

Oxygen-diffusion kinetics in densified, amorphous SiO₂

R. A. B. Devine

Centre National d'Etudes des Télécommunications, Boîte Postale 98, 38243 Meylan, France

J. J. Capponi

Laboratoire de Cristallographie, Centre National de la Recherche Scientifique Boîte Postale 166X, 38041 Grenoble Cedex, France

J. Arndt

Mineralogisch-Petrographisches Institut, Universität Tübingen, D-7400 Tübingen 1, Federal Republic of Germany

(Received 8 October 1986)

The dynamics of intrinsic defect creation and annealing have been studied in dry, amorphous SiO₂ subjected to a hydrostatic pressure of 5 GPa at 600°C. The compaction produced does not in itself result in the creation of significant numbers of paramagnetic defects but it does greatly enhance their creation efficiency when samples are subjected to subsequent γ irradiation. Annealing after irradiation results in a growth in the number of peroxy-radical defects and a correlated decrease in the number of oxygen-vacancy defects. The higher temperatures needed for annealing with respect to the case of undensified, amorphous SiO₂ lead us to conclude that interstitial molecular-oxygen diffusion is severely inhibited in the compacted samples.

It is well known that under certain conditions of temperature and pressure a densification of amorphous silica (*a*-SiO₂) occurs which remains when the compressed sample is returned to ambient conditions. Such a permanent densification may be induced at room temperature using hydrostatic pressures ~ 17 GPa,¹ whereas at higher temperatures the required pressures are lower.² For example, a pressure of 4 GPa at 600°C is adequate to produce a densification $\sim 16\%$. The relative ease with which *a*-SiO₂ can be induced to densify may be related to the flexibility of the structure resulting from the rather small variations in bond energy required to cause significant Si—O—Si bond angle variations³ (0.17 eV/molecule for a $\pm 25^\circ$ angle variation about the mean of 144°). Alternatively, the densification might arise through the creation of large numbers of broken bond defects.¹ Detailed physical information on the densified material is, however, lacking. On a macroscopic scale, measurements² indicate that the refractive index and density of *a*-SiO₂ follow the simple Lorentz-Lorenz relationship⁴ with a curve which fits almost identically that found² for crystalline SiO₂ polymorphs. Although macroscopic, this result suggests that for the greater part,⁵ refractive index variations arise from volume changes in the network rather than bond-length—induced charge transfer, etc.

In interpreting the Brillouin frequency shift and Raman spectra of room-temperature compacted *a*-SiO₂, it has been suggested¹ that the compaction may proceed either by a process of collapse of large network ring structures⁶ into smaller ones, or by a process of bond breaking leaving nonbridging oxygen defects.⁷ In uncompactd *a*-SiO₂ three intrinsic paramagnetic structural defects have been identified and well characterized both experimentally⁸ and theoretically.⁹ They are the oxygen vacancy or E'_1 center ($\equiv \text{Si}^+ \cdot \text{Si} \equiv$), the nonbridging oxygen-hole center

(NBOHC) ($\equiv \text{Si}-\text{O}^\cdot$) and the peroxy radical ($\equiv \text{Si}-\text{O}-\text{O}^\cdot$), the latter being suggested¹⁰ to be the natural partner of the oxygen vacancy center. The creation efficiency, annealing dynamics, and symmetry of these defects have revealed significant information on the basic SiO₄ tetrahedron and the environment around the defect and upon various physical properties of the network such as gaseous diffusion.^{11,12} The present work presents the results of a study of these same defects in compacted *a*-SiO₂, undertaken with the aim of exploring the microscopic properties of the compressed phase.

Samples of Heraeus Suprasil W1 (dry silica, < 10 ppm OH) were cut from commercial 4-mm diam rod in 6-mm lengths and compressed in a belt apparatus¹³ to a pressure of 5 GPa at a temperature of 600°C for 20 min. Al₂O₃ was used as the pressure-transmitting medium. After crushing, samples were irradiated with 115 Mrad of ⁶⁰Co γ rays at room temperature. A series of compacted samples were then given 10-min anneals in flowing Ar gas at temperatures up to 525°C. Uncompactd bulk samples were given identical irradiation and annealing treatment for comparative purposes. Electron-spin resonance (ESR) measurements were carried out using a Bruker ER 200 D spectrometer operating in the region of 9.23 GHz. Oxygen-vacancy centers were studied at room temperature while peroxy-radical and NBOHC resonances were studied at -120°C . All measurements were made using non-saturating microwave power levels and nondistorting field modulation levels.

