


Towards understanding the electronic structure of the simpler members of two-dimensional halide-perovskites

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In this paper, we analyze the band structure of two-dimensional (2D) halide perovskites by considering structures related to the simpler case of the series, $(BA)_2PbI_4$, in which PbI_4 layers are intercalated with butylammonium [$BA = CH_3(CH_2)_3NH_3$] organic ligands. We use density-functional-theory (DFT) based calculations and tight-binding (TB) models aiming to discover a simple description of the bands within 1 eV below the valence-band maximum and 2 eV above the conduction-band minimum, which, including the energy gap, is about a $\Delta E = 5$ eV energy range. The bands in this ΔE range are those expected to contribute to the transport phenomena, photoconductivity, and light emission in the visible spectrum, at room and low temperature. We find that the atomic orbitals of the butylammonium chains have negligible contribution to the Bloch states which form the conduction and valence bands in the above defined ΔE range. Our calculations reveal a rather universal, i.e., independent of the intercalating BA, rigid-band picture inside the above ΔE range characteristic of the layered perovskite “matrix” (i.e., PbI_4 in our example). Besides demonstrating the above conclusion, the main goal of this paper is to find accurate TB models which capture the essential features of the DFT bands in this ΔE range. First, we ignore electron hopping along the c axis and the octahedral distortions and this increased symmetry (from C_2 to C_4) halves the Bravais lattice unit cell size and the Brillouin zone unfolds to a 45° rotated square and this allows some analytical handling of the 2D TB Hamiltonian. The Pb 6s and I 5s orbitals are far away from the above ΔE range and, thus, we integrate them out to obtain an effective model which only includes hybridized Pb 6p and I 5p states. Our TB-based treatment (a) provides a good quantitative description of the DFT band structure, (b) helps us conceptualize the complex electronic structure in the family of these materials in a simple way, and (c) yields the one-body part to be combined with appropriately screened electron interaction to describe many-body effects, such as excitonic bound states.

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I. INTRODUCTION

The discovery and production of semiconducting superlattices has led to a significant advancement in solid-state physics and electronic technology. Further development, however, has been hindered by the infrastructure-demanding time-consuming fabrication procedure for precisely assembling the nanometer-scale structures in such artificial material structures. Rather recently, a new class of superlattices has been discovered, known as Ruddlesden-Popper (RP) two-dimensional (2D) halide-perovskites, that can be self-assembled using wet-chemistry synthesis [1,2]. The interest in this type of materials is not just because of the simplicity of the synthesis process, but more importantly, the interest is due to the fact that they are clean and maintain the periodic structure at the atomic level. These 2D halide-perovskites are analogues to the oxide perovskites described by Ruddlesden and Popper [3], from where they gain their name, and they form self-assembled superlattice structures [1,4–6]. Figure 1 illustrates the relationship of the Ruddlesden-Popper perovskite structures to the three-dimensional (3D) perovskite ABX_3 , where A, B, and X are a small organic cation (such as methylammonium), a metal cation, and a halogen anion, respectively. The bulk perovskite ABX_3 is figuratively sliced apart to form a multilayer sandwich in which amine-terminated organic ligands

A' (such as, butylammonium) are “inserted” in between every m ($m = 1, 2, \dots$) adjacent halide-perovskite layers (see Fig. 1) in the formula $A'_2A_{m-1}B_mX_{3m+1}$. An example of this series is the family [1,2] $(BA)_2(MA)_{m-1}Pb_mI_{3m+1}$ with $BA = CH_3(CH_2)_3NH_3$ and $MA = CH_3NH_3$ ($m = 1, 2, 3, 4, \dots$). These Ruddlesden-Popper hybrid perovskites come with inherent low dimensionality and highly ordered periodic nanostructures. The long amine-terminated organic ligands essentially cleave the 3D perovskite lattices into two-dimensional sheets, forming alternating layers of organic and inorganic supercells along the c -axis crystal orientation. The formation of such layered structures is thermodynamically favorable, which makes it possible to obtain Ruddlesden-Popper hybrid perovskites conveniently using wet-chemistry synthesis.

The family of 3D organic halide-perovskites has received tremendous attention during the past decade due to their promising optoelectronic and solar cell applications. Decreasing the dimensionality from 3D to 2D, more versatile organic cations can be incorporated as templates to produce new structures [7]. The properties of these reduced dimensionality semiconductors are less widely studied but results of many such studies have started to appear in the literature rapidly. Among the many fascinating properties of these materials is white-light emission [8], where the broad emission

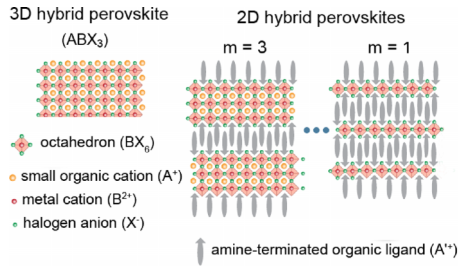


FIG. 1. Illustration of the relationship of the Ruddlesden-Popper perovskite structures to the three-dimensional (3D) perovskite ABX_3 . See text for details.

comes from the transient photo-excited states generated by self-trapped excitons. Furthermore, these materials provide a significant degree of flexibility in tuning their optoelectronic properties by varying the number of perovskite layers and by choosing an appropriate organic ligand. This makes them particularly suitable for several photovoltaic applications [9] and as light emitters [8,10–13]. Fast pump-probe spectroscopies have also revealed useful information about the carrier dynamics and recombination in these RP perovskite materials [14–17], as well as effects of excitonic many-body interactions [18].

Examining them from a different view angle, these RP perovskites can potentially provide a new playground for fundamental physics. Namely, these nearly ideal 2D structures add another interesting family of materials to a growing list of interesting 2D materials, such as, Graphene and its variants and its cousin materials, transition metal dichalcogenides, etc., which can host a variety of interesting phenomena, and potentially new phases of matter. Understanding the physics beyond single-electron phenomena, and in order to make further progress, requires understanding the band structure of these materials at a deeper yet simpler level than the complex multiband picture provided by the DFT calculation. However, the number of atoms in a single unit cell of the Bravais lattice is very large. For example, the unit cell of the Bravais lattice of the $m = 1$ structure, which is the subject of the present paper, contains 156 atoms. Therefore the complexity of the atomic structure of the material, which is reflected in its band structure, might be a reason for not finding them appealing for theoretical investigations of potential novel phenomena. Namely, at first sight, it might seem a hopeless task to try to find a simple picture to describe the electronic structure.

There are numerous publications where DFT and related techniques have been applied in order to understand various aspects of these and related materials. The aim of the present paper is not to add another such study but rather to analyze the known complex band structure [19] of $(BA)_2PbI_4$ and related materials and find a simple (and, if possible, analytical) and accurate way to describe the origin of its features. Interestingly, we find that in the simplest case of $(BA)_2PbI_4$, it is possible to achieve this goal. First, we show that a rigid-band description is accurate for such materials. For example, we find that all the bands within 2 eV below the valence-band maximum (VBM) and 2 eV above the conduction-band minimum (CBM) have negligible contribution from the atomic orbitals contained in the amine-terminated organic ligands.

This is not to say that these larger organic ligands do not play a significant role in various aspects of the crystal formation. For example, the choice of these organic spacers is important in achieving good quality crystallization [20,21] and allows optimization of the film quality [9], and in achieving the crystal orientation and the stability of the system [4,22]. However, once the structure is formed and the positions of all the metal and halogen atoms are given, our calculation illustrates that the Bloch states with energy within 1 eV below the VBM and 2 eV above the CBM, have negligible projection to the atomic orbitals of these large organic ligands. The role of these larger butylammonium organic ligands is simply to act as a charge reservoir which fill completely the highest occupied bands making the material an insulator.

The second part of the present paper is to uncover a simple picture of the band structure responsible for most of the optoelectronic response. We consider a 2D TB model, which ignores the small octahedra distortions and this allows us to reduce the size of the unit cell by a factor of two, a fact that doubles the Brillouin zone (BZ) by unfolding it because of symmetry. This reduces the TB Hamiltonian to a 12×12 matrix for each point in the BZ. The most important conduction and valence bands are obtained as a hybridization of mostly metal-ion p and halogen p orbitals. In the case of our example, $(BA)_2PbI_4$, hybridization occurs between Pb and I p orbitals. To obtain the correct dispersion of the highest valence band, we also need to involve the role of the sp hybridization between the metal-ion s orbital and the p orbital of the halogen atoms that sit at the octahedra corners. Finally, we offer a simple analytic description of the band structure by integrating out this metal-ion s orbital as it is energetically well below the Fermi energy.

The paper is organized as follows. In Sec. II, we present the results of our DFT calculation (including the projection of the Bloch states to atomic orbitals) for those crystalline structures which we think are relevant for the point to be made. In Sec. III, we detail our tight-binding model and how it fits with the DFT results of the bands and the orbital character of the Bloch wave functions. In Sec. IV, we present the analytical model that gives a good approximation to the bands near the Fermi level and we also give the final fit of our TB model to the DFT bands. In Sec. VII, we present our final remarks and conclusions of our study.

II. THE CRYSTAL AND BAND STRUCTURE

A. Crystal structure

Figure 2 illustrates the crystal structure of the low-temperature structure [23] of $(BA)_2PbI_4$ layered material. There are PbI_4 perovskite layers which are intercalated by butylammonium chains. Notice that the atomic positions in two successive PbI_4 perovskite layers are staggered relative to the atomic positions of the same atoms in their nearest neighboring layer. As a result the unit cell of the Bravais lattice is twice as long along the c axis.

