

Structural information from Miedema's theory for the heat of formation of intermetallic phases

T. Rajasekharan and K. Girgis

*Institut für Kristallographie und Petrographie, Eidgenössische Technische Hochschule Zürich,
CH-8092 Zürich, Switzerland*

(Received 6 August 1982)

We show in this paper that, notwithstanding the isotropic nature of Miedema's semi-empirical theory for the heat of formation ΔH of intermetallic phases, and contrary to the present view in literature, the parameters Φ^* and n_{WS} used to predict the signs of ΔH can contribute to our understanding of structural phenomena quantitatively. On a $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map considerable resolution is obtained among the binary systems in which different structure types occur, showing the importance of the energy of formation of an intermetallic compound in deciding the crystal structure that it adopts. The points corresponding to the systems in which the Laves phases (MgCu₂, MgZn₂, and MgNi₂ types), the SiCr₃-type phases, the TiAl₃-type phases, etc., occur, show linear relationships on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map, whereas the binary systems with MoSi₂-type phases are resolved into two lines corresponding to the two groups of MoSi₂-type phases with different c/a ratios and near-neighbor coordinations. The above observations enable us to explain many structural trends that can be noticed from the phase diagrams: for instance, the almost total exclusion of the SiCr₃-type phases from the 250 binary systems in which Laves phases occur and that of Laves phases from the 87 binary systems with SiCr₃-type phases.

I. INTRODUCTION

Since an *ab initio* quantum-mechanical calculation of the crystal structure adopted by intermetallic compounds remains beyond our reach today¹ and since an immense wealth of experimental data exist in the literature in the form of solved structures, it becomes interesting to look for systematics in the available experimental data. Such systematics could form a practical guideline for the production of new members of the same structure type and they could also throw more light on the basic physical processes leading to the occurrence of a particular structure type in a binary system. We examine in this paper the possibility of obtaining structural information on intermetallic compounds using the highly successful thermochemical scheme of Miedema *et al.*²⁻⁴ which predicts the signs of the heat of formation of liquid and solid metallic mixtures with a reliability between 96% to 100%. They have used two variables: Φ^* , which measures the chemical potential (closely related to the experimental work functions or electronegativity), and n_{WS} , which was thought to represent the electron density at the surface of the Wigner-Seitz (WS) cell for each element. A quantum-mechanical interpretation of Miedema's parameters has been proposed by Chelikowsky and Phillips.¹ The equation suggested by Miedema *et al.*^{3,4} for the heat of formation of solid alloys

formed from elements A and B is

$$\Delta H \sim [-(\Delta\Phi^*)^2 + (Q/P)(\Delta n_{WS})^{2/3} - R/P] . \quad (1)$$

Q/P and R/P are constants obtained from the examination of phase-diagram information. $R=0$ for all alloys except for combinations of a d metal with a p element or with divalent Mg or Be. Equation (1) shows that in cases where $R=0$, if $(\Delta\Phi^*)^2/(\Delta n_{WS})^{2/3} > Q/P$ for a particular binary system, the heat of formation is negative and vice versa. Accordingly, Miedema *et al.* have been able to separate "+" signs (representing element combinations with less than 10% mutual solid solubility) and "-" signs (representing element combinations with more than 10% mutual solid solubility and/or having intermediate phases in the phase diagram) on a map using $|\Delta\Phi^*|$ and $|(\Delta n_{WS})^{1/3}|$ as coordinates. From the slope of the straight line that separates the two regions, the value of Q/P was obtained as $9.4 \text{ V}^2/(\text{d.u.})^{2/3}$. [See Refs. 2 and 3 for the definition of Miedema's density units (d.u.)] For the combination of transition metals with p elements the boundary that separates the + and - signs was found from phase-diagram data to be a parabola dependent (slightly) on the valency of the p element. It is this observation that necessitated the introduction of the parameter R in Eq. (1) for the

transition-metal-*p*-element combinations.

Miedema *et al.* have treated only the isotropic (structurally averaged) part of ΔH . So at the first look it is not obvious that the theory will yield any structural information at all. Attempts have been made to separate various structure types occurring in equiatomic phases using the coordinates ($|\Delta\Phi^*|$, $|\Delta n_{WS}^{1/3}|$) without success.⁵ We intend to show in this paper that useful conclusions regarding the structure of intermetallic phases can be obtained if one uses $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ as the coordinates. In the next section we discuss the structure types studied in detail. In addition to those structure types, six other structure types are also treated in this paper.

II. THE STRUCTURE TYPES STUDIED IN DETAIL

(1) The *Laves phases* denote a large group of intermetallic phases AB_2 with one of the following structure types: (a) $MgCu_2$, (b) $MgZn_2$, and (c) $MgNi_2$. The $MgCu_2$ structure is cubic with eight formula units in the elementary cell and belongs to the space group $Fd\bar{3}m$. $MgZn_2$ is hexagonal with

four formula units in the elementary cell and belongs to the space group $P6_3/mmc$. $MgNi_2$ is hexagonal and belongs to the space group $P6_3/mmc$ and has eight formula units in the elementary cell. The above three structure types are closely related and are discussed in detail elsewhere.^{6,7}

(2) The cubic *SiCr₃ type* (A15) belongs to the space group $Pm\bar{3}n$ and has two formula units in the elementary cell. This structure type has evinced a lot of interest because many of the intermetallic compounds which are superconducting with high T_c crystallize in this structure type. A detailed description of this structure type is also available elsewhere.^{8,9}

(3) The *MoSi₂ type* is tetragonal and belongs to the space group I/mmm ; two formula units are in the elementary cell.¹⁰ The description of this structure is based on the recognition of two "ideal" axial ratios: $3\sqrt{2}/\sqrt{3}=2.45$ when close-packed layers are stacked in bcc [110] packing, and $3\sqrt{2}=4.24$ when close-packed layers are packed in close packing. The first group with c/a values ranging from 2.29 to 2.60 can be discussed using $HfAu_2$ as a typical example. The Hf atom lying in the plane $z=\frac{1}{2}$ has ten Au nearest neighbors: eight at equal distances in

