


Electronic structure and magnetism of the triple-layered ruthenate $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ G. Gebreyesus ^{*}*Department of Physics, School of Physical and Mathematical Sciences, College of Basic and Applied Sciences, University of Ghana, Ghana*Prosper Ngabonziza [†]*Max Planck Institute for Solid State Research, Heisenbergstraße 1, 70569 Stuttgart, Germany
and Department of Physics, University of Johannesburg, P.O. Box 524 Auckland Park 2006, Johannesburg, South Africa*Jonah Nagura *Department of Physics, Federal University of Agriculture, Abeokuta, Nigeria*Nicola Seriani *The Abdus Salam ICTP, I-34151, Trieste, Italy*

Omololu Akin-Ojo

*East Africa Institute for Fundamental Research, University of Rwanda, Kigali, Rwanda*Richard M. Martin [‡]*Department of Physics, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801, USA
and Department of Applied Physics, Stanford University, Stanford, California 94305, USA*

(Received 17 November 2021; revised 10 February 2022; accepted 21 March 2022; published 13 April 2022)

We report electronic band structure calculations for $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ which display both ferromagnetic and metamagnetic behavior. The density functional calculations find the ground state to be ferromagnetic in agreement with experiment and we show that the resulting spin polarization has dramatic consequences for the electronic properties. The minority-spin bands are mainly empty and disperse steeply upward at the Fermi energy, whereas the majority-spin bands are full or nearly fully occupied and form narrow bands near the Fermi energy, which could be the electronic origin of the metamagnetism. Inclusion of a Hubbard interaction U applied to the Ru $4d$ states has major effects on the narrow bands, which reveal the role of Coulomb interactions and correlated many-body physics. The results are in qualitative agreement with recent angle-resolved photoemission spectroscopy (ARPES) experiments and show the need for a combined theoretical study and experimental ARPES investigation with better energy resolution to reveal the nature of the narrow bands close to the Fermi level, which is critical for understanding the exotic magnetic properties observed in this material.

DOI: [10.1103/PhysRevB.105.165119](https://doi.org/10.1103/PhysRevB.105.165119)**I. INTRODUCTION**

Strontium ruthenate materials $\text{Sr}_{n+1}\text{Ru}_n\text{O}_{3n+1}$ form a unique class of layered compounds in the Ruddlesden-Popper (RP) series; because the number ($n = 1, 2, 3, \dots, \infty$) of the RuO_6 octahedra in the unit cell represents a critical parameter for obtaining distinct collective phenomena in the ground state, depending on the overlap between Ru $4d$ wave functions through adjacent octahedra [1–3]. The rich array of ground-state properties in this series includes superconductivity, magnetic order and metamagnetism, orbital order, heavy fermions, and other phenomena [1–10]. These materials are

also particular examples of the $4d$ transition metal oxides in which electronic correlations are remarkably strong. The richness of different phenomena in Sr-based layered ruthenates comes from a competition between local and itinerant physics, which is an evidence of strong interplay between charge, spin, orbital, and lattice degrees of freedom [1,11–13].

Among the strontium ruthenates in the RP series, the single-layer Sr_2RuO_4 has been studied extensively, mostly because it is a superconductor with an unconventional order [2,3,14–20], and it is the first example of a noncuprate-layered perovskite superconductor [21]. The electronic properties are due to bands associated with the Ru-O planes which form a square lattice with the same structure as the Cu-O planes in the high-temperature superconductors [22,23]. The other members of the layered strontium ruthenates are formed by stacks of similar layers. The double-layered $\text{Sr}_3\text{Ru}_2\text{O}_7$ has two layers of RuO_6 octahedra per unit cell and it is generally regarded as

^{*}ghagoss@ug.edu.gh[†]p.ngabonziza@fkf.mpg.de[‡]rmartin@illinois.edu

a paramagnetic that is very near a ferromagnetic instability [4]. Due to its interesting diamagnetic features, quantum critical metamagnetism, and nematic fluid behavior with heavy d -electron masses, the $\text{Sr}_3\text{Ru}_2\text{O}_7$ system has also been the subject of many theoretical and experimental investigations [5,6,24–29]. At the other end of this series, the increased tendency toward magnetism culminates in SrRuO_3 , which is a ferromagnetic metal with the cubic perovskite structure that can be viewed as the infinite-layer limit of the series [10,30–32]. Also, SrRuO_3 is a highly conductive metallic oxide with a good lattice match to many oxide materials; thus, it has been widely used as a conductive electrode in diverse oxide heterostructures [10,33–36].

The focus of this paper is the triple-layer system $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, which displays both anisotropic ferromagnetism and intriguing orbital-selective metamagnetic behavior [8,9,11,12,37–41]. The spontaneous ferromagnetic moment is perpendicular to the planes below the Curie temperature of 105 K, with an additional transition at around 60 K observed in magnetic susceptibility [9,11,37,42]. Depending of the sample quality, the reported saturated magnetic moments in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ single crystals range from ~ 1.0 to $\sim 1.7\mu_B/\text{Ru}$ oriented along the c axis [8,9,11,12,43,44]. A spin-polarized neutron diffraction study [44] has shown that the moments are largest on the central layer and are almost completely due to the spin with vanishingly small orbital contributions. For magnetic field applied along the ab plane, there is a decrease in magnetization below the second transition at 60 K, which is accompanied with a metamagnetic transition for fields of about 2.5 T [37,38,43]. A double metamagnetic behavior was also reported in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ with a second in-plane metamagnetic transition at a slightly larger applied field [40].

Electronic band structure calculations of Sr_2RuO_4 [45,46], $\text{Sr}_3\text{Ru}_2\text{O}_7$ [24,25], and SrRuO_3 [47,48] materials have been reported, and compared in detail to angle-resolved photoemission (ARPES) experiments [23,24,26,49–51]. However, it is only recently that the band structure of the $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ system has been studied by ARPES [52]. To our knowledge, there have been no electronic structure calculations for $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ reported so far.

Here, we report density functional calculations of the band structure and Fermi surfaces of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ that reveal the large effects due to the magnetic properties of this system. In agreement with experiment, the ground state of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ is found to be ferromagnetic, so that there are two sets of bands for the two spins. The minority-spin bands have reduced occupation and Fermi surfaces that are mainly in the outer areas of the Brillouin zone (BZ) away from Γ with bands that disperse steeply upward. In contrast, the majority-spin bands are full or nearly fully occupied with the Fermi level near the top of the bands leading to narrow bands near the Fermi energy around the Γ point. As previously found for other Sr-based layered ruthenates [53–56], we find that inclusion of a Hubbard U has dramatic effects on states near the Fermi energy, signifying the role of Coulomb interactions and correlated many-body physics for explaining the exotic ground-state properties, e.g., anomalous metamagnetic behavior in this material.

The calculated bands are in general agreement with ARPES experiments [52], which find holelike and electronlike

pieces of the Fermi surfaces consistent with spin-polarized bands of a ferromagnetic system and very different from the calculated Fermi surfaces if the system is constrained to be non-spin-polarized. However, it will further require combined theoretical study and experimental ARPES investigation with higher-energy resolution to sort out the details of the narrow bands near the Fermi energy.

The remainder of this paper is organized as follows. In Sec. II we discuss the structure of the layered ruthenates with the focus on the three-layered ruthenate materials. Section III contains a discussion about the nature of the bands in these classes of materials, and presents results from a tight-binding model, which are very useful for understanding the DFT results and comparison with the experiment. Section IV contains the technical details of our DFT calculations, and presents the band structures and Fermi surfaces for the spin-polarized and non-spin-polarized calculations. In Sec. V we discuss the comparison of the calculated bands and Fermi surfaces with the experimental data. Section VI discusses the possible consequences for the interesting properties of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ and in Sec. VII are concluding remarks.

II. STRUCTURES OF THE LAYERED RUTHENATES

The structures of the layered strontium ruthenates in the RP sequence $\text{Sr}_{n+1}\text{Ru}_n\text{O}_{3n+1}$, $n = 1, 2, \dots$, consist of Ru-O and Sr-O layers as illustrated in Fig. 1 for $\text{Sr}_4\text{Ru}_3\text{O}_{10}$. At the top is the side view of a triple layer which can be considered as three layers of Ru and O atoms, with four Sr-O layers: two Sr-O layers between the Ru-O layers and two Sr-O layers at the top and bottom of each triple layer. The three-dimensional crystal is formed by stacking the triple layers with the outer Sr-O layers creating an effective insulating barrier between adjacent triple layers, so that the relevant electronic bands are essentially two dimensional. The calculations described in Sec. IV are for the experimentally determined three-dimensional structure [8]. However, the primary results are independent of the stacking and the properties are dominated by the two-dimensional bands of a single triple layer. This is verified by the calculations and is assumed in the interpretations of the ARPES experiments which are analyzed purely in terms of two-dimensional bands [52].

At the bottom of Fig. 1 is a top view of a single Ru-O layer with the rotations of the octahedra indicated by the displacements of the oxygen atoms in the plane. If there are no rotations, as in the single-layer system Sr_2RuO_4 , each Ru atom is equivalent and the primitive unit cell is a square containing one Ru atom. In the double- and triple-layer systems, the rotations double the unit cell so that it is a square containing two Ru atoms and turned by 45° degrees as shown in Fig. 1 (bottom left). Correspondingly, the BZ is a square reduced by $\frac{1}{2}$ and turned by 45° , as depicted at the bottom right in Fig. 1. In the triple layer $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, the rotations of the RuO_6 octahedra in the inner and outer layers are not equivalent; the RuO_6 octahedra in the inner layer rotated by an angle of $\approx 11^\circ$ about the c axis whereas the outer two layers rotated by very nearly half that amount in the opposite direction [8].

