Mobility of Pd and Si in Pd₂Si

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A diffusion-membrane technique is developed and used to determine the mobilities of Pd and Si in Pd₂Si. A multilayer film structure Pd/Pd₂Si/Ge, in which Ge serves as a sink for Pd, is used to monitor the flux of Pd through Pd₂Si. The flux of Si through Pd₂Si is determined by using the thin-film structure Si/Pd₂Si/Cr in which Cr serves as a sink for Si. Values for the mobility of Pd and Si are calculated from the measured flux rates. Good agreement is found between the calculated Pd mobility values and the parabolic rate constant for Pd₂Si growth around 200 °C. The mobility of Si at these temperatures is found to be orders of magnitude lower. The activation energy for Si mobility is found to be 2.1 eV in epitaxial Pd₂Si and 1.7 eV in polycrystalline Pd₂Si.

I. INTRODUCTION

The wide applicability of near-noble-metals and their silicides in microelectronic circuitry has stimulated extensive research into the properties of these materials. A good knowledge and thorough understanding of the silicon/metal interaction are essential for successful utilization of silicon/metal and silicon/silicide contacts. Aluminum, which is commonly used for metallization, may penetrate the metal or silicide layer, causing pitting of the contact. Refractory metals such as Cr, Ti, or W have been recognized to serve as sacrificial diffusion barriers to Al when introduced between the near-noble-metal and the Al layer.^{1,2} The growth kinetics of the silicon/metal and aluminum/metal compounds during interdiffusion of such multilayer films determine the usefulness of such metallization schemes. The Si (and metal) transport properties of the various layers are therefore important parameters necessary to evaluate the applicability and reliability of these metallization schemes.

Although the silicide formation of Pd, one of the nearnoble-metals, has received considerable attention, conclusions regarding the dominant diffusing species during silicide formation are inconsistent.³⁻⁸ The variety of experiments from which conclusions were drawn was carried out under various experimental conditions using different techniques. The results tend to suggest that both Pd and Si diffuse during Pd₂Si growth from Si/Pd couples. The flux ratios of Pd and Si differ from experiment to experiment.

The general (empirical) rule is that when the phase that forms is rich in the metal (or Si) the richer element is also the dominant diffusing species. Both Ni and Pt have been found to be the dominant moving species during Ni₂Si and Pt₂Si formation, respectively.⁶ It is therefore surprising that Pd₂Si should deviate from the empirical rule since Pd, Ni, and Pt all belong to the same group of nearnoble-metals. Further investigation into the mobility of Pd and Si in Pd₂Si is therefore appropriate.

We have developed a diffusion-membrane technique using multilayer films to determine the mobility of Pd and Si in Pd₂Si. A comparison of these values with the parabolic rate constant for Pd₂Si growth should contribute to a better understanding of the diffusion mechanism during its formation.

II. MEASURING PRINCIPLE

Various techniques have been used to determine the intrinsic diffusion coefficients of metal and silicon in metal silicides. Tu made use of implanted Xe markers to determine the intrinsic diffusivity of Ni and Si in Ni₂Si.⁹ Pretorius and co-workers used radioactive Si tracers to determine the Si self-diffusion (intrinsic diffusivity) in PtSi and NiSi.^{10,11} Neither of these measuring techniques has been applied successfully to Pd₂Si.

We propose the following membrane technique to determine the mobility of Si and Pd in Pd₂Si. Consider the Pd₂Si layer as a membrane with a source supply of Si or Pd on the one side. The flux of Si or Pd through Pd₂Si can now be monitored with the aid of a suitably chosen layer on the other side which serves as a sink for the diffusing atoms [see Fig. 1(a)]. The thickness of the Pd₂Si layer is Y and it is assumed that A (Pd or Si) diffuses through Pd₂Si whereas B does not. On diffusing through, A and B form a compound $A_x B$. The flux of A through Pd₂Si, J_A , can be expressed in terms of the chemical potential difference of A across Pd₂Si, $\Delta \mu_A$, and the mobility of A in Pd₂Si, M_A :

$$J_A = -M_A C^A_{\mathrm{Pd},\mathrm{Si}} \Delta \mu_A / Y ,$$

where $C_{Pd,Si}^{A}$ is the atomic concentration of A in Pd₂Si

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FIG. 1. Schematic of multilayer thin-film structures: (a) general, (b) Pd/Si couple, (c) Pd/Pd₂Si/Ge structure for determining Pd mobility, and (d) Si/Pd₂Si/Cr structure for determining Si mobility.

