



## Effects of external pressure on the narrow-gap semiconductor $\text{Ce}_3\text{Cd}_2\text{As}_6$

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Here we report the magnetic and electronic properties of recently discovered  $\text{Ce}_3\text{Cd}_2\text{As}_6$ . At ambient pressure,  $\text{Ce}_3\text{Cd}_2\text{As}_6$  presents a semiconducting behavior with an activation gap of 74(1) meV. At 136 K, a sudden increase of the electrical resistivity and a peak in specific heat are consistent with a charge density wave transition. At low temperatures, antiferromagnetic order of the  $\text{Ce}^{3+}$  ions occurs below  $T_N = 4.0$  K with a magnetic hard axis along the  $c$  axis and a  $\Gamma_6 = |\pm 1/2\rangle$  doublet ground state. The application of external pressure strongly suppresses the charge density wave order, which is completely suppressed above 0.8(1) GPa, and induces a metallic ground state. No evidence for superconductivity is detected above 2 K. Conversely, the antiferromagnetic state is favored by pressure, reaching a transition temperature of 5.3 K at 3.8(1) GPa. Notably, the resistivity anomaly characterizing the antiferromagnetic order changes with increasing pressure, indicating that two different magnetic phases might be present in  $\text{Ce}_3\text{Cd}_2\text{As}_6$  under pressure. This change in ordering appears to be associated to the crossing of the  $T_{\text{CDW}}$  and  $T_N$  lines.

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### I. INTRODUCTION

Emergent phenomena in strongly correlated materials often stem from competing interactions near an order-disorder or a metal-insulator boundary [1]. A demonstrated strategy to uncover emergent phenomena is to tune new materials through these boundaries via chemical doping, magnetic fields, or the application of external pressure. The latter is a particularly important parameter that compresses the material and tunes the electronic structure (e.g., bandwidth, hybridization, and valence) without introducing disorder [2,3].

Notably, materials containing layers of square nets provide a platform for realizing pressure-induced phenomena. At ambient pressure, the low-dimensionality of the square nets makes them susceptible to distortions that may result in a modulation of the conduction electrons, the opening of a semiconducting gap, and the emergence of a charge density wave (CDW) [4–9]. The response of CDWs under external applied pressure is often striking but not yet fully understood. On one hand, the application of pressure may increase the electron-phonon interaction, which favors the formation of CDWs [5,10]. On the other hand, applied pressure

might enhance the interlayer interactions, resulting in a more three-dimensional Fermi surface that suppresses the nesting vectors responsible for the CDWs [10]. The latter effect is much pursued, as the suppression of CDWs typically favors the appearance of conventional superconductivity, due to the increase of the density of states at the Fermi level once the gap is closed and the electron-phonon interaction present in these materials, as observed in  $\text{NbSe}_3$ ,  $\text{Sr}_3\text{Ir}_4\text{Sn}_{13}$ ,  $\text{Ca}_3\text{Ir}_4\text{Sn}_{13}$ , and  $\text{LaPt}_2\text{Si}_2$  [11–13]. However, unconventional superconductivity may also emerge upon suppression of a CDW phase similar to  $\text{CsV}_3\text{Sb}_5$  [14].

In addition, intercalating rare-earth elements with square nets may induce further instabilities, such as magnetism in the localized  $4f$  limit or a heavy Fermi liquid when hybridization between  $4f$  and square net  $p$  bands is in the strong-coupling limit. Once more, emergent phenomena are expected at the localized-delocalized boundary, unconventional superconductivity being the classic example. In particular, several Ce-based compounds display emergent ground states and complex phase diagrams at ambient pressure and under pressure [15]. Finally, the combination of strong electronic correlations and a narrow-gap semiconducting state is a promising route to improve thermoelectricity [16,17].

The large family of “112” materials with general formula  $\text{CeM}_{1-\delta}\text{X}_2$  ( $M$  = transition or post-transition metal,  $X$  = pnictide) crystallizes in the tetragonal  $P4/nmm$  or  $I4/mmm$  structures, which contain layers of pnictide square nets, and hosts many examples of emergent ground states.  $\text{CeAgBi}_2$ , for instance, is antiferromagnetic below 6.4 K and presents several metamagnetic transitions and an anomalous Hall

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contribution for fields applied parallel to the  $c$  axis [18]. Moreover,  $\text{CeAgSb}_2$  presents a ferromagnetic order below 9.8 K, in which a topological magnetic hysteresis of the domains was recently observed [19]; also, it presents a complex evolution of the magnetic phase as a function of applied pressure [20].  $\text{CeAuSb}_2$  is another example of a complex magnetic structure under the application of magnetic fields [21–24] and has been recently proposed to host a nematic state coupled to its antiferromagnetism [25].  $\text{CeAuBi}_2$  also presents several metamagnetic transitions when fields are applied parallel to the  $c$  axis at 2 K [26], which suggests a complex evolution of its magnetic structure with field. Furthermore, this family of materials is well described by a trend in the ground state wave function of the localized moment. For members crystallizing in the  $P4/nmm$  structure, antiferromagnetism is favored by a  $\Gamma_7$  ground state, while ferromagnetism is observed in compounds with a  $\Gamma_6$  ground state [27].

