

Collinear antiferromagnetic order in $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ revealed by neutron diffraction

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The hidden order phase in URu_2Si_2 is highly sensitive to electronic doping. A special interest in silicon-to-phosphorus substitution is due to the fact that it may allow one, in part, to isolate the effects of tuning the chemical potential from the complexity of the correlated f and d electronic states. We investigate the new antiferromagnetic phase that is induced in $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ at $x \gtrsim 0.27$. Time-of-flight neutron diffraction of a single crystal ($x = 0.28$) reveals c -axis collinear $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ magnetic order with localized magnetic moments (~ 2.1 – $2.6 \mu_B$). This points to an unexpected analogy between the (Si,P) and (Ru,Rh) substitution series. Through further comparisons with other tuning studies of URu_2Si_2 , we are able to delineate the mechanisms by which silicon-to-phosphorus substitution affects the system. In particular, both the localization of itinerant $5f$ electrons as well as the choice of \mathbf{q}_m appear to be consequences of the increase in chemical potential. Further, enhanced exchange interactions are induced by chemical pressure and lead to magnetic order, in which an increase in interlayer spacing may play a special role.

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I. INTRODUCTION

The complexity of energetically nearly degenerate electronic states in strongly correlated materials often gives rise to unusual ordering phenomena and exotic physics [1,2]. However, it is challenging to identify the hierarchy of the underlying interactions between charge, orbital, magnetic, and structural degrees of freedom. Such is the case for the hidden order (HO) state that emerges in the heavy-fermion material URu_2Si_2 . Studies using a large set of external parameters to tune this system have revealed a rich phase space of adjacent ordered phases [3–5], many of which are magnetic. Notably, the characteristics and symmetry of the HO state itself are known to be markedly different from conventional spin or charge orders in strongly correlated metals [6]. Even though the onset of HO at $T_0 = 17.5$ K is marked by a second-order symmetry breaking phase transition [7,8], the true symmetry of the associated order parameter remains elusive.

In URu_2Si_2 , strong electronic correlations arise due to the hybridization of localized uranium f electrons with the conduction electrons, as evident from a large single-ion Kondo temperature $T_K = 120$ K [9,10], and the formation of a coherent Kondo lattice at $T^* \approx 70$ K [11–13]. The onset of HO is accompanied by the opening of a charge gap over

about 40% of the Fermi surface (FS), as observed via various methods sensitive to changes in the band structure [7,8,13–18]. This reorganization of the electronic structure below T_0 was originally attributed to the emergence of charge or spin density wave order, which would be typical of itinerant magnetism. More detailed investigations using modern electronic band structure methods revealed a secondary hybridization of a heavy f -like quasiparticle band with a light holelike band at $Q^* = \pm 0.3\pi/a$, which results in the formation of a hybridization gap $\Delta_{Q^*} = 5$ meV [9,10,19]. Both angle-resolved photoemission spectroscopy (ARPES) [20] and neutron spectroscopy [21,22] indicate that even larger parts of the Fermi surface are gapped. Nuclear magnetic resonance (NMR) measurements [23] also indicate the presence of a pseudogap below 30 K.

This dramatic reorganization of the Fermi surface suggests that the HO order is a result of its intricate electronic band structure. However, the complex metallic state of URu_2Si_2 also exhibits strong anisotropy in the spin and charge channels, which is typically associated with localized electronic degrees of freedom [24]. Torque magnetization [25], high-resolution x-ray diffraction [26], elasto-resistance [27], NMR [28], and Raman spectroscopy [6] indicate that the electronic state breaks the tetragonal symmetry of the underlying crystal structure, which led to the proposal that the HO state may be of nematic origin. Even though other x-ray diffraction [29], NMR [30], and thermodynamic [31] studies at ambient pressure have not corroborated this tetragonal

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symmetry breaking, a recent x-ray diffraction study revealed a tetragonal-to-orthorhombic phase transition as a function of pressure [32]. Furthermore, ultrasound measurements have observed an orthorhombic lattice instability due to a volume-conserving strain field with Γ_3 symmetry [33]. Taken together, this suggests that the difference between the various studies may be due to varying crystal quality, resulting in different amounts of internal strain.