B. $(BA)_2PbI_4$ band structure

First, we carried out calculations for the $(BA)_2PbI_4$ structure [19] shown in the left panel of Fig. 2 using the QUANTUM

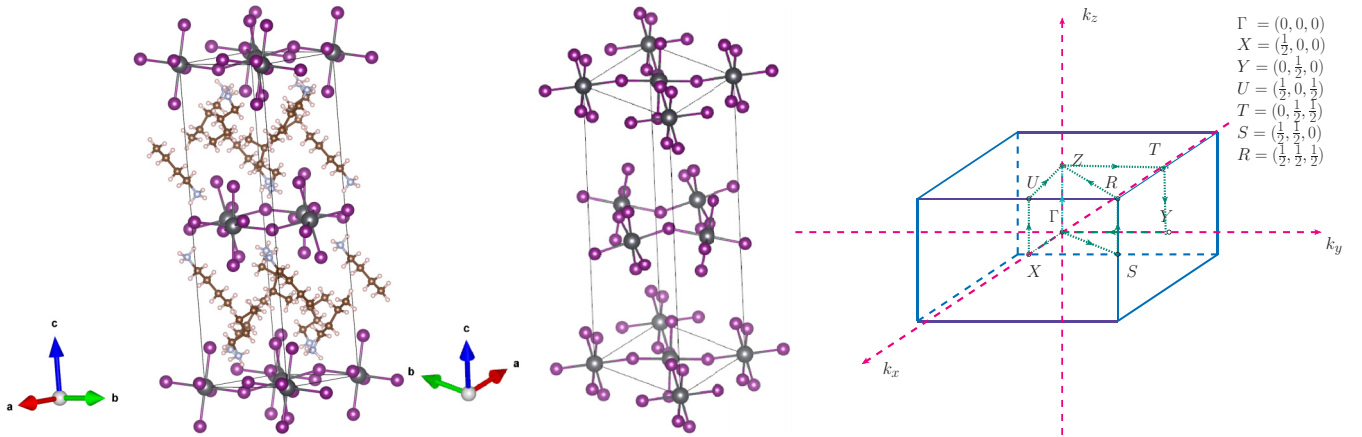


FIG. 2. (Left) Crystal structure of $(BA)_2PbI_4$. Grey (purple) spheres denote the Pb (I) atoms. The gold spheres represent the C atoms bonded to either each other or to small red (or bluish) hydrogen (or nitrogen) atoms and, thus, form the butylammonium chains which intercalate the PbI_4 layers. (Middle) Halide-perovskite matrix. (Right) The Brillouin zone and the path connecting the specific high-symmetry points used to compute the band structure.

ESPRESSO [24] (QE) implementation of the density functional theory (DFT) in the GGA framework. The Perdew-Burke-Ernzerhof (PBE) exchange correlation functional [25] was used with Projector Augmented-Wave [26] pseudopotentials generated with a scalar-relativistic calculation local potential using the “atomic” code by Dal Corso [27]. Our self-consistently converged ground-state calculation used a $8 \times 6 \times 3$ k -point mesh and a 30 Rd energy cutoff. In Appendix A, a convergence study demonstrates that the k -point mesh and the energy cutoff used are accurate for the purpose of the present paper.

Our first objective is to establish that the atomic orbitals of the butylammonium organic ligands, i.e., the orbitals of C, N, and H, have negligible contribution to the Bloch states with energy which falls in the energy window of 1 eV below the top of the VBM and 2 eV above the CBM, which, including the gap is about a 5 eV range. Figure 3 illustrates that there is insignificant projection of each of the Bloch states to the local orbitals of all C, N, and H.

Using the same exchange correlation functional and pseudopotentials as in the case of $(BA)_2PbI_4$, we carried out a self-consistent-field (SCF) DFT ground state calculation for

the “bare-bone” perovskite atomic matrix illustrated in the middle panel of Fig. 2, i.e., without the intercalating butylammonium chains. We then computed the bands along the same crystallographic directions as for the $(BA)_2PbI_4$ for comparison. These bands are shown in Fig. 4 as blue circles and are compared with those of the complete material $(BA)_2PbI_4$ (shown as red lines).

Notice that the agreement between the bands near the Fermi level is very good. The position of the Fermi level is different for the real material $(BA)_2PbI_4$ as compared to the simple PbI_4 matrix because the intercalating butylammonium chains add more electrons to these bands, thus, raising the Fermi level and making it a band insulator. The important conclusion is that the bands near the Fermi level and their Bloch wave functions, assuming the same atomic positions, are determined by the PbI_4 matrix to a good degree of accuracy.

C. Projecting Bloch states near the Fermi energy to atomic orbitals

Next, we wish to determine the orbital character of the bands inside the energy window of our interest. Since we

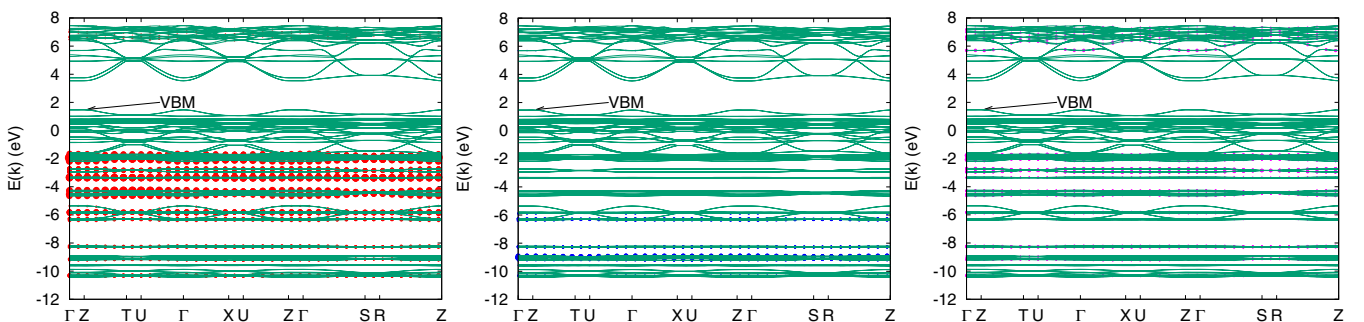


FIG. 3. Projections of the Bloch states to all atomic orbitals of C (left), N (middle), and H (right) present in the butylammonium ligands. All the bands are illustrated with both green lines and red (C), blue (N), or magenta (H) symbols. The size of these symbols is proportional to the magnitude-squared of the projection of each band (for each \vec{k} point) on these orbitals. As a result, because for many bands the projection is negligible, which gives rise to invisibly small size of these symbols, only the lines are visible for these bands. Notice that there is no visible contribution of these orbitals to the bands within the 5 eV energy range of our focus, i.e., from 2 eV above the CBM to 1 eV below the CBM.

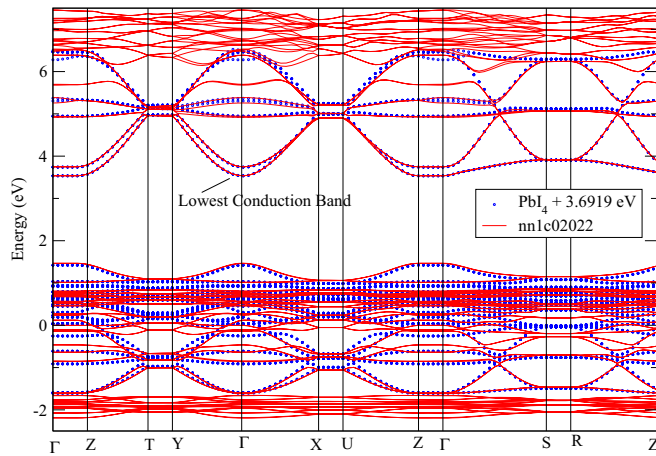


FIG. 4. Comparison of the band structure between the $(BA)_2PbI_4$ material and the PbI_4 matrix. Notice that we have shifted the band energy scale of PbI_4 by 3.6919 eV.

have demonstrated in the previous section that these bands are almost completely determined by the Pb and I atoms we projected the bands in the orbitals of those atoms only. In Figs. 5 and 6, the projection of the bands in the atomic orbitals are shown. The size of the circle is proposal to the contribution of the particular orbital to the given band. The orbitals chosen are those that contribute to the bands and lie within ± 10 eV

from the Fermi level. These are the Pb outer orbitals, i.e., $6s$, $4f$, $5d$, and $6p$ and the I $5s$, $4d$, and $5p$. Notice that the projections for the case of $(BA)_2PbI_4$ (left panels), and those of the PbI_4 halide-perovskite “matrix” (right panels) are very similar. In addition, we note that top valence bands and the lower conduction bands (i.e., with band energy less than about 5 eV and above 0 eV) are mostly made out of p (Pb or I) orbitals. There is a very small amount of s orbital contribution.

D. Removing the organic molecules and adding Cs at the location of N sites: Cs_2PbI_4

In order to clearly demonstrate the almost irrelevance of the butylammonium chains on the band structure near the Fermi level, we calculated the band structure of the a compound obtained by removing all the butylammonium chains and by adding Cs at the N sites of the original $(BA)_2PbI_4$ compound. We choose Cs based on its electronegativity relative to the Fermi level of the PbI_4 matrix, in order to make sure that the energy level of the outer s Cs level falls above the energy of the highest occupied band, such that these Cs s states will be empty inside the compound, thus, resulting in doping the PbI_4 matrix in the same manner as the alkylammonium chains cause doping. The structure is illustrated in the left panel of Fig. 7 and the DFT calculated band structure is compared to the $(BA)_2PbI_4$ compound in right panel of Fig. 7.

