

Superconductivity in the ternary intermetallics YbPd₂Ge₂, LaPd₂Ge₂, and LaPt₂Ge₂

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Superconductivity is found to exist in the ternary intermetallics YbPd₂Ge₂, LaPd₂Ge₂, and LaPt₂Ge₂. Superconductivity-onset temperatures are 1.17, 1.12, and 0.55 K, respectively. Superconductivity was found neither in YbCu₂Ge₂ nor in bulk CeCu₂Si₂, the latter in contrast to previously published work.

I. INTRODUCTION

CeCu₂Si₂ appears to be a mixed-valent compound with the ThCr₂Si₂ structure in which transport and heat-capacity measurements show that the charge carriers have exceptionally large effective masses, two orders of magnitude greater than typically found in *d*-band metals. It has recently been reported to become superconducting,¹ a surprising result for such strongly interactive quasiparticles, and this is suggestive that the pairing interaction might be due to a nonconventional mechanism. We have been unable to reproduce the observed superconductivity in CeCu₂Si₂ above 0.47 K, the limit of our cryostat, and in conjunction with Dynes, have found the indication of only a trace of superconductivity down to 0.06 K, most likely due to the presence of a small fraction of an unidentified minority phase. In the course of this investigation we have discovered superconductivity in related *RPd₂Ge₂*-ordered intermetallic compounds where *R* is a rare earth, and in LaPt₂Ge₂, which possesses the tetragonal ThCr₂Si₂ structure (BaAl₄ type),^{2,3} or a variant of this structure.³ In the ThCr₂Si₂ structure, the Th atoms (or *R* atoms) occupy one set of equivalent sites. There are many ternary phases of the form *RA₂X₂*, where *A* = Pd, Pt, Rh, Ag, Cr, Mn, Fe, Co, Ni, and Cu and *X* = Si and Ge.^{2,3} Of interest to us in this investigation was the occurrence of superconductivity in those phases not containing magnetic *3d* elements. In particular, the lattice constant data² for YbPd₂Ge₂ suggested that Yb is, or nearly is, divalent (a $4f^{14}$ ion) but charge fluctuation or mixed-valent effects ($+2 \rightleftharpoons +3$) perhaps could play an important role in the low-temperature properties.

II. EXPERIMENTAL AND RESULTS

The intermetallic compounds were prepared by induction-melting stoichiometric quantities of the constituents in vitreous carbon crucibles under one atmosphere of Ar. Melting losses were less than 1%.

The as-prepared compound of YbPd₂Ge₂ was always multiphase and it was necessary to perform annealing experiments. Annealing was carried out for all of the compounds in sealed quartz ampoules. For YbPd₂Ge₂, a chip of Yb metal was also placed in the ampoule, but isolated from the compound, to prevent decomposition and to scavenge residual oxygen. Powder x-ray diffractometry (Cr *K* α radiation) was performed to confirm the structure and to ascertain the cleanliness of the samples and metallographic studies were also performed. Superconducting critical temperatures were determined by ac induction measurements and magnetic susceptibility measurements were made using the Faraday method.

III. YbPd₂Ge₂

The superconducting onset temperatures T_c and transition widths, ΔT_c , as well as other pertinent information are shown in Table I. The best material was obtained by annealing for two weeks at 670 °C (sample No. 11). Even so, the diffractometer traces exhibited several very weak lines that could be indexed on the basis of the presence of very small amounts of Yb, Pd₄Yb₃, PdGe, Pd₂Ge, and GeO₂. This was confirmed by metallographic studies which indicated the presence of small amounts of other phases. One extra line, at $d = 2.470 \text{ \AA}$, could not be indexed on the basis of any known phase of Yb, Pd, and Ge. The intensity of this line decreased rather dramatically with annealing, but could not be completely eliminated. In sample No. 11, this line was very weak. DTA measurements showed that YbPd₂Ge₂ undergoes three solid-state transitions before melting (at 732, 878, and 990 °C; evidence of melting obtained at 1425 °C, the limit of the DTA apparatus). We believe that the 2.470- \AA line may be related to the presence of a small amount of one of the untransformed elevated temperature phases.