The above underscores that unraveling the conundrum of HO requires a better understanding of the underlying duality of the itinerant and localized degrees of freedom. To this end, it is helpful to examine the impressive collection of studies in which the HO state has been tuned by various control parameters. This includes the external parameters of high magnetic fields [34,35], pressure and strain [36,37], as well as chemical substitution, on the uranium [38,39] or ruthenium site [40–51].

In all of these cases, it is observed that HO exists in close proximity to magnetic phases, which, in many cases, resemble those found in other tetragonal members of the UT_2Si_2 (T : transition metal) family [52–56]. While this provides some insight, it is often unclear how to disentangle the effects of varying the hybridization (e.g., by varying the degree of delocalization and spin-orbit coupling of the ligand), of the local environment (by variation of bond lengths and angles), and the variation of the Fermi surface. This is exemplified by a recent magnetoresistivity study under combined high pressures and magnetic fields that finds that the effects of both tuning parameters are intertwined [57]. To address this issue, the (Si,P) substitution series has been established [58–62], which is thought to weaken p - f hybridization [63], but affects the spacing and orientation of d -ligands only weakly. With the character of f - d interaction held intact, the donation of one p electron is thought to emphasize the effects of varying the chemical potential in the URu_2Si_2 host. Even though the various consequences of any chemical substitution are necessarily intertwined, the special significance of the (Si,P) series is that it provides a different hierarchy in which the numerous relevant energy scales of URu_2Si_2 are affected. One may therefore hope that this will provide the necessary contrast to disentangle the mechanisms by which the HO state is manipulated in other tuning studies of this material.

An overview of the effects of (Si,P) substitution is given by the schematic phase diagram in Fig. 1. Interestingly, superconductivity (SC) and HO prove to be highly sensitive to very small P doping levels. In particular, the superconducting critical temperature $T_c \approx 1.4$ K weakly increases to a maximum at $x \approx 0.01$, before suppression of SC at $x \approx 0.028$ and suppression of HO at $x \approx 0.035$ [60]. Quantum oscillation measurements in this regime indicate that no significant changes of the Fermi surface are associated with the destruction of the HO phase [64]. Following a paramagnetic Kondo lattice state in the range $0.035 \lesssim x \lesssim 0.26$, antiferromagnetism is abruptly stabilized at $x \gtrsim 0.27$ [59]. To better understand the different roles of the large number of available tuning parameters of HO in URu_2Si_2 , it is of great interest to characterize the order in this new magnetic phase [5]. Here we report on a neutron diffraction study carried out to determine its antiferromagnetic order.

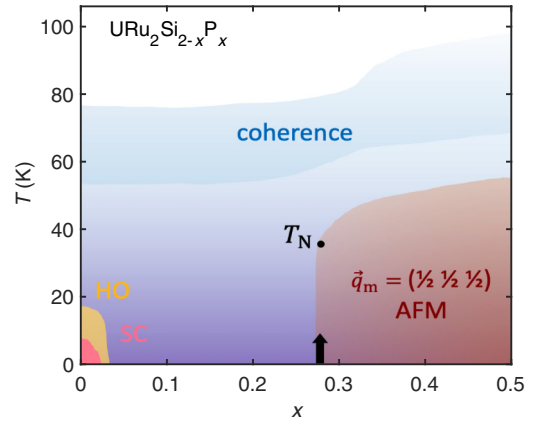


FIG. 1. Simplified phase diagram of $URu_2Si_{2-x}P_x$, adapted from Gallagher *et al.* [59]. Few percents of phosphorus substitution suppress the hidden order (HO) and, with it, superconductivity (SC). The arrow marks the composition of the long-range antiferromagnetically (AFM) ordered sample investigated in this study.