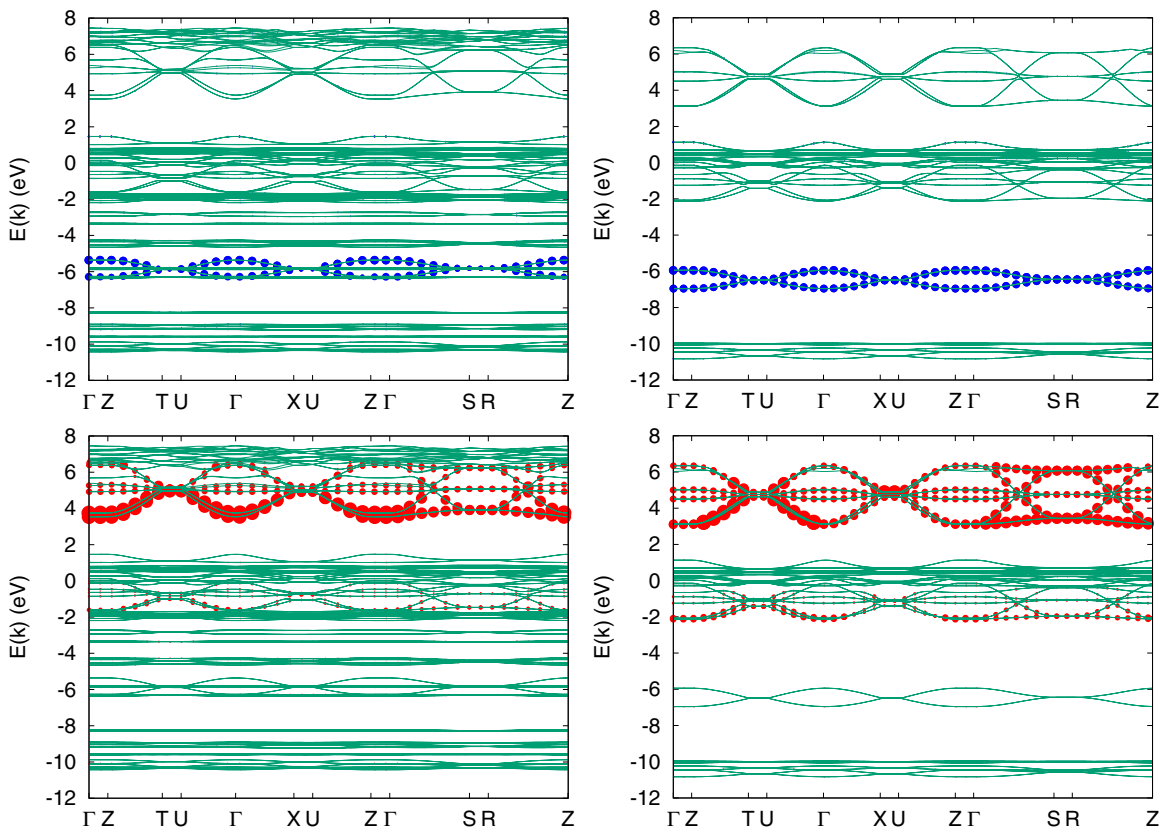


FIG. 5. Projections of the Bloch states to local Pb $6s$ (top) and $6p$ (bottom) orbitals. See caption of Fig. 3 for explanation of the symbol and lines in the plots. The panels on the left correspond to the $(BA)_2PbI_4$, while those on the right correspond to the PbI_4 halide-perovskite “matrix.” The highest occupied band is the one which starts at the Γ point at just above 1 eV on this scale.

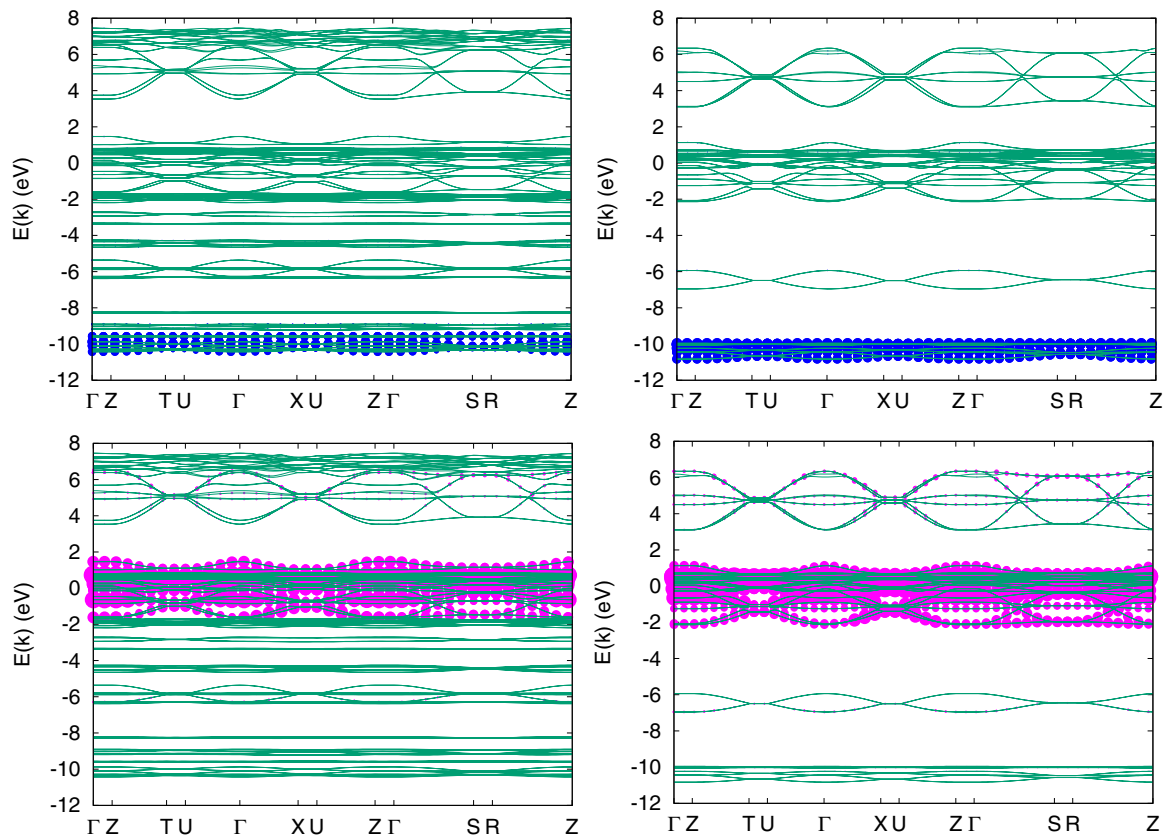


FIG. 6. Projections of the Bloch states to local I $5s$ (top) and I $5p$ (bottom) orbitals. All the bands are illustrated with both lines and blue or magenta symbols. See caption of Fig. 3 for explanation of the symbol and lines in the plots. The panels on the left row correspond to the $(BA)_2PbI_4$, while those on the right correspond to the PbI_4 halide-perovskite “matrix.” The highest occupied band is the one which starts at the Γ point at just above 1 eV on this scale.

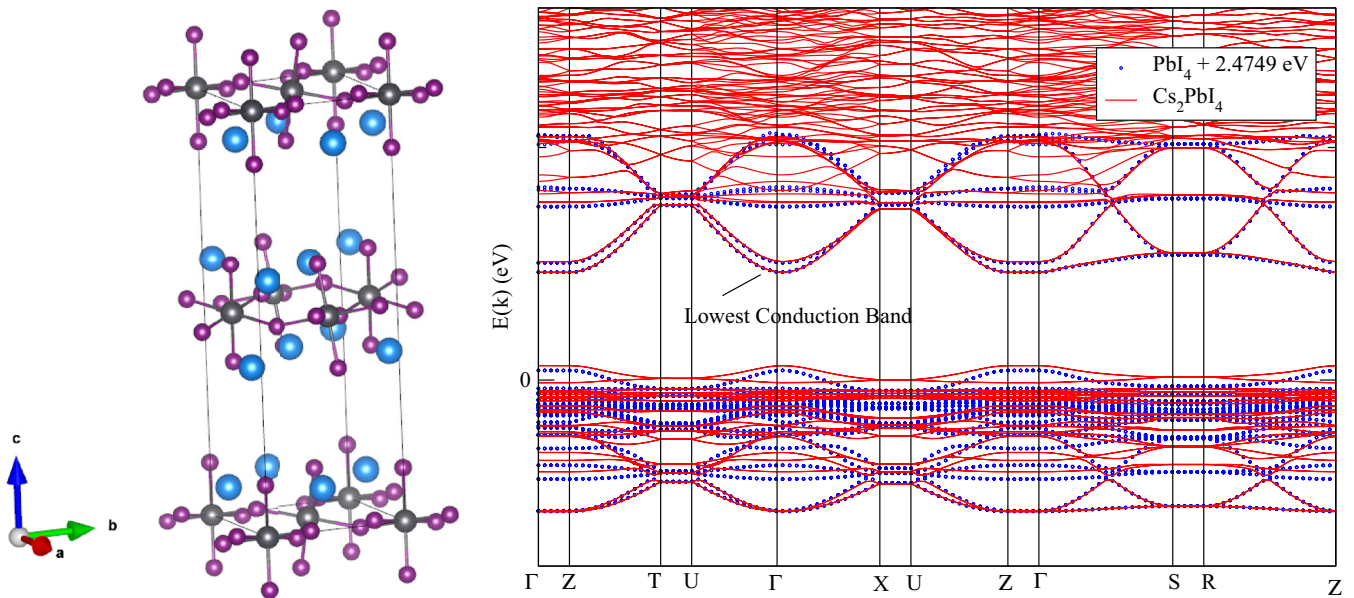


FIG. 7. (Left) The structure of a compound obtained by removing all the butylammonium chains and by adding Cs at the N sites of the original $(BA)_2PbI_4$ compound. (Right) We compare the bands of the PbI_4 (which are very close to those of $(BA)_2PbI_4$ as demonstrated in Fig. 4) with the bands of Cs_2PbI_4 obtained by a DFT calculation for this compound, shifting the energy of all the bands of PbI_4 by a constant amount of 2.4749 eV.