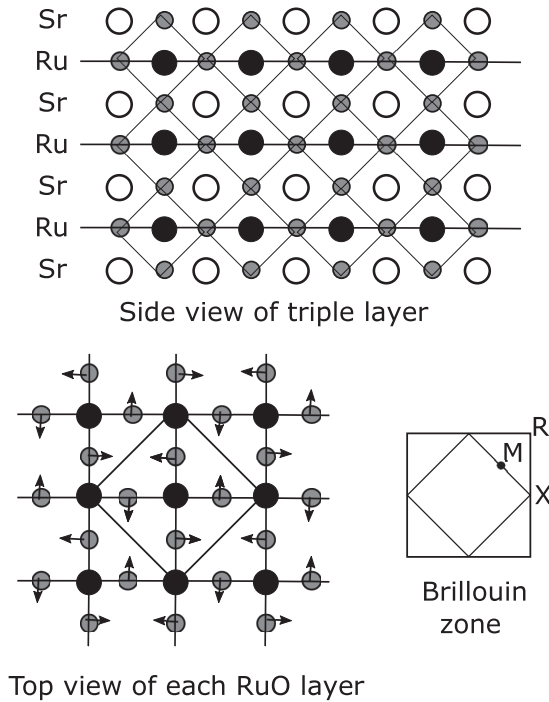


FIG. 1. The structure of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ is based on the triple layer shown at the top, which consists of three layers of Ru atoms surrounded by octahedra of O atoms indicated by the small gray circles. There are layers of Sr between the Ru layers and at the upper and lower sides of a triple layer. The stacking of the layers in the crystal is described in the text but the main features of the electronic structure are determined by the bands of a single triple layer. At bottom left is depicted a layer of Ru and O atoms. If the O atoms were exactly between the Ru atoms, the layer would have one Ru per cell; however, the octahedra are rotated as indicated by the arrows which double the primitive cell as shown by the square. Similarly, the BZ is reduced by a factor of 2 and turned by 45° as indicated at the lower right.

III. NATURE OF THE BANDS

The main features of the band structures of the layered compounds in the RP series can be understood starting from the band structure a single Ru-O layer shown at the bottom left in Fig. 1. In all cases, the Ru d states involved are primarily the 3 t_{2g} states d_{xz} , d_{yz} , and d_{xy} where x , y denote the directions in the plane of the layers and z is perpendicular to the plane. The bands are formed by Ru-Ru hopping via the O p states so that they are mixed Ru and O characters with the states near the Fermi energy composed primarily of Ru d states. There are also oxygen atoms (called apical oxygens) above and below the Ru atoms that have the effect of raising the energy of the d_{xz} and d_{yz} relative to the d_{xy} .

It is very useful for understanding the DFT results and comparison with experiment to have simple models that capture the essential features. The minimal model that captures the main features is a tight-binding model with only one p - d hopping matrix element $t_{pd\pi}$ and one relevant energy difference $\Delta_{pd} = \varepsilon_d - \varepsilon_p$, which is the well-known model for the single-layer Sr_2RuO_4 and used for other materials in the RP series [1–3]. The resulting bands near the Fermi energy

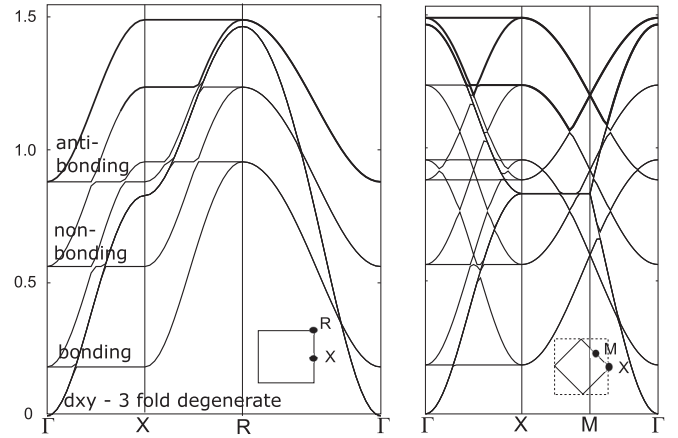


FIG. 2. Characteristic bands of the coupled Ru t_{2g} and O p states found using the tight-binding model described in the text and the structure shown in Fig. 1. Only the bands near the Fermi energy are shown. These have mainly Ru d character; O states are at lower energy. At the left are bands of a triple layer which are like three copies of the single layer taking into account the splitting due to the hopping between neighboring planes in the triple layer. The figure at the right illustrates the effect of doubling the cell, which folds the bands into the smaller BZ and leads to states with maxima at the Γ point. Note the similarity to the bands for the full DFT calculations shown in Figs. 4, 5 7, and 9.

have mainly Ru d character: a d_{xy} band with dispersion that is isotropic in the x , y plane, and a pair of bands d_{xz} and d_{yz} with very anisotropic dispersion. The d_{xz} band disperses along the x direction and is flat in the y direction since the matrix element with the oxygen p states is zero, whereas the d_{yz} band disperses in the y direction and is flat in the x direction. At the Γ point, d_{xz} and d_{yz} are degenerate with energy shifted up relative to the d_{xy} band at Γ since the hopping matrix element $t_{pd\pi}$ mixes the apical oxygen p states with the d_{xz} and d_{yz} bands but does not affect the d_{xy} bands.

The multilayer systems consist of n Ru-O layers as illustrated by the triple-layer system in Fig. 1. Since each n -layer unit is weakly coupled to the adjacent n -layer unit, the final dispersion is essentially two dimensional. However, within n layer each unit, the hopping through the apical oxygen is comparable to the hopping within the planes. This leads to n copies of the bands of a single layer which are split by the interlayer hopping into bonding and antibonding for two layers; bonding, nonbonding, and antibonding for three layers; and so forth. The example of three layers is shown in the left side of Fig. 2, where the three d_{xz} and d_{yz} states per Ru atom form bands with large splitting due to the hopping between the three Ru-O layers. In contrast, the d_{xy} states form three bands, which are degenerate in this model.

At the right of Fig. 2 is the consequence of doubling of the unit cell in real space. This is caused by the rotations of the O octahedra indicated by the arrows in Fig. 1. Since the bands are plotted in the smaller BZ, they appear complicated, and the relation to the unfolded bands is very useful for understanding. Most relevant for our purposes is that the bands at the R point are folded to the Γ point in the reduced BZ. In the original BZ (which applies if there are no rotations of

the octahedra), the maxima in the bands are at the R point and there are no bands near the zone center Γ in the upper energy range. It is only because of the rotations that there are bands with maxima at the zone center Γ that curve downward (holelike). An interesting consequence is that the intensity of the folded bands in a measurement is proportional to the amount of rotation and may appear to be weak in the actual experiments.

There are other effects that are not shown but can be understood qualitatively, e.g., magnetism. In a tight-binding model, this is simply a splitting of the majority- and minority-spin states leading to two sets of bands shifted almost rigidly, like those shown in Figs. 5 and 9. Since there is a large magnetization, the majority-spin bands are almost filled and the Fermi energy must be near the top of the bands. Thus, the majority-spin bands are expected to be holelike bands around Γ with energies at or near the Fermi energy. In contrast, the minority-spin bands have reduced occupation with the Fermi energy in the lower parts of the bands.

IV. DFT CALCULATIONS AND COMPUTATIONAL DETAILS

There is much previous work on the bands of the ruthenates including Sr_2RuO_4 and $\text{Sr}_3\text{Ru}_2\text{O}_7$ which has a structure similar to $\text{Sr}_4\text{Ru}_3\text{O}_{10}$. Density functional calculations have proven very successful in describing the overall features of the observed bands [1–3,23,25,26]. In all cases, the bands near the Fermi energy have mainly Ru t_{2g} d character with characteristic forms determined by only a few parameters, as described in Sec. III. For the nonmagnetic materials, the Fermi surfaces are determined by the band fillings and sensitive only to the relative energies of the d_{xy} and the d_{xz} and d_{yz} bands. There are effects of correlation that cause narrower bandwidths; however, the Fermi surfaces are not greatly affected because the Fermi energy is determined by band filling and the volume is constrained by the Luttinger theorem [57]. However, there is a qualitative difference for ferromagnetic $\text{Sr}_4\text{Ru}_3\text{O}_{10}$. The bands for the two spins have very different fillings and very different Fermi surfaces. The results depend on the magnitude of the spin splitting which is sensitive to the functional.

