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$CeRu_2$ and $CeCo_2$: Superconductors with 4f electrons

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 $CeRu_2$ and $CeCo_2$ have been considered to have tetravalent Ce because of their superconductivity, collapsed volume, and lack of magnetic moments. Resonant photoemission studies show these compounds to be mixed valent, suggesting the need to reexamine the mechanism of their superconductivity.

CeRu₂ and CeCo₂ are superconductors with transition temperatures T_c of 6 and 1 K, respectively.^{1,2} It has generally been assumed that Ce in these materials is tetravalent, i.e., $4f^0$, because studies in which Ce is introduced into various superconducting matrices have shown that a localized 4f electron, either magnetic or nonmagnetic, drives the host T_c to zero for only a few atomic percent Ce.³ CeRu₂ and CeCo₂ also display collapsed volume and a loss of magnetic moments, properties which are common to a group of Ce intermetallic compounds and which are often taken to signal tetravalent Ce for the group.⁴ Alternatively, mixed valence has been proposed for various of these compounds,^{5,6} and, based on a reexamination of the cerium atomic volume and on their 2pcore-level absorption spectra, there is a recent contention that no intermetallic tetravalent cerium compounds exist.^{7,8} This severely challenges the traditional understanding of superconductivity in CeRu₂ and CeCo₂, especially since superconductivity has been reported for only one other material known with certainty to be mixed valent, CeCu₂Si₂, for which the T_c of ~ 0.5 K is less than a tenth that of CeRu₂.^{9, 10}

This Communication describes the results of resonant photoemission measurements on CeRu₂ and CeCo₂, made to search for 4*f* emission. It is convenient to compare with a material known to have 4*f* electrons, and therefore data are also presented for CeAl₃, which is interesting in its own right because it displays 4*f* electron demagnetization below \sim 5 K.

The photoemission experiments reported here were performed at the Stanford Synchrotron Radiation Laboratory using equipment, operating conditions, and experimental procedures exactly as described previously.¹¹ The samples were polycrystalline ingots, fractured *in situ* in a vacuum of 1×10^{-10} Torr, and superconductivity with the correct T_c was verified for the samples of CeRu₂ and CeCo₂. Figures 1, 2, and 3 show the valence-band photoemission spectra of CeAl₃, CeRu₂, and CeCo₂, respectively, for several photon energies $h\nu$ between 80 and 130 eV.

The interpretation of resonant photoemission data for any cerium material, and especially one believed to have no 4f electrons, must take account of the



FIG. 1. CeAl₃ valence-band photoemission spectra for several photon energies. The resolution varies from 0.41 eV at 80 eV to 0.47 eV at 122 eV.

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FIG. 2. $CeRu_2$ valence-band photoemission spectra for several photon energies. The resolution varies from 0.41 eV at 80 eV to 0.47 eV at 122 eV.

fact that the Ce 5d emission resonates with multiplicative gain as large as 0.5 that of the 4f electrons. That this is so is known experimentally from studies of Gd and Tm, where the 4f and 5d emission do not overlap.^{12, 13} It is therefore necessary to have a spectroscopic criterion to identify 4f emission. In CeAl₃, as



FIG. 3. $CeCo_2$ valence-band photoemission spectra for several photon energies. The resolution is 0.83 eV for the 115-120-eV spectra, and is 0.41, 0.45, and 0.46 eV for the 80-, 112-, and 121-eV spectra, respectively.

in cerium itself and many other compounds, the most prominent feature in the on-resonant spectrum (photon energy $h\nu = 121-122$ eV) occurs at ~ 2 -eV binding energy and all workers who have made resonant photoemission studies of cerium compounds agree that this peak has 4f origin.^{11, 14-16} An identifying feature of its resonance behavior is that this peak becomes predominant only for $h\nu > 118$ eV, although the entire valence band resonates for $h\nu > \sim 112$ eV, the Fano minimum.