II. EXPERIMENTAL DETAILS

Single crystals of $URu_2Si_{2-x}P_x$ were synthesized by an indium flux method [58,60]. Aside from the high purity of the resulting crystals, this method also overcomes the issue of the high vapor pressure of phosphorus. The phosphorus concentration $x = 0.28(1)$ of this sample (i.e., $\approx 14\%$ substitution) was determined by energy-dispersive x-ray spectroscopy (EDX). The uncertainty of this value was estimated by performing a number of measurements on different positions of the sample surface. For reference, this composition is marked by a black arrow in the phase diagram in Fig. 1. The magnetic susceptibility $\chi(T)$ of this crystal was measured using a superconducting quantum interference device (SQUID) magnetometer (Quantum Design), in a field of 0.5 T applied either parallel or perpendicular to the c axis. The dimensions of single crystals grown by the molten metal flux technique make magnetic neutron diffraction barely feasible. We selected a crystallite for its large size compared to the average sample yield, with a mass of only ≈ 0.5 mg and dimensions of $0.8 \times 0.8 \times 0.05$ mm³.

The issue of small sample size can be overcome using the latest generation time-of-flight neutron diffractometers, which combine high-brilliance neutron moderators with highly optimized focusing neutron guides. The high flux yield at the sample position enables experiments on single crystals with dimensions of ~ 1 mm and less. In combination with detector banks that collect scattered neutrons over a large solid angle, this allows for experiments that were impossible until recently. Preliminary measurements of the nuclear scattering down to 90 K were carried out at the TOPAZ instrument at SNS (Oak Ridge National Laboratory), which receives neutrons from a decoupled poisoned hydrogen moderator. The investigation of the magnetic order at 2 K was performed at the WISH instrument at the ISIS pulsed neutron source (STFC, Rutherford Appleton Laboratory) [65]. WISH looks onto a solid methane (40 K) moderator, which provides high-brilliance neutron pulses with a broad band of wavelengths from 1 to 10 Å. Neutrons are collected on a detector

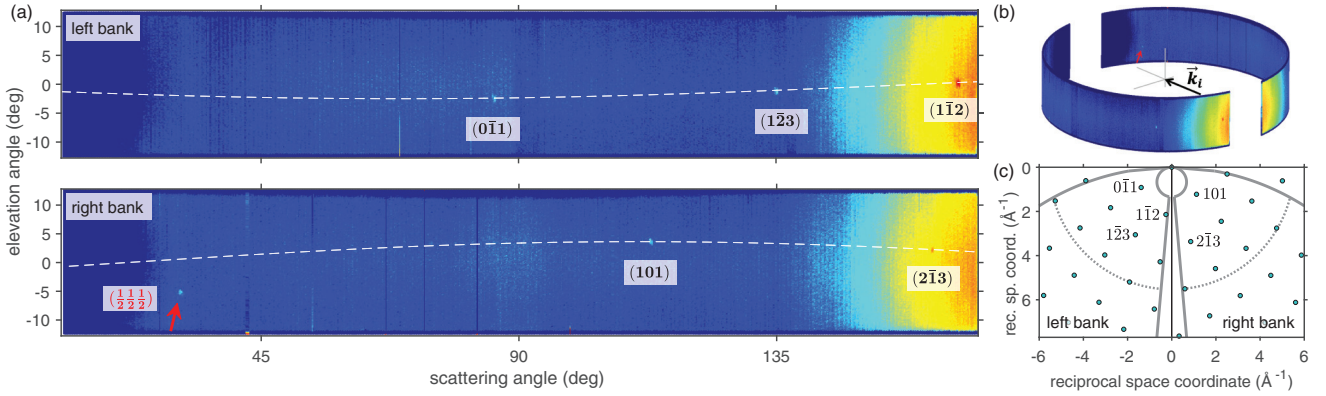


FIG. 2. (a) Overview of Bragg peaks observed at WISH, at 2 K. The upper (lower) panel shows the raw neutron counts detected on the left (right) detector bank, on an arbitrary logarithmic scale. The (101)–(011) plane of reciprocal space is indicated by a white dashed line. The (1,1,1) direction, which features the only observed magnetic Bragg peak, is seen below this plane, at a scattering angle of around 30° . (b) Perspective view of this data, illustrating the layout of the instrument. (c) Schematic view of the (101)–(011) plane of reciprocal space (Note that the magnetic peak at \mathbf{q}_m is observed below this plane). The accessible range of momentum transfers is delineated by a broad gray line and peaks seen in (a) are labeled in analogy. The dotted line indicates the momentum transfer at which the magnetic form factor of uranium has decreased by $1/e$ ($\approx 5.5 \text{ \AA}^{-1}$).