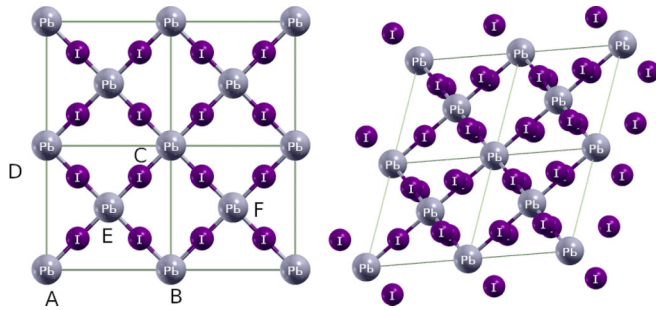


FIG. 8. (Left) The pure two-dimensional structure. (Right) A single layer with all the Pb and I atoms but with undistorted octahedra.

This compound exists in nature but with different lattice constants than those used in this calculation: we wish to keep them the same as in the $(BA)_2\text{PbI}_4$ for direct comparison of the bands. Notice that after shifting the energy of all the bands of PbI_4 by the same amount of 2.4749 eV there is very good agreement between the bands near the Fermi energy. This finding strengthens our conclusion and makes transparent our statement that the bands of the PbI_4 matrix almost solely determine the band structure within 1 eV below the VBM and 2 eV above the bottom of the CBM, which, including the energy gap, is about a 5 eV energy range.

III. TIGHT-BINDING MODEL

From the previous comparison, we conclude that it makes sense to derive a TB model to describe these rigid-band features. In this model, the role of all the interlayer butylammonium chains is to create the structure, and their role in the electronic structure is to provide one additional electron per Pb atom. Therefore, in our TB model which aims at describing these rigid band features for only the bands in the energy window of our interest discussed previously, the intercalating organic chains will be ignored.

We start from a symmetric system without any octahedra distortions and we choose the a and b axes to be of the same length, i.e., the a - b lattice has the symmetry of the square lattice. In the real material, the a and b axes are very close to each other, namely, $|\bar{b}|/|\bar{a}| \sim 0.98$.

First, we ignore the interplane coupling, which reduces the effective unit cell to one with half the number of atoms. Notice that the bands obtained by the DFT calculation have different bandwidths along the c axis. The bands that are mostly made of hydrogen atomic orbitals, i.e., from the butylammonium chains have significant bandwidth of the order of 0.1 eV. The bands of our interest, however, i.e., those within the region of 5 eV near the Fermi level, have remarkably negligible dispersion along the c axis. The bandwidth of these bands is less than 0.01 eV. This justifies our treatment of these bands as two-dimensional.

In addition, as a first step, we are going to ignore the role of the I atoms which are off the plane which reduces the problem to the 2D unit cell shown on the left panel of Fig. 8. We will also discuss how to include these off-planar atoms displayed on the right panel of Fig. 8 as a second stage.

Next, we are going to consider matrix elements of the Hamiltonian between the following twelve (12) states inside

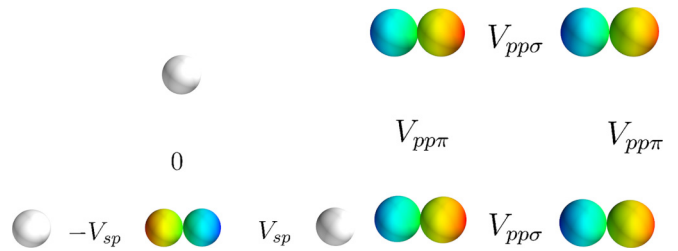


FIG. 9. (Left) Matrix elements of \hat{H} between a p_x orbital and an s -type orbital. The green (yellow) regions of the p orbitals denote negative (positive) sign. As a result these matrix elements can be zero or they can have opposite sign. (Right) The σ - and π -type matrix elements between p orbitals of neighboring atoms. σ ($V_{pp\sigma}$) bonds form when the orbital orientation is along the line connecting the atoms, while π ($V_{pp\pi}$) bonds form when the orbital orientation is perpendicular to the line connecting the atoms. When two orbitals have perpendicular orientations they do not yield a nonzero matrix element.

the unit cell of the reduced Bravais lattice: $|s^{(i)}\rangle$, $|p_x^{(i)}\rangle$, $|p_y^{(i)}\rangle$ and $|p_z^{(i)}\rangle$, where $i = 1, 2, 3$ for the three atoms in the reduced unit cell, i.e., the Pb atom at the origin $(0, 0, 0)$, the I atom at $(a/2, a/2, 0)$, and the I atom at $(-a/2, a/2, 0)$. We note that throughout the rest of the paper, our x and y axes are with respect to the 45° rotated coordinate system, not the original unit cell. The state $|s^{(i)}\rangle$ for $i = 1$ is the Pb $6s$ orbital and for $i = 2, 3$ it is the $5s$ I orbital. The states $|p_x^{(i)}\rangle$ and $|p_y^{(i)}\rangle$ where $i = 1$ corresponds to the Pb $6p$ orbitals, whereas when $i = 2, 3$ it corresponds to the I $5p$ orbitals.

Matrix elements of the form $\langle s^{(i)} | \hat{H} | s^{(j)} \rangle$ when $i = j$ are the on-site energies, which are $E_s^{(1)}, E_s^{(2)} = E_s^{(3)}$, while off-diagonal matrix elements are nonzero only if the atoms are nearest neighbors, in which case they are $V_{ss} \equiv \langle s^{(1)} | \hat{H} | s^{(2)} \rangle = \langle s^{(1)} | \hat{H} | s^{(3)} \rangle$.

Matrix elements between a $6s$ orbital of a Pb atom and a $5p$ -type orbital of any of its nearest neighbor I atoms or vice versa, when nonzero, are given by V_{sp} or V_{ps} . Because of the negative eigenvalue of the p_x orbital with respect to reflections about a plane perpendicular to the x axis, some of these matrix elements are zero and others have a relative minus sign as illustrated in the left panel of Fig. 9.

There two kinds of matrix elements of H between the same type of p orbitals as illustrated in the right panel of Fig. 9. As explained in the figure caption of this figure, they can be either of σ type, i.e., $V_{pp\sigma}$ or of π type, i.e., $V_{pp\pi}$. One has to be careful about their relative sign and when these matrix elements are zero.

First, instead of using the unit cell illustrated in the left panel of Fig. 8 by the ABCD square, which is the one that is necessary to use when the symmetry is broken by the octahedra tilting, we use as unit cell the smaller size square EBFC which is rotated by 45° with respect to the original. This unit cell contains only one Pb and two I atoms and, therefore, 12 states. This doubles the size of the BZ from the orange ABCD square of Fig. 10 to the square labeled $A'B'C'D'$, which is rotated by 45° with respect to the x axis. Therefore, for our convenience, we will solve the problem in this larger BZ (where we have only 12 states for each k) and in order to compare with the results of the original BZ, we will need to

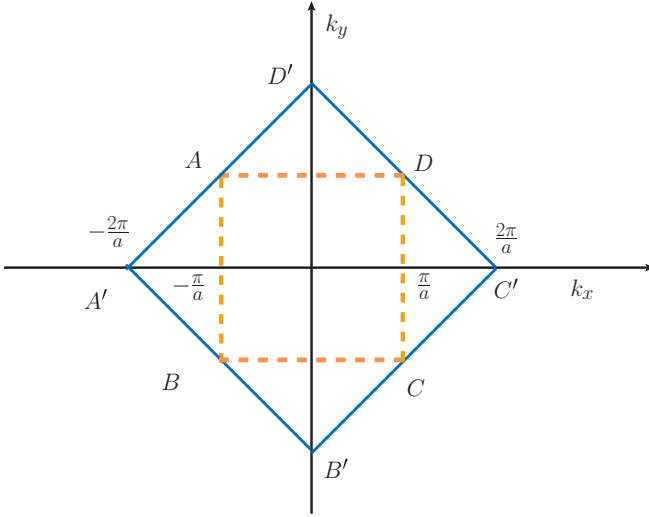


FIG. 10. The Brillouin zones of the original and the reduced simpler structure.

fold this BZ to the smaller size one. This folding will increase the number of bands by a factor of two and, therefore, we will be able to recover the total number of 24 bands, which were present in the original unit cell of the Bravais lattice.

If we include the matrix elements illustrated in Fig. 9 and we transform our Hamiltonian in momentum space, the Hamiltonian matrix becomes momentum-diagonal, i.e.,

$$\hat{H} = \sum_{\vec{k} \in \text{BZ}} \hat{H}_{\vec{k}}, \quad (1)$$

where the sum is over the entire Brillouin zone (blue square of Fig. 10) and $\hat{H}_{\vec{k}}$ is a 12×12 matrix given in Table I, where

$$d_x = 2 \cos\left(\frac{(k_x + k_y)a}{4}\right), \quad d_y = 2 \cos\left(\frac{(k_x - k_y)a}{4}\right), \quad (2)$$

$$d'_x = 2i \sin\left(\frac{(k_x + k_y)a}{4}\right), \quad d'_y = 2i \sin\left(\frac{(k_x - k_y)a}{4}\right). \quad (3)$$

The subscripts x and y in the above \vec{k} -dependent coefficients are labeled according to our x and y axes which are with respect to the 45° rotated coordinate system, not the original