Magnetic susceptibility studies of as-prepared and annealed samples all showed Curie-Weiss contributions to the susceptibility, but the magnitude of this

TABLE I. Superconductivity results for YbPd₂Ge₂.

Sample No.	T_c (K)	ΔT_c	Fraction of sample superconducting (compared to Nb powder)	Remarks
1	1.33	0.31	0.5	As prepared
2	1.30	0.22	0.7	Annealed 24 h at 600 °C; no Yb atmosphere
3	1.23	0.18	0.8	Annealed 1 week at 600 °C; no Yb atmosphere
4	1.30	0.28	0.7	Annealed at 1000 °C for 24 h; no Yb atmosphere
5	1.18	0.17	1	Annealed 650 °C/114 h; no Yb atmosphere
6	1.27	0.30	0.7	As prepared
7	1.17	0.13	0.9	Annealed at 650 °C for 114 h; no Yb atmosphere
8	1.18	0.14	0.9	Annealed at 650 °C for 114 h; Yb atmosphere
9	1.27	0.31	0.6	No anneal
10	1.14	0.12	1	650 °C for 114 h; Yb atmosphere
11	1.17	0.19	1	670 °C for 2 weeks; Yb atmosphere

contribution decreased after annealing. It was the smallest for sample 11. If one assumes that the measured paramagnetic moments arise from Yb³⁺ (reasonable since LaPd₂Ge₂ shows no paramagnetic moments, see below) and that the moment of Yb³⁺ is equal to the free ion value of $4.5\mu_B/\text{Yb}$, the fraction of Yb³⁺ in each sample can be determined. In the as-prepared samples, 50% of the Yb is Yb³⁺, consistent with the large amount of second phase seen in x-ray and metallographic studies of the sample. The lowest Yb³⁺ concentration (~16%) is obtained in the sample annealed for the longest time (No. 11). The susceptibility of the majority phase in this sample may be estimated by the following procedure. The measured susceptibility is fit to the form $\chi = C/(T + \Theta) + \chi_0$ by a least-squares method from 300 to 4.2 K; we assume that the Curie-Weiss term $C/(T + \Theta)$ is due to the small amount of second

phase and that χ_0 is the temperature-independent susceptibility of the majority phase. A good fit to such an expression is obtained (standard deviation $\pm 0.3\%$) making the above assumptions plausible. By this procedure we estimate that the susceptibility of the majority phase is weakly paramagnetic, $\chi_0 = 0.25 \times 10^{-6}$ emu/g.

IV. LaPd₂Ge₂ AND LaPt₂Ge₂

Superconductivity results for these two phases are shown in Table II. No extra diffraction lines were present in the annealed sample of LaPd₂Ge₂, although the metallographic studies revealed the presence of a needlelike precipitate. Basically, the sample was single phase. The as-prepared sample differed little from the annealed and it was essentially

TABLE II. Superconductivity results for LaPd₂Ge₂ and LaPt₂Ge₂.

	T_c (K)	ΔT_c	Remarks
LaPd ₂ Ge ₂	1.12	0.04	Annealed at 670 °C for 5 days
LaPt ₂ Ge ₂	0.55	...	Annealed at 650 °C for 5 days