The conditions of time, temperature, and pressure used to compress our samples lead² to a densification $\sim 16\%$. In "as-compacted" samples (nonirradiated) we found no evidence for the production of measurable numbers of any of the three intrinsic, paramagnetic defects. After subsequent irradiation with 115 Mrad of ⁶⁰Co γ rays, defects

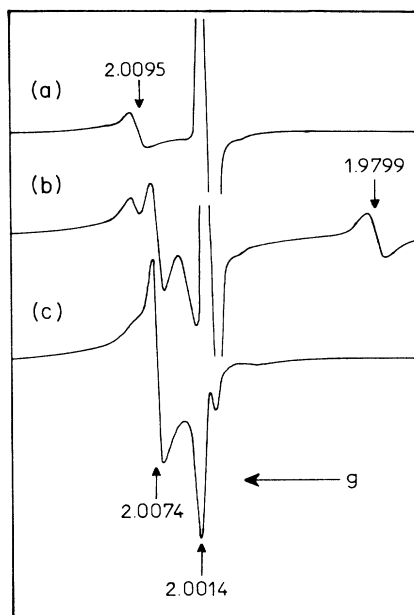


FIG. 1. Observed experimental ESR spectra at -120°C in samples of Suprasil W1 after irradiation with 115 Mrads of ^{60}Co γ rays: (a) sample compacted at 600°C under a pressure of 5 GPa, (b) uncompact sample, and (c) the compacted sample shown in (a) having been annealed for 10 min at 525°C in Ar post irradiation. The resonance at $g = 1.9799$ is due to a Cr in MgO marker. The large resonance shown off scale in (a) and (b) is the E'_1 resonance.

were revealed as demonstrated in the typical resonance spectrum taken at -120°C and shown in Figs. 1(a) and 1(b). The intense central resonance is the E'_1 vacancy center while the curve to the left centered upon a g factor ~ 2.0095 is part of the NBOHC resonance⁸ [Fig. 1(a)], the remaining downward peak is masked by the large E'_1 resonance. Figure 1(b) shows the result obtained at -120°C for an identically irradiated, uncompact sample of Suprasil W1. Noticeably, the peaks associated with the peroxy radical resonance ($g = 2.0074$ and 2.0014) now present, were absent in the irradiated, compacted sample [Fig. 1(a)]. Finally, in Fig. 1(c) we show the resonance spectrum for an irradiated, compacted sample after subsequent anneal at 525°C in flowing Ar gas for 10 min. We note that the E'_1 resonance has almost completely disappeared and that the previously absent peroxy resonance has now appeared.

Prior to the annealing studies we estimated the comparative yields of the different defects in dry $\alpha\text{-SiO}_2$ as a function of compaction. These results are shown in Table I. For the E'_1 resonance, minor differences in line shape were found between the compacted and uncompact states so that numerical integration was performed in order to extract reasonably accurate defect density ratios. Analysis of the line shapes using computer fitting suggested that the differences arose from small variations in the

TABLE I. Relative efficiency of creation of the three intrinsic paramagnetic defects by ^{60}Co γ rays in dry Suprasil W1.

Defect	Ratio of concentrations (compact to uncompact)
E'_1	110
NBOHC	> 17
Peroxy radical	> 43 ^a

^aAfter irradiation plus 10 min at 525°C .

linewidth and g factor distribution widths;¹¹ they disappeared upon annealing suggesting that they were induced by strain and defect density effects. The differences were not consistent with those expected if other forms of the basic E'_1 defect were present. For the other resonances (peroxy and NBOHC) the either incomplete nature of the observed spectrum [e.g., Fig. 1(a)] or the superposition of spectra lead to inherent inaccuracy in determining the defect density ratios. In all cases, however, we note that the three defects are much more readily created by radiation in the compacted $\alpha\text{-SiO}_2$ than in the uncompact. Most striking is the result for the oxygen vacancy center where the process is 2 orders of magnitude more efficient. Note that the peroxy comparison required studies on annealed samples since this resonance was not observed in the irradiated "as-compact" case. The figure quoted for the ratio of compacted to uncompact defect densities was therefore that obtained after annealing.

