

COMMENTS AND ADDENDA

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Shubnikov-de Haas Effect and Magnetic Breakdown in Be[†]

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Magnetic breakdown in the basal plane of Be was investigated in fields up to 150 kG. Geometrical aspects of the open orbits caused by breakdown are discussed. An estimate of 30 kG for the characteristic breakdown field is obtained by applying the theory of Falicov *et al.* Shubnikov-de Haas frequency measurements are reported for the central-cigar cross section. These measurements are relevant to a recent conjecture that the Onsager-Lifshitz quantization scheme may be violated in the presence of magnetic breakdown.

The phenomenon of magnetic breakdown in metals has been investigated experimentally and theoretically, particularly in hexagonal-close-packed metals such as Zn and Mg.^{1,2} The purpose of the present work was to investigate this phenomenon in Be and to compare high-field-magnetoresistance measurements with the theory of Falicov *et al.*² Their calculation of the transverse magnetoresistance of Zn and Mg is applicable to Be because the geometry of the pieces of Fermi surface involved in the breakdown is similar in all three metals.^{2,3} The existence of breakdown in the basal plane in Be has been shown by Alekseevskii and Egorov,⁴ Reed,^{5,6} and Reed and Condon.⁷ However, these authors did not attempt a quantitative comparison of their results with theory. Another motivation for this work was to investigate further the conjecture of Alekseevskii and Egorov⁴ that the Onsager-Lifshitz quantization scheme may be violated, in the presence of magnetic breakdown, in such a way that different frequencies could be observed in the Shubnikov-de Haas (SdH) and de Haas-van Alphen (dHvA) effects.

The measurements were made at the Francis Bitter National Magnet Laboratory in fields up to 150 kG and at $T = 4.2^\circ\text{K}$. Measurements below 4.2°K were not performed because magnetic-interaction effects become predominant there as reported by Reed and Condon.⁷ The electronic techniques

and sample holder were those described previously.^{8,9} The sample, obtained from the Franklin Institute, was a single-crystal rod whose current density \vec{J} was parallel to $[11\bar{2}0]$; its residual-resistance ratio $[\rho(296^\circ\text{K})/\rho(4.2^\circ\text{K})]$ was 100.¹⁰

The effects of magnetic breakdown are demonstrated by the longitudinal-to-transverse magnetoresistance plot of Fig. 1. The field is in the $(1\bar{1}00)$ plane and open orbits induced by breakdown are shown by the dip centered on $[0001]$. The orbits are directed along $[1\bar{1}00]$ and exist for angles of \vec{B} out to $19.2^\circ \pm 0.8^\circ$ away from $[0001]$. This estimate, obtained from the width of the dip at half-amplitude, agrees with the value of 19.5° given by Reed.⁶ The geometry of the open orbit caused by breakdown between the coronet and cigars is shown in Fig. 2. This figure was constructed, approximately to scale, from the work of Tripp *et al.*³ The geometrical factors that control the thickness of the band of open orbits, as the field is tipped away from $[0001]$, are the size of the neck orbit on the coronet and the distance along the hexad axis over which breakdown occurs on the cigar. At a critical angle ψ , corresponding to $\sim 19.2^\circ$, the thickness of the open-orbit band approaches zero as shown in Fig. 2. The implication of this value is that magnetic breakdown occurs over a length $2l \approx 0.063$ a.u., which is approximately 7% of the length of the cigar. Since the energy gap increases

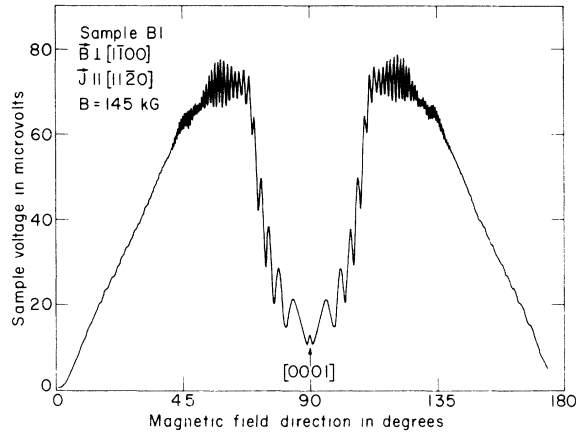


FIG. 1. Longitudinal-to-transverse magnetoresistance-rotation plot at 145 kG. The sample notation is that of Ref. 10.

with distance along the hexad axis, the breakdown length would be expected to increase with field above the minimum breakdown field. Thus the above breakdown length is that corresponding to a field of 145 kG.