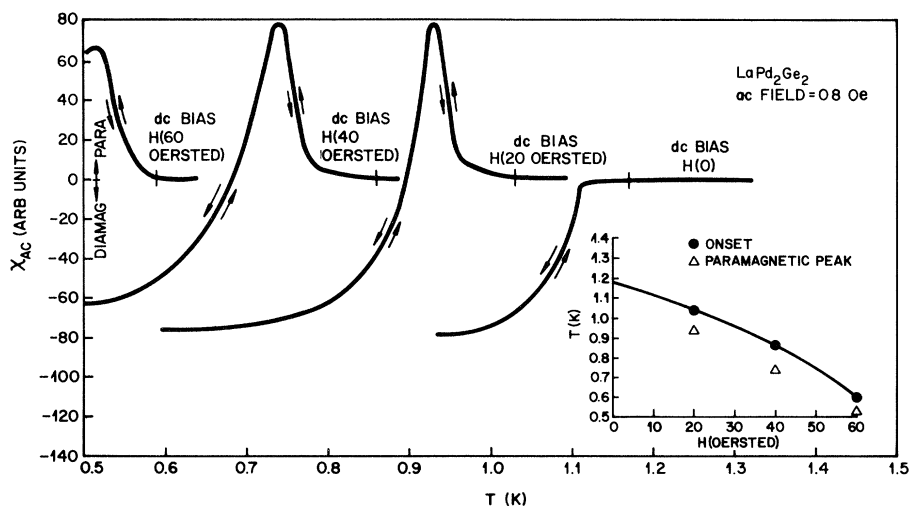


FIG. 1. Low-frequency ac susceptibility of LaPd_2Ge_2 as a function of temperature without and with a dc bias field. Note the development of a reversible differential paramagnetic signal.

single phase. The diffractometer trace of the annealed sample of LaPt_2Ge_2 showed several weak extra lines that could be indexed on the basis of La_2O_3 , Ge, and LaGe_2 . The structure of LaPt_2Ge_2 is tetragonal, but a variant of the ThCr_2Si_2 structure.⁴

The magnetic susceptibility of LaPd_2Ge_2 exhibited a small Curie contribution that is attributed to a low level of paramagnetic impurities (perhaps 20–50 ppm Fe). Otherwise the intrinsic susceptibility is temperature independent and 0.1×10^{-6} emu/g.

LaPt_2Ge_2 exhibits a weakly temperature dependent diamagnetic susceptibility below 300 K. A small Curie-Weiss contribution was observed below 50 K probably caused by a small amount of Fe impurity (20 ppm). The absolute value χ_g for this sample at room temperature is -0.02×10^{-6} emu/g $\pm 20\%$.

The susceptibilities of LaPt_2Ge_2 and of LaPd_2Ge_2 are typical of low density states metals and are slightly smaller than the estimated susceptibility of YbPd_2Ge_2 .

The superconducting transitions as measured by a low-frequency (25 Hz) ac bridge for powdered samples of LaPd_2Ge in low applied fields showed a surprising reversibility as demonstrated by the differential paramagnetic signals shown in Fig. 1. This high degree of reversibility for the entry and exit of small amounts of flux is rarely found in intermetallic compounds. The very low critical field curve shown in the inset suggests that the material may in fact be a type I superconductor.

V. CeCu_2Si_2

Steglich *et al.*¹ recently reported superconductivity in CeCu_2Si_2 ($T_c = 0.5$ K). This material exhibits

low-temperature transport properties suggestive of Ce “unstable 4*f* shell” behavior. The samples Steglich *et al.* studied were prepared by arc melting and their results were obtained on both unannealed and annealed samples. We have not been able to reproduce their superconductivity results. More specifically, a sample of CeCu_2Si_2 annealed at 900 °C for five days (similar to the Steglich *et al.* annealed sample) was normal down to 0.47 K. The sample reported by Steglich *et al.* became superconducting above 0.5 K. Further measurements in a dilution refrigerator down to 0.060 K by Dynes showed no evidence for bulk superconductivity although a noticeable decrease in resistance occurred. This decrease could be due to a smeared out superconducting transition of a minority second phase.

The diffractometer trace for our sample exhibited very weak extra diffraction peaks that could be indexed on the basis of the presence of Cu_3Si , CeCu_2 , CeCu_6 , and CuO as extra phases. The magnetic susceptibility exhibited a temperature-dependent paramagnetic contribution that did not fit Curie-Weiss behavior over any significant temperature range. This may be due to the fact that the crystal-field splittings of Ce^{3+} are frequently on the order of 100 K. On the other hand, CeCu_2 and CeCu_6 , as impurity phases may be contributing to the above susceptibility behavior. CeCu_2 has been shown to exhibit a non-Curie-Weiss temperature-dependent susceptibility⁵ and it is likely that CeCu_6 would exhibit similar behavior.

VI. YbCu_2Ge_2

A sample of YbCu_2Ge_2 , annealed at 670 °C for five days, did not exhibit superconductivity down to 0.47 K.

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