

## Erratum: Witnessing entanglement in quantum magnets using neutron scattering [Phys. Rev. B **103**, 224434 (2021)]

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Two mistakes were made in our paper concerning the definition of the fluctuation-dissipation theorem. The expression given in Sec. II C is  $\chi''(\omega) = \frac{1}{\hbar} \tanh(\hbar\omega\beta/2)\tilde{S}(\omega)$  (suppressing wave-vector dependence). This expression should read  $\chi''(\omega) = \frac{\pi}{\hbar} \tanh(\hbar\omega\beta/2)\tilde{S}(\omega)$  in order to be consistent with the Van Hove formulation of the dynamic structure factor conventionally used

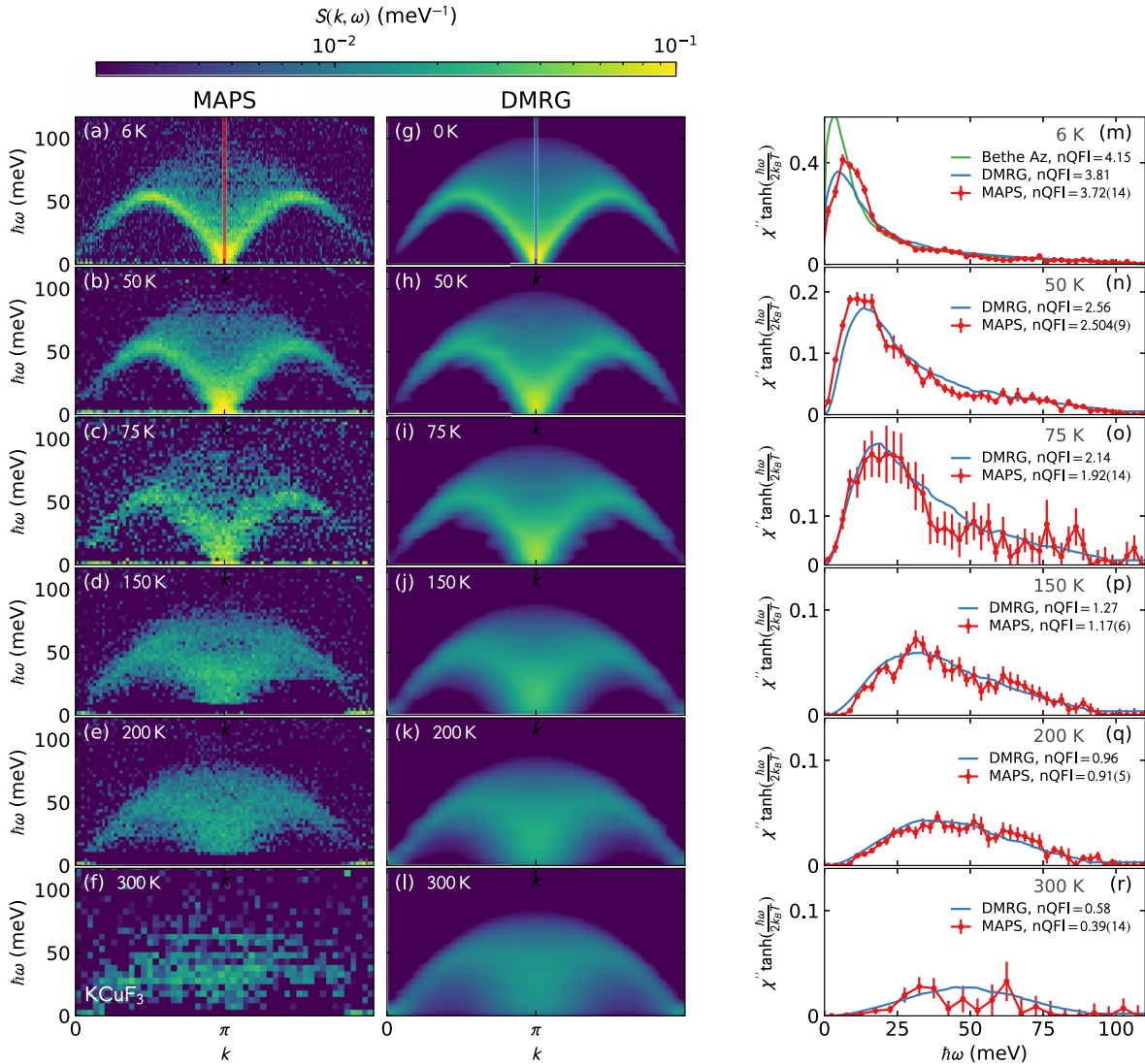


FIG. 1. Spectra and corrected quantum Fisher information. (a)–(f) Neutron scattering of  $\text{KCuF}_3$  measured on MAPS at different temperatures. (g)–(l) DMRG simulated scattering from a 1D HAF with experimental resolution broadening applied. (m)–(r) QFI integrand at  $k = \pi$ , shown with normalized quantum Fisher information ( $\frac{f_Q}{12.5^2}$ ) calculated at that point. At 6 K the data are also compared with the algebraic Bethe ansatz result (m).

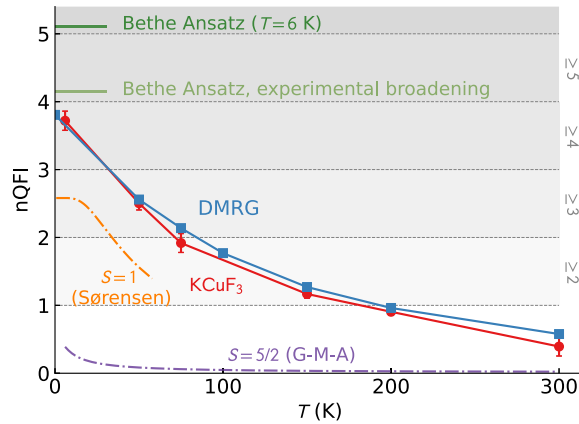


FIG. 2. Corrected normalized quantum Fisher information as a function of temperature. When  $n\text{QFI} > m$ , where  $m > 0$  is an integer, the system is in a state with  $\geq(m+1)$ -partite entanglement. Note that at the lowest temperatures,  $\geq 4$ -partite entanglement is witnessed in  $\text{KCuF}_3$ .

in neutron scattering,  $S(\omega) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dt e^{-i\omega t} \langle \hat{S}\hat{S}(t) \rangle$  (suppressing position dependence), rather than the definition in Eq. (S13) of Hauke *et al.* [1].

