

## New magnetic compounds with Heusler and Heusler-related structures

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Magnetic and structural properties are presented for new ferromagnetic compounds of the form  $Rh_2MnX$  where  $X$  is Al, Ga, In, Tl, Sn, Pb, and Sb. The exchange is described in terms of competing ferromagnetic Mn-Rh-Mn interactions and antiferromagnetic Mn-Mn interactions.

### INTRODUCTION

A number of new magnetic compounds and alloys have been originated from the Mn-Bi-Rh ternary system. Small amounts of Rh in the compound MnBi were reported earlier to be effective in stabilizing the high-temperature phase of MnBi and thus to have application to optically addressable memory.<sup>1</sup> A new group of ferromagnetic compounds of the form  $Mn_3Rh_2Bi_4$  with a new crystal structure was reported to exist in the central region of this ternary<sup>2</sup>; detailed calculations of the magnetic properties were made by Goodenough.<sup>3</sup>

Further exploration of this ternary system, using magnetic separation as a tool for isolating magnetic compositions, revealed that  $Mn_{0.8}Bi_{0.2}Rh$  was a single-phase ferromagnetic compound.<sup>4</sup> The structure of this compound is CsCl with  $Mn_{0.8}Bi_{0.2}$  on one simple cubic sublattice and Rh on the other sublattice. This structure is that of the well-known compound MnRh where the Bi partially substitutes on the Mn sublattice. MnRh is antiferromagnetic,<sup>4</sup> whereas addition of the Bi induces ferromagnetism.

The maximum solubility of Bi in the MnRh lattice is  $Mn_{0.8}Bi_{0.2}Rh$ . This formula can also be written  $Rh_2Mn_{1.6}Bi_{0.4}$ . Both the structure and the composition of this Bi compound bring to mind the Heusler compounds. If another atom could be found with greater solubility than Bi in this lattice then perhaps new Heusler compounds could be generated of the form  $Rh_2TX$ , where  $T$  is a 3d transition-metal atom and  $X$  is a group-B atom. One example of this which has already been reported in  $Rh_2MnGe$ .<sup>5</sup>

We have found and reported earlier that letting  $X = Sn$  allows variation of the  $T$  atom across the entire 3d series including Cu. Some of these  $Rh_2T$ Sn compounds have the Heusler structure and some exhibit a curious tetragonal modification of the Heusler structure.<sup>6</sup> In the present paper we will describe the magnetic and crystallographic properties of the Heusler and Heusler-related compounds generated by variation of atom  $X$  in the manganese compound series  $Rh_2MnX$ .

### SAMPLE PREPARATION AND STRUCTURE

The compounds made were of the formula  $Rh_2MnX$ , where  $X$  is a group-III B, IV B, V B, or VI B element. The compounds were made by mixing together powders of the elements, sealing this mixture in evacuated quartz ampoules and annealing at 250 °C for 7 h, followed by 700 °C for 6 days. The samples were then pulverized and reannealed in vacuum for 5 days at 700 °C. The upper part of Fig. 1 shows an experimental x-ray diffraction pattern taken on a powder sample of  $Rh_2MnPb$  ( $X = Pb$ ). A calculated pattern is shown in the lower part of Fig. 1. This pattern was calculated for the Heusler structure with the aid of the Smith program.<sup>7</sup>

The good agreement between these patterns shows the sample to be of single phase quality and also shows that the Heusler structure is a good representation of the actual structure.

Figure 2 shows the Heusler structure for  $Rh_2MnX$ . The Rh atoms occupy the corners of a simple cubic sublattice. The atoms of the second simple cubic sublattice occupy the body-center positions of the first sublattice. In the fully or-