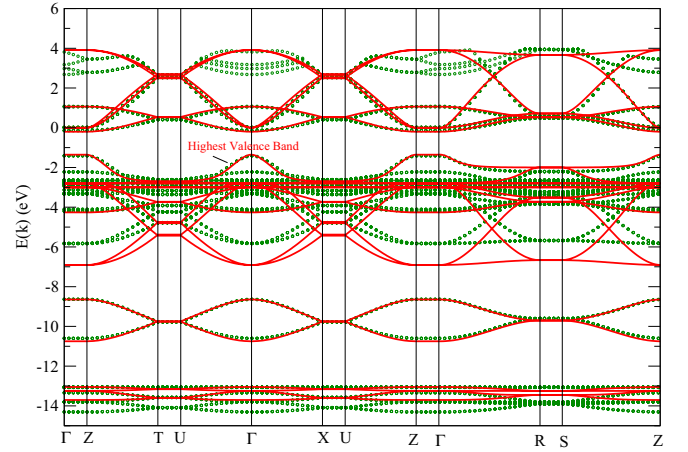


FIG. 11. Comparison of the results obtained by TB with those by DFT for the symmetric PbI_4 matrix.

unit cell. The k_x and k_y , however, are the x and y projections of \vec{k} on the corresponding x and y axes of the original unrotated unit cell. We use the latter for making comparison with the bands obtained by DFT. A numerical diagonalization of the above 12×12 at every \vec{k} point leads to the results illustrated in Fig. 11 which are compared to the results of the DFT of a 3D PbI_4 crystal without octahedral distortions and with $|\vec{a}| = |\vec{b}|$. The results of the fitting parameters is given in Table II. We note that the disagreement for the highest energy bands is due to the fact that these bands are at positive (unbound) energy and, therefore, the DFT calculation has included the effects of unbound electronic states. The DFT calculation involves more states because it includes the off-plane I atoms (see structure illustrated in the right panel of Fig. 8). The TB matrix (a 20×20 matrix) which includes the s and p orbital of these atoms is given in Appendix B. There are several more fitting parameters to use in this case and the agreement can be improved by introducing some of the missing bands away from the Fermi level. Such an approach, however, would lead to significant complication which works against our goal of simplifying the problem and leaving the description near the Fermi energy as accurate as possible. We have implemented this more complex TB Hamiltonian but the fact that it yields

TABLE I. Tight-binding matrix for the truncated 2D case.

	$ s^{(1)}\rangle$	$ p_x^{(1)}\rangle$	$ p_y^{(1)}\rangle$	$ p_z^{(1)}\rangle$	$ s^{(2)}\rangle$	$ p_x^{(2)}\rangle$	$ p_y^{(2)}\rangle$	$ p_z^{(2)}\rangle$	$ s^{(3)}\rangle$	$ p_x^{(3)}\rangle$	$ p_y^{(3)}\rangle$	$ p_z^{(3)}\rangle$
$\langle s^{(1)} $	$E_s^{(1)}$	0	0	0	$V_{ss}d_x$	$V_{sp}d'_x$	0	0	$V_{ss}d_y$	0	$V_{sp}d'_y$	0
$\langle p_x^{(1)} $	0	$E_{p_{xy}}^{(1)}$	0	0	$-V_{ps}d'_x$	$V_{pp\sigma}d_x$	0	0	0	$V_{pp\pi}d_y$	0	0
$\langle p_y^{(1)} $	0	0	$E_{p_{xy}}^{(1)}$	0	0	0	$V_{pp\pi}d_x$	0	$-V_{ps}d'_y$	0	$V_{pp\sigma}d_y$	0
$\langle p_z^{(1)} $	0	0	0	$E_{p_z}^{(1)}$	0	0	0	$V_{pp\pi}d_x$	0	0	0	$V_{pp\pi}d_y$
$\langle s^{(2)} $	$V_{ss}d_x$	$V_{ps}d'_x$	0	0	$E_s^{(2)}$	0	0	0	0	0	0	0
$\langle p_x^{(2)} $	$-V_{sp}d'_x$	$V_{pp\sigma}d_x$	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0	0
$\langle p_y^{(2)} $	0	0	$V_{pp\pi}d_x$	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0
$\langle p_z^{(2)} $	0	0	0	$V_{pp\pi}d_x$	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0
$\langle s^{(3)} $	$V_{ss}d_y$	0	$V_{ps}d'_y$	0	0	0	0	0	$E_s^{(2)}$	0	0	0
$\langle p_x^{(3)} $	0	$V_{pp\pi}d_y$	0	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0	0
$\langle p_y^{(3)} $	$-V_{sp}d'_y$	0	$V_{pp\sigma}d_y$	0	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0
$\langle p_z^{(3)} $	0	0	0	$V_{pp\pi}d_y$	0	0	0	0	0	0	0	$E_{p_z}^{(2)}$

TABLE II. Tight-binding parameters after fitting the DFT bands of the $x-y$ symmetric untilted PbI_4 matrix. The fit is shown in Fig. 11.

$E_s^{(1)}$	$E_{p_{xy}}^{(1)}$	$E_{p_z}^{(1)}$	$E_s^{(2)}$	$E_{p_{xy}}^{(2)}$	$E_{p_z}^{(2)}$	$V_{pp\sigma}$	$V_{pp\pi}$	V_{sp}	V_{ss}	V_{ps}
-9.3	-0.2	-0.2	-13.05	-2.8	-3.0	-2.5	0.8	1.2	-0.6	0.8

16 additional bands (when we fold the BZ) leads to a fitting procedure which does not have a unique and simple solution. For completeness, however, we provide this more complex TB Hamiltonian matrix in the table of Appendix B and we leave out from this paper the ambiguous results of such a fit.

In Sec. IV C, we will use our TB model to fit the results of the actual 3D $(\text{BA})_2\text{PbI}_4$ crystal, where we find that the quality of the fit is equally good. First, however, we would like to discuss an analytical treatment of the problem, which we do in the next section.

IV. ANALYTICAL DESCRIPTION

Here, we describe a simplified analytical treatment of the problem which contains all the essential elements of the original system near the Fermi level.

A. Noninteracting s and p orbitals

First, we define the model subspace, which is all the orbitals that fall inside the energy range of our interest which is 1 eV below the VBM and 2 eV above the CBM. These states are the Pb $6p$ and the I $5p$ orbitals. The Hamiltonian acting inside this model subspace is the following:

$$\hat{H}_M = \hat{H}_s + \hat{H}_{p_x} + \hat{H}_{p_y} + \hat{H}_{p_z}, \quad (4)$$

where \hat{H}_s corresponds to the orbital $|s^{(1)}\rangle$, i.e., the Pb $6s$ orbital, i.e.,

$$\hat{H}_s = E_s^{(1)}|s^{(1)}\rangle\langle s^{(1)}|, \quad (5)$$

and

$$\hat{H}_{p\alpha} = \begin{pmatrix} E_{p\alpha}^{(1)} & V_1^{(\alpha)} & V_2^{(\alpha)} \\ V_1^{(\alpha)*} & E_{p\alpha}^{(2)} & 0 \\ V_2^{(\alpha)*} & 0 & E_{p\alpha}^{(2)} \end{pmatrix} \quad (6)$$

in the basis $|p_\alpha^{(1)}\rangle, |p_\alpha^{(2)}\rangle, |p_\alpha^{(3)}\rangle$, where 1, 2, 3 stand for the Pb and the two p_α ($\alpha = x, y, z$) orbitals of its neighboring I atoms. Here

$$E_{p_x}^{(1,2)} = E_{p_{xy}}^{(1,2)}, \quad V_1^{(x)} = d_x V_{pp\sigma}, \quad V_2^{(x)} = d_y V_{pp\pi}, \quad (7)$$

$$E_{p_y}^{(1,2)} = E_{p_{xy}}^{(1,2)}, \quad V_1^{(y)} = d_x V_{pp\pi}, \quad V_2^{(y)} = d_y V_{pp\sigma}, \quad (8)$$

$$E_{p_z}^{(1,2)} = E_{p_z}^{(1,2)}, \quad V_1^{(z)} = d'_x V_{pp\pi}, \quad V_2^{(z)} = d'_y V_{pp\pi}. \quad (9)$$

Next, we will include the coupling of this subspace to the states which are outside this energy range but not too far away. For example, we will include the coupling between the Pb $6s$ and the I $5p$ because the latter falls inside the subspace of our interest, but we will ignore the coupling V_{ps} between Pb $6p$ and the I $6s$ states as the latter state is too far below (13 eV) the VBM and V_{ps} is only 0.8 eV, namely, $V_{ps}/(E_s^{(1)} - E_s^{(2)})$

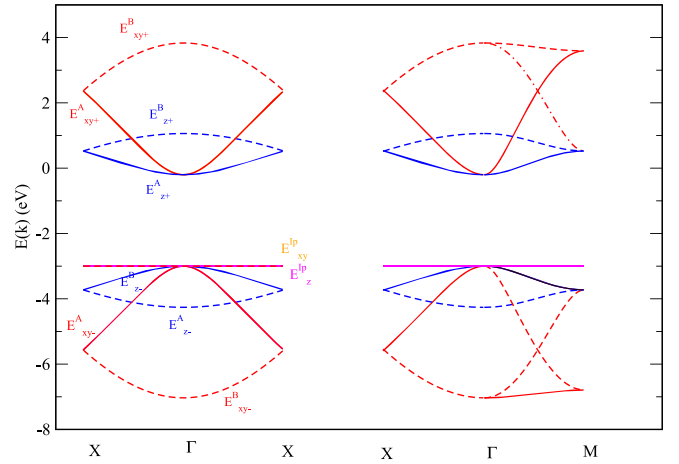


FIG. 12. The bands of the simple analytical 2D tight-binding model along some high symmetry directions without including the coupling between Pb $6s$ and I $5p$.