Here we focus on the newly discovered narrow-gap semiconductor  $\text{Ce}_3\text{Cd}_2\text{As}_6$  [28], whose crystal structure can be viewed as the ordered version of a tetragonal  $\text{CeCd}_{0.67}\text{As}_2$  (112) parent subcell. In this case, the Cd vacancies order within the body-centered structure and give rise to the superstructure  $\text{Ce}_3\text{Cd}_2\text{As}_6$  within a monoclinic  $C2/m$  space group. At ambient pressure,  $\text{Ce}_3\text{Cd}_2\text{As}_6$  and its nonmagnetic analog  $\text{La}_3\text{Cd}_2\text{As}_6$  are both semiconductors wherein a charge density wave phase transition takes place below 136 K and 279 K, respectively [28]. The strong suppression of the charge density wave transition temperature ( $T_{\text{CDW}}$ ) from  $\text{La}_3\text{Cd}_2\text{As}_6$  to  $\text{Ce}_3\text{Cd}_2\text{As}_6$  is an indication that this phase may be completely suppressed by external pressure as  $\text{Ce}_3\text{Cd}_2\text{As}_6$  presents a more compressed crystalline structure than  $\text{La}_3\text{Cd}_2\text{As}_6$ . Also, the activation energy decreases from 105(1) meV to 74(1) meV, going from  $\text{La}_3\text{Cd}_2\text{As}_6$  to  $\text{Ce}_3\text{Cd}_2\text{As}_6$  [28]. By performing experiments under pressure on  $\text{Ce}_3\text{Cd}_2\text{As}_6$  single crystals, one expects to completely suppress the CDW phase transition and to induce a metallic ground state.

To investigate this scenario, we report the electronic and magnetic properties of  $\text{Ce}_3\text{Cd}_2\text{As}_6$  under pressure. At ambient pressure and low temperatures, an antiferromagnetic ordering of the  $\text{Ce}^{3+}$  ions takes place below 4.0 K. The entropy recovered at  $T_N$  is about 75(5) % of the full doublet entropy  $R \ln 2$ , which suggests the presence of either Kondo compensation or short-range interactions. The application of external pressure induces a metallic ground state in  $\text{Ce}_3\text{Cd}_2\text{As}_6$  and completely suppresses the charge density wave order for pressures higher than 0.8(1) GPa. However, superconductivity above 2 K was not detected, possibly due to competition with the magnetically ordered state, as the antiferromagnetic transition temperature is enhanced by pressure, reaching 5.3 K at 3.8(1) GPa. Finally, two different magnetic phases might be present in  $\text{Ce}_3\text{Cd}_2\text{As}_6$  under pressure, due to the change in the resistivity anomaly associated to the antiferromagnetic phase transition.

## II. EXPERIMENTAL DETAILS

Platelike single crystals of  $\text{Ce}_3\text{Cd}_2\text{As}_6$  and  $\text{La}_3\text{Cd}_2\text{As}_6$  were grown by the vapor transport technique. First a polycrystalline seed of  $1\text{La}/\text{Ce}:0.7\text{Cd}:2\text{As}$  was prepared via solid state reaction at 800 °C. Then the polycrystalline powder was

loaded along with iodine in a quartz tube, which was sealed in vacuum. The tube was kept in a temperature gradient from 830 °C to 720 °C. The initial polycrystalline material was kept in the hot zone, and the single crystals precipitated in the cold zone. The synthesized phase was checked by single-crystal x-ray diffraction and energy-dispersive x-ray spectroscopy (EDX) resulting in 1:0.67:2 stoichiometry within experimental error. Magnetic susceptibility was measured in a commercial Quantum Design magnetic property measurement system (MPMS). The specific heat was measured using a commercial small mass calorimeter that employs a quasiadiabatic thermal relaxation technique in a commercial Quantum Design physical property measurement system (PPMS). Electrical resistivity ( $\rho$ ) was performed in a commercial PPMS along with AC bridges (Lakeshore AC372 resistance bridge or Linear Research LR-700 AC resistance bridge) or using the ACT option of the PPMS and measured in the  $ab$  plane in a standard four-probe configuration. Pressures up to 1.9(1) GPa were generated using a self-contained double-layer piston-cylinder-type Cu-Be pressure cell (PCC) with an inner cylinder of hardened NiCrAl [3]. For pressures to 3.8(1) GPa a modified Bridgman anvil cell (mBAC) was used [29]. The pressure-transmitting media used were Daphne oil 7373 and a 1:1 mixture of  $n$ -pentane and isopentane, respectively. Lead served as a low-temperature manometer in both pressure cells [3].

## III. RESULTS

For simplicity, we refer to the tetragonal parent subcell  $I4/mmm$  (Fig. S1 of the Supplemental Material [35]) as the structure of  $\text{Ce}_3\text{Cd}_2\text{As}_6$  in further discussions, without affecting the interpretation of our results. Figure 1(a) presents the magnetic susceptibility ( $\chi = M/H$ ) as a function of temperature at  $\mu_0 H = 1$  T applied parallel and perpendicular to the  $c$  axis of  $\text{Ce}_3\text{Cd}_2\text{As}_6$ . These measurements find that the  $c$  axis is the hard axis. At low temperatures, the  $\text{Ce}^{3+}$  moments order antiferromagnetically below  $T_N = 4.0$  K, as shown in the left inset of Fig. 1(a). The right inset shows the inverse of the magnetic susceptibility for the polycrystalline average ( $\frac{1}{3}\chi_{//c} + \frac{2}{3}\chi_{\perp c}$ ). By performing linear fits in the high-temperature range ( $T \geq 150$  K), a  $\text{Ce}^{3+}$  effective moment of 2.6(1)  $\mu_B$  was obtained, in agreement with the theoretical value for a free  $\text{Ce}^{3+}$  ion. From the Curie-Weiss fit, we also extracted a Weiss temperature of  $-27(1)$  K. Figure 1(b) displays the magnetization ( $M$ ) behavior as a function of applied magnetic field at 2 K. A metamagnetic transition is observed at 1.9 T when fields are applied perpendicular to the  $c$  axis. Further, a saturation starts to develop at 3.4 T. For fields parallel to the hard  $c$  axis, the magnetization linearly increases with increasing magnetic field.