and Y is the thickness of the Pd_2Si layer. We now define the effective mobility of A in Pd_2Si as

$$\overline{M}_A = -M_A \Delta \mu_A ,$$

so that the flux becomes

$$J_A = \overline{M}_A C^A_{\text{Pd-Si}} / Y$$

The rate of $A_x B$ formation $R_{A_x B}$, in terms of the flux of A through Pd₂Si is given by

$$R_{A_xB} = J_A / C_{A_xB}^A$$

or

$$R_{A_{\mathbf{x}}B} = \overline{M}_{A} C^{A}_{\mathrm{Pd}_{2}\mathrm{Si}} / C^{A}_{A_{\mathbf{x}}B} Y , \qquad (1)$$

where $C_{A_xB}^A$ is the atomic concentration of A in A_xB . The effective mobility of A in Pd₂Si is therefore given by

$$\overline{M}_{A} = R_{A_{x}B} Y C^{A}_{A_{x}B} / C^{A}_{Pd_{2}Si} .$$
⁽²⁾

The effective mobility \overline{M}_A is not equivalent to the intrinsic diffusivity of A in Pd₂Si, since A is subject to a "driving force" $\Delta \mu_A$, but it is comparable to the parabolic rate constant $\widetilde{M}_{Pd,Si}$, as determined from the growth kinetics of Pd₂Si:

$$\widetilde{M}_{\rm Pd,Si} = L^2/2t$$
,

where L is the thickness of Pd_2Si formed in time t. The

equivalence between \overline{M} and \widetilde{M} can be shown by replacing A by Pd, B by Si, x by 2, Y by y, where y is the thickness of the growing Pd₂Si layer, and R_{A_xB} by dy/dt so that (2) becomes

$$\overline{M}_{\mathrm{Pd,Si}} = \frac{dy}{dt} y = \frac{y^2}{2t} \bigg|_0^L = \frac{L^2}{2t} = \widetilde{M}_{\mathrm{Pd,Si}} .$$

We have used $\overline{M}_{Pd,Si}$ rather than \overline{M}_{Pd} or \overline{M}_{Si} to indicate that the growth of Pd₂Si could be as a result of both Pd and Si diffusing through it [see Fig. 1(b)].

The following multilayer combinations are suggested for determining Pd and Si mobilities.

(i) Pd mobility. A layer of Ge can serve as a sink for Pd so that Pd_2Ge is formed at the Pd_2Si/Ge interface, as shown in Fig. 1(c).

(ii) Si mobility. The Cr layer which serves as a sink for Si results in $CrSi_2$ formation at the Pd_2Si/Cr interface, as is shown in Fig. 1(d).

The membrane technique as described above allows the determination of the effective mobility of Si or Pd through Pd₂Si, but does not give any other information regarding the diffusion mechanism except the activation energy. It is possible, in the case of the Si/Pd₂Si/Cr structures, that Pd₂Si can dissociate, giving up Si atoms to form CrSi₂ and releasing Pd atoms to diffuse through the Pd₂Si layer to the Si/Pd₂Si interface to form new Pd₂Si. In this case the measured effective mobility will be that of Pd and not Si. We believe that this mechanism is unlikely since in a very similar experiment involving Si transport through Pd₂Si, it was found that Pd₂Si does not dissociate during its oxidation at 750 °C.¹² This is in agreement with our assumption that Si and not Pd diffuses through Pd₂Si during CrSi₂ growth.

Since we have layers other than Pd or Si in contact with Pd_2Si , the driving force $(\Delta\mu)$ across Pd_2Si is influenced, which in turn affects the values obtained for \overline{M}_{Pd} and \overline{M}_{Si} . The driving forces operative in the various multilayer systems are related to the free energy of formation of the growing compound layers. Since the free energies of formation of the silicides are of the same order of magnitude¹³ and since Pd_2Ge is expected to have a free energy of formation of similar magnitude, we believe that the difference in driving forces in the three cases $Pd/Pd_2Si/Si$, $Pd/Pd_2Si/Ge$, and $Si/Pd_2Si/Cr$ will not account for large differences in the measured values of $\widetilde{M}_{Pd,Si}$, \overline{M}_{Pd} , and \overline{M}_{Si} . A comparison of these values is therefore justified.