We carried out first-principles calculations using the pseudopotential, plane-wave implementation of density functional theory in the QUANTUM ESPRESSO distribution [58,59]. All calculations were performed using the PBEsol exchange-correlation functional [60] and ultrasoft pseudopotentials [61] with Sr(4s, 4p, 5s), Ru(4s, 4p, 5s, 4d), and O(2s, 2p) valence states [62,63]. The wave functions were expanded in plane waves with a cutoff of 75 Ry for the kinetic energy and 600 Ry for the charge density. Gaussian smearing with a broadening of 0.007 Ry was used in all our calculations. We have used the experimental lattice parameters and atomic positions taken from Crawford *et al.* [8]. In that reference, the primitive cell contains two triple layers with 68 atoms per cell, but we have used the simplification in which the two triple layers are equivalent (which was suggested as a possibility [8]) in which case it is a base-centered orthorhombic structure with 34 atoms per cell with space group $Cmce$ (64). Given the proportions of a, b, and c of the unit cell (roughly thrice

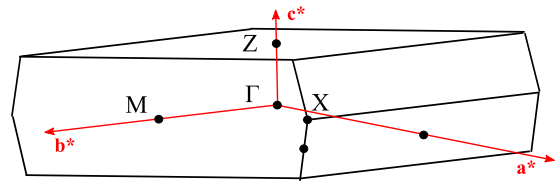


FIG. 3. BZ for $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ structure as described in the text, where a^* , b^* , c^* are the reciprocal lattice vectors for the base-centered orthorhombic structure with space group $Cmce$ (64). Since the dispersion is almost two dimensional, the points are labeled using the notation for a square lattice and the band structure plots are for lines that connect Γ - X - M - Γ points in the $k_z = 0$ plane and Γ - Z on the k_z axis.

longer along c than a and b), integrations over the BZ of this lattice were determined to be $8 \times 8 \times 3$ Monkhorst-Pack [64] k -point meshes. The grids were $32 \times 32 \times 8$ for the Fermi surface calculations. Structural relaxations were performed with atomic forces converged to within 1.0×10^{-5} eV/Å, while energies were converged to within 1.0×10^{-11} eV. The results presented here are for the optimized structure with the above criteria.

In all cases, the lowest energy is a ferromagnet with large spin polarization in general agreement with the experiment. The moment calculated for the optimized structure is $1.77\mu_B/\text{Ru}$, which is in agreement with the experimental values in the range ~ 1.0 to $\sim 1.7\mu_B/\text{Ru}$ [8,9,11,12,43,44]. The moments are largest on the central layer in agreement with a spin-polarized neutron diffraction study [44]. Since there is large spin polarization, there is a large difference in occupation of the majority- and minority-spin bands and we expect qualitative effects on the states at the Fermi energy.

The BZ for the actual orthorhombic crystal structure is shown in Fig. 3. The calculated bands are plotted for the lines Γ - X - M - Γ - Z connecting the points shown in the figure. However, since there is only small dispersion in the direction perpendicular to the layers, the bands are almost two dimensional and we have used the notation for the points in the BZ zone appropriate for a two-dimensional square lattice. This is the notation used for the ARPES experiments [52] where the data are analyzed in terms of a two-dimensional square BZ. We have investigated the k_z dispersion in other parts of the BZ and we find that the dispersion at all places studied is less than that along the Γ to Z line which is shown in the following figures for the band structures.

A. Non-spin-polarized calculations

We first present the calculated band structures and Fermi surfaces of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ for the case where the system is constrained to have no spin polarization, as shown in Fig. 4 and the left side of Fig. 6, respectively. The conclusion in Sec. V is that the results do not agree with the experiment [52]. Nevertheless, this is important to show the large difference from the spin-polarized results and it provides a simpler picture of the bands that are a guide to understanding the more complicated case where there is no spin degeneracy and there are twice as many bands.

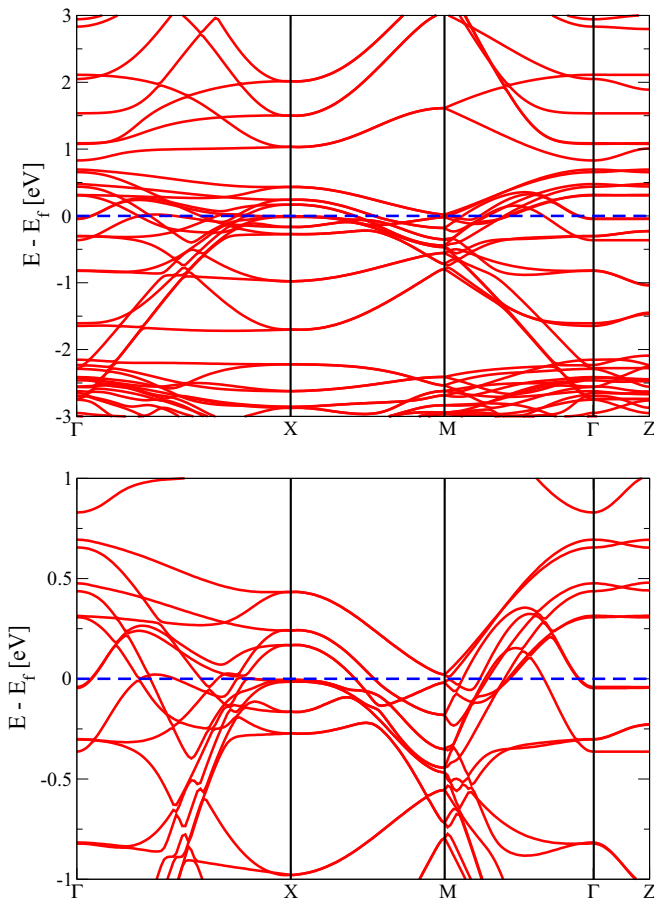


FIG. 4. Non-spin-polarized bands of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ plotted along the lines connecting the points shown in Fig. 3. The Fermi energy is set to zero and indicated by the dashed blue line. The bands at the bottom of the top figure have mainly oxygen in character; those from -2 to 1 eV are Ru t_{2g} d bands mixed with O p states. The steeply increasing bands are due to the d_{xy} states, whereas the d_{xz} and d_{yz} form pairs of states degenerate at Γ with one almost flat and the other with large dispersion along the Γ - X direction. These states in different layers are strongly coupled leading to bonding, non-bonding, and antibonding bands at approximately -1.6 , -0.8 , and 0.3 eV at Γ .

The position of the bands relative to the Fermi energy is determined by the fact that there are four electrons in the six t_{2g} states per Ru atom, i.e., $\frac{2}{3}$ filling of the t_{2g} bands. As we can see from the upper part of Fig. 4, the bands below around -2 eV are primarily oxygen in character and the bands above approximately 1 eV are e_g Ru d states. The main focuses here are the t_{2g} bands in the range -2 to 1 eV and the bands that extend to below -2 eV in the range of the oxygen bands. The t_{2g} bands are mainly Ru d character and they look remarkably like the tight-binding bands shown in Fig. 2.

This is shown most clearly by the flat bands along Γ to X that identify the bonding, nonbonding, and antibonding d_{xz} and d_{yz} bands and the images that are folded into the small BZ very much like that shown in the right side of Fig. 2. The magnitude of the splitting of the bonding and antibonding bands is approximately 1.3 eV which is very similar to the value of approximately 1.2 eV calculated for $\text{Sr}_3\text{Ru}_2\text{O}_7$ [25]. The bonding and nonbonding d_{xz} and d_{yz} bands are evident as

shown by the almost flat bands; the antibonding bands are not as clear since they are mixed with other bands that are folded into the BZ. The bonding and nonbonding combination of the d_{xz} and d_{yz} states in the three layers are mainly full and the Fermi energy is in the range of the antibonding bands, where there are many bands.

The bands that increase steeply upward going away from Γ both toward X and toward M have d_{xy} character. There are three bands that are almost decoupled from one another with one band for the central layer and two forming a degenerate pair at a different energy since the potentials at the central layer and outer layers are different.

The calculated Fermi surfaces shown in the left side of Fig. 6 are qualitatively similar to those for $\text{Sr}_3\text{Ru}_2\text{O}_7$ [24]; the shapes of the contours are similar, but the number of bands and the sizes of the pieces of Fermi surfaces are different for three layers compared to two in $\text{Sr}_3\text{Ru}_2\text{O}_7$. The flat pieces of the Fermi surface are due to the d_{xz} and d_{yz} bands, which are almost flat as a function of k_x for d_{yz} and k_y for d_{xz} . These flat pieces of the Fermi surface show how the lattice is oriented since they are perpendicular to the rows and columns of Ru atoms. Along the boundaries of the BZ, the Fermi surfaces are complicated due to many bands crossing the Fermi energy, as can be seen in Fig. 4. We have not tried to interpret these since they depend upon the details and are not relevant to the spin-polarized ferromagnetic system considered next.

B. Spin-polarized calculations

The calculated band structures and Fermi surfaces for the same crystal structure using the same PBEsol functional, but allowing for spin polarization, are shown in Figs. 5 and 6, respectively.

The ground state is found to be ferromagnetic so that there are two sets of bands for the two spins. The bands have very nearly the same form as in the non-spin-polarized calculation shown in Fig. 4, except shifted relative to one another. As expected from the large polarization, the shift of the bands is large which changes the picture qualitatively from a non-magnetic state. The majority-spin bands are shifted down and are almost filled so that they have holelike character near the Fermi energy. These are the dominant bands near the zone center and there are several bands near the Fermi energy that are quite flat, especially along the line from Γ to X .

In contrast, the minority-spin bands shifted up so that they are approximately $\frac{1}{3}$ filled and have electronlike character dispersing upward across the Fermi energy as shown in Fig. 5. There are three d_{xy} bands for the three layers that disperse steeply upward across the Fermi energy: two on the outer layers and very nearly degenerate, and one on the central layer with slightly different energy. These give rise to the nearly circular pieces of the Fermi surface shown in the right side of Fig. 6. The nonbonding combination of d_{xz} and d_{yz} are the pair of bands that are degenerate at -0.4 eV, and which disperse upward to form the square-shaped piece of Fermi surface. The character of the minority-spin bands is essentially the same for all the spin-polarized calculations reported here and shown in Figs. 5 and 9.

