

Frequencies, eigenvectors, and single-crystal selection rules of $\mathbf{k}=0$ phonons in $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$: Theory and experiment

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We have determined the components of the Raman tensor of single-crystal superconducting $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ for the phonons at ~ 500 , 440, 330, and 150 cm^{-1} . They all correspond to the totally symmetric modes both in the tetragonal (A_{1g}) and orthorhombic (A_g) structures except for the 330-cm^{-1} mode which is B_{1g} tetragonal. We have also performed lattice dynamical calculations with parameters extracted from those of perovskite and other metallic oxides. The frequencies and eigenvectors obtained agree well with the experimental results. This work settles discrepancies found in the previous literature.

Many reports on vibrations of high- T_c superconductors have appeared.¹⁻⁵ The disagreement between assignment of phonons by various authors have spurred measurements on single crystals.⁶⁻⁹ In one case⁹ incorrect experimental selection rules were obtained. Little theoretical work on phonon frequencies and eigenvectors has been published so far.^{10,11} While in Ref. 10 the semiconducting O_6 compound was treated, Ref. 11 has adjusted the force constants to fit the experimental data. Here we present theoretical results from parameters not adjusted but taken over from known oxide perovskite and other metallic oxide potentials. We also present an unambiguous experimental determination of the Raman selection rules for single crystals and the details of the Raman tensors. The results agree with the theoretical predictions. This fact enables us to propose an assignment of the ir phonons observed in ceramic samples.

Single crystals of $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ were grown from a nonstoichiometric melt by slow cooling following Ref. 12. Tetragonal platelets (maximum dimensions $2.5 \times 2.5 \times 0.08\text{ mm}$) formed and were separated from the flux; their composition was checked with an EDAX microprobe. With a single-crystal diffractometer we found the crystal studied here to be oriented with the c axis perpendicular to the large surface of the platelet. The lattice constants were $a = 3.869\text{ Å}$, $b = 3.872\text{ Å}$, and $c = 11.71\text{ Å}$, values typical for unannealed crystals.¹² These values for a , b , and c indicate δ values slightly smaller than that of the tetragonal phase ($\delta \approx 0.5$).¹³ By measuring the shielding in a weak magnetic field, T_c (onset) was found to be 40 K (at 10% of maximum shielding). The complete shielding effect was 77% while the Meissner effect was found to be 6%. The details of the Raman setup are given in Ref. 2.

The potentials for the lattice dynamical calculation were taken from least-squares fits to BaTiO_3 ,¹⁴ SrO ,¹⁵

and NiO ,¹⁶ with the only modification that the polarizability in the z direction of the Cu-O bonds was set equal to their longitudinal polarizability. This change is justified by the structural differences between the perovskites and the superconductor, as shown in the $\delta \approx 1$ case.¹⁰

Let us first consider the case where the wave vector of the incident light \mathbf{k}_i is parallel to the scattered one \mathbf{k}_s and to the c axis of the crystal (curves 1-6 of Fig. 1). In curve 1 both the (linear) polarization of the incident light ($\hat{\mathbf{e}}_i$) and the orientation of the (linear) analyzer ($\hat{\mathbf{e}}_s$) are along the x axis. (No differences were found between the x and y axes since the crystal is nearly tetragonal. Even for an orthorhombic crystal no difference would be expected because of twinning.¹²) Curve 2 shows no detectable depolarized scattering to within noise. From the absence of all phonon features under cross polarization we conclude that all observed peaks have A_g symmetry in the orthorhombic structure, i.e., $a_{ij} = 0$ for $i \neq j$ where a_{ij} are the components of the Raman tensor.