In order to verify that the growth of $A_x B$ is limited by the diffusion of A through Pd₂Si, the following properties of the $A_x B$ layer grown on Pd₂Si need to be investigated. (i) It should grow linear with time since $t^{1/2}$ kinetics would imply growth controlled by diffusion through a layer which grows with time, that being $A_x B$. (ii) Equation (1) predicts $R_{A_x B} \propto 1/Y$ so that the rate of $A_x B$ growth should be inversely proportional to the thickness of the interposed Pd₂Si layer. If these two conditions are not satisfied, in which case one of the interfaces or diffusion through $A_x B$ is controlling $A_x B$ growth, then this membrane technique cannot be applied.

III. EXPERIMENTAL PROCEDURE

The various films were sequentially deposited onto chemically cleaned substrates by electron-beam evaporation of the metals at pressures below 10^{-6} Torr. Thin films of Ge, Si, and Pd were deposited on SiO₂ substrates whereas Pd and Cr were deposited on Si(111) and Si(100) substrates. Annealing was carried out in a quartz tube vacuum furnace at pressures below 2×10^{-7} Torr. The Ge/Si/Pd samples were pre-annealed at 220°C for an hour to obtain the structure SiO₂/Ge/Pd₂Si/Pd. In cases where the Si layer was very thin or very thick, it was necessary to adjust the pre-annealing temperature or time accordingly to ensure complete Pd₂Si formation without any detectable Pd₂Ge. The formation of Pd₂Ge was obtained by further annealing the samples at temperatures between 200 and 240°C for various periods of time.

By annealing the Si(111)/Pd/Cr samples at 380 °C for 10 min, Si(111)/Pd₂Si/Cr structures were obtained. The annealing time for Si(100)/Pd/Cr samples with approximately 1- μ m-thick Pd layers, was increased to 50 min at 380 °C. No CrSi₂ formation was detected after such a heat treatment whereas Pd₂Si formation was complete. The samples were further annealed at temperatures between 410 and 500 °C for various periods of time during which CrSi₂ was formed on Pd₂Si. In order to investigate Pd₂Ge and CrSi₂ growth without an interposed Pd₂Si layer, SiO₂/Ge/Pd and Si(100)/Cr structures were prepared and annealed following the procedure outlined above.

Rutherford backscattering spectrometry employing mega-electron-volt ${}^{4}\text{He}{}^{2+}$ ions was used to determine the thickness of the various layers and the composition of the compounds formed. The phases of the various compounds were identified by means of glancing-angle x-ray spectroscopy using the Read and Guinier cameras. It is found that Pd₂Si grows epitaxially on Si(111), which was verified from the absence of the characteristic Pd₂Si lines on the x-ray micrographs.

IV. RESULTS AND DISCUSSION

A. $Si/Pd_2Si/Cr$ interdiffusion

Figure 2 shows the Rutherford backscattering spectra of Si(111)/Pd/Cr samples as deposited, after preannealing at 380°C for 10 min and after further annealing at 470°C for 20 min. For clarity, the spectrum of the sample annealed at 470°C is shown separately. When comparing the two spectra in the top part of the figure, it can be seen that the Cr signal has remained unchanged which implies that no CrSi₂ was formed during the preannealing process. The broadening of the Pd signal and the appearance of the Si step indicate that palladium silicide has formed. The phase of this silicide could not be identified by means of glancing angle x-ray diffraction since the characteristic lines of Pd₂Si were absent due to the epitaxy of Pd₂Si grown on Si(111). The ratio of the Pd-signal height to the step height of the Si signal was used to establish the composition as Pd₂Si. The step in the Cr (low-energy side) and the second step in the Si signal



FIG. 2. Rutherford backscattering spectra of Si(111)/Pd/Cr samples: (top) as-deposited and after pre-annealing at 380 °C, and (bottom) after further annealing at 470 °C. The arrows indicate the energies of ${}^{4}\text{He}^{2+}$ ions scattered from the elements when located on the surface of the sample.