The Fermi surfaces for the majority and minority spins are shown in the center and right figures in Fig. 6. The surfaces are

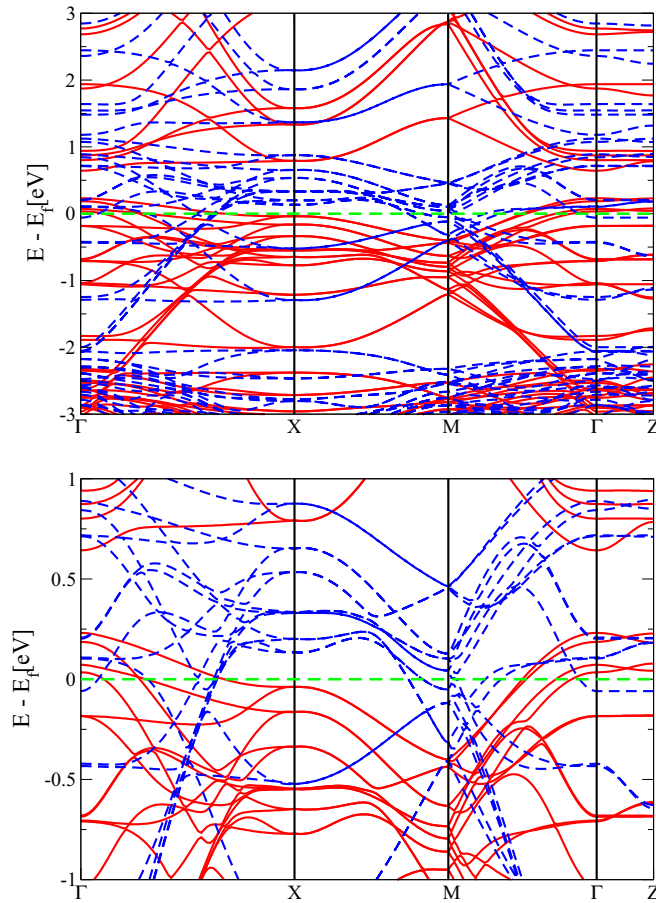


FIG. 5. Spin-polarized bands of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ plotted along the lines connecting the points shown in Fig. 3. The Fermi energy is set to zero and indicated by the dashed green line. There are two sets of bands for the majority (red solid lines) and minority bands (blue dashed lines) spin. The bands from the two spins are shifted almost rigidly, except that lower-energy parts of the majority-spin bands are strongly mixed with oxygen bands. As discussed in the text, the bands for each spin are very similar to the bands for the non-spin-polarized case in Fig. 4. The Fermi surfaces for the two spin states are shown in the center and right figures in Fig. 6.

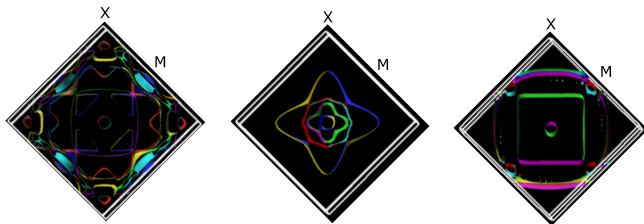


FIG. 6. Fermi surfaces for $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ resulting from two different calculations. At the left is the result for a calculation constrained to have no spin polarization corresponding to the bands in Fig. 4. The two figures in the center and right are the Fermi surfaces for the majority and minority spins in the spin-polarized calculation corresponding to the bands in Fig. 5. As described in the text, the spin-polarized calculations explain the major features of the ARPES results shown in Fig. 11.

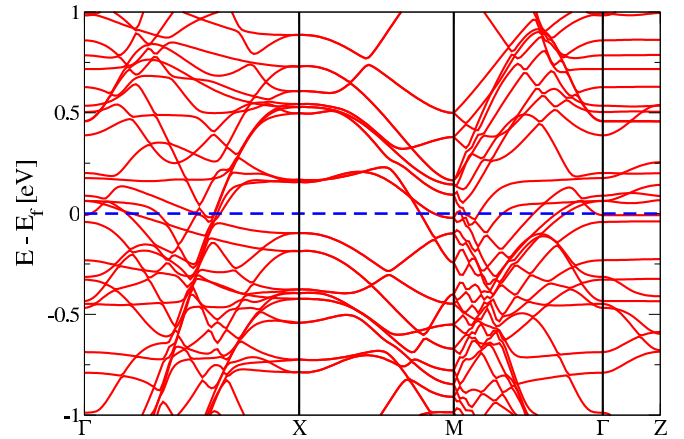


FIG. 7. Bands of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ with the inclusion of spin-orbit coupling plotted along the lines connecting the points shown in Fig. 3. The Fermi energy is set to zero and indicated by the dashed blue line. The changes are mainly small shifts of bands and small gaps; however, this has large effect on the narrow bands near the Fermi energy, which can be seen in the Fermi surfaces in Fig. 8.

qualitatively different from those for the non-spin-polarized system shown at the left in Fig. 6. The surfaces for each spin state are much simpler and have more straightforward interpretations. Because of the large polarization, pieces of the Fermi surface near the zone center are mainly due to the majority-spin bands, whereas pieces nearer than the BZ boundary are mainly due to the minority spins.

The pieces of the Fermi surface shaped like extended ovals with the long direction pointing toward the X points at the BZ corners are due to the majority-spin d_{xz} and d_{yz} bands, and they are very similar to t_{2g} bands found in many two-dimensional systems, such as oxide interfaces in [65]. However, in the present case, the position of the majority-spin bands relative to the Fermi energy is very sensitive to the functional used in the calculation. This is shown in the calculations described below where the Hubbard U correction leads to a reduction of the size of the ovals until they disappear completely. Comparison of the bands with ARPES experiments is presented in Sec. V.

C. Spin-orbit coupling

The calculated band structures and Fermi surfaces for the same crystal structure using the same PBEsol functional with the inclusion of the spin-orbit couplings are shown in Figs. 7 and 8, respectively. Since the spin-orbit effects are small compared to the bandwidths, the bands are still primarily majority and minority spin. By comparing the bands in Figs. 4 and 7, we can see that the effects are mainly to open small gaps and remove band crossings. The effects are most significant in cases where there are band crossings for narrow bands near the Fermi energy, which is illustrated by the Fermi surfaces shown in Fig. 8. Note that in all the calculations, there are large square parts of the Fermi surfaces, which is the most prominent feature found in the ARPES experiments as described in Sec. V.

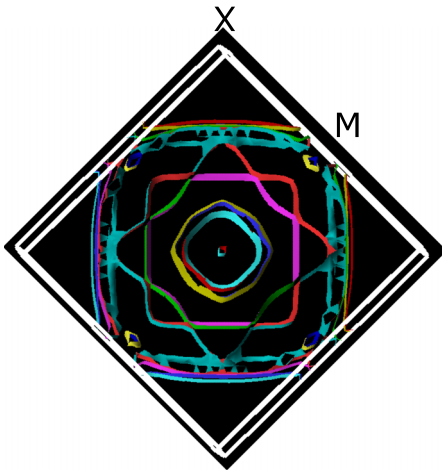


FIG. 8. Fermi surfaces for $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ including spin-orbit coupling corresponding to the bands in Fig. 7; the surfaces are close to a sum of those for the majority and minority spins shown, respectively, in the center and right figures in Fig. 6, with changes mainly at the places where there are narrow bands.

D. Including a Hubbard U correction

We also performed DFT + U calculation by applying a Hubbard U correction on the Ru $4d$ states [66]. Calculations with similar functionals with and without an added U have been reported for Sr_2RuO_4 [53,54], $\text{Sr}_3\text{Ru}_2\text{O}_7$ [55], and SrRuO_3 [56]. The values of U used in the previous studies [1,53–56,67–69] were in the region 1.5–4.0 eV. Here we present results for two moderate values, $U = 1.0$ and $= 2.0$ eV, which show the large effects on the electronic properties of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$.

The calculated band structures using the PBEsol + U functional with the inclusion of Hubbard $U = 1.0$ and 2.0 eV are shown in Fig. 9. The crystal structure is the same as in the calculations for $U = 0$ eV. As expected, the addition of a repulsive interaction U for electrons in the Ru $3d$ orbitals leads to a larger splitting of the bands. For $U = 1.0$ eV the highest states of the majority-spin band just touch the Fermi energy. This is the dividing point where the nature of the Fermi surface changes. For $U < \approx 1.0$ eV, there are pieces of Fermi surface which are holelike around the zone center for the majority-spin bands. For $U > \approx 1.0$ eV the majority-spin bands are completely filled and there are only minority-spin bands at the Fermi energy. As U increases above this value, there are small quantitative changes in the shapes of the bands and the Fermi surfaces since the filling of the minority-spin bands is fixed at $\frac{1}{3}$. This is illustrated in Fig. 9 where we can see that the majority-spin bands shift downward relative to the Fermi energy for $U = 2.0$ eV, but the minority-spin bands stay almost the same.