Figure 2(a) displays the specific heat divided by temperature ( $C_p/T$ ) per mole of lanthanide for both  $\text{Ce}_3\text{Cd}_2\text{As}_6$  and  $\text{La}_3\text{Cd}_2\text{As}_6$  as a function of temperature. A lambda-type peak defines the antiferromagnetic transition temperature at  $T_N = 4.0$  K, in agreement with susceptibility measurements. By subtracting the phonon contribution given by the specific heat of its nonmagnetic analog  $\text{La}_3\text{Cd}_2\text{As}_6$ , we obtain the magnetic specific heat ( $C_{\text{mag}}/T$ ), displayed in Fig. 2(b) from 0.33 K to 30 K. We note that due to the small mass of

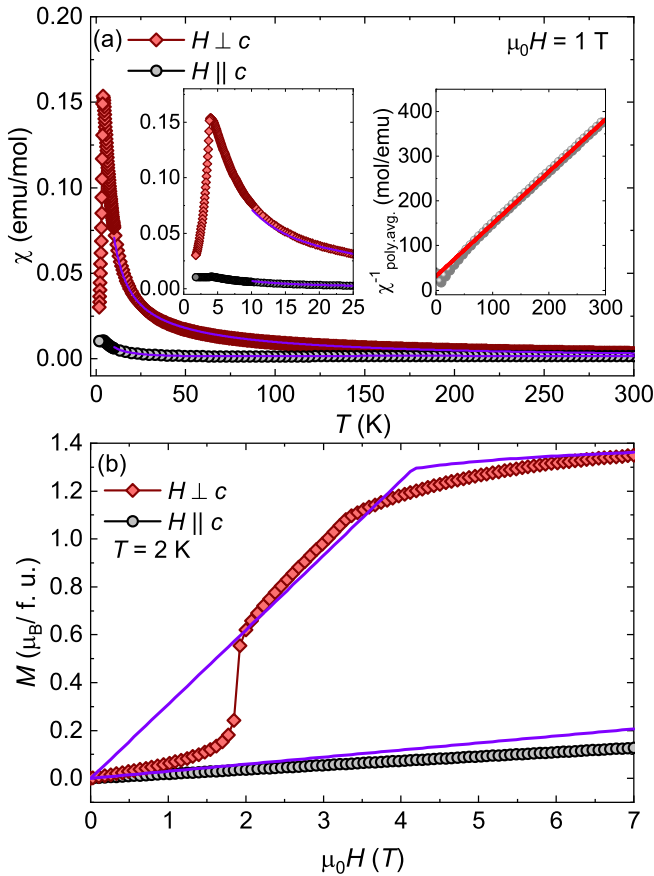


FIG. 1. (a) Magnetic susceptibility ( $\chi = M/H$ ) as a function of temperature at  $\mu_0 H = 1$  T. The left inset presents a closer view of the magnetic susceptibility at low temperatures. The right inset shows the inverse of the polycrystalline average magnetic susceptibility as a function of temperature. The solid red line is an extrapolation of a linear fit performed at high temperatures ( $T \geq 150$  K). (b) Magnetization as a function of applied magnetic field at 2 K. The solid purple lines are fits using a CEF mean-field model.

the single crystals, specific heat measurements of  $\text{La}_3\text{Cd}_2\text{As}_6$  were collected on a mosaic of 10 single crystals of total mass around 1 mg, which could lead to errors in the phonon subtraction. The bump centered around 90 K may be associated with a Schottky anomaly due to first crystal field excited state.

The solid purple lines in Fig. 1 and in Fig. 2(b) are fits using a crystalline electric field (CEF) mean-field model considering competing nearest-neighbor interactions and the tetragonal CEF Hamiltonian:  $\mathcal{H} = g_J \mu_B \mathbf{H} \cdot \mathbf{J} + z_i J_i^{\text{ex}} \cdot \langle J^{\text{ex}} \rangle + B_2^0 O_2^0 + B_4^0 O_4^0 + B_4^4 O_4^4$ , where  $g_J$  is the Landé  $g$  factor,  $\mu_B$  is the Bohr magneton,  $\mathbf{H}$  is the applied magnetic field, and  $\mathbf{J}$  is the total angular momentum.  $z_i J_i^{\text{ex}}$  represents the  $J_i$  mean-field exchange interactions ( $i = \text{AFM, FM}$ ) between the  $z_i$  nearest neighbors.  $B_n^m$  are the CEF parameters and the  $O_n^m$  are the Steven's operators [30]. By simultaneously performing fits to  $\chi(T)$ ,  $M(H)$ , and  $C_{\text{mag}}/T$  data, we extract the CEF scheme and two exchange parameters for this compound. For the CEF parameters we obtained the following values:  $B_2^0 \approx 22.8$  K,  $B_4^0 \approx -0.3$  K, and  $B_4^4 \approx -4.3$  K. These parameters imply a ground state composed of a  $\Gamma_6 = |\pm 1/2\rangle$  doublet, a first excited state  $\Gamma_7^2 = 0.40|\pm 5/2\rangle + 0.91|\mp$