In curve 3 we have rotated $\hat{\mathbf{e}}_i$ and $\hat{\mathbf{e}}_s$ by 45° so that they are both parallel to either $[110]$ or $[\bar{1}\bar{1}0]$. We find the same spectrum for orientations along either direction and thus display only one of them. In the spectrum taken under crossed polarization ($\hat{\mathbf{e}}_i \perp \hat{\mathbf{e}}_s \parallel [110]$, curve 4) we see that the peak at $\sim 330\text{ cm}^{-1}$ has recovered the strength of that in curve 1, while the other peaks continue to be weak or absent. From this observation we conclude that the Raman polarizability must have components $a_{xx} = -a_{yy}$. Furthermore, from the absence of the peak at 330 cm^{-1} in curve 8, taken for $\hat{\mathbf{e}}_i \parallel \hat{\mathbf{e}}_s \parallel z$, we infer that $a_{zz} \approx 0$. We have thus specified all Raman tensor components of this mode. They are consistent with the measurement with circularly polarized light (curves 5 and 6). The absence of the peak at 330 cm^{-1} for equal angular momenta of incident and scattered light (in an absolute

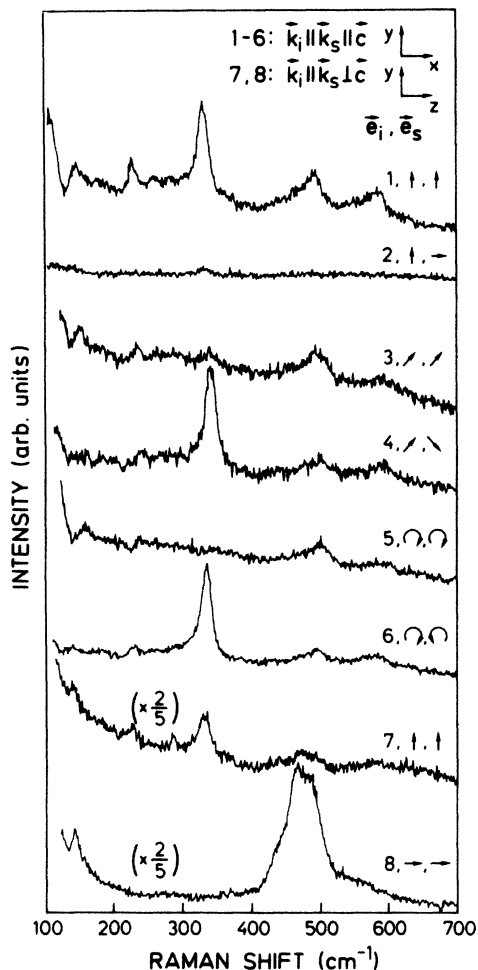


FIG. 1. Polarization-dependent Raman spectra of single-crystalline $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$. All curves were obtained in the back-scattering geometry ($\mathbf{k}_i \parallel \mathbf{k}_s$); the orientations of the polarization vectors with respect to the crystal axes are indicated in the figure.

frame; opposite if we take the direction of propagation $\mathbf{k}_{i,s}$ as the quantization axis) shows that $\alpha_{xx} = -\alpha_{yy}$ to good accuracy.

The connection between this peak and the superconducting mechanism of these compounds has been shown to lie in an anomalous temperature dependence.^{17,18} When the ceramic samples of $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ are cooled below the superconducting transition temperature this peak lowers its frequency (softens), whereas all other Raman peaks are observed to stiffen. In the nonsuperconducting compound $\text{YBa}_2\text{Cu}_3\text{O}_6$ where this mode is also present (with tetragonal B_{1g} symmetry), no anomalous temperature dependence has been found.¹⁰ Closest in frequency in our lattice dynamical calculation for $\delta=0$ lies an O(II)-O(III) (z) out-of-phase mode (347 cm^{-1} , Fig. 2) with the O(II)-O(III) motion required for the Raman tensor; we thus assign it to the experimental peak at 330 cm^{-1} . There is hence an anomaly related to superconductivity for vibrations in the so-called "planes" of $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$. We have also searched for this anomaly

in the single crystals studied here, but found it to be weaker than in high-oxygen-content ceramic samples. The reason for that is found in curve 8 of Fig. 1, where $\mathbf{k}_i \parallel \mathbf{k}_s \perp c$ axis. The broad, large peak centered at $\sim 475 \text{ cm}^{-1}$ corresponds⁴ to a distribution of O contents between 6.0 and 6.6. A distribution like this is characteristic for those of our single crystals which were not oxygen treated after preparation. We also note that the peak at 330 cm^{-1} has resonant behavior;¹⁹ when the spectra in Fig. 1 were taken with the 4880-\AA Ar^+ laser line, this peak decreased by a factor of 5–10.