The minority-spin Fermi surface for $U = 2.0$ eV is shown in Fig. 10. It is essentially the same for $U = 1.0$ eV (not shown) and, in fact, the minority-spin Fermi surface is qualitatively the same for $U = 0$, which can be seen by comparing Fig. 10 with the Fermi surface at the right side in Fig. 6. Thus, a definitive result of our calculations is that for any reasonable functional, the result for the minority-spin Fermi surface has

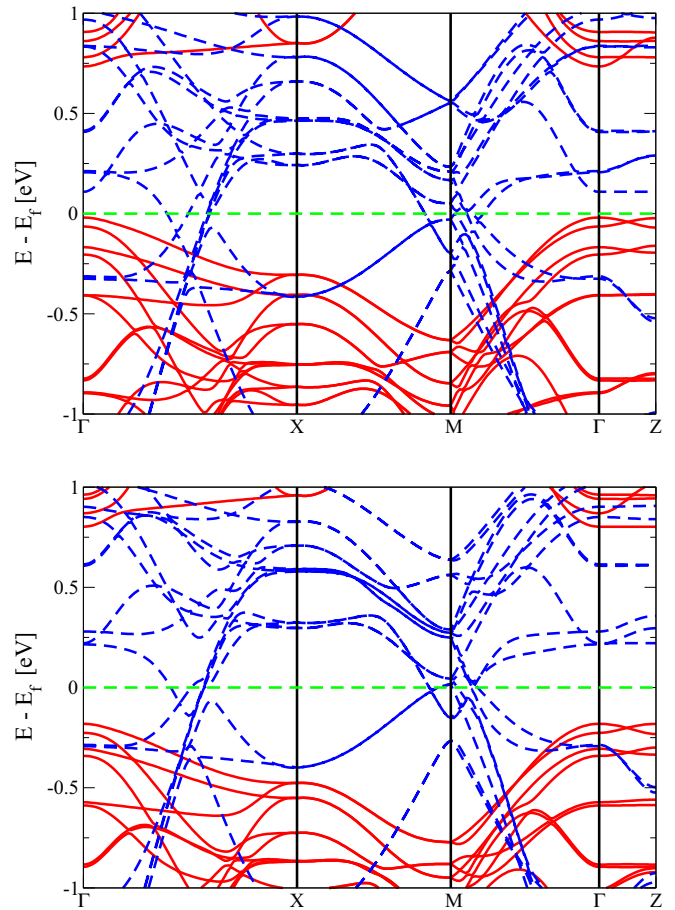


FIG. 9. Spin-polarized bands of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ with the inclusion of the Hubbard U interaction on the Ru atoms plotted along the lines connecting the points shown in Fig. 3. The Fermi energy is set to zero and indicated by the dashed green line. The upper figure is for $U = 1.0$ eV and the lower figure is for $U = 2.0$ eV. As U increases, the spin splitting increases and at $U = 1.0$ eV the majority-spin band is filled and just touches the Fermi energy. For larger U the system is fully polarized, a half-metallic system with only minority-spin states at the Fermi energy. The Fermi surfaces for the minority spin for $U = 2.0$ eV are shown in Fig. 10; the surfaces for $U = 1.0$ eV are very similar and are not shown.

qualitatively the same form. As discussed in the following section, this is the major feature seen in the experimental ARPES data.

V. COMPARISON WITH ARPES EXPERIMENTS

In this section we compare the theoretical calculations presented above with recent ARPES experiments [52] which are focused on the Fermi surface and the bands near the Fermi energy. Since the experiments do not resolve the spin, there is no direct evidence of the ferromagnetic moments nor of the difference between majority and minority spins. Thus, the conclusions regarding the effects of spin polarization are prediction of the theory. Nevertheless, the experimental results [52] provide definitive tests of the theory because the energies of the bands are greatly modified by the polarization, and we

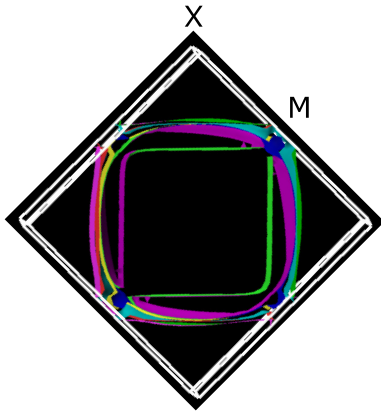


FIG. 10. The Fermi surface for the minority-spin bands calculated with the inclusion of Hubbard U . The results shown are for $U = 2.0$ eV and the Fermi surface is essentially the same for $U = 1.0$ eV. There are no bands for the majority spin that cross the Fermi energy, but there are bands near the Γ point that are close to the Fermi energy as shown in Fig. 9.

can compare the predicted and measured energies of the bands and shapes of the Fermi surfaces.

In the comparison with the ARPES experiments, it is important to keep in mind that the experiments may be affected by changes of the structure at the surface, and there may be surface states like those found, for example, in ARPES for Sr_2RuO_4 [22]. However, careful studies to distinguish bulk and surface states are not reported in the experimental paper [52].

A. Fermi surfaces

First, we consider the Fermi surfaces because this is the information that shows most clearly the primary result of this work: the large effect of the ferromagnetism on the electronic spectrum. In Fig. 11 is shown the intensity at the Fermi energy as a function of the momentum in two dimensions parallel to the layers. The figure is a reproduction of Fig. 1(a) in [52], where the data are presented with the square BZ oriented as shown (with the sides vertical and horizontal) which corresponds to the way the data was acquired. In order to compare with the present calculations which are presented in the conventional orientation of the BZ, the image must be turned by 45° ; this is only a choice of presentation that does not affect the conclusions.¹ The bright areas indicate states at or very near the Fermi energy.

The dominant feature in Fig. 11 is the square-shaped region indicated as α_1 and α_2 . This is in good agreement (when turned by 45°) with the sets of bands in the Fermi surfaces for the minority spin calculated for the system with ferromagnetic order: the right side of Fig. 6 for the calculation with $U = 0$ or Fig. 8 where the spin-orbit coupling is included, or Fig. 10,

¹In the experimental paper [52] it was stated that at their measurement temperature the structure was modified so that the BZ was turned by 45° ; however, this is not needed and the analysis of the experiments is consistent with the structure we use here.

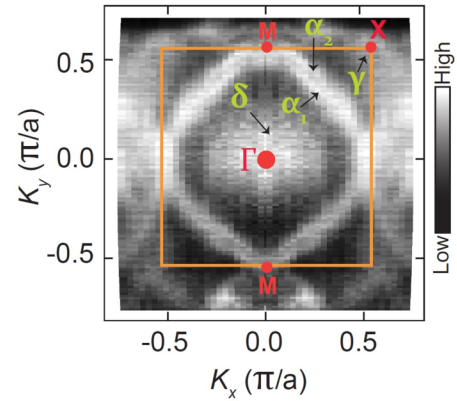


FIG. 11. The experimental Fermi surface reported in [52]. This is taken from Fig. 1(a) of [52] with no changes. As explained in the text, it must be turned by 45° to compare with the present theoretical calculations. The white areas correspond to the large intensity at the Fermi energy, and it is immediately apparent that the primary feature, the large roughly square shape, agrees well with the calculated Fermi surface for the ferromagnetic calculations (Fig. 8 or 10) and it is qualitatively different from the calculated Fermi surface if the system is forced to have no spin polarization (the left side of Fig. 6).

which shows the Fermi surface calculated with $U = 2.0$ eV, which is essentially the same for all $U > 1.0$ eV where there is only the minority-spin Fermi surface. As shown most clearly in Fig. 10, the calculations find this to be a combination of the nearly circular pieces of Fermi surface shown due to the three nearly degenerate d_{xy} bands in the three Ru-O layers and the square-shaped piece of Fermi surface due to the non-bonding combination of d_{xz} and d_{yz} states in a triple layer. The character of the bands called α_1 and α_2 is supported by the other results in the ARPES experiments, which show that these are wide bands with a large slope in agreement with the calculated minority-spin bands. These are robust results found in all the spin-polarized calculations and clear support for the interpretation in terms of spin-polarized bands.

The observation that these bands have strong intensity in the measurements is additional evidence for the theoretical conclusion that they are due to minority-spin bands. This can be understood following the description of the bands in the pedagogical model in Sec. III, where it is clear that in the lower-energy range the bands are properties of the lattice of Ru-O octahedra independent of whether the octahedra are rotated or not. This is the energy range applicable for the minority-spin bands in the DFT calculations, which have reduced occupation (approximately $\frac{1}{3}$) so that the Fermi energy is in the lower part of the minority-spin t_{2g} bands. In contrast, the upper parts of the bands are “folded” into the smaller BZ as indicated in right-hand side of Fig. 2. This is the energy range relevant for the majority-spin bands which are fully or nearly fully occupied in the DFT calculations. Thus, the intensity is nonzero only because of the rotations of the octahedra, and the intensity for these bands may be weak in the experimental measurements for momenta in the first BZ.

Thus, the major features in the ARPES data are well explained by the theoretical calculations with spin polarization and are quantitatively different from the Fermi surfaces

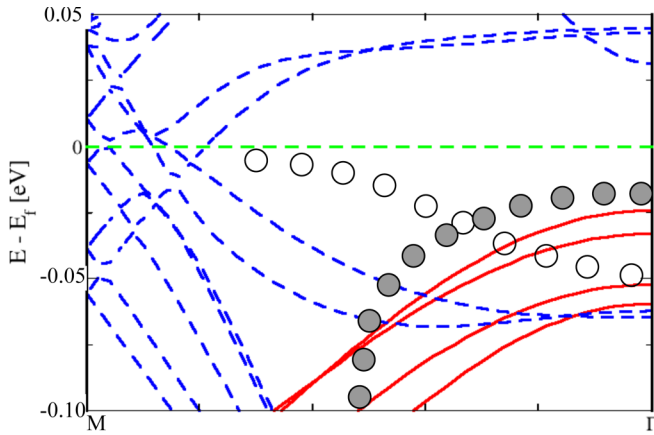


FIG. 12. Comparison of theory and experiment for narrow bands near the Fermi energy in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ along the line from Γ to M . The experimental results from [52] are represented by open and gray solid circle curves as described in the text. The theoretical bands are calculated with Hubbard $U = 1.5$ eV and they have been scaled by a factor of 5, which is similar to the renormalization factors found in the previous works of ~ 6 for $\text{Sr}_3\text{Ru}_2\text{O}_7$ [26] and 3–5 for Sr_2RuO_4 [23].

predicted by the non-spin-polarized calculations shown in the left side of Fig. 6. In that case, the occupation of the bands is $\frac{2}{3}$, the Fermi energy is in the upper parts of the bands, and there are a host of pieces of Fermi surface (left side of Fig. 6) near the zone boundaries that are not seen in the experimental data. They are analogous to results of the non-spin-polarized calculations and the experimental results for $\text{Sr}_3\text{Ru}_2\text{O}_7$ [24].