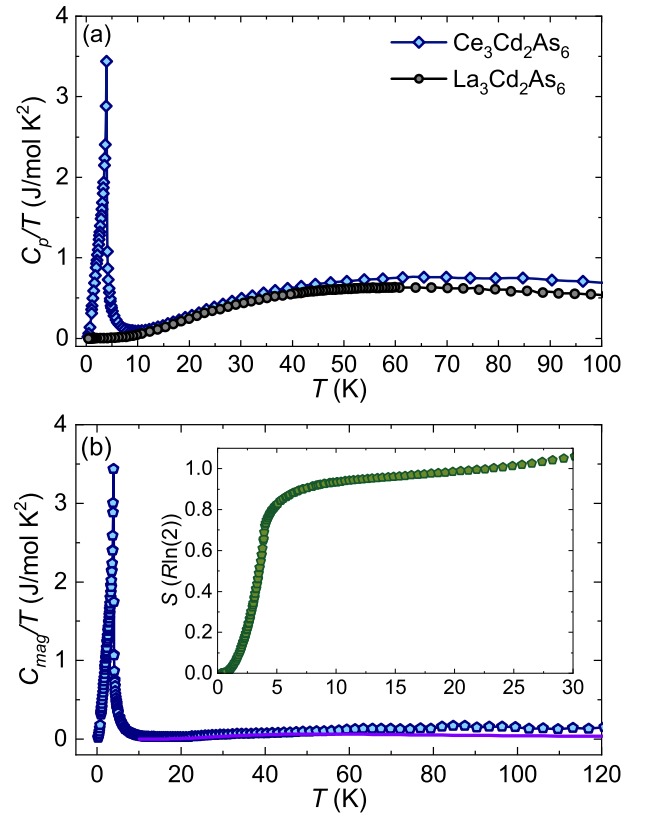


FIG. 2. (a) Specific heat measurements for  $\text{Ce}_3\text{Cd}_2\text{As}_6$  and  $\text{La}_3\text{Cd}_2\text{As}_6$ . (b) Magnetic specific heat and entropy of  $\text{Ce}_3\text{Cd}_2\text{As}_6$ . The solid purple line is a fit using a CEF mean-field model. The inset shows the entropy obtained from the magnetic specific heat as a function of temperature.

$3/2\rangle$  doublet at 160 K, and a second excited state of a  $\Gamma_7^1 = 0.91|\pm 5/2\rangle - 0.40|\mp 3/2\rangle$  doublet at 480 K. Interestingly, in spite of its  $\Gamma_6$  ground state,  $\text{Ce}_3\text{Cd}_2\text{As}_6$  is antiferromagnetic and does not follow the general trend of the “112” family of compounds, possibly due to its body-centered  $I4/mmm$  subcell and superstructure modification. For the exchange parameters we obtain  $z_{\text{AFM}} J_{\text{AFM}}^{\text{ex}} = 0.8$  K and  $z_{\text{FM}} J_{\text{FM}}^{\text{ex}} = -1.3$  K. The presence of ferromagnetic and antiferromagnetic exchange interactions is required to simulate the field-induced transition when  $H||ab$ . We note that the CEF parameters acquired from the fits of macroscopic data may not be fully accurate and must be treated with caution. The fact that this simple model does not account for the sharp feature around 1.9 T in the magnetization for  $H||ab$  indicates that more exchange parameters, or anisotropic terms such as  $J_x$ ,  $J_y$ , and  $J_z$ , may be needed to fully fit the macroscopic data of this material. Nevertheless, this model yields important information about  $\text{Ce}_3\text{Cd}_2\text{As}_6$ , such as the  $\Gamma_6$  ground state and that the  $\text{Ce}^{3+}$  ions are much localized. The integral of  $C_{\text{mag}}/T$  vs  $T$  gives the magnetic entropy presented in the inset of Fig. 2(b) that reaches 75(5) % of the doublet entropy  $R \ln(2)$  at  $T_N$ . A partial Kondo effect or short-range magnetic correlations could be responsible for the presence of entropy above  $T_N$ .

Figure 3(a) displays the electrical resistivity ( $\rho$ ) of  $\text{Ce}_3\text{Cd}_2\text{As}_6$ , which reveals semiconducting behavior. At 136 K, a sharp increase of the resistivity [displayed in the

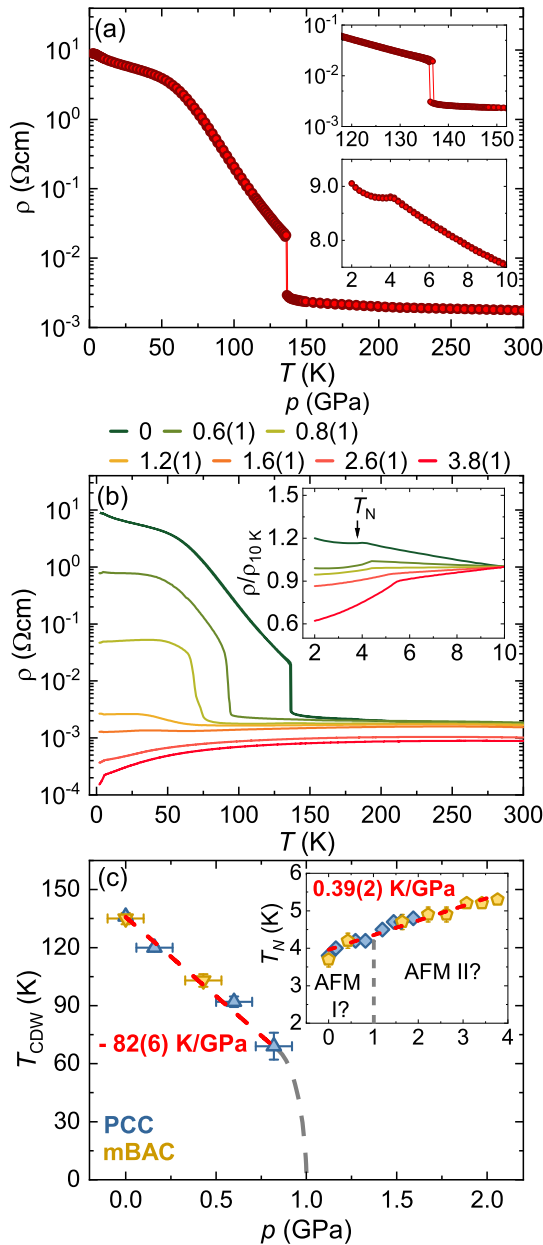


FIG. 3. (a) Electrical resistivity as a function of temperature for  $\text{Ce}_3\text{Cd}_2\text{As}_6$ . The top inset shows a closer view to the charge density wave transition at high temperatures. The bottom inset displays a zoomed-in view of the low-temperature behavior. (b) Electrical resistivity as a function of temperature at several pressures. The data up to 1.6(1) GPa were taken using a PCC. The 2.6(1) GPa and 3.8(1) GPa measurements were performed in a mBAC. We note that resistivity values for the mBAC results were obtained using the expected resistivity value for  $\text{Ce}_3\text{Cd}_2\text{As}_6$  at room temperature. The inset displays the low-temperature range ( $T \leq 10$  K) of the electrical resistivity normalized by its 10 K value. The arrow points to  $T_N$ . (c) Temperature-pressure phase diagram for  $\text{Ce}_3\text{Cd}_2\text{As}_6$ . The dashed red lines are linear fits and the dashed gray lines are a guide to the eye.