(high-energy side) in the bottom part of the figure indicate that CrSi_2 has formed. The width of both these steps are measures for the thickness of the CrSi_2 layer. Both Si and Cr signals of CrSi_2 formed on thick Pd_2Si (~1 μ m) are superimposed on the Pd signal. Since the Cr signal is larger than the Si signal, it was used for determination of the CrSi₂ layer thickness.

No Cr diffusion into the Pd or Pd₂Si layers is observed in any of these spectra or spectra of similar samples. The Pd₂Si/CrSi₂ interface (high-energy side of the Pd signal in Fig. 2) has remained sharp, although Si has diffused through Pd₂Si to form CrSi₂. Degradation of the Si(111)/Pd₂Si interface (long tail of the low-energy edge of Pd signal) is observed when comparing the Pd signals of the spectra of the samples annealed at 380 and 470 °C. We believe that the nonuniform Si(111)/Pd₂Si interface as indicated by the tail in the Pd signal arises from the release of Si from the substrate. The mechanism of Si release is still unknown but could involve Pd₂Si dissociation at the Si(111)/Pd₂Si interface. The Si(100)/Pd₂Si interface of similar samples (same metal-layer thickness and heat treatment) is found to remain uniform during CrSi₂ growth and is therefore not affected by the Si release or transport process.

In simple metal/silicon or metal/metal phase-forming interactions, it is relatively easy to conclude from the growth kinetics whether the interaction is controlled by an interface reaction (t kinetics) or by diffusion of the reactants ($t^{1/2}$ kinetics). The interpretation of multilayer growth kinetics is complicated by the fact that a large number of interfaces and layers can control the layer



FIG. 3. Growth kinetics of $CrSi_2$ at 470 °C, illustrating the effect of interposed Pd₂Si on the growth rate.

growth. In such cases linear growth kinetics are no longer a sufficient condition for the interaction to be controlled by an interfacial reaction.

Figure 3 shows a plot of $CrSi_2$ thickness formed at 470 °C against annealing time. The growth of $CrSi_2$ on Si(100) using Si/Cr thin-film couples is compared with CrSi₂ growth on 2000-Å-thick epitaxial Pd₂Si using Si(111)/Pd/Cr multilayer films. Since the CrSi₂ formation on Si is linear with time, we conclude that it is limited by an interface reaction. The linear kinetics for the



FIG. 4. The rate of $CrSi_2$ growth on Pd_2Si is shown to be inversely proportional to the Pd_2Si layer thickness.

case of interposed Pd₂Si does not necessarily imply that it is also reaction limited. Equation (1) predicts that the CrSi₂ formation will be linear with time (constant rate) and inversely proportional to the interposed Pd₂Si layer thickness when it is controlled by diffusion through the Pd₂Si layer. The lower growth rate (gradient) for CrSi₂ on Pd₂Si in Fig. 3 is a direct result of the diffusion limiting effect of the Pd₂Si layer.

Figure 4 shows the variation of $CrSi_2$ growth rate with reciprocal Pd₂Si thickness at 450 °C for both epitaxial and nonepitaxial Pd₂Si. It is obvious that the growth rate of CrSi₂ is inversely proportional to the interposed Pd₂Si layer, confirming the prediction of (1). The more thickness of nonepitaxial Pd₂Si necessary to limit Si transport for CrSi₂ growth is related to the greater mobility of Si in polycrystalline Pd₂Si compared to epitaxial Pd₂Si. The effect of the Pd₂Si layer thickness on CrSi₂ growth rate is discussed in greater detail in Ref. 14.

B. Pd/Pd₂Si/Ge interdiffusion

The backscattering spectra of Ge/Si/Pd samples are shown in Fig. 5. The spectrum of the sample annealed at 220°C for 3 h is shown separately. On comparing the two spectra in the top of the figure, a broadening of the Si signal as well as a step in the Pd signal are observed. This indicates that Pd₂Si, which was confirmed by x-ray diffraction, has formed. The Ge signal around channel 430 has remained unchanged, indicating that no Pd₂Ge has formed on pre-annealing at 220°C. After further annealing the sample at 220°C for 3 h, a change in the Ge signal is observed (see bottom of Fig. 5). The decrease of the un-



FIG. 5. Backscattering spectra of SiO₂/Ge/Si/Pd structures: (top) as-deposited and after pre-annealing at 220 °C, and (bottom) after further annealing at 220 °C, Pd₂Ge is formed.

reacted Pd was determined and from that the Pd_2Ge thickness was calculated.