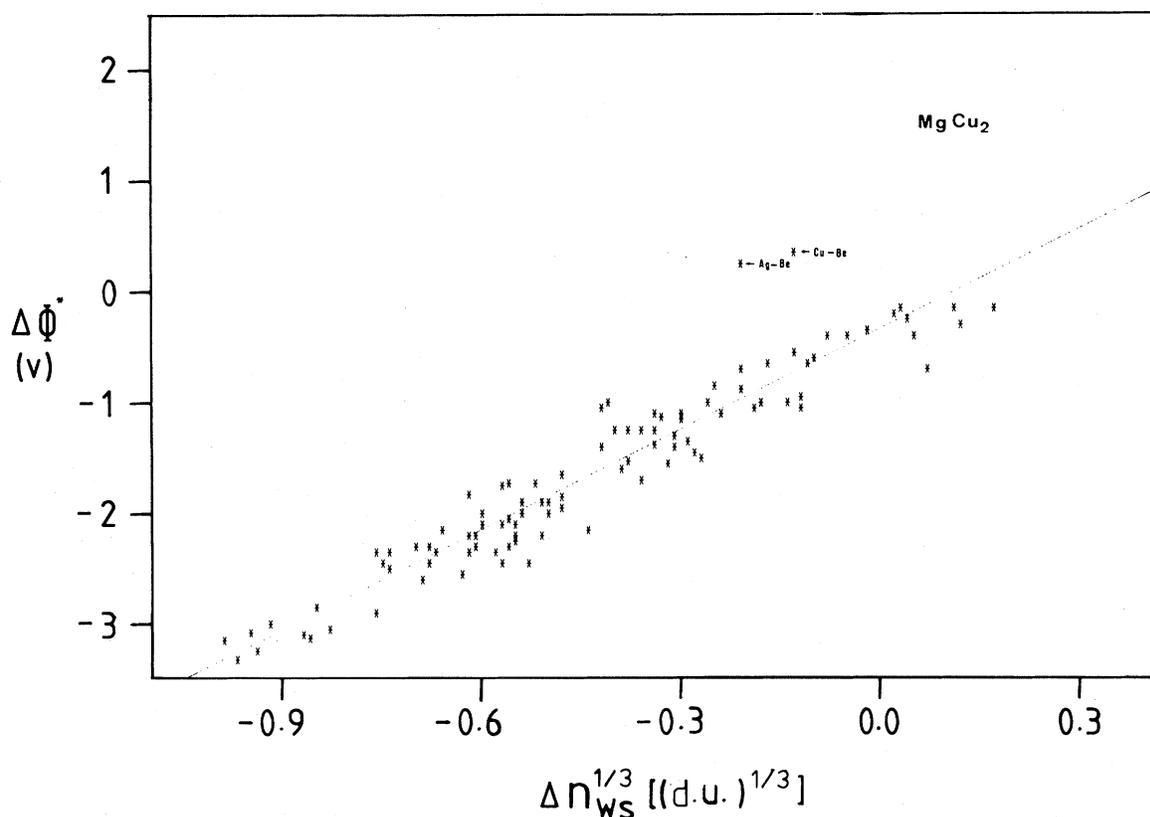


FIG. 1. A plot of $\Phi_A^* - \Phi_B^*$ vs $(n_{WS}^A)^{1/3} - (n_{WS}^B)^{1/3}$ for the binary systems with $MgCu_2$ -type phases. *A* is the minority element.

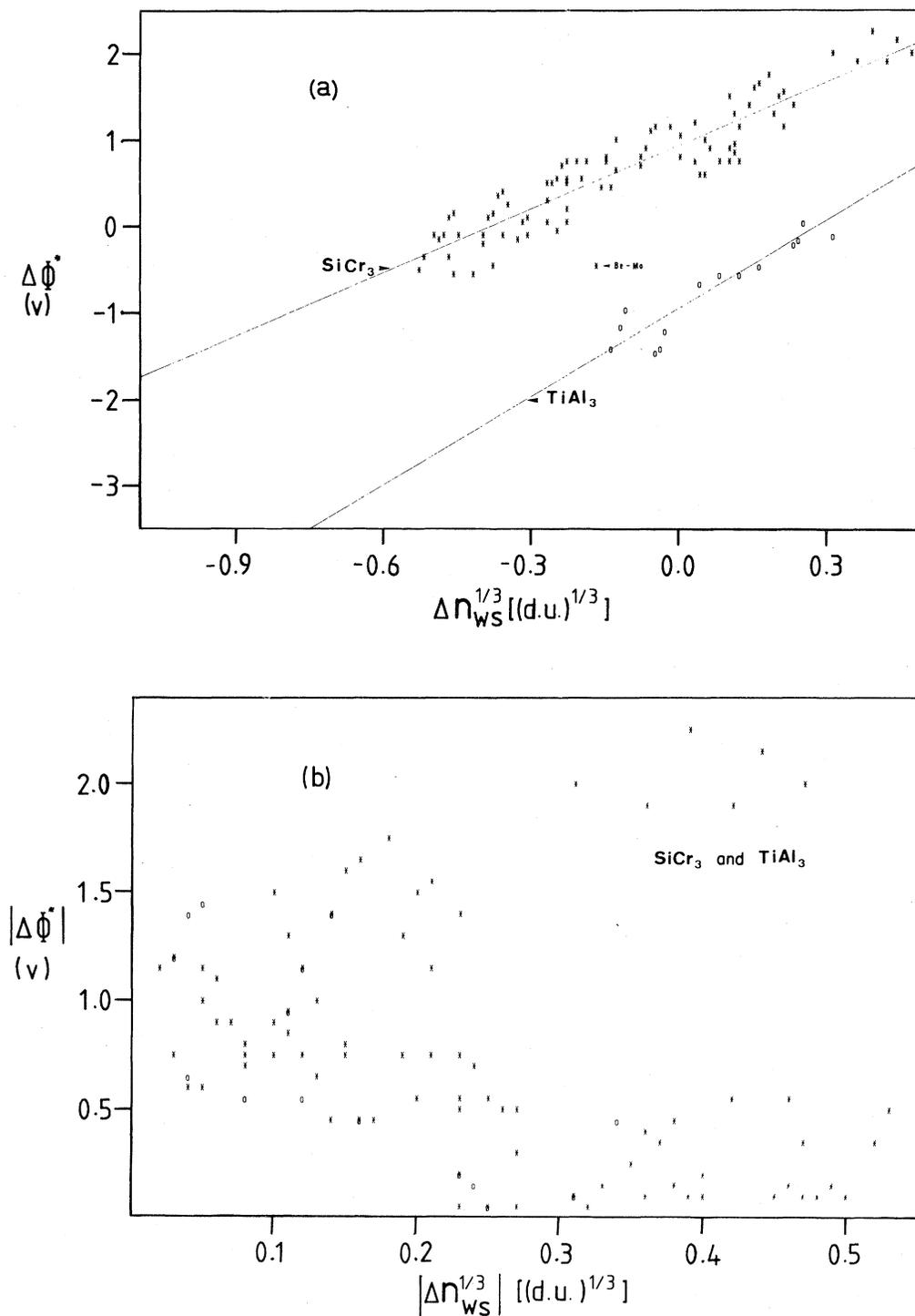


FIG. 2. (a) Plots of $\Phi_A^* - \Phi_B^*$ vs $(n_{WS}^A)^{1/3} - (n_{WS}^B)^{1/3}$ for the binary systems with SiCr₃-type (crosses) and TiAl₃-type (circles) phases. A is the minority element. (b) Binary systems with SiCr₃-type (represented by crosses) and TiAl₃-type (represented by circles) compounds are depicted on a $(|\Delta\Phi^*|, |\Delta n_{WS}^{1/3}|)$ map. The linear dependence between $\Delta\Phi^*$ and $(\Delta n_{WS})^{1/3}$ observed in (a) and also the resolution of the binary systems with SiCr₃-type and TiAl₃-type phases are lost in this diagram.