The annealing dynamics of E'_1 centers and peroxy radicals in irradiated compacted Suprasil W1 were measured following 10-min anneals; the results are presented in Figs. 2(a) and 2(b). The solid curve in Fig. 2(a) shows the annealing of E'_1 defects in irradiated, uncompact Suprasil W1 while the solid triangles show the results obtained for the compacted samples. For simplicity we have normalized our signal amplitudes to unity (corrected for line shape variation) corresponding to 4.7×10^{17} E'_1 defects per cm^3 for the compacted samples and 4.3×10^{15} E'_1 defects per cm^3 for uncompact samples. The precision in the absolute value determination is $\lesssim 50\%$. The peroxy-radical signal amplitude in compacted Suprasil W1 is shown by the solid circles in Figs. 2(b) while uncompact results, taken from Ref. 14, are shown by the solid stars.

It has been suggested¹¹ that peroxy radicals might form by the trapping of diffusing, molecular oxygen at E'_1 centers, thereby annihilating them. Results previously obtained in uncompact Suprasil W1 support this tendency, but we believe that the results on compacted Suprasil W1 shown in Figs. 2(a) and 2(b) confirm this hypothesis beyond reasonable doubt. Furthermore, when one estimates the peroxy-radical density at saturation in the compacted samples one obtains the figure of 3.3×10^{17} peroxy radicals per cm^3 as compared to an E'_1 density prior to annealing of 4.7×10^{17} per cm^3 . Given the anticipated inaccuracy in absolute-value determination previously mentioned, we consider the agreement between these two numbers to be better than satisfactory. The annealing re-

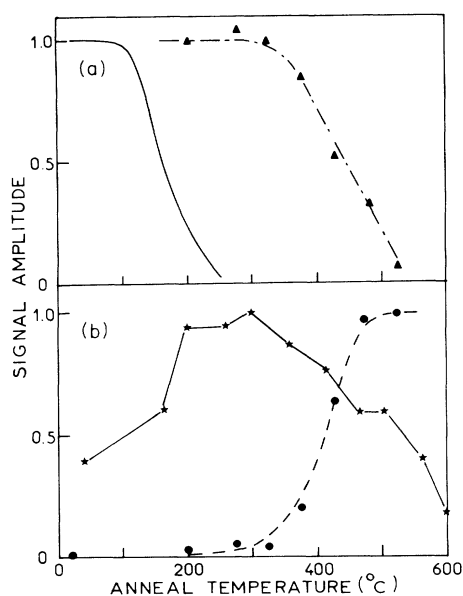


FIG. 2. Results of isochronal annealing in compacted and uncompact a -SiO₂ samples: (a) solid line, E_1 defects in uncompact Suprasil W1; solid triangles, E_1 defects in Suprasil W1 compacted by 16% and (b) solid stars peroxy-radical defects in uncompact Suprasil W1 (Ref. 14); solid circles, peroxy-radical defects in Suprasil W1 compacted by 16%.

sults [Figs. 2(a) and 2(b)] and the results for the initial defect creation by room-temperature irradiation then confirm the hypothesis that peroxy-radical defects form by the transformation of E_1 defects by the trapping of diffusing O₂. The results further suggest, by comparison with the behavior observed in uncompact Suprasil W1, that the diffusion of O₂ in compacted silica is severely inhibited as compared to the uncompact case. Within the framework of the Waite model¹⁵ for correlated defect or interstitial annealing one can demonstrate that the results shown in Fig. 2(a) suggest that the diffusion coefficient of O₂ in our compacted silica is between 10⁷ and 10⁹ times smaller at 450°C than in uncompact silica. Assuming a constant diffusion coefficient preexponential factor this suggests that the activation energy in compacted silica must increase from the uncompact value¹⁶ of 1.17 eV to 2 eV < E_A < 2.4 eV. Such a reduction in diffusion related

to compaction is consistent with a reduction found in crystalline polymorphs. In α quartz,¹⁷ whose density is approximately that of our compacted samples, the diffusion coefficients at 600°C is between 10⁴ and 10⁷ times less than that in uncompact a -SiO₂.

The restriction in the ease with which gas molecules diffuse through compacted silica is consistent with the picture of a network composed of a distribution of interstitial sizes,¹⁸ themselves resulting from a multiple ring structure in the amorphous lattice.^{6,19} Compression of the structure should lead to a reduced "pore" size (via reduction in the mean Si—O—Si bond angle) which would certainly account for a decreased diffusion. Such a modification may indeed have been confirmed by recent Raman spectroscopy studies²⁰ on a -SiO₂ at high pressure which have been interpreted as suggesting that both the width of the bond angle distribution and its mean value decrease under applied pressure.