top inset of Fig. 3(a)] might be associated to a charge density wave transition.  $\text{Ce}_3\text{Cd}_2\text{As}_6$  presents Cd vacancies ordered in a stripe pattern, that breaks the parent compounds' four-fold symmetry and are accompanied by several structural

modifications [28]. The displacement of Cd ions also affects their ligands (As1, As2, and As3), leading to distorted As square nets [28]. These As square nets resemble the Te nets that host similar superstructures to  $\text{Ce}_3\text{Cd}_2\text{As}_6$  and also present CDW phases, such as  $\text{K}_{0.33}\text{Ba}_{0.67}\text{AgTe}_2$  [31]. Around 60 K, an inflection in electrical resistivity may indicate the presence of additional conduction channels prone to exist at low temperatures in narrow-gap semiconductors due to disorder [32]. At low temperatures, there is a well-defined kink around 4 K that reflects the loss of spin-disorder scattering below the antiferromagnetic transition, as displayed in the bottom inset of Fig. 3(a). The transition temperature values for the CDW and AFM transitions were inferred from the temperature-dependent resistivity data by identifying the local extrema in  $d(\ln\rho)/dT$  [33] and  $d\rho/dT$  [34] data, respectively. For  $T_{\text{CDW}}$ , the error bars were estimated based on Gaussian fits of the peaks observed in the temperature derivative of the natural logarithm of the resistivity, and for  $T_N$ , the broadness of the anomalies was used to estimate the error bars. More details about the criteria used for the transition temperatures can be found in the Supplemental Material [35].

Figure 3(b) displays the electrical resistivity as a function of temperature at several pressures. Overall the application of external pressure suppresses the semiconducting behavior and induces a metallic ground state, i.e., a positive  $d\rho/dT$ , as pressures increase above  $\approx 2$  GPa. Further increasing pressure enhances the metallic behavior of the material, which presents a residual resistance ratio (RRR)  $\approx 6$  and a room temperature resistivity of the order of 0.9 m $\Omega$  cm at 3.8(1) GPa. Furthermore, external pressure strongly suppresses the charge density wave transition temperature ( $T_{\text{CDW}}$ ) at an initial rate of  $-82(6)$  K/GPa, which cannot be identified anymore above 0.8(1) GPa (see Fig. 3(c) and Fig. S2 of the Supplemental Material [35]). We note that the critical pressure of suppression of the CDW order cannot be unambiguously determined by resistivity measurements alone. Therefore, the dashed gray lines in Fig. 3(c) represent only one possibility for the critical pressure. A suppression of the CDW order with increasing pressure was expected, as discussed before, because external pressure can destroy the nesting vectors responsible for the CDW phase, closing the energy gaps created by this phase and favoring a metallic ground state [10]. For  $\text{Ce}_3\text{Cd}_2\text{As}_6$ , the nesting vector responsible for the putative CDW phase is unknown; therefore electronic band structure calculations and angle-resolved photoemission spectroscopy (ARPES) are needed to investigate the Fermi surface of this material. A metallic state was observed at high pressures in many semiconducting materials in which the gap is created by CDW order, e.g.,  $\text{CeRhAs}$  [36–38],  $\text{Sr}_3\text{Ir}_4\text{Sn}_{13}$ ,  $\text{Ca}_3\text{Ir}_4\text{Sn}_{13}$  [12],  $\text{La}_3\text{Co}_4\text{Sn}_{13}$  [39], and many transition metal dichalcogenides [40–43]. This behavior is in contrast with the observed gap evolution as a function of pressure for Kondo insulators such as  $\text{Ce}_3\text{Bi}_4\text{Pt}_3$  wherein the gap is created by the hybridization between the conduction electrons and  $4f$  electrons. Applying external pressure enhances this hybridization, leading to larger energy gaps at high pressures [37,44]. Nevertheless, x-ray diffraction experiments under pressure are desired to investigate the lattice evolution as a function of pressure.

Conversely, the antiferromagnetic transition temperature is gradually increased by increasing pressure, reaching a  $T_N$  of

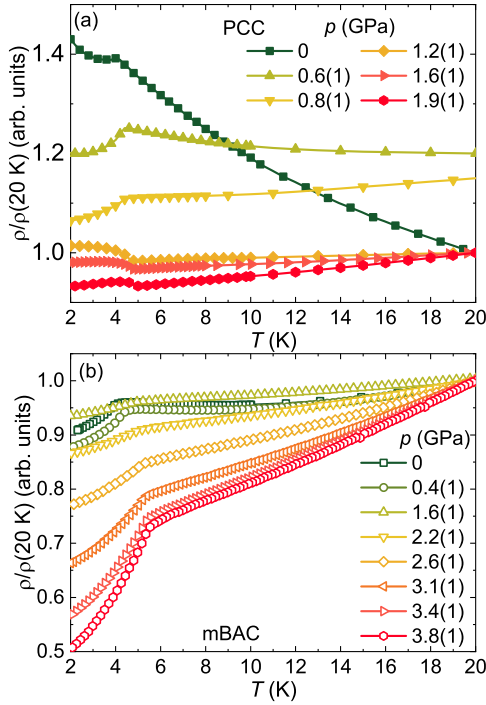


FIG. 4. Normalized electrical resistivity as a function of temperature for  $\text{Ce}_3\text{Cd}_2\text{As}_6$  for measurements using a PCC (a) and a mBAC (b) at selected pressures. The curves at 0.6(1) GPa and 0.8(1) GPa in panel (a) were vertically shifted to improve visualization.