~ 0.06 . We will consider $V_{ss} = 0$ as this couples the Pb $6s$ and I $5s$ orbitals which are far from the above energy range and both states fall outside the energy domain of our interest. The resulting Hamiltonian is given by

$$\hat{H} = \hat{H}_M + \hat{H}_{sp}. \quad (10)$$

The \hat{H}_{sp} couples the Pb s (i.e., $|s^{(1)}\rangle$) with $|p_x^{(2)}\rangle$ and $|p_y^{(3)}\rangle$.

$$\hat{H}_{sp} = \begin{pmatrix} E_s^{(1)} & V_{sp}d'_x & V_{sp}d'_y \\ -V_{sp}d'_x & E_{p_{xy}}^{(2)} & 0 \\ -V_{sp}d'_y & 0 & E_{p_{xy}}^{(2)} \end{pmatrix}, \quad (11)$$

Diagonalization of the 3×3 matrices $\hat{H}_{p_x}, \hat{H}_{p_y}, \hat{H}_{p_z}$, of the model subspace yields the following eigenvalues:

$$E_{\pm}^{\alpha}(\vec{k}) = \frac{E_{p\alpha}^{(1)} + E_{p\alpha}^{(2)}}{2} \pm R_{\alpha} \quad E_0^{\alpha}(\vec{k}) = E_{p\alpha}^{(2)}, \quad (12)$$

$$R_{\alpha} \equiv \sqrt{\left(\frac{E_{p\alpha}^{(1)} - E_{p\alpha}^{(2)}}{2}\right)^2 + |V_1^{(\alpha)}|^2 + |V_2^{(\alpha)}|^2}. \quad (13)$$

The corresponding eigenstates for each of the above cases are given by the following form:

$$|\psi_{\pm}^{(\alpha)}\rangle = \frac{1}{\xi_{\pm}^{(\alpha)}} \left[(E_{p\alpha}^{(2)} - E_{\pm}^{(\alpha)}) |p_{\alpha}^{(1)}\rangle + V_1^{(\alpha)} |p_{\alpha}^{(2)}\rangle + V_2^{(\alpha)} |p_{\alpha}^{(3)}\rangle \right], \quad (14)$$

$$|\psi_0^{(\alpha)}\rangle = \frac{V_2^{(\alpha)} |p_{\alpha}^{(2)}\rangle - V_1^{(\alpha)} |p_{\alpha}^{(3)}\rangle}{\eta^{(\alpha)}}, \quad (15)$$

$$\xi_{\pm}^{(\alpha)} = \sqrt{(E_{p\alpha}^{(2)} - E_{\pm}^{(\alpha)})^2 + |V_1^{(\alpha)}|^2 + |V_2^{(\alpha)}|^2}, \quad (16)$$

$$\eta^{(\alpha)} \equiv \sqrt{|V_1^{(\alpha)}|^2 + |V_2^{(\alpha)}|^2}. \quad (17)$$

The unit cell doubling and 45° rotation leads to the folding of the BZ which doubles the number of eigenvalues to 18 (some of which are degenerate in this higher symmetry Hamiltonian). They are illustrated in Fig. 12 and are compared to the DFT calculation in Fig. 13 using the same values of the

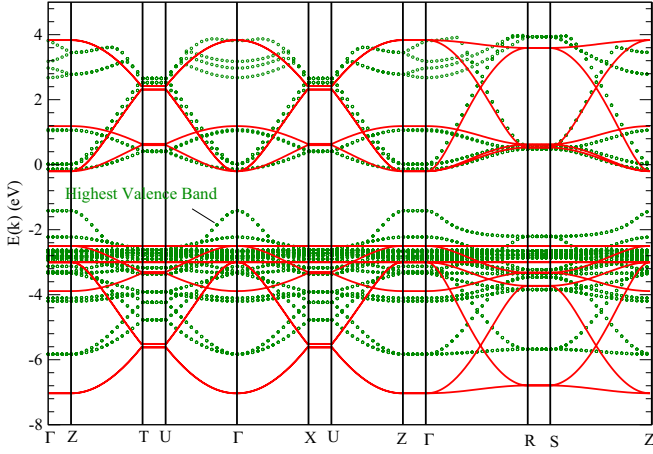


FIG. 13. Comparison with the DFT results of the 2D PbI₄ crystal with the results of our analytical 2D tight-binding model without including the coupling between Pb 6s and I 5p.

parameters of Table II but ignoring the *s* orbitals completely, i.e., $V_{ss} = V_{sp} = V_{ps} = 0$. Notice that the occupied bands are described well, notice, however, that the top of the valence band (i.e., the I 5p state) is flat and the model does not describe its dispersion. The main reason for this discrepancy is its coupling to the Pb 6s and it is corrected in the next section.

B. Integrating out the Pb *s* orbitals

The I 5p bands E_x^0 or E_y^0 are flat in this approximation. Next we include the role of the Pb 6s and I 5p hybridization which will account for the dispersion of the top valence band (which corresponds to the flat band illustrated by the green line in Fig. 12. The hybridization of the Pb 6s orbital with the I 5p, i.e., the matrix element V_{sp} , couples the $|s^{(1)}\rangle$ and $|p_x^{(2)}\rangle$ and $|p_y^{(2)}\rangle$ as in Eq. (11).

This interaction is crucial when there is an exact or almost degeneracy at specific \vec{k} points between the $E_{\pm}^{\alpha}(\vec{k})$ and $E_0^{\alpha}(\vec{k})$ bands given by Eq. (13). This happens at the Γ point for the four bands $E_{\pm}^{\alpha}(\vec{k})$ and $E_0(\vec{k})$ when they are folded at $(k_x, k_y) = (2\pi/a, 0)$. If we include the state $|s^{(1)}\rangle$ this

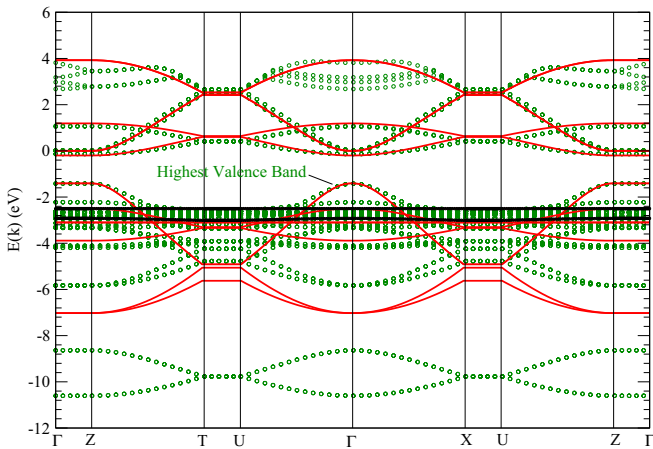


FIG. 14. Comparison of the bands with the DFT results after including the role of the V_{sp} interaction in our simple model.

TABLE III. TB parameters after fitting the DFT bands of the (BA)₂PbI₄. The fit is shown in Fig. 15.

$E_s^{(1)}$	$E_{p_{xy}}^{(1)}$	$E_{p_z}^{(1)}$	$E_s^{(2)}$	$E_{p_{xy}}^{(2)}$	$E_{p_z}^{(2)}$	$V_{pp\sigma}$	$V_{pp\pi}$	V_{sp}	V_{ss}	V_{ps}
-6.0	-3.3	-4.9	-10.0	1.1	-0.5	-1.8	0.5	1.8	0.6	0.8

becomes a 5×5 matrix, which cannot be analytically diagonalized. However, because the state $|s^{(1)}\rangle$ is well below the Fermi level such that $V_{sp}/(E_{p_{xy}}^{(2)} - E_s^{(1)}) \sim 0.1$, we can apply perturbation theory in V_{sp} . More precisely, we will apply quasi-degenerate stationary perturbation theory [28] to integrate out the *s* orbital in second order. Using Eq. (17) for the eigenstates, the second-order-corrected diagonal matrix elements (i.e., which include the effect of the virtual transition from the $|\psi^{(-)}\rangle$ state to $|s^{(1)}\rangle$ state and back) are given as follows:

$$\mathcal{E}_{-}^{(x)}(\vec{k}) = E_{-}^{(x)}(\vec{k}) + \frac{|V_1^{(x)}V_{sp}d'_x|^2}{|\xi_{-}^{(x)}|^2(E_{-}^{(x)} - E_s^{(1)})}, \quad (18)$$

$$\mathcal{E}_{-}^{(y)}(\vec{k}) = E_{-}^{(y)}(\vec{k}) + \frac{|V_1^{(y)}V_{sp}d'_y|^2}{|\xi_{-}^{(y)}|^2(E_{-}^{(y)} - E_s^{(1)})}, \quad (19)$$

$$\mathcal{E}_0^{(x)}(\vec{k}) = E_0^{(x)}(\vec{k}) + \frac{|V_1^{(x)}V_{sp}d'_x|^2}{|\eta^{(x)}|^2(E_0^{(x)} - E_s^{(1)})}, \quad (20)$$

$$\mathcal{E}_0^{(y)}(\vec{k}) = E_0^{(y)}(\vec{k}) + \frac{|V_1^{(y)}V_{sp}d'_y|^2}{|\eta^{(y)}|^2(E_0^{(y)} - E_s^{(1)})}. \quad (21)$$

The off-diagonal matrix elements vanish along the high-symmetry directions $\Gamma \rightarrow Z \rightarrow T \rightarrow U \rightarrow \Gamma \rightarrow X \rightarrow U \rightarrow \Gamma$ in our plot of Fig. 14. Along other directions, the off-diagonal elements do not necessarily vanish. In this case, we would need to diagonalize the 4×4 matrix. However, this cannot be done analytically, and if we need to resort to a numerical diagonalization we might as well diagonalize the full matrix to obtain the more exact nonperturbative solution. The purpose of this subsection was to demonstrate that the origin of the dispersion of the upper valence band is from the V_{sp} coupling which, we feel, has been achieved.

C. Fitting the full crystal with the same tight-binding model

Now that we have an analytical understanding of the origin of the band character near the Fermi level, we can proceed and modify the parameters of our TB model (Table I) listed in Table II in order to provide a good fit of the material of our focus, i.e., (BA)₂PbI₄.