However, the above fluctuation-dissipation theorem is only valid for  $S(\omega)$  symmetrized over  $\omega$ ,  $\tilde{S}(\omega) = S(\omega) + S(-\omega)$ , which our data were not (no negative  $\omega$  data were collected in the experiment). The correct fluctuation-dissipation theorem to use for nonsymmetrized data is [2]

$$\chi''(\omega) = \pi(1 - e^{-\hbar\omega\beta})S(\omega), \quad (1)$$

which applies to dynamic susceptibilities defined as

$$\chi''(\omega) = -i \int_0^{\infty} dt e^{i\omega t} \phi(t), \quad \phi(t) \equiv \frac{i}{\hbar} \langle [\hat{S}(t), \hat{S}] \rangle, \quad (2)$$

i.e., where the time integral runs from 0 to  $\infty$ , as is assumed by both Hauke *et al.* [1] and Lovesey [2].

Due to an error in the code used to calculate  $\chi''$  (and hence QFI) we inadvertently instead used  $\chi''(\omega) = \frac{\pi}{2}(1 - e^{-\hbar\omega\beta})S(\omega)$ . As a result, all reported  $\chi''$  and QFI values are smaller than their real values by a factor of two. This implies that QFI is a more powerful probe of solid state entanglement than previously realized. In particular, we now conclude that at least quadripartite entanglement is witnessed in  $\text{KCuF}_3$  at the lowest temperatures, and that bipartite entanglement is present up to at least 150 K. Corrected versions of Figs. 2 and 5 of our original paper are provided in Figs. 1 and 2 of this Erratum, respectively.

Recently, we became aware of the need for this correction through work reported in Ref. [3], and thank Varun Menon for originally raising the issue. We note that this incorrect factor of 1/2 also affects the results of Refs. [4,5], the  $\chi''$  and QFI values of which will be corrected in the same fashion.

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