Figure 1 shows also some rather large-amplitude quantum oscillations which are due to breakdown. These oscillations show up clearly in field sweeps for \vec{B} near $[0001]$, where their amplitude is $\sim 50\%$ of the dc-background magnetoresistance. The SdH frequency of the oscillations at $[0001]$ is 9.4×10^6

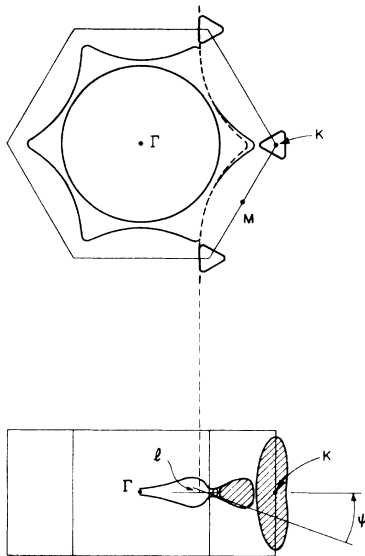


FIG. 2. Fermi-surface model of Be showing the second-zone coronet and third-zone cigars. The open orbit at the critical angle ψ is shown in dashed lines in the upper view. The model is based on the work of Tripp *et al.* (Ref. 3).

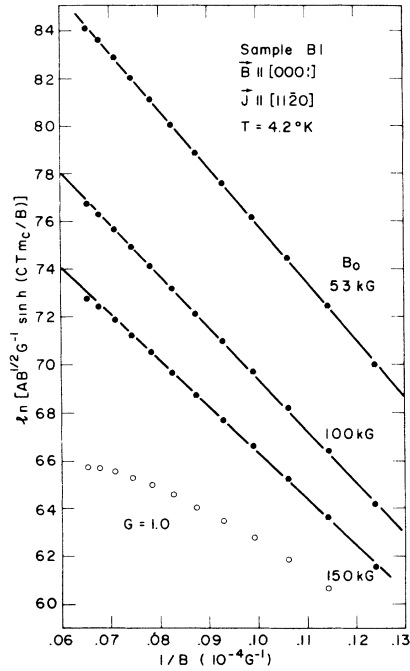


FIG. 3. Scattering-temperature plot showing the effect of breakdown through the probability G for completing a complete cigar orbit.

G , which corresponds to the central-cigar cross section. Figure 3 shows the amplitude of the oscillations, as a function of B^{-1} , on the usual scattering-temperature plot.¹¹ The circles are the experi-

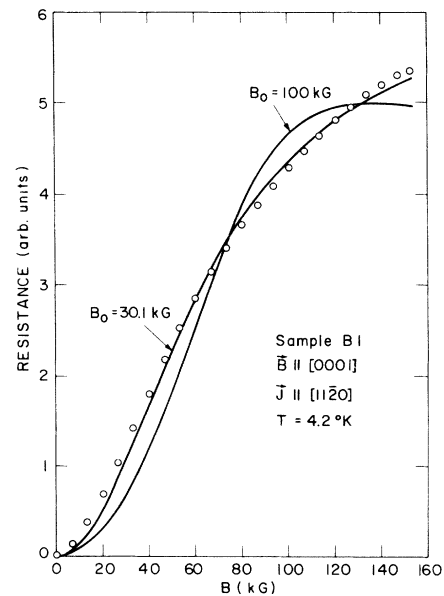


FIG. 4. Experimental data (open circles) and fits of the theory of Falicov *et al.* (Ref. 2) to the nonoscillatory magnetoresistance. Fits with two values of B_0 are shown.