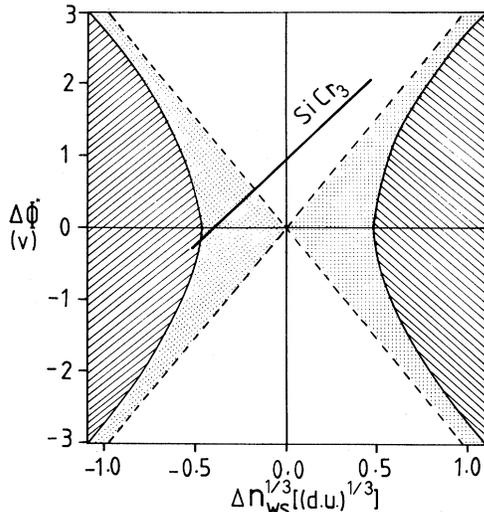


FIG. 3. Line corresponding to the binary systems with SiCr_3 -type compounds is shown on a $(\Delta\Phi^*, (\Delta n_{\text{ws}})^{1/3})$ map. The unshaded region has t - t and t - p combinations, the dotted region has only t - p combinations, and there are no compounds in the region shaded with lines. It is the negative-energy contribution from the R term that enables the t - p compounds to occur in the dotted region, though the positive-energy contribution from the $(\Delta n_{\text{ws}})^{2/3}$ term is more than the negative-energy contribution from $(\Delta\Phi^*)^2$ term in this region.

the planes $z \sim \frac{1}{3}$ and $\frac{2}{3}$, and two almost at the same distance along the c direction in the planes $z \sim \frac{1}{6}$ and $\frac{5}{6}$. The Au atom has essentially a coordination of the same form, with five atoms of the like kind and five atoms of the unlike kind. In the other group where the axial ratio varies from 3.18 to 4.68, only the atoms lying in the planes $z \sim \frac{1}{3}$ and $\frac{2}{3}$ are equidistant from the central atom. The distance of

the two atoms in the c direction is 20–30% greater, and four more atoms in the a direction are at about 10% greater distance. CuTi_2 is a typical member of this group. The compounds AlCr_2 , MoU_2 , and AlTc_2 have axial ratios that lie midway between the ranges of these two groups and do not fit into the above classification.

(4) The TiAl_3 type is tetragonal with two formula units in the elementary cell and belongs to the space group $I4/mmm$.¹¹

III. OBSERVATIONS AND DISCUSSION

We see from Eq. (1) that in cases where $R = 0$,

$$\Delta H < 0 \quad \text{if} \quad (\Delta\Phi^*)^2 / (\Delta n_{\text{ws}})^{2/3} > Q/P, \quad (2)$$

i.e.,

$$\Delta H < 0 \quad \text{if} \quad |\Delta\Phi^*| / |(\Delta n_{\text{ws}})^{1/3}| > (Q/P)^{1/2}. \quad (3)$$

It is reasonable to assume that in cases where one looks for structural information in a binary alloy in which the two types of atoms occupy different and distinct point positions, the actual direction of charge transfer and hence the signs of $\Delta\Phi^*$ and $(\Delta n_{\text{ws}})^{1/3}$ are of significance. In fact, inequality (2) offers a more general solution than inequality (3), i.e., ΔH is negative if $-(Q/P)^{1/2} < \Delta\Phi^* / (\Delta n_{\text{ws}})^{1/3} < (Q/P)^{1/2}$. $\Delta H = 0$ would correspond to $\Delta\Phi^* / (\Delta n_{\text{ws}})^{1/3} = \pm(Q/P)^{1/2}$.

In Fig. 1 we show a plot of $\Phi_A^* - \Phi_B^*$ vs $(n_{\text{ws}}^A)^{1/3} - (n_{\text{ws}}^B)^{1/3}$ for the binary systems in which the Laves phases with the MgCu_2 -type structure occur: A is the minority element. The list of the binary systems with phases isostructural with MgCu_2 is from Ref. 11. Out of a total of 173 such binary systems only 126 systems formed by the 62 elements whose Φ^* and $(n_{\text{ws}})^{1/3}$ values are avail-

TABLE I. Slopes (m) and intercepts on the $(\Phi_A^* - \Phi_B^*)$ axis (c) of the straight lines obtained in plotting $\Phi_A^* - \Phi_B^*$ vs $(n_{\text{ws}}^A)^{1/3} - (n_{\text{ws}}^B)^{1/3}$ are listed. A is the minority element. R is the regression factor. The structural data are taken from Ref. 11 except for the SiCr_3 structure type, in which case the data are taken from Ref. 13.

| Number | Structure type | Pearson symbol | Number of representatives | m [V/(d.u.) ^{1/3}] | c (V) | R |
|--------|--------------------------|----------------|---------------------------|-----------------------------------|------------|------|
| 1 | MgCu_2 | $cF24$ | 126 | 3.02 | -0.33 | 0.95 |
| 2 | MgZn_2 | $hP12$ | 67 | 3.30 | -0.20 | 0.96 |
| 3 | MgNi_2 | $hP24$ | 10 | 2.05 | -0.58 | 0.92 |
| 4 | SiCr_3 | $cP8$ | 87 | 2.44 | 0.94 | 0.92 |
| 5 | TiAl_3 | $tI8$ | 14 | 3.20 | -0.94 | 0.93 |
| 6 | CeCu_2 | $oI12$ | 11 | 2.02 | -0.80 | 0.96 |
| 7 | CuTi_3 | $tP4$ | 10 | 3.62 | 0.12 | 0.89 |
| 8 | Ni_2Al_3 | $hP5$ | 11 | 2.85 | 0.03 | 0.96 |
| 9 | AsNa_3 | $hP8$ | 16 | 2.34 | 0.73 | 0.98 |
| 10 | TiNi_3 | $hP16$ | 10 | 4.12 | -0.64 | 0.97 |
| 11 | CoAs_3 | $cI32$ | 6 | 2.51 | -0.33 | 0.89 |

able^{3,12} could be considered. The points lie on a straight line on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map with a regression factor $r=0.95$. The only two exceptions among the 126 binary systems considered are the Cu-Be and Ag-Be systems, and these exceptions will be discussed later.

67 binary systems having MgZn₂-type phases and 10 binary systems with MgNi₂-type phases form straight lines on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map with $r=0.96$ and 0.92 , respectively.

Figure 2(a) shows plots of $\Delta\Phi^*$ vs $(\Delta n_{WS})^{1/3}$ for 87 binary systems with SiCr₃-type compounds and for 14 binary systems with TiAl₃-type compounds. They form two separate straight lines with $r=0.92$ and 0.93 , respectively. Figure 2(b) illustrates that the binary systems with SiCr₃- and TiAl₃-type phases get mixed up in a $(|\Delta\Phi^*|, |(\Delta n_{WS})^{1/3}|)$ map. In Figure 3, we show schematically the distribution of points corresponding to the transition-metal-transition-metal (*t-t*) combinations ($R=0$), and the transition-metal-*p*-metal (*t-p*) combinations ($R\neq 0$).