B. Narrow bands near the Fermi energy

The other prominent feature in Fig. 11 is a broad area of intensity in the region around the zone center labeled δ . Even though there is no direct evidence in this figure for a Fermi surface, this indicates narrow bands at or near the Fermi energy within the energy resolution over a large area of the BZ around the zone center. This also is in qualitative agreement with the present calculations for the ferromagnetic system where the majority-spin bands are filled or almost filled with the maxima at the zone center. The theory predicts that these bands near the top of the majority-spin bands should have weaker intensity in the experiment because they have nonzero intensity near the center of the BZ only because of the rotations of the octahedra. The observations support this interpretation and are very different from the predictions of the non-spin-polarized calculation shown in Fig. 4 where there are no holelike bands around the zone center at the Fermi energy.

In addition to the Fermi-surface map reproduced in Fig. 11, the experimental paper [52] also reports bands around the zone center. We can compare directly with the two bands along the Γ - M direction analyzed most carefully in [52]. For measurements with linear horizontal polarization (LHP) there is a band that curves upward with a bandwidth of order 40 meV; the dispersion shown in Fig. 4(c) of [52] is reproduced here as the open-circle curve in Fig. 12. For measurements with linear vertical polarization (LVP) there

is a band that has a maximum at Γ approximately 20 meV below the Fermi energy and curving downward; the dispersion reported in Fig. S2c in the Supplemental Material of [52] is shown here as the gray solid-circle curve in Fig. 12.

These results are compared with the theoretical calculations in Fig. 12, where the energies of the calculation bands have been scaled by a factor of 5. We have chosen to present the results for calculations with $U = 1.5$ eV. The theoretical results for other values of U would be similar except the majority-spin bands would be shifted upward or downward. The renormalization factor of 5 is very reasonable compared with the previous works of factors ~ 6 for $\text{Sr}_3\text{Ru}_2\text{O}_7$ [26] and 3–5 for Sr_2RuO_4 [23].

There are other features evident in the experimental ARPES results that clearly indicate intensity near the Fermi energy near the X point, labeled γ in Fig. 11 and described in the text of [52] along the M to X lines. These are also present in the calculated bands as indicated by the narrow bands along the Γ to X and M to X lines which are in the same energy range as those shown in Fig. 12. However, these features are not presented in as much detail as the results shown in Fig. 12. We do not attempt to analyze the narrow bands further. The theory is not at a stage where the bands can be predicted in detail and further work will need a combination of high-resolution experiments in conjunction with theoretical calculations. In addition, we should bear in mind various caveats: the dispersion of the bands perpendicular to the layers, the resolution reported in the paper is or order 20 meV, and there may be other effects like impurity scattering and surface states.

VI. CONSEQUENCES FOR OTHER PROPERTIES

The most novel properties of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ relate to magnetism. In addition to the ferromagnetic moment along the c axis there is metamagnetic behavior for applied magnetic fields parallel to the layers. Such a large susceptibility can be explained by narrow bands at the Fermi energy, which has been proposed as the explanation for the observed metamagnetic behavior in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ [8,9,11,12,37–41] and studied in more detail in $\text{Sr}_3\text{Ru}_2\text{O}_7$ [5,6,24–29]. Narrow bands in the right energy range are found in this work for DFT calculations using the PBEsol functional and DFT+ U for values of U up to around $U = 2.0$ eV. However, we have concluded that the states of the art for the theory and the experimental work done so far are not definitive enough to establish the nature of the narrow bands in detail.

A result of this work is that for Hubbard interaction $U > 1.0$ eV, $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ is predicted to be a half-metal, with only minority-spin bands at the Fermi energy. Such materials are potentially of technological interest for spintronics applications [70] where the transport properties involve interesting many-body effects [71]. The present calculations present interesting possibilities of a nearly-half-metal with almost filled majority-spin bands that have large susceptibility due to narrow bands at the Fermi energy. This would combine effects related to metamagnetism and half-magnetism that could be controlled by applied fields. For larger interactions U , there would be no narrow bands near the Fermi energy and the calculations would find behavior more like other half-metals.

VII. CONCLUSIONS

The main conclusion of this work is the effect on the electronic bands of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ due to the ferromagnetic order. Because the bands are spin polarized with a large ferromagnetic moment (known experimentally and found in the theoretical calculations), the majority-spin bands are full or almost full and the minority-spin bands have an occupancy of approximately $\frac{1}{3}$. The result is that there are wide bands crossing the Fermi energy due to the minority spin, whereas there are narrow bands at or near the Fermi energy for the majority spin. When the spin-orbit coupling is included, there is still the distinction of two types of bands and there is additional structure in the narrow bands near the Fermi energy. The fact that there are several narrow bands leads to a large density of states that varies rapidly near the Fermi energy, which is a way to understand metamagnetism as well as more than one metamagnetic transition as has been observed in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ [40].

The experimental ARPES measurements [52] are in good agreement with the theoretical predictions for spin-polarized bands and qualitatively different from the calculated bands if the system is constrained to be non-spin-polarized. As discussed in Sec. V, the wide bands for the minority spins explain the dominant features of the ARPES data that are roughly square and circular Fermi surfaces, whereas the narrow bands of the majority spins are in qualitative agreement with the observations of narrow bands around the zone center. The features due to the minority spins remain the same in all the spin-polarized calculations. However, the position of the majority-spin bands relative to the Fermi energy is very sensitive to the choice of the functional, which is illustrated by the large changes as a function of an added Hubbard U interaction on the Ru d states.

Finally, it is important to put this work in perspective in the big picture of electronic interactions and correlation in the layered strontium ruthenates and, more generally, transition metal oxides. The ARPES measurements indicate

narrow bands roughly a factor of 5 narrower than those found in the present DFT calculations. This is comparable to the other layered strontium ruthenates [23,26] and indicates similar effects of correlations. In addition, however, there is another effect of interactions, the magnitude of the spin splitting in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ due to the ferromagnetic order. Our results show that there are narrow majority-spin bands for a range of added interactions U on the Ru d states, which appear to be qualitatively consistent with the experimental data. But, our knowledge of functionals and the energy resolution of the ARPES experiments done so far are not sufficient to pin down details. Exploring the nature of the narrow bands in future work can provide insight into the role of interactions and correlations in the entire range of these layered systems and transition metal oxides in general.

Note added. Recently, we learned of previous ARPES measurements that are reported in talks in 2005 and 2008, but were never published. To our knowledge, they are consistent with the data in [52]. See Ref. [72] for more details.

ACKNOWLEDGMENTS

This work was supported by the 2019 Innovation Fund of the American Physical Society through the U.S.-Africa Initiative in Electronic Structure and by the Abdus Salam International Centre for Theoretical Physics. We gratefully acknowledge the African School for Electronic Structure Methods and Applications (ASESMA) without which this work could not have been done. We thank E. Carleschi, J. Chang, M. Cuoco, B. P. Doyle, and M. Horio for fruitful discussions. This work used the Extreme Science and Engineering Discovery Environment (XSEDE), which is supported by National Science Foundation Grant No. ACI-1548562, and the Centre for High Performance Computing (CHPC), South Africa. Finally, we are indebted to A. Fodorov, J. Denlinger, and J. Allen for bringing to our attention unpublished ARPES measurements on $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ referred to in [72].

-
- [1] M. Malvestuto, E. Carleschi, R. Fittipaldi, E. Gorelov, E. Pavarini, M. Cuoco, Y. Maeno, F. Parmigiani, and A. Vecchione, Electronic structure trends in the $\text{Sr}_{n+1}\text{Ru}_n\text{O}_{3n+1}$ family ($n = 1, 2, 3$), *Phys. Rev. B* **83**, 165121 (2011).
 - [2] C. Bergemann, A. P. Mackenzie, S. R. Julian, D. Forsythe, and E. Ohmichi, Quasi-two-dimensional Fermi liquid properties of the unconventional superconductor Sr_2RuO_4 , *Adv. Phys.* **52**, 639 (2003).
 - [3] A. P. Mackenzie and Y. Maeno, The superconductivity of Sr_2RuO_4 and the physics of spin-triplet pairing, *Rev. Mod. Phys.* **75**, 657 (2003).
 - [4] S.-I. Ikeda and Y. Maeno, Magnetic properties of bilayered $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Phys. B (Amsterdam)* **259**, 947 (1999).
 - [5] S. A. Grigera, R. S. Perry, A. J. Schofield, M. Chiao, S. R. Julian, G. G. Lonzarich, S. I. Ikeda, Y. Maeno, A. J. Millis, and A. P. Mackenzie, Magnetic field-tuned quantum criticality in the metallic ruthenate $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Science* **294**, 329 (2001).
 - [6] R. A. Borzi, S. A. Grigera, J. Farrell, R. S. Perry, S. J. S. Lister, S. L. Lee, D. A. Tennant, Y. Maeno, and A. P. Mackenzie, Formation of a nematic fluid at high fields in $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Science* **315**, 214 (2007).
 - [7] P. Rivero, V. Meunier, and W. Shelton, Half-metallic ferromagnetism in $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Phys. Rev. B* **95**, 195106 (2017).
 - [8] M. K. Crawford, R. L. Harlow, W. Marshall, Z. Li, G. Cao, R. L. Lindstrom, Q. Huang, and J. W. Lynn, Structure and magnetism of single crystal $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ a ferromagnetic triple-layer ruthenate, *Phys. Rev. B* **65**, 214412 (2002).
 - [9] F. Weickert, L. Civale, B. Maiorov, M. Jaime, M. B. Salamon, E. Carleschi, A. M. Strydom, R. Fittipaldi, V. Granata, and A. Vecchione, Missing magnetism in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$: Indication for antisymmetric exchange interaction, *Sci. Rep.* **7**, 3867 (2017).
 - [10] G. Koster, L. Klein, W. Siemons, G. Rijnders, J. S. Dodge, C.-B. Eom, D. H. A. Blank, and M. R. Beasley, Structure, physical properties, and applications of SrRuO_3 thin films, *Rev. Mod. Phys.* **84**, 253 (2012), and references therein.
 - [11] G. Cao, L. Balicas, W. H. Song, Y. P. Sun, Y. Xin, V. A. Bondarenko, J. W. Brill, S. Parkin, and X. N. Lin, Competing ground states in triple-layered $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ verging on itinerant