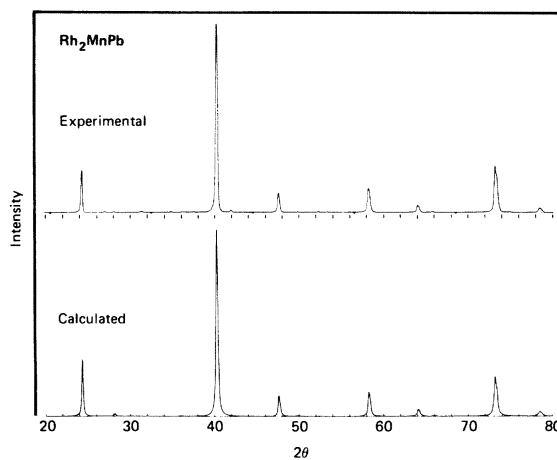
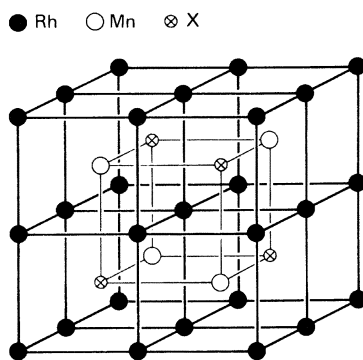


FIG. 1. Experimental and calculated x-ray diffraction pattern for a powder sample of  $Rh_2MnPb$ .

FIG. 2. Fully ordered Heusler structure,  $Rh_2TX$ .

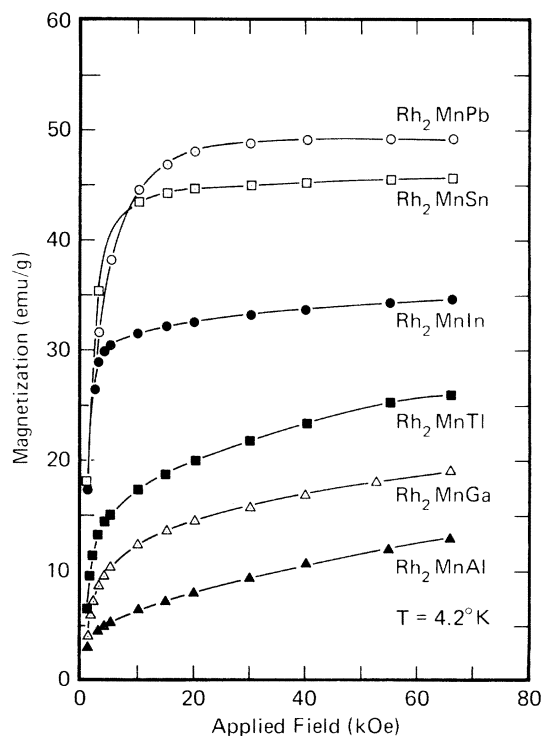
dered Heusler structure the second sublattice is occupied by Mn and X atoms (e.g.,  $X = Pb$ ) ordered as shown.

Three variations of the fully ordered Heusler structure have been observed. The first, already mentioned, is with Mn and Bi on the second sublattice as in  $Mn_{0.8}Bi_{0.2}Rh$ . Another variation is where the Mn and X atoms of the second sublattice are completely disordered on this second sublattice (e.g.,  $Rh_2MnAl$ ). A third variation is where the unit cell of Fig. 2 is stretched along a cube edge, say the  $c$  direction, to produce a unit cell which is tetragonal.<sup>6</sup> The cell shrinks along the  $a$  and  $b$  directions to maintain a nearly constant volume. The tetragonal distortion is large ( $c/a \sim 1.2$ ).  $Rh_2MnSb$  is the only known example of this structure among the Mn compounds but the structure is seen often when Mn is replaced by other  $3d$  transition metal elements.<sup>6</sup>

#### MAGNETIC MEASUREMENTS

Magnetization measurements taken at 4.2 °K on powder samples are shown in Fig. 3. These data were taken with a vibrating sample magnetometer, using a superconducting solenoid with a maximum field of about 70 kOe. In the compounds shown in Fig. 3,  $X$  belongs to group IIIB (Al, Ga, In, Tl) or group IVB (Pb, Sn). Not only is the magnetization lower for the group-III B compounds, but the approach to saturation is more gradual. The group-III B compounds are not really saturated even at 65 kOe.

Figure 4 shows inverse susceptibility versus temperature data taken on the same samples. These data are taken also with a vibrating-sample magnetometer but the field is provided by a 12-in. electromagnet. The system calibrations were checked by making measurements of the susceptibility and Curie temperature of pure nickel. The paramagnetic Curie temperatures  $\Theta$  given by the horizontal axis intercepts in Fig. 4 are much low-

FIG. 3. Magnetization vs applied magnetic field at 4.2 °K for  $Rh_2MnX$  compounds.

er for the group III B compounds. The  $Rh_2MnAl$  compound shows significant curvature in its  $1/\chi$ - $T$  curve.