The behavior, on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map, of intermetallic compounds with six more different structure types was studied and they also form straight lines on the map. The results are summarized in Table I. Among the binary systems with the compounds listed as having an AsNa₃-type crystal structure in Ref. 11, the system Ir-Si (with the compound IrSi₃) was observed to lie far from the straight line. This compound is exceptional, with a value for c/a of 1.52, which is unusual for AsNa₃-type compounds whose c/a values lie around 1.8. Hence Pearson has expressed doubts about the correct assignment of the structure type.¹¹ The value of R quoted in Table I for this structure type is excluding the Ir-Si system. ($R=0.94$, if the Ir-Si system is also included.)

A $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map for the binary systems with MoSi₂-type phases [Fig. 4(a)] shows two straight lines with some interesting exceptions. The different behavior of the binary systems with MoSi₂-type phases compared to those of the binary systems with Laves phases, SiCr₃- and TiAl₃-type phases, and also the other structure types listed in Table I, can be understood if one recalls from Sec. II that the intermetallic compounds belonging to this structure type can be classified into two groups with different near-neighbor coordinations based on their c/a ratios. A plot of c/a vs R_A/R_B for the MoSi₂-type phases is given in Fig. 4(b). A comparison of Figs. 4(a) and 4(b) clearly shows that the two lines on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map correspond to the two groups of compounds with two different c/a ratios. The phases AlCr₂, AlTc₂, and MoU₂ which cannot be classified into either of the two groups on the

basis of the c/a ratios are exceptions on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map too. The only exceptions to the above-mentioned correlation are the compounds CdTi₂ and AuMn₂ [Nos. 23 and 28 in Figs. 4(a) and 4(b)]. CdTi₂ has an exceptionally high c/a ratio for a MoSi₂-type compound and hence its structure has been considered doubtful in Ref. 11. We cannot offer an explanation for the exceptional behavior of AuMn₂ at present. The geometry of the structure and the assumption of simultaneous contacts between spherical atoms in *A-A* and *A-B* directions would necessitate an ideal radius ratio $R_A/R_B=1$ for the MoSi₂-type phases.¹⁰ The following can be seen from Figs. 4(a) and 4(b): The MoSi₂-type compounds with high c/a ratios have $R_A/R_B < 1$ and positive $\Delta\Phi^*$; the MoSi₂-type compounds with low c/a ratios have $R_A/R_B > 1$ and negative $\Delta\Phi^*$. Therefore one can conclude that the direction of the difference in chemical potential is such that the transfer of electrons between the atoms would make the R_A/R_B closer to the ideal value; for instance, electrons are transferred from *B* to *A* for the higher- (c/a) group.

It is seen from Figs. 1-4(a) and Table I that *it is possible to define a definite region on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map where points corresponding to the binary systems with a particular structure type are situated.* There are three binary systems which are exceptions: The Cu-Be and Ag-Be systems among the 250 binary systems with Laves phases and the Mo-Be system among the 87 binary systems with SiCr₃-type phases. These alloy systems of Be are exceptions probably because of the small size of the Be atoms. The above observation enables us to conclude, at least for the structure types that we have considered, that a proper relationship between $\Delta\Phi^*$ and $(\Delta n_{WS})^{1/3}$ is a *necessary* condition for the occurrence of a particular structure type in a binary system. To elaborate the above condition means that even if two elements satisfy the "geometrical conditions" for the formation of a structure type such as favorable R_A/R_B ratio, etc., the binary system will not have a particular structure type unless it is properly located on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map.

An examination of the phase-diagram data leads to an interesting observation that *certain structure types do not occur together in the same binary system.* For instance, in the 250 binary systems in which Laves phases occur, the SiCr₃-type phases do not occur and in 87 binary systems in which SiCr₃-type phases occur, the Laves phases do not occur: The only two exceptions are the Mo-Be system and the Ir-Zr system. BeMo₃ and IrZr₃ are of SiCr₃ type, ZrIr₂ is of MgCu₂ type, and MoBe₂ is of MgZn₂ type. Savitskii and Gribulya¹⁴ have also made a similar observation which they propose as a