The result of the fit is illustrated in Fig. 15 and the values of the parameters that produce this TB fit are given in Table III. We notice that the TB model describes very well all the bands which are due to the orbitals in PbI₄. However, these orbitals are the only ones that have non-negligible contribution to the bands in the energy range of our interest, which is discussed in the abstract.

V. SPIN-ORBIT COUPLING

Once we have established our TB model, it is straightforward to include the contribution of spin-orbit coupling (SOC).

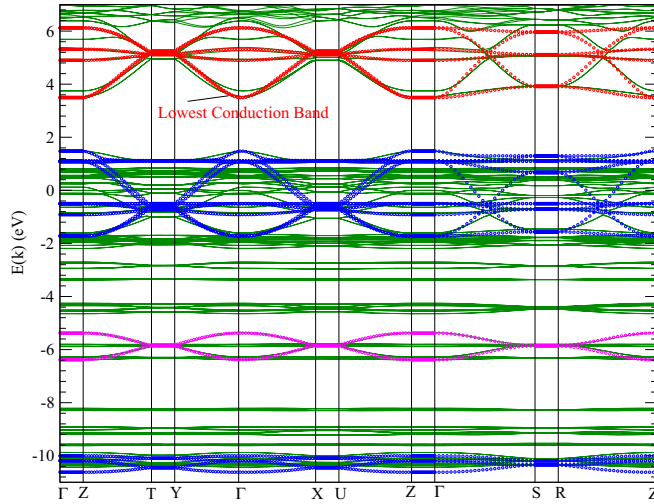


FIG. 15. Comparison of the bands of the real material $(BA)_2PbI_4$ obtained with the DFT (green solid lines) with the result of the fit using our TB model (open circles).

The $6s$ orbitals of Pb and the $5s$ orbital of the two I atoms in the unit cell give no contribution. The nonzero contribution comes from the $6p$ Pb orbital and the $5p$ orbital of the two I atoms (atoms 2 and 3 in our unit cell). We need to add the following part to the TB Hamiltonian and carry out the diagonalization in a space of double dimension.

$$\hat{H}_{\text{SOC}} = J_{\text{Pb}}\vec{\sigma}_1 \cdot \vec{l}_1 + J_I\vec{\sigma}_2 \cdot \vec{l}_2 + J_I\vec{\sigma}_3 \cdot \vec{l}_3, \quad (22)$$

where J_{Pb} (and J_I) is the SOC of the Pb $6p$ (and of the I $5p$) orbital with the electron spin. Each of the three terms above in the basis $|p_x^{(i)} \uparrow\rangle, |p_y^{(i)} \uparrow\rangle, |p_z^{(i)} \uparrow\rangle, |p_x^{(i)} \downarrow\rangle, |p_y^{(i)} \downarrow\rangle, |p_z^{(i)} \downarrow\rangle$, for $i = 1, 2$, and 3 is given in Table IV in units of J_i ($J_1 = J_{\text{Pb}}$, $J_2 = J_I$, $J_3 = J_I$).

In order to add the SOC in the DFT calculation we need fully relativistic pseudopotentials (FRP) and we will utilize those provided in Ref. [29]. First, we carry out a SCF calculation on the same size k -point mesh as in our previous calculations and the highest energy-cutoff suggested in the pseudopotential files. In Fig. 16, we illustrate that our TB model fits with the same level of accuracy the results using these FRP without SOC. We need this step in order to carry out the fit of the calculation with SOC. We note that the values of the TB parameters are only slightly off using this FRP. Next, without changing any of these TB parameters, we simply add the SOC Hamiltonian described in the previous paragraph and we fit the results by using only the

TABLE IV. The SOC Hamiltonian in units of J_i .

	$ p_x^{(i)} \uparrow\rangle$	$ p_y^{(i)} \uparrow\rangle$	$ p_z^{(i)} \uparrow\rangle$	$ p_x^{(i)} \downarrow\rangle$	$ p_y^{(i)} \downarrow\rangle$	$ p_z^{(i)} \downarrow\rangle$
$\langle p_x^{(i)} \uparrow $	0	$-i$	0	0	0	1
$\langle p_y^{(i)} \uparrow $	i	0	0	0	0	$-i$
$\langle p_z^{(i)} \uparrow $	0	0	0	-1	i	0
$\langle p_x^{(i)} \downarrow $	0	0	-1	0	i	0
$\langle p_y^{(i)} \downarrow $	0	0	$-i$	$-i$	0	0
$\langle p_z^{(i)} \downarrow $	1	i	0	0	0	0

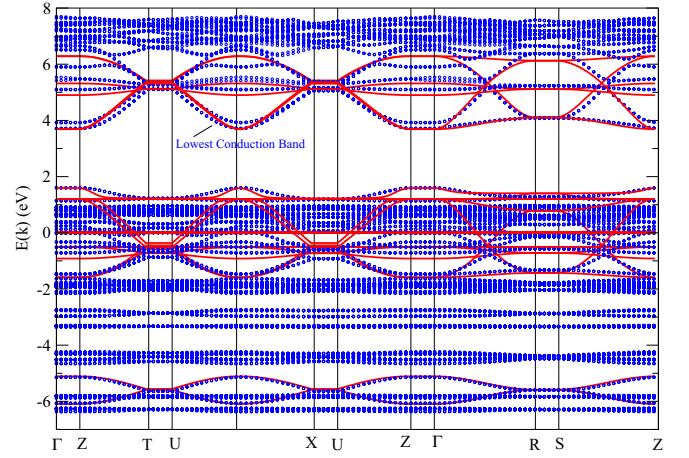


FIG. 16. Calculation with fully relativistic pseudopotentials provided in Ref. [29] but without SOC.

two fitting parameters J_{Pb} and J_I . The results are shown in Fig. 17. The effect of the SOC coupling is large as expected for Pb and I, however, we also know from previous work that DFT tends to overestimate the effects of the SOC (see related discussion in Refs. [30,31]). However, the goal of the present paper is to provide a reasonable starting one-electron model Hamiltonian without the inclusion of SOC and other effects, such as the effects of correlations. As discussed, the model can be the starting point for calculations to include these effects more accurately. For example, the effect of the SOC could be more accurately included by fitting the value of the SOC (using the above model independent form) to the results of a quasi-particle self-consistent GW calculation [32] or to the experimental results for the gap or other experimentally determined parameters.

The simple calculation presented above demonstrates the value of the present paper where the TB and the effective models were determined. Namely, without changing the

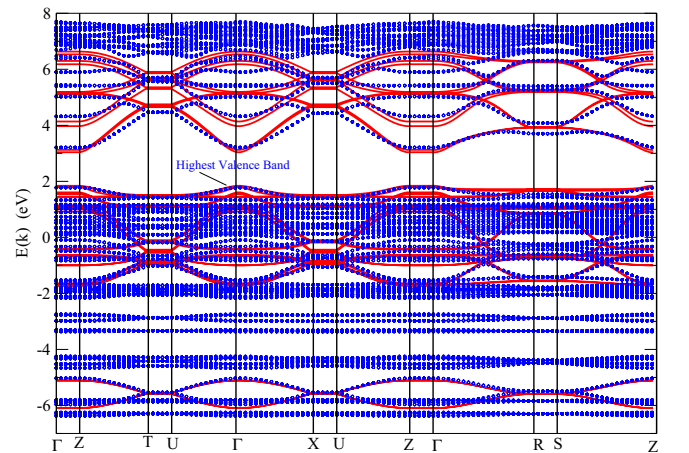


FIG. 17. Calculation with fully relativistic pseudopotentials provided in Ref. [29] and with SOC. The red lines are the fit using the TB model without changing the values of the parameter used in Fig. 16, i.e., the parameters used without SOC, and $\hbar J_1 = 0.9$ eV and $\hbar J_2 = 0.7$ eV.

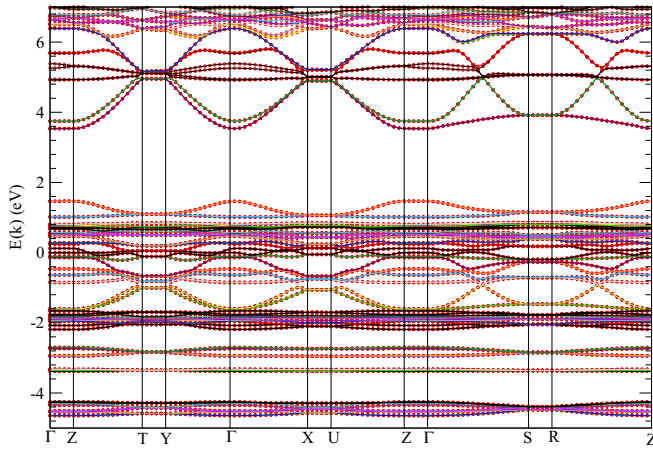


FIG. 18. Demonstration of convergence with k -point mesh size. The red open circles are the results obtained on a $8 \times 8 \times 3$ k -point mesh while the solid lines are the results using the QE default $4 \times 3 \times 1$ size.

parameters of the TB Hamiltonian, i.e., using their values determined without the inclusion of the SOC, we were able to accurately include the effects of the SOC by simply adding the SOC to our model. Similarly, other effects, such as the effects of correlations, the Jahn-Teller effect, optical response, etc., can be included starting from the present model.

VI. OTHER TERMS

There are other important physical effects which yield corrections to the above treatment, such as the octahedra distortions.

The octahedral distortions break the C_4 symmetry and that folds the BZ back to its observed form. In addition, they open gaps at high-symmetry points and lift band-degeneracies along high-symmetry directions.