Table I gives crystallographic data and Curie temperatures for the  $Rh_2MnX$  compounds. For the fully ordered Heusler structure the lattice constant  $a$  is given, whereas for the disordered Heusler structure  $B2$ , the double lattice constant is given to facilitate direct comparison with the  $L2_1$  compounds. The paramagnetic Curie temperature  $\Theta$  was taken from  $1/\chi$ - $T$  plots, and the ferromagnetic Curie temperature  $T_C$  was measured as the temperature of the abrupt decrease in magnetization with increasing temperature while the sample is maintained in a small constant field. Data shown for  $X = Bi$  and  $Te$  are actually for the compositions  $Rh_2Mn_{1.6}Bi_{0.4}$  and  $Rh_2Mn_{1.6}Te_{0.4}$ .

Table II gives the results of magnetization measurements on the  $Rh_2MnX$  compounds. The moment data are in terms of a formula unit rather than per Mn atom, since some moment will probably reside on the Rh atoms as in the case of  $FeRh$ . The determination of the Mn moment must await neutron diffraction measurements.

#### DISCUSSION

The data on  $X = Bi$  and  $X = Te$  (Table I) were reported earlier<sup>4</sup> and are included for completeness.

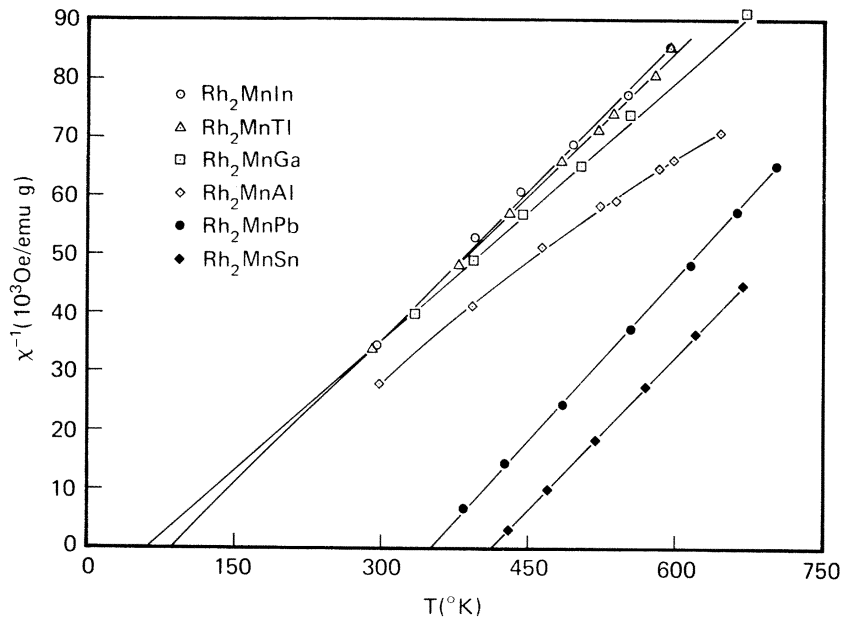


FIG. 4. Inverse susceptibility vs temperature for  $\text{Rh}_2\text{MnX}$  compounds.

$\text{Rh}_2\text{MnGe}$  reported earlier by Hames and Crangle<sup>5</sup> and our data (Tables I and II) are in substantial agreement. The remaining compounds are believed to be reported here for the first time. For  $X = \text{Sb}$  the structure is the same tetragonal modification of the fully ordered Heusler structure exhibited by and reported earlier for some of the  $\text{Rh}_2\text{TlSn}$  compounds.<sup>6</sup> Table I shows that for  $X$  in group

IVB ( $\text{Rh}_2\text{MnPb}$ ,  $\text{Rh}_2\text{MnSn}$ , and  $\text{Rh}_2\text{MnGe}$ ) the structure is the fully ordered Heusler structure  $L2_1$  (Figure 2), whereas for  $X$  in group IIIB ( $\text{Rh}_2\text{MnTl}$ ,  $\text{Rh}_2\text{MnIn}$ ,  $\text{Rh}_2\text{MnGa}$ , and  $\text{Rh}_2\text{MnAl}$ ) the structure is the disordered Heusler structure  $B2$  where the Mn and X atoms are disordered. The driving force for ordering in this system is a strong function of the valence of element X. The