5.3 K at 3.8(1) GPa, as shown in the inset of Fig. 3(c). The enhancement of  $T_N$  can be associated with two intertwined factors. First, the application of external pressure leads to an increase in the density of states at the Fermi level, as the material becomes more metallic with increasing pressure, favoring higher transition temperatures. Second, the exchange interaction between the localized moments might also be enhanced by pressure, favoring the magnetic order. This behavior indicates that  $\text{Ce}_3\text{Cd}_2\text{As}_6$  is in the localized limit of the Doniach diagram [45] and that experiments at higher pressures are in need to elucidate whether the Kondo effect can be tuned to overcome the antiferromagnetism, suppressing  $T_N$  and maybe inducing a quantum critical point. We note that the exchange interaction responsible for the magnetism in  $\text{Ce}_3\text{Cd}_2\text{As}_6$  is still an open question. Figure 3(c) summarizes the evolution of the transition temperatures of  $\text{Ce}_3\text{Cd}_2\text{As}_6$  as a function of pressure using both the piston cylinder cell (PCC) and the Bridgman cell (mBAC). A detailed summary of the results obtained with both cells is presented in the Supplemental Material [35].

Notably, under pressure the signature associated with magnetic order changes at  $T_N$  in the PCC measurements, as shown in Fig. 4(a). At first, a sudden decrease in the resistivity below  $T_N$  is observed due to the loss of spin disorder scattering. However, at pressures of 1.2(1) GPa and higher, this feature changes to an upturn at  $T_N$  possibly due to the formation of a superzone gap. This change is not observed in the mBAC measurements [Fig. 4(b)], in which the antiferromagnetic order leads to a sudden decrease in the resistivity for all studied pressures. We speculate that this difference is caused by

different current directions. Although the currents in both pressure cells were in the  $ab$  plane, fortunately, the current direction in the PCC allowed us to detect this change of scattering. Moreover, this difference indicates that two different magnetic phases might be present under pressure, which seems to be related with the suppression of the CDW phase in  $\text{Ce}_3\text{Cd}_2\text{As}_6$  as shown in Fig. 3(c). Another possible scenario is a change in the Fermi surface as the CDW is completely suppressed without any changes in the magnetic order. Nevertheless, the determination of the origin of this scattering type change requires more experiments with well-defined current directions and is beyond the scope of this paper.

#### IV. CONCLUSIONS

In summary, we report the electronic and magnetic properties of the newly discovered compound  $\text{Ce}_3\text{Cd}_2\text{As}_6$  [28]. At ambient pressure,  $\text{Ce}_3\text{Cd}_2\text{As}_6$  is a narrow-gap semiconductor and presents a charge density wave phase transition at 136 K. Specific heat measurements show a lambda-type peak at 4.0 K associated with the antiferromagnetic ordering of the  $\text{Ce}^{3+}$  ions in  $\text{Ce}_3\text{Cd}_2\text{As}_6$ . Furthermore, CEF fits of  $\chi(T)$ ,  $M(H)$ , and  $C_{\text{mag}}/T$  yielded a CEF scheme with the following parameters:  $B_2^0 \approx 22.8$  K,  $B_4^0 \approx -0.3$  K, and  $B_4^4 \approx -4.3$  K. These parameters result in a ground state composed of a  $\Gamma_6 = |\pm 1/2\rangle$  doublet. Moreover, two competing exchange parameters were estimated with the following values:  $z_{AFM}J_{AFM}^{\text{ex}} = 0.8$  K and  $z_{FM}J_{FM}^{\text{ex}} = -1.3$  K.

The application of external pressure on  $\text{Ce}_3\text{Cd}_2\text{As}_6$  single crystals induced a metallic ground state and strongly suppressed the charge density wave phase, which was not observed for pressures above 0.8(1) GPa. Coincident with the suppression of the CDW phase, a change of the anomaly in the resistivity associated to the antiferromagnetic order was observed with increasing pressure, indicating that  $\text{Ce}_3\text{Cd}_2\text{As}_6$  might host two different antiferromagnetic orders under pressure. The details of the ordering wave vectors of these phase will have to be studied in future. The antiferromagnetic state was favored by pressure as the system became metallic, reaching a  $T_N$  of 5.3 K at 3.8(1) GPa. Experiments at higher pressures are needed to investigate whether  $T_N$  can be suppressed maybe leading to a quantum critical point.