It is well documented that impurities in a film can influence the diffusion of atoms through it.^{15–17} The extent to which impurities enhance or inhibit diffusion depends on their interaction with the matrix and diffusing atoms and on the diffusion mechanism. Experimentally obtained values for diffusion coefficients could therefore be perturbed by foreign atoms in the diffusing medium. The bottom spectrum in Fig. 5 indicates detectable amount of Ge diffusion into the Pd₂Si layer. Germanium does not differ much from silicon with respect to its chemistry and atomic volume. We therefore believe that the presence of Ge in Pd₂Si will not have an appreciable influence on the diffusion of Pd through the Pd₂Si layer.

The growth of Pd₂Ge at 200 °C is illustrated in Fig. 6. Linear thickness and time scales are used in the figure in order to compare the two cases (with and without interposed Pd₂Si). The parabolic growth (linear when the thickness is plotted against the square root annealing time) of Pd₂Ge from Pd/Ge couples is characteristic of diffusion-limited growth and in agreement with reported studies on Pd₂Ge formation.¹⁸ The interposed Pd₂Si layer is seen to have a dramatic effect on the Pd₂Ge kinetics. It was pointed out in Sec. II that linear kinetics in a multilayer-compound-forming system is not necessarily indicative of interfacial-reaction-controlled growth. Although the growth of Pd₂Ge on Pd₂Si is seen to be linear with annealing time, it is in fact limited by diffusion of Pd through the Pd₂Si layer, the thickness of which does not increase with time.



FIG. 6. Growth kinetics of Pd_2Ge at 200 °C, showing parabolic growth for Pd/Ge couples and linear growth for $Ge/Pd_2Si/Pd$ structures. Note that different time scales are used for the two curves.

C. Mobility of Si and Pd

The effective mobilities of Si and Pd in Pd₂Si were calculated from Eq. (2), and the values of the parameters used are tabulated in Table I. These mobility values are plotted against the reciprocal temperature in Fig. 7. Curves representing the parabolic rate constant for Pd₂Si growth, obtained from the literature,^{8,19} are also shown in this figure. Good agreement between the effective mobility of Pd and the parabolic rate constant is obvious. This would suggest that the Pd mobility is sufficiently high to account for the Pd₂Si growth.

Calculated values for Si mobility in epitaxial and polycrystalline Pd₂Si are seen to be significantly different. The slopes of the two curves represent activation energies of 2.1 and 1.7 eV associated with Si diffusion through epitaxial and polycrystalline Pd₂Si, respectively. The higher mobility of Si in polycrystalline Pd₂Si could be attributed to Si grain-boundary diffusion. Furthermore, grain boundaries can serve as sources of point defects, thereby enhancing the Si vacancy concentration in polycrystalline Pd₂Si, resulting in a high Si mobility. The high activation energy of 2.1 eV for epitaxial Pd₂Si is characteristic of a lattice diffusion process, e.g., vacancy diffusion. [Similar values were obtained for Si self-diffusion in PtSi (2.1 eV) and NiSi (2.3 eV).^{10,11}] The activation energy associated with Pd₂Si growth (epitaxial and polycrystalline) was reported to be approximately 1.0 ± 0.1 eV.⁸ It is therefore obvious that Si is not the only diffusing species during Pd₂Si growth. This is supported by the results of Si marker experiments in Ref. 6. A probable diffusion mechanism for Pd₂Si growth is a diffusion process in

TABLE I. Data used in calculating the effective mobility through а Pd₂Si membrane using the relation $\overline{M} = R_{A_x B} Y C_{A_x B}^A / C_{Pd_2 Si}^A$, where A is the diffusant, B is the sink, the phase formed is $A_x B$, Y is the membrane thickness, and R the growth rate of $A_{\mathbf{x}}B.$ $C_{\rm CrSi_2}^{\rm Si} / C_{\rm CPd_2Si}^{\rm Si} = 2.35;$ $C_{\rm Pd_{2}Ge}^{\rm Pd}/C_{\rm Pd_{2}Si}^{\rm Pd} = 0.96.$