TABLE I. Crystal structure and Curie temperature for  $\text{Rh}_2\text{MnX}$  compounds.

| IIIB ( $s^2p$ )                        | IVB ( $s^2p^2$ )                      | VB ( $s^2p^3$ )                       | VIB ( $s^2p^4$ )                      |
|--|---------------------------------------|---------------------------------------|---------------------------------------|
| Al                                     | Si                                    | P                                     | S                                     |
| $B2$                                   | multiphase                            |                                       |                                       |
| $2a = 6.005 \text{ \AA}$               |                                       |                                       |                                       |
| $\Theta = \text{nonlinear}$            |                                       |                                       |                                       |
| $T_C = 85-105 \text{ }^\circ\text{K}$  |                                       |                                       |                                       |
| Ga                                     | Ge                                    | As                                    | Se                                    |
| $B2$                                   | $L2_1$                                | multiphase                            |                                       |
| $2a = 6.056 \text{ \AA}$               | $a = 6.030 \text{ \AA}$               |                                       |                                       |
| $\Theta = 60 \text{ }^\circ\text{K}$   | $\Theta = 460 \text{ }^\circ\text{K}$ |                                       |                                       |
| $T_C = 75-85 \text{ }^\circ\text{K}$   | $T_C = 450 \text{ }^\circ\text{K}$    |                                       |                                       |
| In                                     | Sn                                    | Sb                                    | Te                                    |
| $B2$                                   | $L2_1$                                | tetragonal (new)                      | limited solubility                    |
| $2a = 6.287 \text{ \AA}$               | $a = 6.252 \text{ \AA}$               | $a = 5.898 \text{ \AA}$               | $a = 6.184 \text{ \AA}$               |
| $\Theta = 81 \text{ }^\circ\text{K}$   | $\Theta = 413 \text{ }^\circ\text{K}$ | $c = 6.987 \text{ \AA}$               | $\Theta = 150 \text{ }^\circ\text{K}$ |
| $T_C = 100-110 \text{ }^\circ\text{K}$ | $T_C = 412 \text{ }^\circ\text{K}$    | $\Theta = 312 \text{ }^\circ\text{K}$ | $T_C = -$                             |
|  |                                       | $T_C = 304 \text{ }^\circ\text{K}$    |                                       |
| Tl                                     | Pb                                    | Bi                                    | Po                                    |
| $B2$                                   | $L2_1$                                | limited solubility                    |                                       |
| $2a = 6.324 \text{ \AA}$               | $a = 6.332 \text{ \AA}$               | $a = 6.220 \text{ \AA}$               |                                       |
| $\Theta = 90 \text{ }^\circ\text{K}$   | $\Theta = 348 \text{ }^\circ\text{K}$ | $\Theta = 185 \text{ }^\circ\text{K}$ |                                       |
| $T_C = -$                              | $T_C = 338 \text{ }^\circ\text{K}$    | $T_C = -$                             |                                       |

TABLE II. Magnetization data for  $Rh_2MnX$  compounds.  $p$  is the effective moment per formula unit calculated from the slope of  $1/\chi$  vs  $T$ ;  $\mu_1$  is the moment calculated from  $p = 2[\frac{1}{2}\mu_1 + (1 + \frac{1}{2}\mu_1)^{1/2}]$ ;  $\sigma$  is the magnetization;  $H_c$  is the coercive force; and  $\mu$  is the moment per formula unit calculated from  $\sigma$ .

| X  | Al                | Ga                | In                | Tl   | Ge   | Sn   | Pb   |
|--|-------------------|-------------------|-------------------|------|------|------|------|
| $p(\mu_B)$                                 | ~4.1              | 4.23              | 4.28              | 4.75 | 4.64 | 4.19 | 4.49 |
| $\mu_1(\mu_B)$                             | ~3.3              | 3.35              | 3.40              | 3.86 | 3.75 | 3.30 | 3.60 |
| $\sigma(\text{emu/g})$ ,<br>5 °K, 6 kOe    | 13.1 <sup>a</sup> | 19.8 <sup>a</sup> | 34.5 <sup>a</sup> | 26.6 | 77.5 | 45.6 | 49.2 |
| $\mu(\mu_B)$ ,<br>5 °K, 66 kOe             | >0.7              | >1.2              | >2.3              | >2.2 | 4.62 | 3.10 | 4.12 |
| $\sigma(\text{emu/g})$ ,<br>300 °K, 19 kOe | ...               | ...               | ...               | ...  | 64.2 | 33.6 | 28.2 |
| $H_c(\text{Oe})$ ,<br>300 °K               | ...               | ...               | ...               | ...  | 70   | 38   | ...  |