For the spectra of CeCo₂ and CeRu₂ the broad emission at ~ 2.6-eV binding energy behaves in exactly the same fashion and is therefore assigned as 4femission. Although in these cases the 4f state is more difficult to pick out because it is broad and located on a large background of Co or Ru d states, the similarity to the behavior in CeAl₃ is unmistakable. The finding of 4f emission in CeRu₂ and CeCo₂, and its greater width than in CeAl₃, is qualitatively consistent¹⁷ with all reported core-level photoemission and photoabsorption spectra, which show both $4f^0$ and $4f^1$ structure for CeCo₂ (Refs. 8 and 18) and CeRu₂,¹⁹ but only $4f^1$ structure for CeAl₃.²⁰

There is not complete agreement about the assignment of the resonating emission nearer to the Fermi level in Ce valence-band photoemission. Some work ers^{11} assign it as 5d emission, and this assignment is supported by the fact that strongly resonating 5demission is expected in the spectrum,^{12, 13} and by the behavior pointed out above, that its resonance begins and peaks at slightly lower photon energies than for the 2-eV feature. Such a difference for n = 4 and 5 shell emission was predicted theoretically²¹ and observed experimentally²² for the 4d, 5s, 5p, and 5d lanthanum emission of LaB₆. Other workers²³⁻²⁵ have asserted that, in some cerium materials, the emission near the Fermi level is also of 4f origin. Based on previous work, then, the sharp feature near E_F in the CeRu₂ spectrum could be assigned either as 4f or 5d. We prefer the 5d assignment for the reasons given above, but at present this is an unresolved issue.

Since the resonance technique has been used here to identify 4f emission well away from the Fermi level, it should be mentioned that Baer and his collaborators²⁰ have given a number of arguments that in several cerium materials, including γ -Ce and CeAl₃, the 4f emission lies very close to the Fermi level, with the remainder of the spectrum being of 5d origin. Of necessity these arguments are all indirect, having been undertaken to overcome the fact that the small spectral weight and lack of spectral structure of the $4f^1 \rightarrow 4f^0$ transition have made it impossible to identify by direct spectroscopic means in conventional x-ray photoemission (XPS). Other XPS workers^{26, 27} have also confronted this problem and have found cause to place the 4f emission at 2 eV in γ -Ce, in agreement with the conclusions from resonant photoemission.

Considering now these materials' superconductivity, at present we can only point out certain relevant aspects of the problem. Translational symmetry is the most obvious new ingredient capable of superceding the impurity behavior described in the opening paragraph, because it implies coherent 4f states with vanishing quasiparticle scattering as $T \rightarrow 0$. The details of the 4f coherence have yet to be elucidated. Although a simple 4f band model could account for the rather temperature-independent magnetic susceptibilities² of CeCo₂ and CeRu₂, such a model is inappropriate if the 4f Coulomb energy U is $\sim 5-6$ eV, as claimed from Bremsstrahlung isochromat spectroscopy on several cerium materials.²⁰ Presumably some sort of Fermi-liquid picture applies at low temperature and T_c is then material dependent, being governed by the quasiparticles' residual interactions, about which little is known, and by their coupling to phonons, possibly strong in a mixed valent system. As has been emphasized for CeCu₂Si₂ (Ref. 9) a Fermi-liquid enhancement of γ may be important; the values of γ for CeRu₂ and the isostructural $4f^0$ compound LaRu₂ are,²⁸ respectively, 23.3 ± 0.61 and

13.6 \pm 0.13 mJ/mole K², in roughly the same ratio as the values of T_c , which are 5.92 and 3.2 K, respectively. The Debye temperatures of CeRu₂ and LaRu₂ are nearly equal, implying similar phonon spectra, although the possibility is not ruled out that a particular feature of the phonon spectrum may account for the higher T_c of CeRu₂. Finally, the fact that rareearth substitutions in CeRu₂ initially raise T_c , even when the rare-earth ions carry localized moments,¹ is motivation to speculate that the response of the cerium lattice Fermi liquid to magnetic impurities is different than that of other Fermi liquids.