bank that continuously covers a wide range of scattering angles ($10^\circ \leq 2\theta \leq 170^\circ$) with 1-m-tall position-sensitive ^3He detectors. WISH also employs an oscillating radial collimator that defines a cylindrical collimated area in the center of the sample tank, which provides the low background required for studies with such small samples. The crystal was mounted in a dedicated low background cryostat (Oxford Instruments) with the [110] and [101] directions in the horizontal plane. The broad margin of accessible out-of-plane momentum transfers covered a volume equivalent to (more than) one full Brillouin zone, which was crucial for the identification of the magnetic propagation vector. An illustration of the accessible range of reciprocal space resulting from this configuration is given in Fig. 2.

III. RESULTS

Figure 3 shows the temperature dependence of the magnetic susceptibility $\chi(T)$ in a magnetic field of 0.5 T, applied

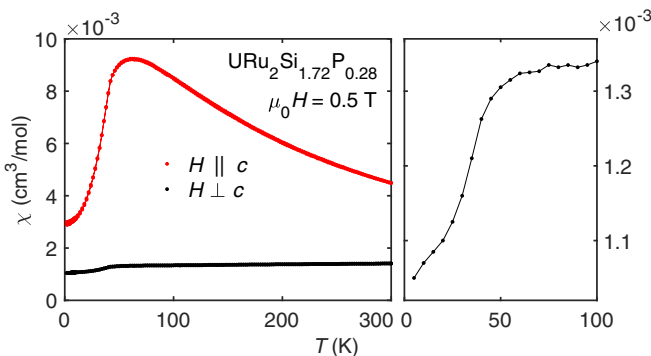


FIG. 3. Magnetic susceptibility of the $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ ($x = 0.28$) crystal investigated by neutron diffraction, in a field of 0.5 T applied either parallel or perpendicular to the c axis. The right panel shows a detailed view of the $\mathbf{H} \perp \mathbf{c}$ data. As in the parent compound, the characteristics are dominated by the onset of Kondo screening around $T_{\text{coh}} \approx 80 \text{ K}$, as well as a strong c -axis single-ion anisotropy.

parallel and perpendicular to the c axis. The characteristics are dominated by the strong single-ion magnetic anisotropy, which indicates the same c -axis Ising character known of the parent compound [7]. The broad maximum around $T_{\text{coh}} \approx 80 \text{ K}$ marks the onset of Kondo screening of the magnetic moments. The magnetic phase transition is associated with a marked decrease in $\vec{H} \perp \vec{c}$ susceptibility. These observations are also consistent with recent NMR measurements, which revealed a commensurate internal field $H_{\text{int}} \approx 0.85 \text{ kOe}$ oriented along the c direction in the antiferromagnetic state of $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ [61].

With the sample orientation and time-of-flight range of the WISH experiment illustrated in Fig 2, the accessible nuclear Bragg peaks were indexed in the $I4/mmm$ unit cell of the parent compound (lattice parameters $a = 4.12 \text{ \AA}$ and $c = 9.57 \text{ \AA}$). The scale factor of the nuclear intensities was refined with the Rietveld method using FULLPROF [66] after a single-crystal Lorentz correction had been performed in MANTID [67]. Variables of this fit included the vertical position of Si/P ions [$z = 0.38(1)$] and a parameter controlling the extinction correction [66]. The (Si,P) stoichiometry was fixed to the value determined by EDX [$x = 0.28(1)$]. A comparison of measured and calculated intensities is shown in Fig. 4. Numerical values and a detailed account of this fit is provided in the Supplemental Material [68].

The intensity of these reflections was tracked between 2 and 80 K. At low temperatures, our measurements reveal a magnetic Bragg peak at momentum transfer $\mathbf{Q} = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$, which corresponds to the magnetic propagation vector $\mathbf{q}_m = \mathbf{Q}$. The widths of this peak in reciprocal space were the same as for nuclear peaks, which indicates that the range of the order is not limited by the coherence length of magnetic correlations. Other instances of scattering from the \mathbf{q}_m vector were also identified in higher-order Brillouin zones, although these intensities are increasingly suppressed by the magnetic form factor.