<sup>a</sup>Not saturated.

ordering is not a function of lattice parameter since the lattice constants shown in Table I overlap between the X group-III $B$  and X group-IV $B$  columns. Therefore we infer that the valence of element X influences the atomic ordering primarily by variation of electron concentration. The ordering does not depend much on cooling rate since X group-III $B$  samples were very slowly cooled without inducing  $L2_1$  order, and X group-IV $B$  samples were quenched without inducing  $B2$ -type disorder.

Also a strong function of the valence of element X are the magnetic properties. For X in group IV $B$  the compounds are all ferromagnetic with relatively high Curie temperatures. The Curie temperatures decrease with increasing lattice constant (Table I). For X in group III $B$  the Curie temperatures are a factor of 4 to 5 smaller, and the Curie temperatures increase with increasing lattice constant.

It is believed that the principal exchange interaction in these compounds is a ferromagnetic Mn-Rh-Mn exchange. A second exchange mechanism which may occur is an antiferromagnetic Mn-Mn exchange between nearest-neighbor Mn. For X in group IV $B$  there are no nearest-neighbor Mn and the compounds exhibit a high Curie temperature and rapid magnetic saturation (Fig. 3). For X in group III $B$  there are many antiferromagnetic Mn-Mn interactions present due to atomic disorder and therefore the Curie temperature is reduced. This may explain the slow approach to saturation of the X group-III $B$  compounds.

The Mn-Mn exchange will fall off more rapidly with distance than the Mn-Rh-Mn exchange due to the larger  $d$  orbitals of the Rh. This causes the Curie temperature to increase with increasing separation for X in group III $B$  (rapid reduction

in Mn-Mn exchange) and to decrease with increasing distance for X in group IV $B$  (slow reduction in Mn-Rh-Mn exchange).

One consequence of this model is that the magnetic moment at low temperatures should be low for the X group-III $B$  compounds due to the presence of antiparallel spins, and Fig. 3 shows that this is the case. For the X group-IV $B$  compounds the low-temperature moments are comparable to the moments calculated from high-temperature susceptibility.

## CONCLUSION

A number of new magnetic compounds have been found. These compounds have the formula  $Rh_2MnX$  where X = Al, Ga, In, Tl, Sn, Pb, and Sb. For X in group IV $B$  the compounds are ferromagnetic with Curie temperatures above room temperature and exhibit large magnetic moments at low temperature and large effective moments above the Curie temperature. For X in group III $B$  the compounds predominantly are ferromagnetic with low Curie temperatures and low moments at low temperature. The difference in atomic order between X group-III $B$  and X group-IV $B$  compounds is attributed to the variation of electron concentration by atom X, and the difference in magnetic behavior between X group III $B$  and X group-IV $B$  compounds is ascribed to the presence of antiferromagnetic Mn-Mn exchange interactions due to variation of this atomic order.

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<sup>1</sup>Kenneth Lee, J. C. Suits, and G. B. Street, *Appl. Phys. Lett.* 26, 27 (1975).

<sup>2</sup>G. Bryan Street, J. C. Suits, and Kenneth Lee, *Solid State Commun.* 14, 33 (1974).

<sup>3</sup>J. C. Suits, G. B. Street, Kenneth Lee, and J. B. Goodenough, *Phys. Rev. B* 10, 120 (1974).

<sup>4</sup>J. C. Suits, *IBM J. Res. Dev.* 19, 422 (1975).

<sup>5</sup>F. A. Hames and J. Crangle, *J. Appl. Phys.* 42, 1336 (1971).

<sup>6</sup>J. C. Suits, *Solid State Commun.* 18, 423 (1976).

<sup>7</sup>D. K. Smith, Lawrence Radiation Laboratory Report No. UCRL-50264 (unpublished).