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- ¹M. Wilhelm and B. Hillenbrand, J. Phys. Chem. Solids <u>31</u>, 559 (1970).
- ²T. F. Smith, H. L. Luo, M. B. Maple, and I. R. Harris, J. Phys. F 1, 896 (1971).
- ³See, for example, M. B. Maple, Appl. Phys. <u>9</u>, 179 (1976), and references therein.
- ⁴H. J. van Daal and K. H. J. Buschow, Phys. Status Solidi (a) <u>3</u>, 853 (1970).
- ⁵F. R. de Boer, W. H. Dijkman, W. C. M. Mattens, and A. R. Miedema, J. Less Common Met. <u>64</u>, 241 (1979).
- ⁶J. M. Lawrence, P. S. Riseborough, and R. D. Parks, Rep. Prog. Phys. <u>44</u>, 1 (1981).
- ⁷D. K. Wohlleben, in *Valence Fluctuations in Solids*, edited by L. M. Falicov, W. Hanke, and M. B. Maple (North-Holland, Amsterdam, 1981), p. 1.
- ⁸K. R. Bauschspiess, W. Kboksch, E. Holland-Moritz, H. Launois, R. Pott, and D. K. Wohlleben, in Ref. 7, p. 417.
- ⁹F. Steglich, J. Aarts, C. D. Bredl, W. Lieke, D. Meschede, W. Franz, and H. Schafer, Phys. Rev. Lett. <u>43</u>, 1892 (1979).
- $^{10}\alpha$ -Ce under ~ 40-kbar pressure near the α - α' phase boundary is superconducting with $T_c \sim 50$ mK [C. Probst and J. Wittig, in *Handbook on the Physics and Chemistry of Rare Earths*, edited by K. A. Gschneidner, Jr., and L. Eyring (North-Holland, Amsterdam, 1978), Vol. I, Chap. 10], but traditional arguments would suggest near tetravalence for this case.
- ¹¹J. W. Allen, S.-J. Oh, I. Lindau, J. M. Lawrence, L. I.

Johansson, and S. B. Hagstrom, Phys. Rev. Lett. <u>46</u>, 1100 (1981).

- ¹²W. F. Egelhoff, Jr., G. G. Tibbetts, M. H. Hecht, and I. Lindau, Phys. Rev. Lett. <u>46</u>, 1071 (1981).
- ¹³F. Gerken, J. Barth, K. L. I. Kobayashi, and C. Kunz, Solid State Commun. <u>35</u>, 179 (1980).
- ¹⁴M. Croft, J. H. Weaver, D. J. Peterman, and A. Franciosi, Phys. Rev. Lett. <u>46</u>, 1104 (1981).
- ¹⁵M. Campagna, W. Gudat, W. Eberhardt, R. Rosei, J. H. Weaver, E. Kaldis, and F. Hulliger, Physica (Utrecht) 102B, 367 (1980).
- ¹⁶N. Martensson, B. Reihl, and R. D. Parks, Solid State Commun. 41, 573 (1982).
- ¹⁷References 18 and 19 conclude that CeRu₂ and CeCo₂ are tetravalent from an erroneous interpretation of their *L*edge absorption spectra, in which a small (\sim 2-eV) chemical shift is mistaken for the large (\sim 7-eV) valenceinduced shift (see Ref. 8).
- ¹⁸C. N. R. Rao, D. D. Sharma, P. R. Sarode, R. Vijayaraghavan, S. K. Dhar, and S. K. Malik, J. Phys. C <u>14</u>, L451 (1981).
- ¹⁹P. R. Sarode, D. D. Sarma, C. N. R. Rao, E. V. Sampathkumaran, L. C. Gupta, and R. Vijayaraghavan, Mat. Res. Bull. <u>16</u>, 175 (1981).
- ²⁰Y. Baer, H. R. Ott, J. C. Fuggle, and L. E. DeLong, Phys. Rev. B 24, 5384 (1981), and references therein.
- ²¹A. Zangwill and P. Soven, Phys. Rev. Lett. <u>45</u>, 204 (1980).
- ²²M. Aono, T.-C. Chiang, F. J. Himpsel, and D. E. Eastman, Solid State Commun. <u>37</u>, 471 (1981).
- ²³A. Franciosi, J. H. Weaver, N. Martensson, and M. Croft, Phys. Rev. B <u>24</u>, 3651 (1981).
- ²⁴D. M. Wieliczka, C. G. Olson, D. W. Lynch, and J. H.

Weaver, Bull. Am. Phys. Soc. <u>27</u>, 276 (1982).

- ²⁵R. D. Parks, N. M. Martensson, and B. Reihl, in *Proceedings of the International Conference on Valence Instabilities, Zürich, April 13-16, 1982, edited by P. Wachter (North-Holland, Amsterdam, in press).*
- ²⁶G. Crecelius, G. K. Wertheim, and D. N. E. Buchanan, Phys. Rev. B <u>18</u>, 6519 (1978).
- ²⁷A. Platau and S.-E. Karlson, Phys. Rev. B <u>18</u>, 3820 (1978).
- ²⁸R. R. Joseph and K. A. Gschneidner, Jr., in Proceedings of the Seventh Rare Earth Research Conference, Coronado, California, October 28-30, 1968 (unpublished), pp. 7-16.