Figure 5 shows the temperature dependence of the integrated intensity of this peak, on a scale of the estimated

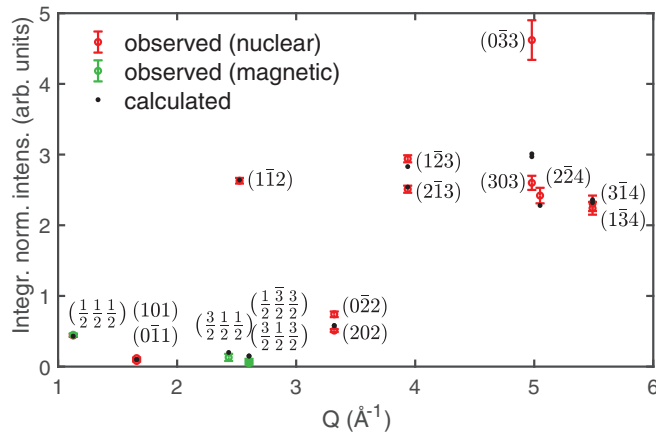


FIG. 4. Comparison of Bragg intensities calculated for $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ ($x = 0.28$) with those measured at WISH at 2 K. The overall scale factor was inferred from a refinement of nuclear intensities (red). The magnitude of the ordered magnetic moment was then fitted separately to reproduce the intensities of magnetic reflections (green). Numerical values of these fits are given in the Supplemental Material [68].

ordered magnetic moment M per uranium ion (see below). A fit of the temperature dependence of the ordered moment via $M(T) \propto (1 - \frac{T}{T_N})^\beta$ yields a Néel temperature of $T_N = 32.5(1.3)$ K and a critical exponent of $\beta = 0.24(6)$. The large uncertainty notwithstanding, this is in line with the three-dimensional (3D) Ising character ($\beta_{\text{th}} = 0.32$) [69] and the strong magnetic anisotropy evident from Fig. 3.

IV. ANALYSIS AND INTERPRETATION

The present study represents a favorable case in which the magnetic structure is fully determined by the observation of a single magnetic peak, given the constraints inferred

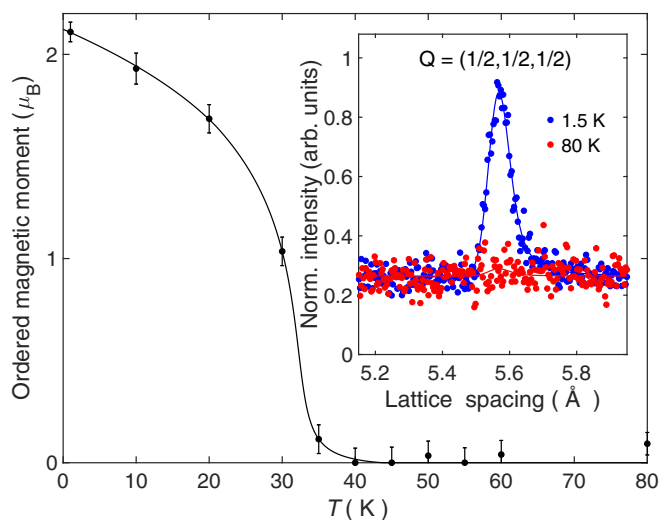


FIG. 5. Temperature dependence of the ordered magnetic moment in $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ ($x = 0.28$), with an ordering temperature of $T_N = 32.6(7)$ K and critical exponent $\beta = 0.31(4)$. The inset illustrates the emergence of a magnetic Bragg reflection at momentum transfer $\mathbf{Q} = \mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$.