- ferromagnetism with critical fluctuations, *Phys. Rev. B* **68**, 174409 (2003).
- [12] Z. Q. Mao, M. Zhou, J. Hooper, V. Golub, and C. J. O'Connor, Phase Separation in the Itinerant Metamagnetic Transition of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, *Phys. Rev. Lett.* **96**, 077205 (2006).
- [13] G. Cao, S. McCall, and J. E. Crow, Observation of itinerant ferromagnetism in layered $\text{Sr}_3\text{Ru}_2\text{O}_7$ single crystals, *Phys. Rev. B* **55**, R672 (1997).
- [14] K. Ishida, H. Mukuda, Y. Kitaoka, K. Asayama, Z. Q. Mao, Y. Mori, and Y. Maeno, Spin-triplet superconductivity in Sr_2RuO_4 identified by ^{17}O knight shift, *Nature (London)* **396**, 658 (1998).
- [15] A. P. Mackenzie, R. K. W. Haselwimmer, A. W. Tyler, G. G. Lonzarich, Y. Mori, S. Nishizaki, and Y. Maeno, Extremely Strong Dependence of Superconductivity on Disorder in Sr_2RuO_4 , *Phys. Rev. Lett.* **80**, 161 (1998).
- [16] G. M. Luke, Y. Fudamoto, K. M. Kojima, M. I. Larkin, J. Merrin, B. Nachumi, Y. J. Uemura, Y. Maeno, Z. Q. Mao, Y. Mori, H. Nakamura, and M. Sigrist, Time-reversal symmetry-breaking superconductivity in Sr_2RuO_4 , *Nature (London)* **394**, 558 (1998).
- [17] Y. Maeno, T. M. Rice, and M. Sigrist, The intriguing superconductivity of strontium ruthenate, *Phys. Today* **54**(1), 42 (2001).
- [18] S.-I. Ikeda, Y. Maeno, S. Nakatsuji, M. Kosaka, and Y. Uwatoko, Ground state in $\text{Sr}_3\text{Ru}_2\text{O}_7$: Fermi liquid close to a ferromagnetic instability, *Phys. Rev. B* **62**, R6089 (2000).
- [19] M. W. Haverkort, I. S. Elfimov, L. H. Tjeng, G. A. Sawatzky, and A. Damascelli, Strong Spin-Orbit Coupling Effects on the Fermi Surface of Sr_2RuO_4 and Sr_2RhO_4 , *Phys. Rev. Lett.* **101**, 026406 (2008).
- [20] C. N. Veenstra, Z.-H. Zhu, M. Raichle, B. M. Ludbrook, A. Nicolaou, B. Slomski, G. Landolt, S. Kittaka, Y. Maeno, J. H. Dil, I. S. Elfimov, M. W. Haverkort, and A. Damascelli, Spin-Orbital Entanglement and the Breakdown of Singlets and Triplets in Sr_2RuO_4 Revealed by Spin- and Angle-Resolved Photoemission Spectroscopy, *Phys. Rev. Lett.* **112**, 127002 (2014).
- [21] Y. Maeno, H. Hashimoto, K. Yoshida, S. Nishizaki, T. Fujita, J. G. Bednorz, and F. Lichtenberg, Superconductivity in a layered perovskite without copper, *Nature (London)* **372**, 532 (1994).
- [22] A. Damascelli, D. H. Lu, K. M. Shen, N. P. Armitage, F. Ronning, D. L. Feng, C. Kim, Z.-X. Shen, T. Kimura, Y. Tokura, Z. Q. Mao, and Y. Maeno, Fermi Surface, Surface States, and Surface Reconstruction in Sr_2RuO_4 , *Phys. Rev. Lett.* **85**, 5194 (2000).
- [23] A. Tamai, M. Zingl, E. Rozbicki, E. Cappelli, S. Riccò, A. de la Torre, S. McKeown Walker, F. Y. Bruno, P. D. C. King, W. Meevasana, M. Shi, M. Radović, N. C. Plumb, A. S. Gibbs, A. P. Mackenzie, C. Berthod, H. U. R. Strand, M. Kim, A. Georges, and F. Baumberger, High-Resolution Photoemission on Sr_2RuO_4 Reveals Correlation-Enhanced Effective Spin-Orbit Coupling and Dominantly Local Self-Energies, *Phys. Rev. X* **9**, 021048 (2019).
- [24] A. Tamai, M. P. Allan, J. F. Mercure, W. Meevasana, R. Dunkel, D. H. Lu, R. S. Perry, A. P. Mackenzie, D. J. Singh, Z.-X. Shen, and F. Baumberger, Fermi Surface and Van Hove Singularities in the Itinerant Metamagnet $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Phys. Rev. Lett.* **101**, 026407 (2008).
- [25] D. J. Singh and I. I. Mazin, Electronic structure and magnetism of $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Phys. Rev. B* **63**, 165101 (2001).
- [26] M. P. Allan, A. Tamai, E. Rozbicki, M. H. Fischer, J. Voss, P. D. C. King, W. Meevasana, S. Thirupathiah, E. Rienks, J. Fink, D. A. Tennant, R. S. Perry, J. F. Mercure, M. A. Wang, J. Lee, C. J. Fennie, E.-A. Kim, M. J. Lawler, K. M. Shen, A. P. Mackenzie *et al.*, Formation of heavy d-electron quasiparticles in $\text{Sr}_3\text{Ru}_2\text{O}_7$, *New J. Phys.* **15**, 063029 (2013).
- [27] J. Lee, M. P. Allan, M. A. Wang, J. Farrell, S. A. Grigera, F. Baumberger, J. C. Davis, and A. P. Mackenzie, Heavy d-electron quasiparticle interference and real-space electronic structure of $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Nat. Phys.* **5**, 800 (2009).
- [28] S. A. Grigera, P. Gegenwart, R. A. Borzi, F. Weickert, A. J. Schofield, R. S. Perry, T. Tayama, T. Sakakibara, Y. Maeno, A. G. Green, and A. P. Mackenzie, Disorder-sensitive phase formation linked to metamagnetic quantum criticality, *Science* **306**, 1154 (2004).
- [29] B. Binz and M. Sigrist, Metamagnetism of itinerant electrons in multi-layer ruthenates, *Europhys. Lett.* **65**, 816 (2004).
- [30] H. T. Dang, J. Mravlje, A. Georges, and A. J. Millis, Electronic correlations, magnetism, and Hund's rule coupling in the ruthenium perovskites SrRuO_3 and CaRuO_3 , *Phys. Rev. B* **91**, 195149 (2015).
- [31] L. Klein, J. S. Dodge, C. H. Ahn, G. J. Snyder, T. H. Geballe, M. R. Beasley, and A. Kapitulnik, Anomalous Spin Scattering Effects in the Badly Metallic Itinerant Ferromagnet SrRuO_3 , *Phys. Rev. Lett.* **77**, 2774 (1996).
- [32] P. B. Allen, H. Berger, O. Chauvet, L. Forro, T. Jarlborg, A. Junod, B. Revaz, and G. Santi, Transport properties, thermodynamic properties, and electronic structure of SrRuO_3 , *Phys. Rev. B* **53**, 4393 (1996).
- [33] P. Ngabonziza, Y. Wang, P. A. van Aken, J. Maier, and J. Mannhart, Inelastic electron tunneling spectroscopy at high-temperatures, *Adv. Mater.* **33**, 2007299 (2021).
- [34] Y. Ding, X. Xu, A. Bhalla, X. Yang, J. Chen, and C. Chen, Switchable diode effect in BaZrO_3 thin films, *RSC Adv.* **6**, 60074 (2016).
- [35] V. Sampath Kumar and M. K. Niranjan, Effect of surface structure on workfunction and Schottky-barrier height in $\text{SrRuO}_3/\text{SrTiO}_3$ (001) heterojunctions, *J. Appl. Phys.* **115**, 173705 (2014).
- [36] S. Lee, B. A. Apgar, and L. W. Martin, Strong visible-light absorption and hot-carrier injection in $\text{TiO}_2/\text{SrRuO}_3$ heterostructures, *Adv. Energy Mater.* **3**, 1084 (2013).
- [37] Y. J. Jo, L. Balicas, N. Kikugawa, E. S. Choi, K. Storr, M. Zhou, and Z. Q. Mao, Orbital-dependent metamagnetic response in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, *Phys. Rev. B* **75**, 094413 (2007).
- [38] D. Fobes, T. J. Liu, Z. Qu, M. Zhou, J. Hooper, M. Salamon, and Z. Q. Mao, Anisotropy of magnetoresistivities in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$: Evidence for an orbital-selective metamagnetic transition, *Phys. Rev. B* **81**, 172402 (2010).
- [39] Y. Liu, J. Yang, W. Wang, H. Du, W. Ning, L. Ling, W. Tong, Z. Qu, Z. Yang, M. Tian, G. Cao, and Y. Zhang, Size effect on the magnetic phase in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, *New J. Phys.* **18**, 053019 (2016).
- [40] E. Carleschi, B. P. Doyle, R. Fittipaldi, V. Granata, A. M. Strydom, M. Cuoco, and A. Vecchione, Double metamagnetic transition in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, *Phys. Rev. B* **90**, 205120 (2014).
- [41] W. Schottenhamel, M. Abdel-Hafeez, R. Fittipaldi, V. Granata, A. Vecchione, M. Hücker, A. U. B. Wolter, and B. Büchner,