Depending on the problem at hand to address, these terms can be important to include, which can be added on top the

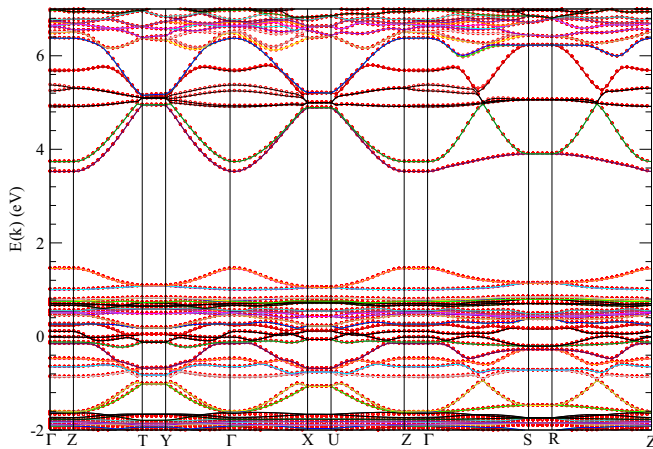


FIG. 19. Demonstration of convergence with respect to the energy cutoff E_c . The red-open circles are the results obtained on a $8 \times 8 \times 3$ k -point mesh using $E_c = 30$ Rd while the solid lines are the results using $E_c = 40$ Rd.

TB Hamiltonian considered in the present work. There are problems, however, where the TB-treatment of the present paper can be a good starting point [33]. For example, and this is one of our future projects, starting from this TB model we can include the role of electron interactions to study exciton bound states. These involve bound states of electron/hole pairs excited from near the VBM to near the CBM where the TB description is reasonably good.

The goal of the present paper was to provide an as simple as possible and yet accurate analytic and semianalytic description of the complex band structure of the simplest member of the series of the 2D halide-perovskite materials. Future work should extend this treatment to $m > 1$ members of this family and should include the role of the above smaller effects. Another direction should be to include the role of electron interactions in many-body phenomena, such as the role of excitons in these systems.

VII. DISCUSSION AND SUMMARY

Simplifying the very complex band structure of the 2D Ruddlesden-Popper perovskite materials and providing a simple model which accurately reproduces its main features and, which quantitatively describes it, is the main goal of the present paper. Such a simplified picture, not only allows us to grasp the physics of the electronic structure of these materials, which might help our thinking forward, but it can also provide a simple description of the one-body part of an effective many-body Hamiltonian to use to carry out many-body calculations.

First, we have illustrated that, in the simplest case of the series $(BA)_2(MA)_{m-1}Pb_mI_{3m+1}$ with $m = 1$, i.e., for $(BA)_2PbI_4$, the bands in the energy range: 1 eV below the VBM to 2 eV above the CBM, a 5 eV range covering the range of the photoelectric response, have negligible contribution from the atomic orbitals contained in the BA ligands. This conclusion is not to diminish the significance of organic chains in general. As an example of their significance, we would like to mention that the small organic chains, such as the MA, (which are absent in the $m = 1$ case) reduce dielectric screening within a monolayer of RP perovskite materials, which helps generate stable excitons at room temperature with binding energies of the order of hundreds of meV [6,34,35]. In addition, their physical properties are influenced by the number of layers that affects the exciton binding energy. However, in the $m = 1$ case, the role of the atoms of these BA chains is solely to stabilize the structure and to act as a charge reservoir which fills completely the highest occupied bands making the material an insulator. We demonstrate that other materials, which share the same halide-perovskite core matrix or even just the “matrix” formed by the same halide-perovskite layer (keeping the structure and all the atomic distances the same) have very similar band structure in the above defined energy range.

Further, we analyzed the complex band structure of this class of materials for $m = 1$ and found a simple 2D TB model which can accurately reproduce the band structure as obtained by DFT in the above energy window. We were also able to simplify this TB model in such a way that it allows an analytical, transparent, and accurate way to describe the origin of all the features of the band structure. As a consequence,

TABLE V. TB matrix when the pair off-the-plane I atoms are included.

	$s^{(1)}$	$p_x^{(1)}$	$p_y^{(1)}$	$p_z^{(1)}$	$s^{(2)}$	$p_x^{(2)}$	$p_y^{(2)}$	$p_z^{(2)}$	$s^{(3)}$	$p_x^{(3)}$	$p_y^{(3)}$	$p_z^{(3)}$	$s^{(4)}$	$p_x^{(4)}$	$p_y^{(4)}$	$p_z^{(4)}$	$s^{(5)}$	$p_x^{(5)}$	$p_y^{(5)}$	$p_z^{(5)}$
$s^{(1)}$	$E_s^{(1)}$	0	0	0	$V_{ss}d_x$	$V_{sp}d'_x$	0	0	$V_{ss}d_y$	0	$V_{sp}d'_y$	0	V_{ss}	0	0	V_{sp}	V_{ss}	0	0	V_{sp}
$p_x^{(1)}$	0	$E_{p_{xy}}^{(1)}$	0	0	$-V_{ps}d'_x$	$V_{pp\sigma}d_x$	0	0	0	$V_{pp\pi}d_y$	0	0	0	$V_{pp\pi}$	0	0	0	$V_{pp\pi}$	0	0
$p_y^{(1)}$	0	0	$E_{p_{xy}}^{(1)}$	0	0	0	$V_{pp\pi}d_x$	0	$-V_{ps}d'_y$	0	$V_{pp\sigma}d_y$	0	0	0	$V_{pp\pi}$	0	0	0	$V_{pp\pi}$	0
$p_z^{(1)}$	0	0	0	$E_{p_z}^{(1)}$	0	0	0	$V_{pp\pi}d_x$	0	0	0	$V_{pp\pi}d_y$	$-V_{ps}$	0	0	$V_{pp\sigma}$	$-V_{ps}$	0	0	$V_{pp\sigma}$
$s^{(2)}$	$V_{ss}d_x$	$V_{ps}d'_x$	0	0	$E_s^{(2)}$	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
$p_x^{(2)}$	$-V_{sp}d'_x$	$V_{pp\sigma}d_x$	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0	0	0	0	0	0	0	0	0	0
$p_y^{(2)}$	0	0	$V_{pp\pi}d_x$	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0	0	0	0	0	0	0	0	0
$p_z^{(2)}$	0	0	0	$V_{pp\pi}d_x$	0	0	0	$E_{p_z}^{(2)}$	0	0	0	0	0	0	0	0	0	0	0	0
$s^{(3)}$	$V_{ss}d_y$	0	$V_{ps}d'_y$	0	0	0	0	0	$E_s^{(2)}$	0	0	0	0	0	0	0	0	0	0	0
$p_x^{(3)}$	0	$V_{pp\pi}d_y$	0	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0	0	0	0	0	0
$p_y^{(3)}$	$-V_{sp}d'_y$	0	$V_{pp\sigma}d_y$	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0	0	0	0	0	0
$p_z^{(3)}$	0	0	0	$V_{pp\pi}d_y$	0	0	0	0	0	0	$E_{p_z}^{(2)}$	0	0	0	0	0	0	0	0	0
$s^{(4)}$	V_{ss}	0	0	V_{ps}	0	0	0	0	0	0	0	0	$E_s^{(2)}$	0	0	0	0	0	0	0
$p_x^{(4)}$	0	$V_{pp\pi}$	0	0	0	0	0	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0	0
$p_y^{(4)}$	0	0	$V_{pp\sigma}$	0	0	0	0	0	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0	0	0	0	0
$p_z^{(4)}$	$-V_{sp}$	0	0	$V_{pp\pi}$	0	0	0	0	0	0	0	0	0	0	0	$E_{p_z}^{(2)}$	0	0	0	0
$s^{(5)}$	V_{ss}	0	0	V_{ps}	0	0	0	0	0	0	0	0	0	0	0	0	$E_s^{(2)}$	0	0	0
$p_x^{(5)}$	0	$V_{pp\pi}$	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0	0
$p_y^{(5)}$	0	0	$V_{pp\sigma}$	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	$E_{p_{xy}}^{(2)}$	0
$p_z^{(5)}$	$-V_{sp}$	0	0	$V_{pp\pi}$	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	$E_{p_z}^{(2)}$

a simple band picture emerges out of the complexity of the bands as obtained by straightforward application of DFT. By considering a two-dimensional TB model, which ignores the small octahedra distortions, it allows us to reduce the size of the unit cell by a factor of two, a fact that doubles the Brillouin zone by unfolding it because of symmetry. This reduces the TB Hamiltonian to a smaller matrix. The most important conduction and relevant conduction and valence bands are obtained as a hybridization of mostly Pb $6p$ and I $5p$ orbitals in the case of our example $(BA)_2\text{PbI}_4$. To obtain the correct dispersion of the highest valence band, we also needed to involve the role of the sp hybridization between the Pb $6s$ orbital and the $5p$ orbital of the I atoms that form the octahedra corners. Finally, we offer a simple analytic description of the band structure by integrating out this Pb $6s$ orbital as it sits energetically well below the above mentioned energy range.

As already discussed in the previous section, there are several directions where the results of the present work can be useful and also ways in which other effects can be incorporated depending on the problem at hand. For example, the simple TB model uncovered in this paper will be useful in carrying out many-body calculations to describe excitonic properties of these materials [36] which is our future goal.

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APPENDIX A: DFT CONVERGENCE STUDY

Our self-consistently converged ground-state calculation used a $4 \times 3 \times 1$ and a $8 \times 4 \times 3k$ -point-mesh size. The results are compared in Fig. 18 and we conclude that the $8 \times 6 \times 3$ size is large enough for the purpose of the present paper. In Fig. 19, we compare the results of our DFT calculation for $8 \times 6 \times 3$ size for energy cutoff of 30 Rd (red circles) and 40 Rd (solid lines) to show that using 30 Rd as the energy cutoff is large enough for the purpose of the present paper.

APPENDIX B: MORE COMPLEX TIGHT-BINDING MATRIX

Generalization of our TB model of Table I for the model complete 2d-halide-perovskite illustrated in the right panel of Fig. 8 leads to the 20×20 TB matrix given in Table V. This model includes the pair of I atoms per Pb atom which are off the plane and complete the octahedra.

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