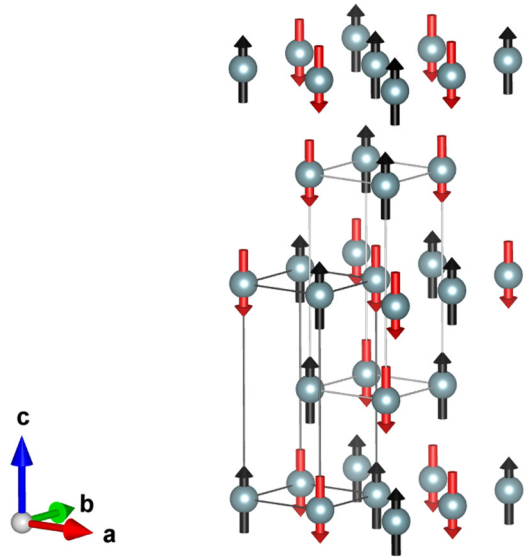


FIG. 6. c -axis collinear $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ magnetic structure of $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ ($x = 0.28$), described by the Shubnikov group I_{c4_1}/acd (No. 142.570). For clarity, only uranium ions are shown and opposite spins are drawn in different colors.

from magnetometry and symmetry. Representational analysis was performed using the ISODISTORT program [70,71]. The propagation vector $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ represents the P point (“k12”) in the Brillouin zone of space group $I4/mmm$. For magnetic moments at the uranium site, this yields two magnetic irreducible representations (irreps), $\text{mk}12\text{t}2$ and $\text{mk}12\text{t}5$ (Kovalev notation), with magnetic moments along the c axis and in plane, respectively.

Each irrep provides a choice of order parameter directions, corresponding to magnetic space groups. The resulting magnetic structures are illustrated in Fig. S1 of the Supplemental Material [68]. Since magnetometry clearly indicates c -axis Ising anisotropy (cf. Fig. 3), any magnetic space group corresponding to $\text{mk}12\text{t}5$ (for which all moments are confined to the a - b plane) can be ruled out. Of the three choices for $\text{mk}12\text{t}2$ ($P1$, $P3$, and $C1$), $P3$ is not physical as it forces one U site to be nonmagnetic, and $C1$ represents the unlikely case that the size of the ordered magnetic moment varies between the two U ions (of the same Wyckoff site, i.e., in an environment of the same point symmetry). There is then indeed only one possible solution, $\text{mk}12\text{t}2$ $P1$, which corresponds to the magnetic space group I_{c4_1}/acd (No. 142.570).

The resulting magnetic structure is illustrated in Fig. 6. The basis and origin of the I_{c4_1}/acd magnetic cell, specified in terms of lattice vectors of the parent paramagnetic space group $I4/mmm$, are $(\bar{1}, 1, 0)$, $(\bar{1}, \bar{1}, 0)$, $(0, 0, 2)$, and $(\frac{1}{4}, \frac{1}{4}, \frac{7}{4})$, respectively. It is the $\bar{4}m'2'$ symmetry of the Wyckoff site $8a$ which constrains the magnetic dipole moments to align with the c axis.

By comparing the integrated magnetic and nuclear intensities, we obtain an ordered magnetic moment of $\mu \sim 2.1$ – $2.6 \mu_B$. This range reflects the systematic uncertainty of the intensity scale factor (cf. Fig. 4 and Ref. [68]). By comparison, the uncertainty due to the choice of the neutron magnetic form factor [72] is not significant (U^{3+}

vs U^{4+} , $d\mu \approx 0.04 \mu_B$). This result is consistent with local-moment-like antiferromagnetic phases in related UA_2B_2 compounds [73,74].

V. DISCUSSION

Our neutron diffraction study demonstrates that the magnetically ordered state observed in $URu_2Si_{2-x}P_x$ for phosphorus concentrations $x \gtrsim 0.27$ is described by a c -axis collinear antiferromagnetic structure with propagation vector $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ and an ordered moment of ~ 2.1 – $2.6 \mu_B$. The commensurate character of this state suggests that it arises from local uranium magnetic moments. This is corroborated by a recent NMR study of an $x = 0.33$ single crystal that evidenced a homogeneous antiferromagnetic state emerging due to the localization of $5f$ electrons at higher P concentrations [61], as well as measurements of the Sommerfeld coefficient of the specific heat, which is reduced sharply once magnetic order emerges for $x \gtrsim 0.27$ [59]. In this section, we highlight the most relevant similarities and differences between our findings and other tuning studies of URu_2Si_2 [5].