A second peak associated with the planes in the superconductor is usually observed in ceramic materials around 435 cm^{-1} .^{2,6,20} In Fig. 1 it is only seen in curve 8 as a shoulder. Its Raman tensor must hence be of the form $\alpha_{xx} = \alpha_{yy} = 0$; only α_{zz} is nonzero. This peak corresponds to the O(II)-O(III) in-phase vibration along z (373 cm^{-1} , Fig. 2).

At low frequencies in Fig. 1 we find a peak at 145 cm^{-1} in curves 1, 3, 5, and 7 with roughly equal amplitude and in curve 8 with 1.5 times that amplitude. The peak is absent in curves 2, 4, and 6. The opposite selection rule of this peak compared to the peak at 330 cm^{-1} for circular polarizations implies $\alpha_{xx} = \alpha_{yy}$, and the ratio of the peak amplitudes for z^2 and x^2 geometry (curve 8) sets the magnitude of the α_{zz} component to $|\alpha_{zz}| \cong 1.2 |\alpha_{xx}|$. In our lattice dynamical calculation this peak corresponds to the Cu(2) vibration (159 cm^{-1} , Fig. 2). The Ba (z) mode below that (132 cm^{-1} in Fig. 2) is not clearly seen in Fig. 1, but it has been reported⁶ to lie near 116 cm^{-1} . The peak with the largest amplitude in (single-phase) ceramic material^{2,6,20} and also in the single-crystal data of Fig. 1 is centered here at 475 cm^{-1} . It appears only in curve 8 (except through leakage in some of the other curves), hence it also has the Raman tensor components $\alpha_{xx} = \alpha_{yy} \cong 0$. This mode has been associated, rather conclusively and by many groups,^{1,2,5,6,8,10} with the O(IV)-Cu(1) A_g mode calculated at 515 cm^{-1} (Fig. 1).

There has been controversy²¹ about Raman features with frequencies $> 550 \text{ cm}^{-1}$. The selection rules of Fig. 1 can help identify impurity phases that do not have the same symmetry as the superconducting crystal or are randomly oriented. We conclude from arguments such as those used above that the Raman tensors of the broad peak at $\sim 580 \text{ cm}^{-1}$ and the feature at 220 cm^{-1} are not compatible with the crystal symmetry. While diagonal, according to the data for the polarizations parallel to the crystal axes, the data for circular polarization and those for polarizations parallel to the base diagonals ($x+y$, $x-y$) require that the α_{xx} and α_{yy} components be *incoherent*. The data for polarizations along z show that $\alpha_{zz} \cong 0$. Incoherent components suggest defect-induced scattering. Thus one possible assignment of the 220- and 580-cm^{-1} modes is to the x - and y -polarized TO modes calculated to be at 130 and 576 cm^{-1} . O(I) defects, which exist to a great deal for $\delta \cong 0.5$, may induce the Raman scattering. The defect-induced nature and the twinning would cause the incoherence of α_{xx} and α_{yy} .

Having discussed all the observed Raman peaks we now propose an assignment of the corresponding ir Davydov partners. Detailed far-infrared data on single crystals are