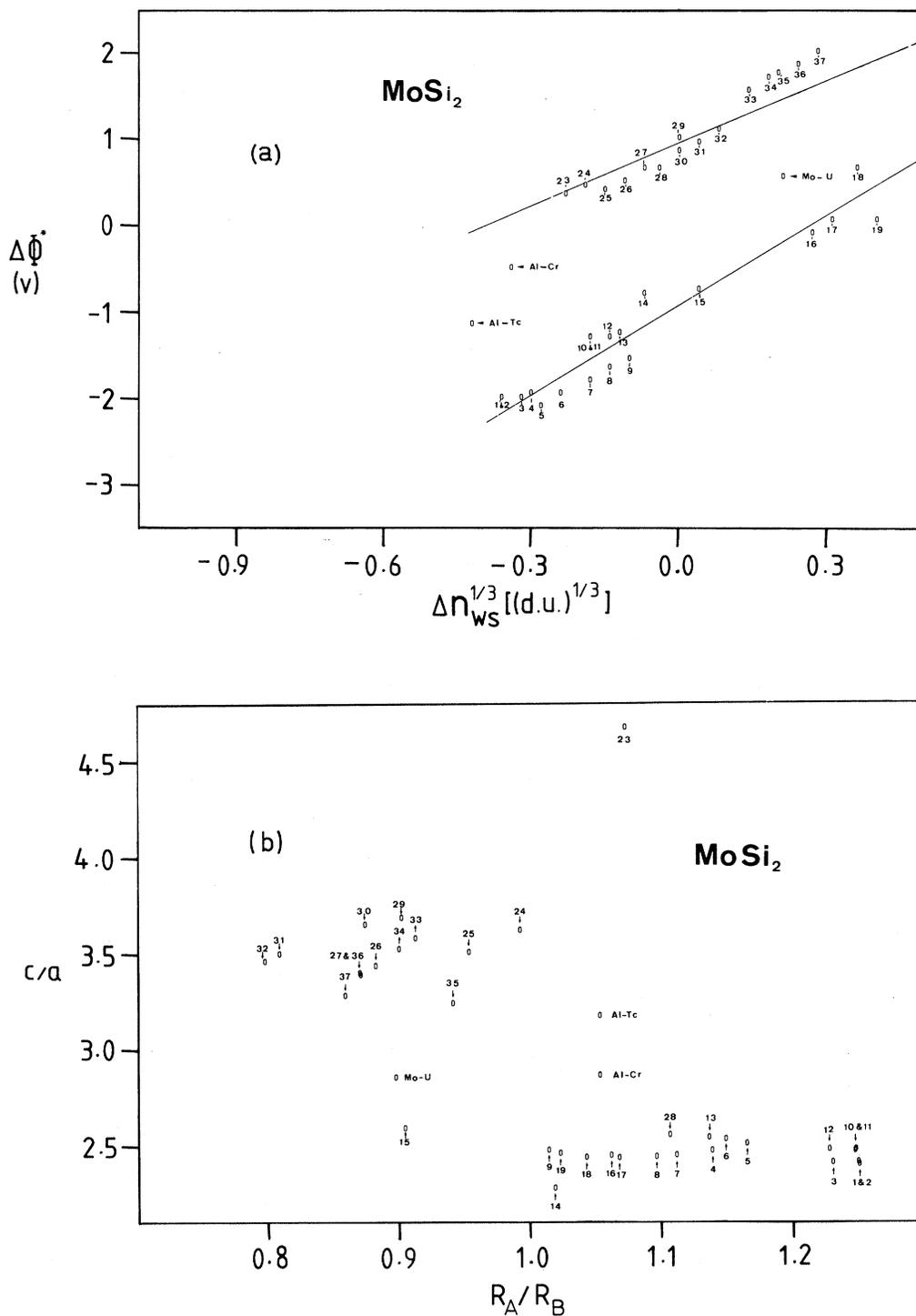


FIG. 4. (a) A plot of $\Phi_A^* - \Phi_B^*$ vs $(n_{WS}^A)^{1/3} - (n_{WS}^B)^{1/3}$ for the binary systems in which MoSi₂-type phases occur. The numbers in the figure correspond to the serial numbers used in Table II. Comparison with Fig. 4(b) shows that the two lines in the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map correspond to the two groups of compounds with the two different axial ratios. The compounds MoU₂, AlTc₂, and AlCr₂ are exceptions in both Figs. 4(a) and 4(b). (b) A plot of c/a vs R_A/R_B for the MoSi₂-type compounds. The serial numbers of Table II are used for referring to the compounds.

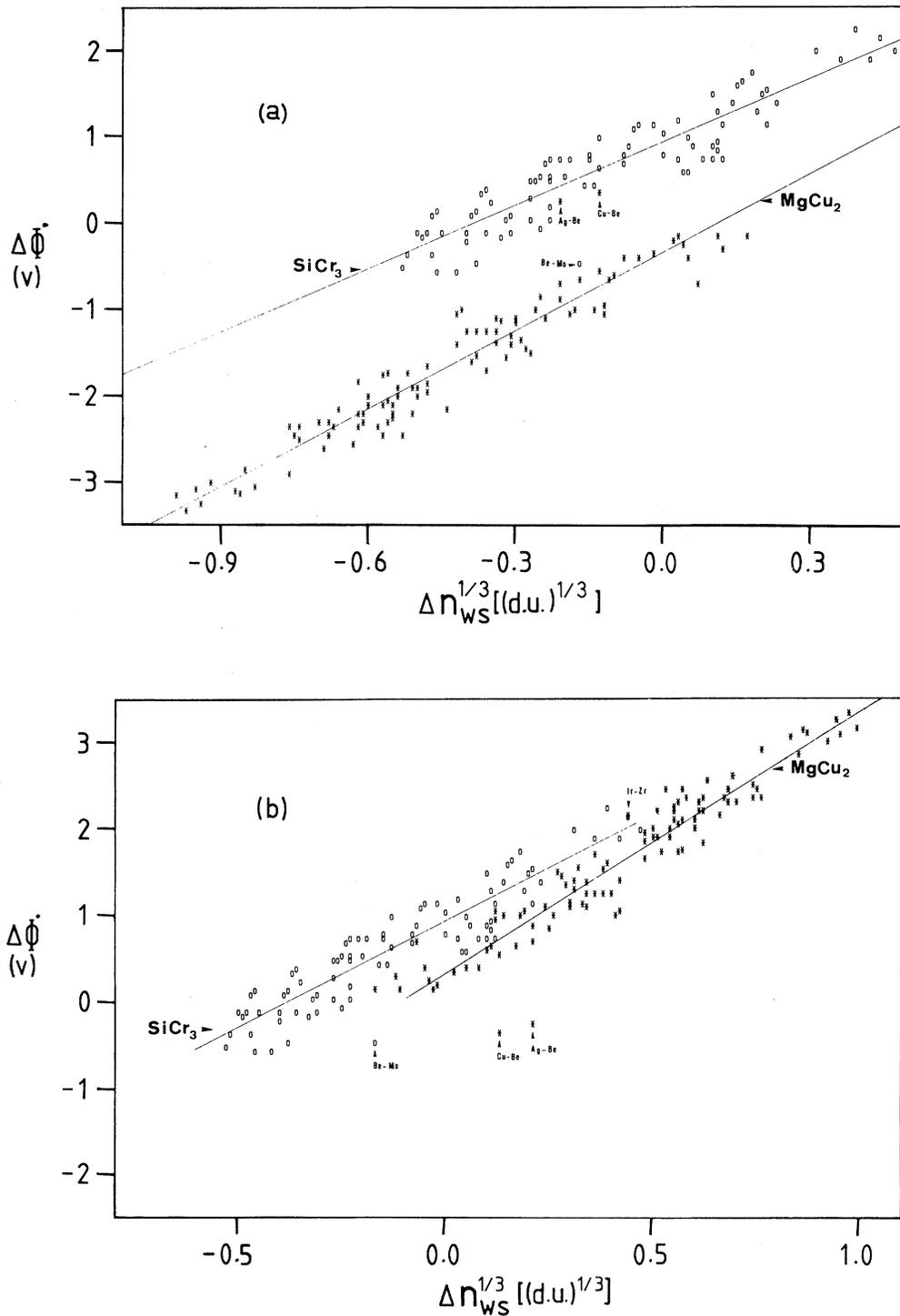


FIG. 5. (a) Plots of $\Phi_A^* - \Phi_B^*$ vs $(n_{ws}^A)^{1/3} - (n_{ws}^B)^{1/3}$ for the binary systems with MgCu_2 -type phases (represented by crosses) and SiCr_3 -type phases (represented by circles). The good resolution between the two lines on the $(\Delta\Phi^*, (\Delta n_{ws})^{1/3})$ map explains the absence of a SiCr_3 -type phase at composition AB_3 in a binary system with a MgCu_2 -type phase at composition AB_2 . (b) Plots of $\Phi_B^* - \Phi_A^*$ vs $(n_{ws}^B)^{1/3} - (n_{ws}^A)^{1/3}$ for the binary systems with MgCu_2 -type phases and $\Phi_A^* - \Phi_B^*$ vs $(n_{ws}^A)^{1/3} - (n_{ws}^B)^{1/3}$ for the systems with SiCr_3 -type phases.