mental points and the dots were obtained by correcting the data by a factor $G \equiv (1 - e^{-B_0/B})^{3/2}$, which is the probability that an electron *completes* a cyclotron orbit on the cigar without breakdown.²

Three values of the characteristic breakdown field B_0 were tried, all of order 10^5 G. These values were chosen because (a) the magnetoresistance deviates markedly from quadratic above ~ 50 kG, (b) Alekseevskii and Egorov estimated $B_0 \sim 150$ kG from magnetoresistance measurements,⁴ and (c) Tripp *et al.*³ estimated $B_0 = 120$ kG from their pseudopotential calculation. However, as shown in Fig. 3, the quality of the fit is very insensitive to the value of B_0 , so that this method appears to be a poor one for estimating this field. A similar behavior was found by Falicov *et al.*,² who fitted dHvA amplitude data in Zn essentially equally as well with B_0 values of 2.7 or 6 kG.

In view of this problem an attempt was made to fit the field dependence of the nonoscillatory part of the magnetoresistance, using the theory of Falicov *et al.*² This theory for Zn and Mg divided the conductivity into a classical part for the large orbits in the second zone plus a quantum-mechanical part responsible for the switching probabilities, at the junctions in the third zone, of the classical network. It is assumed that the impurity-scattering probability on the third-zone orbits is so small that it can be neglected in comparison with the probability for interband transitions due to magnetic breakdown. While this assumption is only of marginal validity in our sample, application of the theory to our sample can be justified to some extent since $(\omega_c \tau)_{\text{cigar}} \gg 1$ at the highest fields,¹⁰ and also because the amplitude of the oscillations ($\sim 0.5\rho_{\text{nonosc}}$) is large enough to indicate that the switching probabilities play a major role in the field dependence. The nonoscillatory part of the magnetoresistance is shown by the points in Fig. 4. Equation (3.7) of Falicov *et al.*² was used, along with¹

$$\rho_{11} = \sigma_{11}(\sigma_{11}^2 + \sigma_{12}^2)^{-1},$$

to fit the field dependence of the magnetoresistance. The full expression for ρ_{11} contained two parameters, in addition to the breakdown field B_0 to be determined. Various values of B_0 were chosen and a nonlinear-least-squares fit was performed to find the best values of the other parameters. The best fit of the data was obtained with $B_0 = 30.1$ kG as shown in Fig. 4. In order to check this fitting procedure for sensitivity to B_0 , the fit obtained with $B_0 = 100$ kG is also shown; it is clearly inferior to the $B_0 = 30.1$ kG fit. It is interesting to note that in the recent work of Reed and Condon⁷ on much "higher" purity samples than ours, the effects of magnetic breakdown are already quite obvious at $B \sim 20$ kG. Thus

our estimate of $B_0 \approx 30$ kG as well as this latter work suggests that B_0 for Be is considerably less than earlier estimates of ~ 150 kG.⁴ For a more detailed comparison with theory, measurements of ρ_{11} and ρ_{12} , on a sample of higher purity (e. g., that of Ref. 7), would be desirable.

A final point concerns the result of Alekseevskii and Egorov,⁴ who found that their SdH frequency results deviated from the dHvA results of Watts by $\sim 15\%$, which was more than the estimated experimental uncertainties. On this basis they speculated that: "It is not excluded that under the conditions of magnetic breakdown the Lifshitz-Onsager quantization scheme may be violated and the electron motion, strictly speaking, will not be described by the orbits observed in the dHvA effect." Figure 5 shows our SdH frequency data [$F(\theta)$] in the $(1\bar{1}00)$ plane, along with the dHvA data of Tripp *et al.*³ $F(\theta)$ was obtained from the data of Fig. 1, using the constant-field rotation method. Our results show no discrepancy with the dHvA results of Tripp *et al.*³ Except perhaps for systematic frequency-calibration or sample-orientation errors, it is difficult to understand the source of the disagreement in the earlier work. Since the frequency is controlled by geometrical factors alone, it is reasonable that the central-cigar frequency is the same in both the dHvA and SdH effects.