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- [1] E. Morosan, D. Natelson, A. H. Nevidomskyy, and Q. Si, Strongly correlated materials, *Adv. Mater.* **24**, 4896 (2012).
- [2] A. K. McMahan, Pressure-induced changes in the electronic structure of solids, *Physica B+C* **139**, 31 (1986).
- [3] M. Nicklas, Pressure probes, in *Strongly Correlated Systems*, Springer Series in Solid-State Sciences, Vol. 180 (Springer, Berlin, 2015), pp. 173–204.
- [4] P. Monceau, Electronic crystals: An experimental overview, *Adv. Phys.* **61**, 325 (2012).
- [5] G. Grüner, The dynamics of charge-density waves, *Rev. Mod. Phys.* **60**, 1129 (1988).
- [6] J. A. Wilson, F. J. Di Salvo, and S. Mahajan, Charge-density waves and superlattices in the metallic layered transition metal dichalcogenides, *Adv. Phys.* **24**, 117 (1975).
- [7] C.-W. Chen, J. Choe, and E. Morosan, Charge density waves in strongly correlated electron systems, *Rep. Prog. Phys.* **79**, 084505 (2016).
- [8] G. A. Papoian and R. Hoffmann, Hypervalent bonding in one, two, and three dimensions: Extending the Zintl-Klemm concept to nonclassical electron-rich networks, *Angew. Chem., Int. Ed.* **39**, 2408 (2000).
- [9] S. Klemenz, A. K. Hay, S. M. L. Teicher, A. Topp, J. Cano, and L. M. Schoop, The role of delocalized chemical bonding in square-net-based topological semimetals, *J. Am. Chem. Soc.* **142**, 6350 (2020).
- [10] R. H. Friend, D. Jerome, R. F. Frindt, A. J. Grant, and A. D. Yoffe, Electrical conductivity and charge density wave formation in  $4\text{H}_b\text{-TaS}_2$  under pressure, *J. Phys. C: Solid State Phys.* **10**, 1013 (1977).
- [11] A. M. Gabovich, A. I. Voitenko, and M. Ausloos, Charge- and spin-density waves in existing superconductors: Competition between Cooper pairing and Peierls or excitonic instabilities, *Phys. Rep.* **367**, 583 (2002).
- [12] L. E. Klintberg, S. K. Goh, P. L. Alireza, P. J. Saines, D. A. Tompsett, P. W. Logg, J. Yang, B. Chen, K. Yoshimura, and F. M. Grosche, Pressure- and Composition-Induced Structural Quantum Phase Transition in the Cubic Superconductor  $(\text{Sr}, \text{Ca})_3\text{Ir}_4\text{Sn}_{13}$ , *Phys. Rev. Lett.* **109**, 237008 (2012).
- [13] B. Shen, F. Du, R. Li, A. Thamizhavel, M. Smidman, Z. Y. Nie, S. S. Luo, T. Le, Z. Hossain, and H. Q. Yuan, Evolution of charge density wave order and superconductivity under pressure in  $\text{LaPt}_2\text{Si}_2$ , *Phys. Rev. B* **101**, 144501 (2020).
- [14] F. H. Yu, D. H. Ma, W. Z. Zhuo, S. Q. Liu, X. K. Wen, B. Lei, J. J. Ying, and X. H. Chen, Unusual competition of superconductivity and charge-density-wave state in a compressed topological kagome metal, *Nat. Commun.* **12**, 3645 (2021).
- [15] B. D. White, J. D. Thompson, and M. B. Maple, Unconventional superconductivity in heavy-fermion compounds, *Physica C* **514**, 246 (2015).
- [16] J. M. Tomczak, Thermoelectricity in correlated narrow-gap semiconductors, *J. Phys.: Condens. Matter* **30**, 183001 (2018).
- [17] T. Takabatake, T. Sasakawa, J. Kitagawa, T. Suemitsu, Y. Echizen, K. Umeo, M. Sera, and Y. Bando, Thermoelectric properties of Ce-based Kondo semimetals and semiconductors, *Physica B* **328**, 53 (2003).
- [18] S. M. Thomas, P. F. S. Rosa, S. B. Lee, S. A. Parameswaran, Z. Fisk, and J. Xia, Hall effect anomaly and low-temperature metamagnetism in the Kondo compound  $\text{CeAgBi}_2$ , *Phys. Rev. B* **93**, 075149 (2016).
- [19] R. Prozorov, S. L. Bud'ko, and P. C. Canfield, Topological magnetic hysteresis in single crystals of  $\text{CeAgSb}_2$  ferromagnet, *J. Phys.: Condens. Matter* **34**, 145802 (2022).
- [20] V. A. Sidorov, E. D. Bauer, N. A. Frederick, J. R. Jeffries, S. Nakatsuji, N. O. Moreno, J. D. Thompson, M. B. Maple, and Z. Fisk, Magnetic phase diagram of the ferromagnetic Kondo-lattice compound  $\text{CeAgSb}_2$  up to 80 kbar, *Phys. Rev. B* **67**, 224419 (2003).
- [21] L. Zhao, E. A. Yelland, J. A. N. Bruin, I. Sheikin, P. C. Canfield, V. Fritsch, H. Sakai, A. P. Mackenzie, and C. W. Hicks, Field-temperature phase diagram and entropy landscape of  $\text{CeAuSb}_2$ , *Phys. Rev. B* **93**, 195124 (2016).
- [22] J. Park, H. Sakai, A. P. Mackenzie, and C. W. Hicks, Effect of uniaxial stress on the magnetic phases of  $\text{CeAuSb}_2$ , *Phys. Rev. B* **98**, 024426 (2018).
- [23] G. G. Marcus, D.-J. Kim, J. A. Tutmaher, J. A. Rodriguez-Rivera, J. O. Birk, C. Niedermeyer, H. Lee, Z. Fisk, and C. L. Broholm, Multi- $q$  Mesoscale Magnetism in  $\text{CeAuSb}_2$ , *Phys. Rev. Lett.* **120**, 097201 (2018).
- [24] S. Seo, V. A. Sidorov, H. Lee, D. Jang, Z. Fisk, J. D. Thompson, and T. Park, Pressure effects on the heavy-fermion antiferromagnet  $\text{CeAuSb}_2$ , *Phys. Rev. B* **85**, 205145 (2012).
- [25] S. Seo, Xiaoyu Wang, S. M. Thomas, M. C. Rahn, D. Carmo, F. Ronning, E. D. Bauer, R. D. dos Reis, M. Janoschek, J. D. Thompson, R. M. Fernandes, and P. F. S. Rosa, Nematic State in  $\text{CeAuSb}_2$ , *Phys. Rev. X* **10**, 011035 (2020).
- [26] M. M. Piva, R. Tartaglia, G. S. Freitas, J. C. Souza, D. S. Christovam, S. M. Thomas, J. B. Leão, W. Ratcliff, J. W. Lynn, C. Lane, J.-X. Zhu, J. D. Thompson, P. F. S. Rosa, C. Adriano, E. Granado, and P. G. Pagliuso, Electronic and magnetic properties of stoichiometric  $\text{CeAuBi}_2$ , *Phys. Rev. B* **101**, 214431 (2020).
- [27] P. F. S. Rosa, R. J. Bourg, C. B. R. Jesus, P. G. Pagliuso, and Z. Fisk, Role of dimensionality in the Kondo  $\text{CeTX}_2$  family: The case of  $\text{CeCd}_{0.7}\text{Sb}_2$ , *Phys. Rev. B* **92**, 134421 (2015).
- [28] M. M. Piva, M. C. Rahn, S. M. Thomas, B. L. Scott, P. G. Pagliuso, J. D. Thompson, L. M. Schoop, F. Ronning, and P. F. S. Rosa, Robust narrow-gap semiconducting behavior in square-net  $\text{La}_3\text{Cd}_2\text{As}_6$ , *Chem. Mater.* **33**, 4122 (2021).
- [29] E. Colombier and D. Braithwaite, Simple adaptation of the Bridgman high pressure technique for use with liquid media, *Rev. Sci. Instrum.* **78**, 093903 (2007).