- Dilatometric study of the metamagnetic and ferromagnetic phases in the triple-layered $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ system, *Phys. Rev. B* **94**, 155154 (2016).
- [42] Z. Xu, X. Xu, R. S. Freitas, Z. Long, M. Zhou, D. Fobes, M. Fang, P. Schiffer, Z. Mao, and Y. Liu, Magnetic, electrical transport, and thermoelectric properties of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$: Evidence for a field-induced electronic phase transition at low temperatures, *Phys. Rev. B* **76**, 094405 (2007).
- [43] M. Zhou, J. Hooper, D. Fobes, Z. Mao, V. Golub, and C. O'Connor, Electronic and magnetic properties of triple-layered ruthenate $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ single crystals grown by a floating-zone method, *Mater. Res. Bull.* **40**, 942 (2005).
- [44] F. Forte, L. Capogna, V. Granata, R. Fittipaldi, A. Vecchione, and M. Cuoco, Suppression of the orbital magnetic moment driven by electronic correlations in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, *Phys. Rev. B* **100**, 104440 (2019).
- [45] T. Oguchi, Electronic band structure of the superconductor Sr_2RuO_4 , *Phys. Rev. B* **51**, 1385 (1995).
- [46] D. J. Singh, Relationship of Sr_2RuO_4 to the superconducting layered cuprates, *Phys. Rev. B* **52**, 1358 (1995).
- [47] A. Herklotz and K. Dörr, Characterization of tetragonal phases of SrRuO_3 under epitaxial strain by density functional theory, *Eur. Phys. J. B* **88**, 60 (2015).
- [48] N. Miao, N. C. Bristowe, B. Xu, M. J. Verstraete, and P. Ghosez, First-principles study of the lattice dynamical properties of strontium ruthenate, *J. Phys.: Condens. Matter.* **26**, 035401 (2013).
- [49] A. Damascelli, Z. Hussain, and Z.-X. Shen, Angle-resolved photoemission studies of the cuprate superconductors, *Rev. Mod. Phys.* **75**, 473 (2003), and references therein.
- [50] H. Ryu, Y. Ishida, B. Kim, J. R. Kim, W. J. Kim, Y. Kohama, S. Imajo, Z. Yang, W. Kyung, S. Hahn, B. Sohn, I. Song, M. Kim, S. Huh, J. Jung, D. Kim, T. W. Noh, S. Das, and C. Kim, Electronic band structure of (111) SrRuO_3 thin films: An angle-resolved photoemission spectroscopy study, *Phys. Rev. B* **102**, 041102(R) (2020).
- [51] D. E. Shai, C. Adamo, D. W. Shen, C. M. Brooks, J. W. Harter, E. J. Monkman, B. Burganov, D. G. Schlom, and K. M. Shen, Quasiparticle Mass Enhancement and Temperature Dependence of the Electronic Structure of Ferromagnetic SrRuO_3 thin films, *Phys. Rev. Lett.* **110**, 087004 (2013).
- [52] P. Ngabonziza, E. Carleschi, V. Zabolotnyy, A. Taleb-Ibrahimi, F. Bertran, R. Fittipaldi, V. Granata, M. Cuoco, A. Vecchione, and B. P. Doyle, Fermi surface and kink structures in $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ revealed by synchrotron-based ARPES, *Sci. Rep.* **10**, 21062 (2020).
- [53] H.-L. Huang and H.-T. Jeng, Orbital ordering and magnetism in layered perovskite ruthenate Sr_2RuO_4 , *Sci. Rep.* **10**, 7089 (2020).
- [54] C. Noce and M. Cuoco, Energy bands and Fermi surface of $\text{Sr}_3\text{Ru}_3\text{O}_7$, *Phys. Rev. B* **59**, 2659 (1999).
- [55] A. Putatunda, G. Qin, W. Ren, and D. J. Singh, Competing magnetic orders in quantum critical $\text{Sr}_3\text{Ru}_2\text{O}_7$, *Phys. Rev. B* **102**, 014442 (2020).
- [56] O. Grånäs, I. Di Marco, O. Eriksson, L. Nordström, and C. Etz, Electronic structure, cohesive properties, and magnetism of SrRuO_3 , *Phys. Rev. B* **90**, 165130 (2014).
- [57] J. M. Luttinger, Fermi surface and some simple equilibrium properties of a system of interacting fermions, *Phys. Rev.* **119**, 1153 (1960).
- [58] P. Giannozzi, S. Baroni, N. Bonini, M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, G. L. Chiarotti, M. Cococcioni, I. Dabo, A. D. Corso, S. de Gironcoli, S. Fabris, G. Fratesi, R. Gebauer, U. Gerstmann, C. Gougousis, A. Kokalj, M. Lazzeri, L. Martin-Samos *et al.*, QUANTUM ESPRESSO: a modular and open-source software project for quantum simulations of materials, *J. Phys.: Condens. Matter* **21**, 395502 (2009).
- [59] P. Giannozzi, O. Andreussi, T. Brumme, O. Bunau, M. B. Nardelli, M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, M. Cococcioni, N. Colonna, I. Carnimeo, A. D. Corso, S. de Gironcoli, P. Delugas, R. A. DiStasio, A. Ferretti, A. Floris, G. Fratesi, G. Fugallo *et al.*, Advanced capabilities for materials modelling with QUANTUM ESPRESSO, *J. Phys.: Condens. Matter* **29**, 465901 (2017).
- [60] J. P. Perdew, A. Ruzsinszky, G. I. Csonka, O. A. Vydrov, G. E. Scuseria, L. A. Constantin, X. Zhou, and K. Burke, Restoring the Density-Gradient Expansion for Exchange in Solids and Surfaces, *Phys. Rev. Lett.* **100**, 136406 (2008).
- [61] D. Vanderbilt, Soft self-consistent pseudopotentials in a generalized eigenvalue formalism, *Phys. Rev. B* **41**, 7892 (1990).
- [62] A. Dal Corso, Ultrasoft pseudopotentials from the pslibrary, PSLibrary version.
- [63] A. Dal Corso, Pseudopotentials periodic table: From h to pu, *Comput. Mater. Sci.* **95**, 337 (2014).
- [64] H. J. Monkhorst and J. D. Pack, Special points for Brillouin-zone integrations, *Phys. Rev. B* **13**, 5188 (1976).
- [65] Z. Zhong, A. Tóth, and K. Held, Theory of spin-orbit coupling at $\text{LaAlO}_3/\text{SrTiO}_3$ interfaces and SrTiO_3 surfaces, *Phys. Rev. B* **87**, 161102(R) (2013).
- [66] V. I. Anisimov, J. Zaanen, and O. K. Andersen, Band theory and Mott insulators: Hubbard U instead of Stoner I, *Phys. Rev. B* **44**, 943 (1991).
- [67] M. Cuoco, C. Noce, and A. Romano, Interplay between Hund coupling and Hubbard interaction in Sr_2RuO_4 , *Phys. Rev. B* **57**, 11989 (1998).
- [68] A. Liebsch and A. Lichtenstein, Photoemission Quasiparticle Spectra of Sr_2RuO_4 , *Phys. Rev. Lett.* **84**, 1591 (2000).
- [69] Z. V. Pchelkina, I. A. Nekrasov, T. Pruschke, A. Sekiyama, S. Suga, V. I. Anisimov, and D. Vollhardt, Evidence for strong electronic correlations in the spectra of Sr_2RuO_4 , *Phys. Rev. B* **75**, 035122 (2007).
- [70] I. Žutić, J. Fabian, and S. Das Sarma, Spintronics: Fundamentals and applications, *Rev. Mod. Phys.* **76**, 323 (2004).
- [71] M. I. Katsnelson, V. Y. Irkhin, L. Chioncel, A. I. Lichtenstein, and R. A. de Groot, Half-metallic ferromagnets: From band structure to many-body effects, *Rev. Mod. Phys.* **80**, 315 (2008).
- [72] See F. Wang, J. W. Allen, J. D. Denlinger, X. N. Lin, and G. Cao, Three-dimensional band structure of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$, Bull. Am. Phys. Soc., 2005 March Meeting (<https://meetings.aps.org/link/BAPS.2005.MAR.J25.8>) and R. S. Singh, F. Wang, J. W. Allen, J. D. Denlinger, X. N. Lin, and G. Cao, Fermi surface mapping of $\text{Sr}_4\text{Ru}_3\text{O}_{10}$ using angle resolved photoemission, Bull. Am. Phys. Soc. 2008 March Meeting (<http://meetings.aps.org/link/BAPS.2008.MAR.X31.12>).