In summary, the open-orbit region at 145 kG was estimated as approximately 7% of the length of the third-zone cigars. This estimate relied heavily on the Fermi-surface model of Tripp *et al.*³ It was shown that the simple method of applying a correction factor for breakdown to the SdH amplitude data was a poor method for estimating the breakdown field. However, a fit of the nonoscillatory part of the magnetoresistance to the theory of Falicov *et al.*² gave $B_0 \approx 30$ kG. This value appears to be more

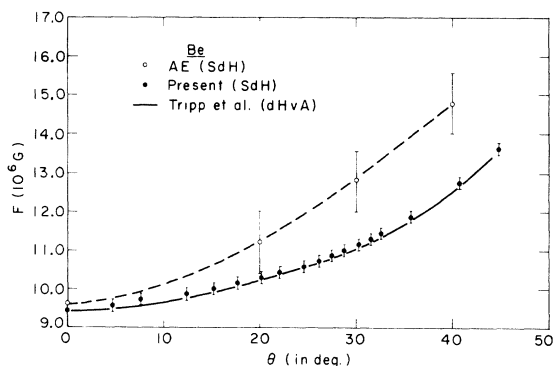


FIG. 5. Comparison of the SdH frequency data with those of Alekseevskii and Egorov (AE) (Ref. 4), and with the dHvA data of Tripp *et al.* (Ref. 3).

reasonable, when compared with the recent data of Reed and Condon,⁷ than earlier estimates of $B_0 \geq 100$ kG. Finally, the orientation dependence of

the SdH frequency was found to agree well with dHvA results, in contradiction to recent results in the literature.

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Comments on Pseudopotential Form Factors for White Tin

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While preparing for an experimental investigation of electron-phonon mass enhancement and relaxation times in white tin, we were confronted with large inconsistencies in the published form factors for this metal.¹⁻⁵ In this paper we are presenting the results of calculations that should clarify this matter a great deal.

The experiments on which our calculations are based are Fermi-surface calipers obtained from very accurate rf size-effect measurements.⁶ With a plane-wave matrix, large enough to assure convergence with respect to the number of reciprocal lattice vectors, we fitted the four form factors $V(K_n)$, $K_n < 2k_F$, and the Fermi energy (or mass) at 14 selected points, representative of the Fermi surface. The best fit is obtained for $E_F = 2.11$ (throughout this paper, we use $\hbar^2/2ma^2$ as unit of energy) [Fig. 1(a)]. Compared with the nearly-free-electron value $E_{NFE} = 2.312$, this suggests a Fermi-surface-band effective mass of about $0.92m_0$. For the best energy value the fit shows a rms deviation in the calipers of about $0.004k_0$ ($k_0 = \Gamma L = 2\pi/a$), which is slightly larger than the experimental inaccuracies. In Fig. 1(a) one sees that this Fermi energy can be determined to within 0.01 (our units).

Our fitted form factors show a large energy de-

pendence [Fig. 1(b)]. We are confident that these curves are representative of the entire Fermi surface, rather than of the specific sample of data points.

In the following we shall compare our results with previous calculations. The work of Weisz¹ was based on Gantmakher's pioneering rf size-effect data. Due to some uncertainties in the line shape the inaccuracy of Gantmakher's calipers amounts to several percent. Weisz found a best Fermi energy of about 2.24 (our units). For this energy we find a fit which has a rms deviation seven times larger than that of our best fit. We feel that Weisz's deviation from our results is due to the large limitation in accuracy and in the number of his data points.

Stafleu and de Vroomen² took E_F equal to the free-electron value and used a "few-plane-wave" matrix. In view of their aim to obtain a rough agreement with their de Haas-van Alphen (dHvA) experiments the discrepancies with our results are reasonable.

Craven's results³ are of more serious concern to us. He claims to describe to within 1% the accurate dHvA measurements of Craven and Stark⁷ by form factors, given in Fig. 1(b), at an energy $E_F = 2.34$ (our units). However, his interpolated