The present results clearly demonstrate the validity of the anticipated E_1 —peroxy-radical interrelationship and underline the important role of oxygen diffusion in the creation of peroxy radicals even at room temperature. These results undermine the idea that the peroxy radicals originate from intrinsic peroxy bridges (Si—O—O—Si) which open under irradiation. The effect of limited O₂ diffusion in densified a -SiO₂ on the creation of peroxy radicals may hold the key to the lack of observation of this defect in crystalline quartz, and this suggests that a search for the defect in damaged quartz after high-temperature annealing might be instructive.

Finally, it should be added that the observation of enhanced radiation sensitivity in compacted a -SiO₂ may be relevant in the area of device technology. Various authors^{21,22} have identified the presence of radiation-induced oxygen-vacancy centers in the Si-SiO₂ interfacial region, and their importance in the degradation of devices has been underlined.²² The possibility that oxides of Si may grow with large residual stress²³ indicates that they may be significantly more radiation sensitive than unstressed forms so that care must be exercised to minimize the appearance of possible sources of stress.

ACKNOWLEDGMENT

The Laboratoire de Cristallographie is associe au Centre National de la Recherche Scientifique.

¹M. Grimsditch, Phys. Rev. Lett. **52**, 2379 (1984).

²J. Arndt and D. Stöfler, Phys. Chem. Glasses **10**, 117 (1969).

³A. K. Revesz and G. V. Gibbs, in *The Physics of MOS Insulators*, edited by G. Lucovsky, S. T. Pantelides, and F. L. Galeener (Pergamon, New York, 1980), p. 92.

⁴M. Born and E. Wolf, *Principles of Optics: Electromagnetic Theory of Propagation, Interference and Diffraction of Light*, 6th ed. (Pergamon, New York, 1983), p. 88.

⁵Although it has been argued that bond polarizability plays an important role, measured values suggest that its variation with density accounts for only 20% of the total refractive in-

dex variation with pressure; see A. K. Revesz, Phys. Rev. Lett. **27**, 1578 (1971).

⁶R. J. Bell and P. Dean, Philos. Mag. **25**, 1381 (1972).

⁷J. C. Phillips, Solid State Phys. **37**, 93 (1982).

⁸D. L. Griscom, J. Non-Cryst. Solids **31**, 241 (1978).

⁹A. H. Edwards, MRS Research Symposium Proceedings, Vol. **61**, 1982.

¹⁰A. H. Edwards and W. Beall Fowler, Phys. Rev. B **26**, 6649 (1982).

¹¹D. L. Griscom, *Proceedings of the International Symposium on Structure and Bonding in Non-Crystalline Solids*, Reston,

- 1982, edited by G. E. Walrafen and A. K. Revesz (Plenum, New York, 1986).
- ¹²R. A. B. Devine, Nucl. Instrum. Methods B **1**, 378 (1984).
- ¹³J. J. Capponi, Ph.D. thesis, Université Scientifique et Médicale, Grenoble, 1973 (unpublished).
- ¹⁴M. Stapelbroek, D. L. Griscom, E. J. Friebele, and G. H. Sigel, J. Non-Cryst. Solids **32**, 313 (1979).
- ¹⁵T. R. Waite, Phys. Rev. **107**, 463 (1957).
- ¹⁶G. Hetherington and K. H. Jack, Phys. Chem. Glasses **5**, 147 (1964).
- ¹⁷B. J. Giletti and R. A. Yund, J. Geophys. Res. **89**, 4039 (1984).
- ¹⁸J. F. Shackelford and J. S. Masaryk, J. Non-Cryst. Solids **30**, 127 (1978).
- ¹⁹F. Galeener, Solid State Commun. **44**, 1037 (1982).
- ²⁰R. J. Hemley, H. K. Mao, P. M. Bell, and B. O. Mysen, Phys. Rev. Lett. **57**, 747 (1986).
- ²¹P. M. Lenahan and P. V. Dressendorfer, J. Appl. Phys. **55**, 3495 (1984).
- ²²F. J. Grunthaner, B. F. Lewis, J. Maserjian, and A. Madhukar, J. Vac. Sci. Technol. **20**, 747 (1982).
- ²³C. Camelin, G. Demazeau, A. Straboni, and J. L. Buevoz, Appl. Phys. Lett. **48**, 1211 (1986).