strain field scattering depend on the values chosen for  $\beta_1$  and  $\beta_2$ .

(5) On p. 4863, second column, the last four lines should read:

However, the phonon scattering width (cross section per unit length) of dislocations deduced from the present measurements, using Leibfried's approach, is quite close to the value expected from physical arguments.

Compton Scattering and Electron Momentum Density in Beryllium, R. Currat, P. D. DeCicco, and Roy Kaplow [Phys. Rev. B 3, 243 (1971)]. We regret that, in mentioning previous positron annihilation work which has been done, we inadvertently omitted reference to the work of S. Berko [Phys. Rev. 128, 2166 (1962)].

Far-Infrared Properties of Lattice Resonant Modes. V. Second-Order Stark Effect, B. P. Clayman, R. D. Kirby, and A. J. Sievers [Phys. Rev. B 3, 1351 (1971)]. A type-setting error was introduced during the galley-correction procedure and resulted