Most importantly, the magnetic state discovered in the (Si,P) system reveals an unexpected parallel to the phase diagram of $URu_{2-x}Rh_xSi_2$ [45,46]. In this series, the HO vanishes for $x \gtrsim 0.08$. As in the (Si,P) system, this is followed by a paramagnetic heavy Fermi liquid region, for $0.08 \lesssim x \lesssim 0.18$ [45,46,48]. Finally, for $0.18 \lesssim x \lesssim 0.7$, a similar long-range $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ (or $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, L)$ [46]) antiferromagnetic state appears, with $\mu \approx 2 \mu_B$ and T_N up to 44 K, as in the present case.

An equivalence of (Si,P) and (Ru,Rh) substitution is far from obvious. For example, one may expect that the main effect of substituting U or Ru (e.g., by Np, Fe, Os, Ru, Rh) may be to alter the d - f hybridization. It was only pointed out very recently that the case of (Si,P) may actually have similar consequences, given that the radial contraction of the p -orbitals weakens the p - f hybridization [63].

On the other hand, both the (Si,P) and (Ru,Rh) series are markedly set apart from isoelectronic chemical substitutions (i.e., by Fe, Os, and Ge, respectively). The latter have a stronger impact on bond lengths and angles, as well as on spin-orbit coupling (SOC) [75]. By comparison, structural modifications in the (Si,P) series are more gentle and the increase in SOC is negligible [59]. On the other hand, a considerable consequence of adding p electrons to the system must be the rise in chemical potential.

Taken together, this suggests that (Si,P) and (Ru,Rh) substitution affect the HO state of URu_2Si_2 in two ways: by moderate chemical pressure and electron donation. It is then interesting to trace the roles of these two effects in stabilizing $\vec{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ large moment magnetic order.

On the one hand, the increase of the chemical potential binds the f -electron states well below the Fermi energy, favoring their localization. On the other hand, at higher x , the decrease in unit cell volume becomes more relevant, as it increases the exchange integrals between neighboring $5f$ orbitals, promoting long-range magnetic order. This increase of correlations at higher substitution levels also enhances the coherence temperature, as observed in electrical transport and

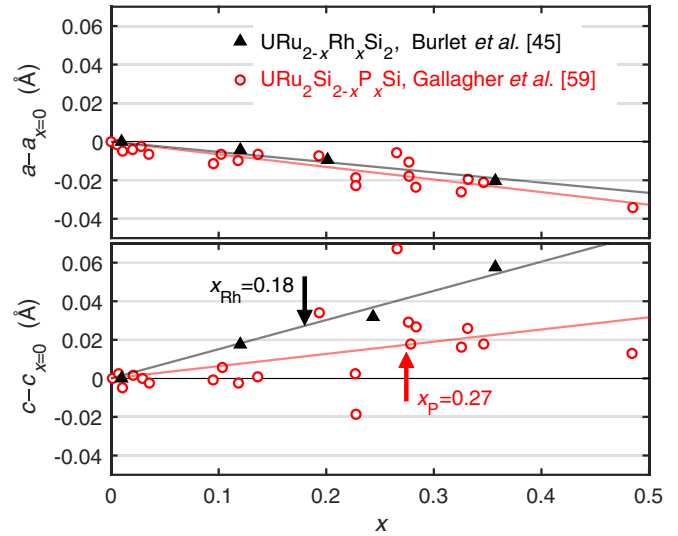


FIG. 7. Changes of the lattice parameters a (top) and c (bottom) in the (Ru,Rh) and (Si,P) substitution series. The data are adapted from studies by Burlet *et al.* [45] and Gallagher *et al.* [59]. In the lower panel, the critical compositions at which local moment $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ magnetic order emerges in either series are marked by arrows.

magnetic susceptibility measurements (discussed below in the context of Fig. 8) [59].