- [30] P. G. Pagliuso, D. J. Garcia, and E. Miranda, Evolution of the magnetic properties and magnetic structures along the  $R_m M \text{In}_{3m+2}$  ( $R = \text{Ce, Nd, Gd, Tb}$ ;  $M = \text{Rh, Ir}$ ; and  $m = 1, 2$ ) series of intermetallic compounds, *J. Appl. Phys.* **99**, 08P703 (2006).
- [31] X. Zhang, J. Li, B. Foran, S. Lee, H.-Y. Guo, T. Hogan, C. R. Kannewurf, and M. G. Kanatzidis, Distorted square nets of tellurium in the novel quaternary polytelluride  $\text{K}_{0.33}\text{Ba}_{0.67}\text{AgTe}_2$ , *J. Am. Chem. Soc.* **117**, 10513 (1995).
- [32] K. Seeger, *Semiconductor Physics: An Introduction*, 9th ed. (Springer, Berlin, 2004).
- [33] R. A. Craven and S. F. Meyer, Specific heat and resistivity near the charge-density-wave phase transitions in  $2\text{H-TaSe}_2$  and  $2\text{H-TaS}_2$ , *Phys. Rev. B* **16**, 4583 (1977).
- [34] M. E. Fisher and J. S. Langer, Resistive Anomalies at Magnetic Critical Points, *Phys. Rev. Lett.* **20**, 665 (1968).
- [35] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.105.094443> for a detailed summary of the electrical resistivity data obtained with both piston cylinder and Bridgman pressure cells.
- [36] T. Sasakawa, T. Suemitsu, T. Takabatake, Y. Bando, K. Umeo, M. H. Jung, M. Sera, T. Suzuki, T. Fujita, M. Nakajima, K. Iwasa, M. Kohgi, Ch. Paul, St. Berger, and E. Bauer, Successive phase transitions and energy-gap formation in  $\text{CeRhAs}$ , *Phys. Rev. B* **66**, 041103(R) (2002).
- [37] S. Yoshii, M. Kasaya, H. Takahashi, and N. Mori, Effect of pressure on the size of energy gap in semiconducting mixed-valent rare-earth compounds, *Physica B* **223-224**, 421 (1996).
- [38] K. Umeo, K. Masumori, T. Sasakawa, F. Iga, T. Takabatake, Y. Ohishi, and T. Adachi, Pressure effects on the phase transitions and energy gap in  $\text{CeRhAs}$ , *Phys. Rev. B* **71**, 064110 (2005).
- [39] L. Mendonça-Ferreira, F. B. Carneiro, M. B. Fontes, E. Baggio-Saitovitch, L. S. I. Veiga, J. R. L. Mardegan, J. Stremper, M. M. Piva, P. G. Pagliuso, R. D. dos Reis, and E. M. Bittar, Pressure effects on the structural and superconducting transitions in  $\text{La}_3\text{Co}_4\text{Sn}_{13}$ , *J. Alloys Compd.* **773**, 34 (2019).
- [40] B. Sipos, A. F. Kusmartseva, A. Akrap, H. Berger, L. Forró, and E. Tutiš, From Mott state to superconductivity in  $1\text{T-TaS}_2$ , *Nat. Mater.* **7**, 960 (2008).
- [41] T. Kiss, T. Yokoya, A. Chainani, S. Shin, T. Hanaguri, M. Nohara, and H. Takagi, Charge-order-maximized momentum-dependent superconductivity, *Nat. Phys.* **3**, 720 (2007).
- [42] A. F. Kusmartseva, B. Sipos, H. Berger, L. Forró, and E. Tutiš, Pressure Induced Superconductivity in Pristine  $1\text{T-TiSe}_2$ , *Phys. Rev. Lett.* **103**, 236401 (2009).
- [43] A. H. Castro Neto, Charge Density Wave, Superconductivity, and Anomalous Metallic Behavior in 2D Transition Metal Dichalcogenides, *Phys. Rev. Lett.* **86**, 4382 (2001).
- [44] J. C. Cooley, M. C. Aronson, and P. C. Canfield, High pressures and the Kondo gap in  $\text{Ce}_3\text{Bi}_4\text{Pt}_3$ , *Phys. Rev. B* **55**, 7533 (1997).
- [45] S. Doniach, in *Valence Instabilities and Related Narrow Band Phenomena*, edited by R. D. Parks (Plenum, New York, 1977), p. 169.