Т	$Y(Pd_2Si)$	$A_x B$ rate R	Mobility \overline{M}
(°C)	(Å)	(Å/s)	$(10^{-13} \text{ cm}^2/\text{s})$
	Si through	epitaxial Pd ₂ Si to fr	om CrSi ₂
450	1100	0.3	0.77
470	1000	0.75	1.75
485	1420	1.14	3.8
500	1420	2.9	7.6
	Si through pol	ycrystalline Pd ₂ Si to	from CrSi ₂
420	9500	0.12	2.7
425	9500	0.16	3.6
445	9500	0.45	9.9
445	15 500	0.27	9.8
455	9500	0.52	11.6
470	9500	0.85	19.0
470	15 500	0.47	17.2
480	9500	1.1	25.0
	Pd through poly	verystalline Pd ₂ Si to	form Pd ₂ Ge
200	1600	0.022	0.034
220	1600	0.083	0.132



FIG. 7. Arrhenius plot of the effective mobilities of Pd and Si in Pd₂Si. Also shown are the parabolic rate constants for Pd₂Si from Pd/Si couples from Refs. 8 and 17.

which the Pd and Si fluxes are coupled, which results in a significantly lower activation energy than that measured for Si mobility. Although the activation energy for Si mobility in polycrystalline Pd_2Si is lower than that through epitaxial Pd_2Si due to the presence of grain boundaries, it is still significantly higher than the activation energy for Pd_2Si growth.

When the Si mobility curves are extrapolated down to approximately 200 °C, the values are found to be orders of magnitude lower than the parabolic rate constant for both epitaxial and polycrystalline Pd_2Si growth. The mobility of Si is therefore far too low to account for the Pd_2Si growth. This confirms the earlier statement that Si is not the only moving species during Pd_2Si growth.

The validity of the comparison of the effective mobility obtained by means of the membrane technique to the parabolic rate constant as obtained from the growth kinetics of Pd₂Si is based on the assumption that the chemical potential difference $\Delta\mu$ across the Pd₂Si layer is not very different in the three cases Si/Pd₂Si/Pd, Si/Pd₂Si/Cr, and Pd/Pd₂Si/Ge. Based on free-energy-of-formation considerations, we believe the driving forces in the three cases to be of the same order.

V. SUMMARY AND CONCLUSIONS

A diffusion-membrane technique has been developed and successfully applied to determine the mobility of Pd and Si respectively in Pd₂Si. A multilayer film consisting of Ge/Si/Pd was used to determine the mobility of Pd. Initial annealing of this system produced Ge/Pd₂Si/Pd. During subsequent annealing the flux of Pd through Pd₂Si was monitored with the aid of Ge which served as a sink for Pd by forming Pd₂Ge. The mobility of Si in Pd₂Si was determined in a similar manner using Si(111)/Pd/Cr and Si(100)/Pd/Cr structures. In this case Cr served as a sink for Si and formed CrSi₂ with the Si which diffused through the Pd₂Si layer in Si/Pd₂Si/Cr.

The mobilities of Pd and Si were calculated by making use of the rate of Pd₂Ge and CrSi₂ formation, respectively. The value for Pd mobility determined around 200°C agrees with the parabolic rate constant for Pd₂Si growth. The mobility of Si in both polycrystalline and epitaxial Pd₂Si at this temperature is found to be orders of magnitude lower. The higher Si mobility in polycrystalline Pd₂Si is attributed to the presence of grain boundaries. An activation energy of 2.1 eV has been determined for the Si mobility in epitaxial Pd₂Si, whereas the value of 1.7 eV has been found for polycrystalline Pd₂Si. Because of the relatively low Si mobility we conclude that Si is not the only diffusing species during Pd₂Si growth. Radioactive marker experiments indicate that both Pd and Si diffuse with comparable fluxes during Pd₂Si formation.⁶ This suggests that the atomic transport mechanism in the diffusion membrane (Pd₂Si) is not necessarily identical to the diffusion mechanism during the membrane compound (Pd_2Si) growth. Care should therefore be taken when comparing the mobilities to the parabolic rate constant for drawing conclusions regarding the dominant diffusion species during thin-film compound formation.

The successful application of the diffusion membrane technique to the Pd_2Si system demonstrates its use as a method for determining the mobility of components in a compound. The technique can be applied to other systems provided suitable combinations of multilayer films can be found.

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