guideline for looking for new representatives of the σ phase. They point out that "in systems with a σ phase there also exist phases of the SiCr_3 , Laves (MgZn_2), and CsCl type which may be called concomitants of the forecast structure." We demonstrate below the utility of the $(\Delta\Phi^*, (\Delta n_{\text{WS}})^{1/3})$ maps for predicting the "concomitants" of a particular structure type. Figure 5(a) is a plot of $\Phi_A^* - \Phi_B^*$ vs $(n_{\text{WS}}^A)^{1/3} - (n_{\text{WS}}^B)^{1/3}$ for the systems with MgCu_2 -type crystal structure and for those with SiCr_3 -type structure: A is the minority element. Since the two lines are without overlap we can conclude that if a MgCu_2 -type compound occurs at composition AB_2 in a binary system, a SiCr_3 -type compound cannot exist in the system at composition AB_3 . In Fig. 5(b) we plot $\Phi_B^* - \Phi_A^*$ vs $(n_{\text{WS}}^B)^{1/3} - (n_{\text{WS}}^A)^{1/3}$ for the binary systems with MgCu_2 -type phases along with a plot of $\Phi_A^* - \Phi_B^*$ vs $(n_{\text{WS}}^A)^{1/3} - (n_{\text{WS}}^B)^{1/3}$ for the systems with SiCr_3 -type compounds. These two lines are nearly resolved. However, the partial overlap suggests that at least in some systems with a MgCu_2 -type compound at composition AB_2 , a compound with SiCr_3 -type

structure can exist with the minority element of the MgCu_2 -type compound as its majority element. Indeed, the only system which is an exception to the mutual exclusion of MgCu_2 phases and SiCr_3 phases from the same binary system is in conformity with the above conclusion. ZrIr_2 is MgCu_2 type whereas IrZr_3 has a SiCr_3 -type structure. The mutual exclusion of the other Laves phases (MgZn_2 and MgNi_2) and SiCr_3 -type phases from the same binary system can also be similarly explained.

Figure 2(a) shows that a plot of $\Phi_A^* - \Phi_B^*$ vs $(n_{\text{WS}}^A)^{1/3} - (n_{\text{WS}}^B)^{1/3}$ resolves the binary systems with the AB_3 phases with SiCr_3 - and TiAl_3 -type structures. This figure would enable us to say whether a SiCr_3 - or TiAl_3 -type compound could exist in a certain binary system with a particular element as the minority element. Nevertheless a plot of $\Phi_B^* - \Phi_A^*$ vs $(n_{\text{WS}}^B)^{1/3} - (n_{\text{WS}}^A)^{1/3}$ for the binary systems with TiAl_3 phases on the same map would show that SiCr_3 -type phases and TiAl_3 -type phases could coexist in the same binary system at 1:3 and 3:1 compositions, respectively. AlNb_3 , AlV_3 , GaNb_3 , and PtV_3 are SiCr_3 -type compounds whereas NbAl_3 , VAl_3 ,

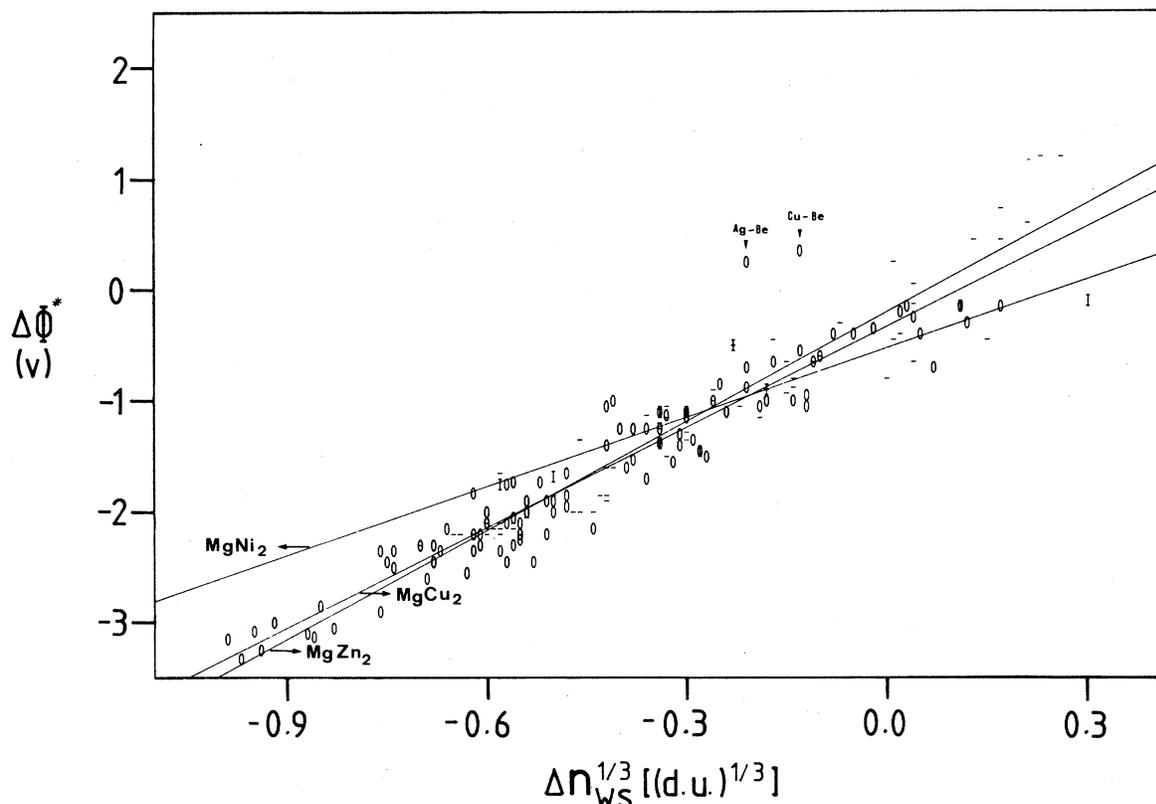


FIG. 6. Plots of $\Phi_A^* - \Phi_B^*$ vs $(n_{\text{WS}}^A)^{1/3} - (n_{\text{WS}}^B)^{1/3}$ for the binary systems with the Laves phases MgCu_2 , MgZn_2 , and MgNi_2 types. Systems with MgCu_2 -type phases are represented by "0," those with MgZn_2 -type phases with "-" and those with MgNi_2 -type phases with "I." A combination of two or more symbols represents a binary system with two or more of the Laves phases occurring at different temperatures.

NbGa₃, and VPt₃ are TiAl₃-type phases. Another relationship between structure types that coexist in the same binary system is the following. When an intermetallic compound with SiCr₃-type structure coexists with a MoSi₂-type phase at compositions 1:3 and 2:1, respectively, the MoSi₂-type compound belongs to the lower-axial-ratio group. Examples are the SiCr₃-type compounds AuTi₃, SiMo₃, and GeMo₃ which coexist with the MoSi₂-type compounds Au₂Ti, Si₂Mo, and Ge₂Mo of the lower-axial-ratio group. But, if the MoSi₂-type compounds exist at the composition 1:2 they belong to the larger-axial-ratio group. AgZr₃ and AuZr₃ which are SiCr₃ type coexist with the MoSi₂-type phases AgZr₂ and AuZr₂ which belong to the group of compounds with larger c/a values. The resolution of the binary systems with MoSi₂-type phases into two groups based on their c/a values on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map in Fig. 4(a) enables us to explain the above behavior. A plot of $(\Phi_A^* - \Phi_B^*)$ vs $(n_{WS}^A)^{1/3} - (n_{WS}^B)^{1/3}$ for the SiCr₃-type phases on Fig. 4(a) would coincide with the upper line corresponding to the binary systems with larger-axial-ratio compounds, showing that, in the same binary system, if the SiCr₃- and MoSi₂-type compounds

have the same minority element, the MoSi₂-type compound will belong to the larger-axial-ratio group. A plot of $(\Phi_B^* - \Phi_A^*)$ vs $(n_{WS}^B)^{1/3} - (n_{WS}^A)^{1/3}$ for the SiCr₃-type compounds will coincide with the lower line corresponding to the lower-axial-ratio group compounds showing that, when SiCr₃- and MoSi₂-type compounds exist with different minority elements in the same binary system, the MoSi₂-type compounds will belong to the lower-axial-ratio group, e.g., AuTi₃ (SiCr₃ type) and Au₂Ti (MoSi₂ type).