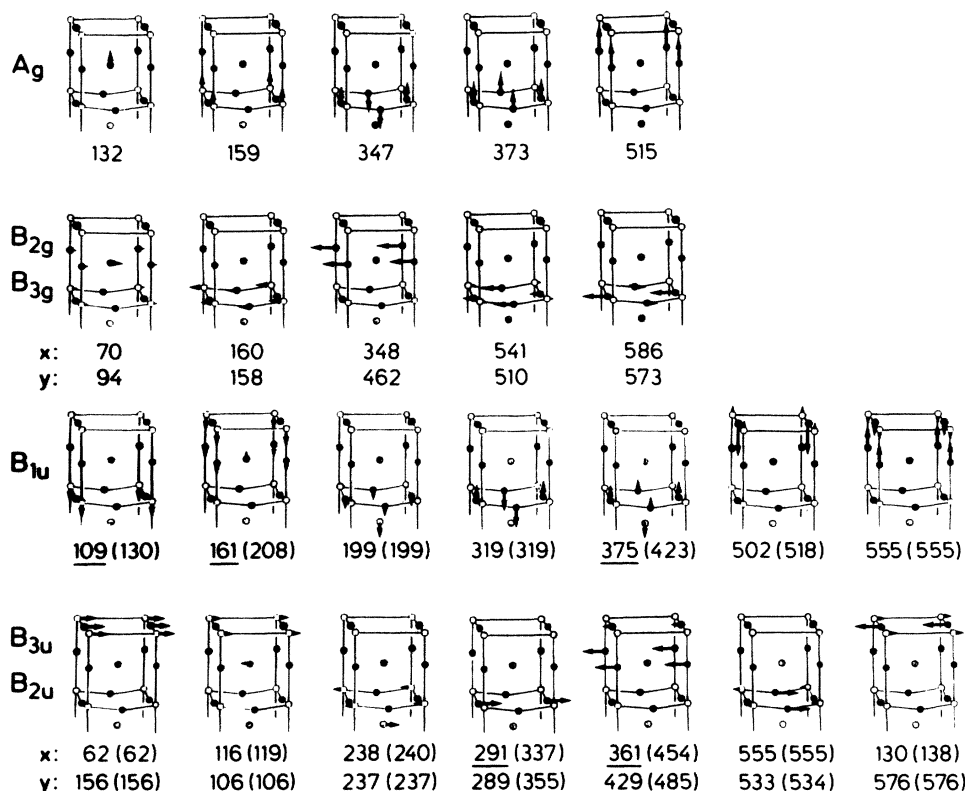


FIG. 2. Eigenvectors of all 36 optical modes of $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ in the upper half on the unit cell. The displacements for the lower part can be simply obtained from the upper ones. The frequencies are below each mode; for the odd modes the TO (LO) frequencies are given outside (inside) parentheses. (Filled circles—O; open circles—Cu; hatched circles—Ba; dotted circles—Y.) In the text, O(II), O(III), and Cu(2) refer to atoms in the Cu-O planes adjacent to Y; O(I) and Cu(1) are atoms in the chains of the superconductor.

not available so we turn to measurements on ceramic materials.^{2,18} Here we find peaks at 115, 150, 193, 275, 320, and 575 cm^{-1} . The three lowest frequencies have been suggested to belong to Ba and Y vibrations; the latter assignment is evidenced by the shift in frequency when Y is replaced by a rare earth.² We suggest that the weak peak at 115 cm^{-1} corresponds to the mode at 109 cm^{-1} (Fig. 2), while 150 and 193 cm^{-1} are the Ba (B_{1u}) frequency and the Y (B_{1u}), calculated at 161 and 199 cm^{-1} , respectively. The next two peaks in the doublet at 275 and 320 cm^{-1} , also possess an anomalous temperature dependence,^{18,22} and we find five likely frequencies in Fig. 2: the B_{3u} and B_{2u} in-plane movements (290, 361, and 429 cm^{-1}), and two B_{1u} out-of-plane vibrations (319 and 375 cm^{-1}). We assign the 290-cm^{-1} in-plane vibration to the experimental peak at 275 cm^{-1} . In the tetragonal material, the mode at 319 cm^{-1} (Fig. 2) is exactly silent since it has no net dipole moment: it is indeed not seen in ir measurements of $\text{YBa}_2\text{Cu}_3\text{O}_6$,¹⁰ and is probably too small to be detected in ir measurements of $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$. The calculated eigenmode at 375 cm^{-1} (Fig. 2) shows a large shift upon the replacement of Y by a rare earth which is not observed in the experiment. Thus, we assign the peak at 320 cm^{-1} to the remaining mode (B_{2u} and B_{3u} symmetries; 361 and 429 cm^{-1}) which has a large x - y anisotropy. The two frequencies may be smeared out in polycrystalline materi-

al and not observed separately in spite of their rather large TO-LO splitting. The peak at 575 cm^{-1} could correspond to the bond stretching O(I) vibration which has already appeared, defect induced, in the Raman spectra. A definitive assignment, however, must come from single-crystal ir measurements.

In conclusion, we have determined the Raman tensor elements of four of the phonons of superconducting $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ by polarization-dependent Raman scattering on single crystals. We have also presented a calculation of the eigenvectors and frequencies of all phonons obtained in a lattice dynamical calculation. The agreement between theory and experiment is good. The phonons with the temperature anomaly observed below T_c have been shown to be related to the CuO planes in the superconductor.

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