

Phenomenological derivation of nonmonotonic temperature dependences in antiferromagnetic superconductors using a two-fluid model

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We demonstrate that the nonmonotonic temperature dependences of the upper critical field and the superconducting order parameter of an antiferromagnetic superconductor can be deduced phenomenologically from a modified version of the well-known two-fluid model.

In a recent letter, Nass *et al.*¹ have obtained a nonmonotonic temperature (T) dependence for the thermodynamical critical field [$H_c(T)$] and gap parameter $\Delta(T)$ of an antiferromagnetic superconductor,² using a rigorous mean-field theory which assumes phonon-mediated pairing in the system. Without prejudice to possible interplay of nonphononic mechanisms,³ which might either favor or counteract superconductivity in this complex system, we suggest that it may be possible to deduce a nonmonotonic temperature dependence phenomenologically as exemplified below by considering a modified version of the well-known two-fluid model.⁴

In zero magnetic field, we can relate the superfluid concentration $X_0(T)$ of an ordinary superconductor to its thermodynamic critical field $H_{c0}(T)$ through,⁴

$$H_{c0}^2(T) = H_0^2 [2(T^2/T_c^2)[1 - X_0(T)]^{1/2} + X_0(T) - 2T^2/T_c^2] \quad (1)$$

where $X_0(T) = 1 - T^4/T_c^4$ and H_0 is the critical field at zero temperature. For an antiferromagnetic superconductor which manifests a molecular field^{1,2} $H_Q(T)$ below the Néel temperature T_N , we assume that the modified superfluid concentration $X(T)$ can be approximately described by

$$X(T) = X_0(T)F(H_Q/H_{c2}) \quad (2)$$

with the function $F(H_Q/H_{c2})$ chosen such that $F(0) = 1$ and $F(1) = 0$. Here, $H_{c2}(T) = \sqrt{2}K(T)H_c(T)$ is the upper critical field^{1,2} of the antiferromagnetic superconductor and the parameter $K(T)$ in the two-fluid model can be approximated⁵ by

$$K(T) = K(0)(1 + T^2/T_c^2)^{-1/2}(1 + T/T_c)^{-1/2} .$$

Following Nass *et al.*,¹ we assume $T_N = T_c/2$, but with an alternative description $H_Q(T) = H_Q(0) \times (1 - T^2/T_N^2)^{1/2}$. Further we choose $F(H_Q/H_{c2}) = (1 - H_Q/H_{c2})^2$ as a trial function. Under these modifications, the thermodynamical critical field $H_c(T)$ of the antiferromagnetic superconductor will follow from Eq. (1) if we replace $X_0(T)$ by the modified value $X(T)$.

In terms of the reduced variables, $t = T/T_c$, $h(t) = H_c(T)/H_0$ and $h_q(0) = H_Q(0)/[\sqrt{2}K(0)H_0]$, the expression for the critical field of an antiferromagnetic superconductor would become

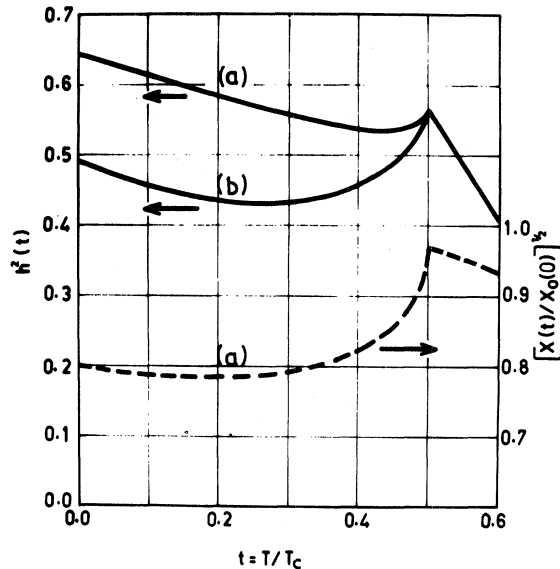


FIG. 1. Temperature dependence of $h(t)$ and $X(t)/X_0(0)$ of an antiferromagnetic superconductor for $h_q(0) = 0.16$ (curves a) and 0.21 (curve b).

$$h^2(t) = (1-t^4) \left[1 - \frac{h_q(0)}{h(t)} (1+t^2)^{1/2} (1+t)^{1/2} (1-4t^2)^{1/2} \right]^2 + 2t^2 \left[1 - (1-t^4) \left[1 - \frac{h_q(0)}{h(t)} (1+t^2)^{1/2} (1+t)^{1/2} (1-4t^2)^{1/2} \right]^2 \right]^{1/2} - 2t^2 . \quad (3)$$

For $T \geq T_N$, $H_Q(T) = 0$ and the above equation will simply reduce to the familiar parabolic form.⁴

In Fig. 1, we depict the variation of $h^2(t)$ with t for $h_q(0) = 0.16$ and 0.21 as obtained from Eq. (3) and also the temperature dependence of $[X(t)/X_0(0)]^{1/2}$ as obtained from Eq. (2) for $h_q(0) = 0.16$. We find that in spite of the crudeness of the various approximations used, the nonmonotonic temperature dependence^{1,2} of the thermodynamical critical field and of the superfluid order parameter is also qualitatively

borne out by this simple two-fluid picture, irrespective of the pairing and depairing mechanisms that might be operative in an antiferromagnetic superconductor.

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