The MgCu₂-, the MgZn₂-, and the MgNi₂-type compounds are closely related structurally and Laves has suggested that the major factor controlling the formation of these phases is the relative size of the A and B atoms.¹⁵ On the basis of a hard-sphere-packing model, ideal space filling is achieved when the ratio of the atomic radii $R_A/R_B = 1.225$. It has been convincingly demonstrated by the work on ternary alloy systems that the occurrence of one or the other Laves phases, once the geometrical conditions are fulfilled, is determined by Brillouin-zone effects.⁶ Figure 6 shows that these closely related structure types occur in binary systems lying close together on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map.

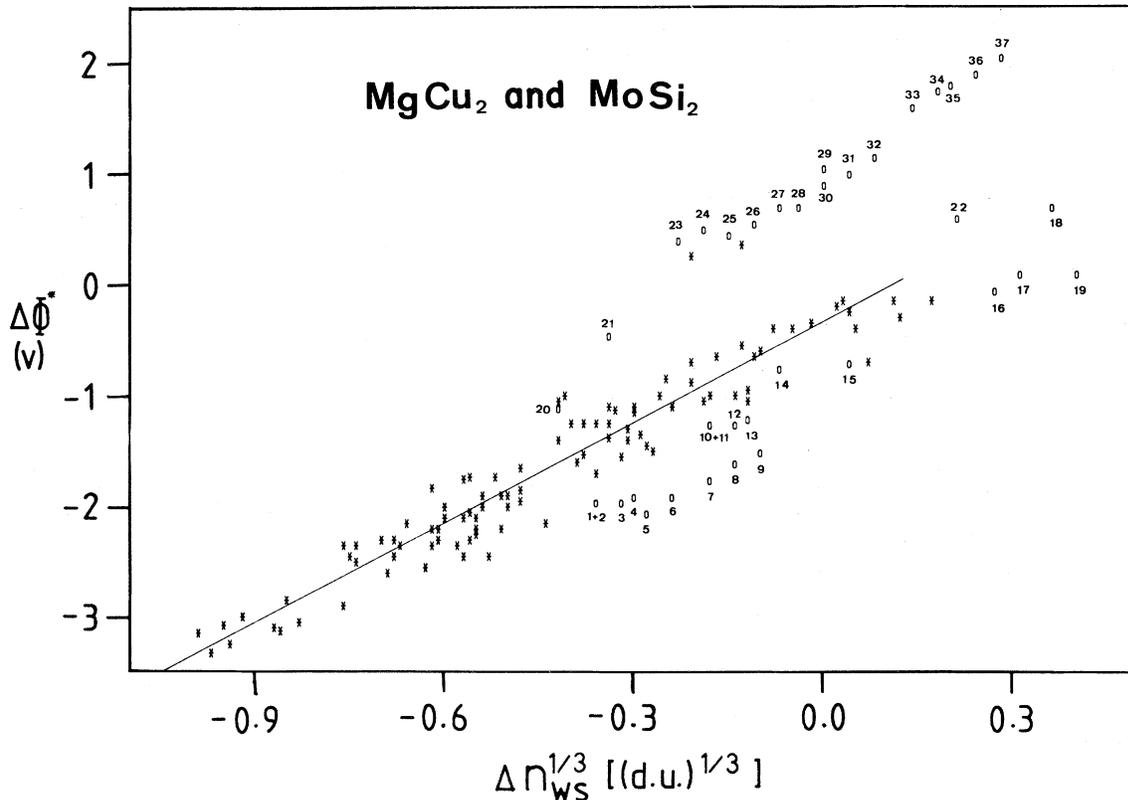


FIG. 7. Plot of $\Phi_A^* - \Phi_B^*$ vs $(n_{WS}^A)^{1/3} - (n_{WS}^B)^{1/3}$ for the binary systems with MgCu₂ phases (crosses) and MoSi₂-type phases (circles). The numbers refer to the serial numbers used in Table II.

In Fig. 7 the binary systems having 1:2 phases with MgCu_2 - and MoSi_2 -type crystal structures are shown on the same $(\Delta\Phi^*, (\Delta n_{\text{WS}})^{1/3})$ map. The binary systems with the MoSi_2 phases of the larger c/a values are well separated from those with the MgCu_2 -type phases. The resolution of the binary systems with MoSi_2 -type phases of smaller c/a values is not very good: Nevertheless we can explain at least partially a "curious anomaly" noted by Nevitt.¹⁰ He has pointed out that though many combinations of the elements of the Cu group with A components (i.e., minority elements) from the long periods have nearly the proper radius ratio (1.225) for crystallizing in one of the Laves phases, they crystallize with MoSi_2 -type structure. Table II lists the MoSi_2 -type compounds with their radius ratios. The combinations YAu_2 , GdAu_2 , DyAu_2 , ScAu_2 , ZrPd_2 , and HfPd_2 (Nos. 1–6 and 8 in the figure) have radius ratios which are favorable for the formation of Laves phases, but are not located in the region where systems with MgCu_2 phases are located on the $(\Delta\Phi^*, (\Delta n_{\text{WS}})^{1/3})$ map. The combinations YAg_2 , GdAg_2 , DyAg_2 , and ScAg_2 (Nos. 10–13), also with favorable radius ratios for the formation of Laves phases, do not crystallize with the MgCu_2 -type structure probably for the same reason; however, they are not very well resolved from the systems with MgCu_2 -type phases. One can also demonstrate the same for the other Laves phases (MgZn_2 and MgNi_2 types).

Though the difference in equilibrium energies of a particular intermetallic compound in closely related structures is known to be small, it becomes evident on inspection of structural data compilations^{11,16–18} that some structural information can be grossly deduced from the Periodic Table. For instance, a rela-

tionship between the position of elements in the Periodic Table and their equipoint occupation in a structure can be noticed.¹³ Pearson¹⁹ has pointed out that this is possible because in structures which are quite different from one another, the near-neighbor coordinations that determine a large part of the enthalpy are different. However, there has been no quantitative way of understanding structure types on the basis of the energy of formation of the binary systems in which they exist. The $(\Delta\Phi^*, (\Delta n_{\text{WS}})^{1/3})$ maps give us a means of doing the same. For instance, we can see from Figs. 5(a) and 5(b), respectively, that the resolution of the binary systems with MgCu_2 and SiCr_3 structure types on the $(\Delta\Phi^*, (\Delta n_{\text{WS}})^{1/3})$ map results from (i) the specific direction of the chemical potential difference ($\Delta\Phi^*$) and electron density difference $[(\Delta n_{\text{WS}})^{1/3}]$ between particular point positions of a structure type, and (ii) the fact that the magnitude of the chemical potential differences ($|\Delta\Phi^*|$) in binary systems with a SiCr_3 -type compound is generally greater than that of a system with a MgCu_2 -type phase at a given $|(\Delta n_{\text{WS}})^{1/3}|$ value. This shows that the energy of formation of a binary system with a SiCr_3 -type compound is greater than that of a system with a MgCu_2 -type compound.

IV. CONCLUSIONS

It has been shown in this paper that, at least for the structure types we have considered, the structure of an intermetallic compound depends on the energy of formation of the binary system in which it occurs, and that one can separate those structure types that occur in systems with considerably different energy of formation, and hence with very dif-

TABLE II. Ratios of the radii of the minority element to those of the majority element in MoSi_2 -type phases are listed. The serial numbers used here are the same as the one referring to the compounds in Figs. 4(a), 4(b), and 7. The table shows that some MoSi_2 -type phases have nearly the ideal radius ratio (1.225) for crystallizing in one of the Laves phases. Figure 7 and the text explain as to why they do not crystallize as a Laves phase.

| Number | Compound | R_A/R_B | Number | Compound | R_A/R_B | Number | Compound | R_A/R_B |
|--------|-----------------|-----------|--------|-----------------|-----------|--------|-----------------|-----------|
| 1 | YAu_2 | 1.25 | 14 | CuZr_2 | 1.02 | 26 | ZnHf_2 | 0.88 |
| 2 | GdAu_2 | 1.25 | 15 | MnAu_2 | 0.87 | 27 | ZnZr_2 | 0.87 |
| 3 | DyAu_2 | 1.23 | 16 | MoSi_2 | 1.06 | 28 | AuMn_2 | 1.11 |
| 4 | ScAu_2 | 1.14 | 17 | WSi_2 | 1.07 | 29 | AgZr_2 | 0.90 |
| 5 | ZrPd_2 | 1.16 | 18 | ReSi_2 | 1.04 | 30 | CuTi_2 | 0.87 |
| 6 | HfPd_2 | 1.15 | 19 | MoGe_2 | 1.02 | 31 | CuHf_2 | 0.81 |
| 7 | ZrAu_2 | 1.07 | 20 | AlTc_2 | 1.05 | 32 | CuZr_2 | 0.80 |
| 8 | HfAu_2 | 1.10 | 21 | AlCr_2 | 1.05 | 33 | AuHf_2 | 0.91 |
| 9 | TiAu_2 | 1.01 | 22 | MoU_2 | 0.90 | 34 | AuZr_2 | 0.90 |
| 10 | YAg_2 | 1.25 | 23 | CdTi_2 | 1.07 | 35 | PdTi_2 | 0.94 |
| 11 | GdAg_2 | 1.25 | 24 | CdHf_2 | 0.99 | 36 | PdHf_2 | 0.87 |
| 12 | DyAg_2 | 1.23 | 25 | ZnTi_2 | 0.95 | 37 | PdZr_2 | 0.86 |
| 13 | ScAg_2 | 1.14 | | | | | | |

ferent average energy in the near-neighbor bonds, on a $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map. Differentiating closely related structure types lying close together on the $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ map might require the consideration of effects such as Fermi-surface effects which are weaker in comparison with the isotropic volume-dependent forces contributing to the energy of formation of a solid and described by Miedema's theory for the heat of formation of intermetallic phases. The reason for the occurrence of two groups of compounds with different axial ratios in the MoSi_2 structure type is found to be the difference in

the direction of $\Delta\Phi^*$ between the point positions in these two groups. The $(\Delta\Phi^*, (\Delta n_{WS})^{1/3})$ maps have been demonstrated to be quite effective in explaining in a quantitative fashion the structural trends exhibited in the phase diagrams and in accounting for structures which are concomitant with a particular structure type.

ACKNOWLEDGMENT

This work was supported by the Swiss National Science Foundation.

-
- ¹J. R. Chelikowsky and J. C. Phillips, *Phys. Rev. B* **17**, 2453 (1978); *Phys. Rev. Lett.* **39**, 1687 (1977).
- ²R. Boom, F. R. de Boer, and A. R. Miedema, *J. Less-Common Metals* **45**, 237 (1976); **46**, 271 (1976).
- ³A. R. Miedema, R. Boom, and F. R. de Boer, *J. Less-Common Metals* **41**, 283 (1975).
- ⁴A. R. Miedema, P. F. de Châtel, and F. R. de Boer, *Physica (Utrecht)* **100B**, 1 (1980).
- ⁵A. Zunger, in *Structure and Bonding in Crystals*, edited by M. O'Keeffe and A. Navrotsky (Academic, New York, 1981), Vol. I, p. 129.
- ⁶J. H. Wernik, in *Intermetallic Compounds*, edited by J. H. Westbrook (Wiley, New York, 1967), pp. 198–204.
- ⁷W. B. Pearson, *The Crystal Chemistry and Physics of Metals and Alloys* (Wiley, New York, 1972), pp. 654–660.
- ⁸M. V. Nevitt, in *Intermetallic Compounds*, Ref. 6, pp. 217–220.
- ⁹W. B. Pearson, *The Crystal Chemistry and Physics of Metals and Alloys*, Ref. 7, pp. 669–673.
- ¹⁰M. V. Nevitt, in *Intermetallic Compounds*, Ref. 6, pp. 224–227.
- ¹¹W. B. Pearson, *Handbook of Lattice Spacings and Structures of Metals and Alloys* (Pergamon, Oxford, 1967), Vol. II.
- ¹²A. Borseese, G. Borzone, and R. Ferro, *J. Less-Common Metals* **70**, 213 (1980).
- ¹³P. Villars, K. Girgis, and F. Hulliger, *J. Solid State Chem.* **42**, 89 (1982).
- ¹⁴E. M. Savitskii and V. B. Gribulya, *Dokl. Akad. Nauk SSSR* **223**, 1383 (1975) [*Sov. Phys.—Dokl.* **223**, 414 (1975)].
- ¹⁵F. Laves and H. Witte, *Metallwirtsch. Metallwiss. Metalltech.* **14**, 645 (1935).
- ¹⁶M. Hansen, *Constitution of Binary Alloys* (McGraw-Hill, New York, 1958).
- ¹⁷E. Parthé, *Crystal Chemistry of Tetrahedral Structures* (Gordon and Breach, New York, 1970).
- ¹⁸*Treatise on Solid State Chemistry*, edited by N. B. Hannay (Plenum, New York, 1973–77), Vols. 1–7.
- ¹⁹W. B. Pearson, *The Crystal Chemistry and Physics of Metals and Alloys*, Ref. 7, p. 126.