The amount of chemical pressure exerted by (Si,P) and (Ru,Rh) substitution is similar. In Fig. 7, we compare the variation of the lattice parameters in the two doping series, based on data reported by Burlet *et al.* [45] and Gallagher *et al.* [59]. The contraction of the basal plane (parameter a , corresponding to the nearest-neighbor spacing of uranium ions) is, in fact, identical within the uncertainties of the measurements. On the other hand, the c -axis expansion (which likely acts as a handle on interlayer correlations) due to (Ru,Rh) substitution

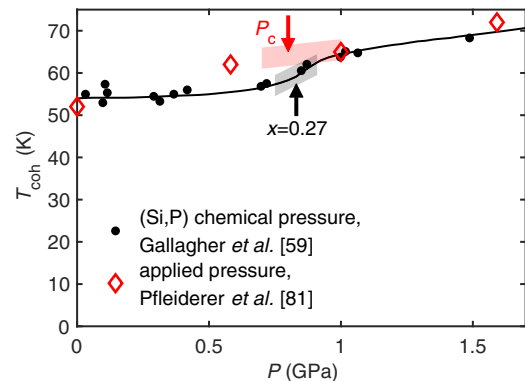


FIG. 8. Pressure dependence of the heavy-fermion coherence scale T_{coh} , inferred from a broad maximum in magnetic susceptibility $\chi(T)$ curves (cf. Fig. 3). Data for chemical pressure in the present doping series, adapted from Ref. [59], is compared to measurements of the parent compound under applied pressure, reported by Pfeleiderer *et al.* [81]. The arrows and shaded margins indicate the regimes where long-range magnetic order is induced by hydrostatic (red) and chemical (black) pressure.

is significantly larger than in the (Si,P) series. Notably, the values of c at which long-range order sets in are similar in both compounds (cf. arrows in Fig. 7). Even though the number of outliers in the $c - c_{x=0}$ data by Gallagher *et al.* calls for caution, this makes for a tentative explanation for the difference in the corresponding critical substitution levels.

In this context, studies of URu₂Si₂ under applied hydrostatic pressure are of special interest because they allow one to single out the effects attributable to structural variations. Applied pressure drives a quantum phase transition from the HO state to an antiferromagnetic phase with propagation vector $\mathbf{q}_m = (0, 0, 1)$ at a critical pressure $P_c = 0.7\text{--}1.0$ GPa [76] (see, also, Ref. [77] and references therein). At ambient pressure, neutron scattering originally observed a similar magnetic phase, however, with a much reduced magnetic moment of about $0.01 \mu_B$ [78]. One important conclusion of the extensive efforts to investigate URu₂Si₂ under pressure is that this small moment antiferromagnetism (SMAF) at ambient pressure is likely a parasitic effect, induced locally by remnant strain [79]. The closing of a spin gap at \mathbf{q}_m , as observed in neutron scattering [80], emphasizes that the HO state is fundamentally different from the large moment pressure-induced antiferromagnetic phase. It is also consistent with the view that pressure enhances exchange interactions.

Chemical and applied pressure can be compared quantitatively using the Birch-Murnaghan equation of state [$P_{\text{chem}} \approx 9 \text{ MPa} (\Delta V/V)$], as reported by Gallagher *et al.* [59]. In Fig. 8, we use this relation to compare the variation of T_{coh} in the (Si,P) series (inferred from the broad maximum in $\chi(T)$, data adapted from Ref. [59]) to the corresponding results of a high-pressure URu₂Si₂ study by Pfeleiderer *et al.* [81] (the same effect is also observed in the resistivity measurements [82]). The two means of compressing the lattice indeed increase the coherence temperature by similar amounts. This suggests that even though the resulting antiferromagnetic structures differ, similar physics is at play in stabilizing the local moment order parameter. This is also supported by the observation that for the P concentration at which the antiferromagnetic phase sets in, the lattice contraction corresponds to a chemical pressure of ≈ 0.75 GPa [59], which is comparable to P_c [76,77].

To recapitulate, (Si,P) and (Ru,Rh) substitution affect the lattice in a similar way, and both induce a $(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ large moment magnetic order. The equivalent amount of applied pressure similarly induces a large moment antiferromagnetic state in URu₂Si₂, albeit with propagation vector (0,0,1). The fact that the donation of one electron is the common difference between these substitution series and applied pressure suggests that the variation of the chemical potential has the role of selecting the different magnetic symmetry.

In the Supplemental Material [68], we also provide a discussion of the effect of applied magnetic fields (with reference to the studies [83–88]). This is another tuning parameter that may serve to reveal similarities and differences between different ground states of URu₂Si₂ derivatives. Here, we merely note that as in URu₂Si_{2-x}P_x, an antiferromagnetic state can be induced in the paramagnetic regime of URu_{2-x}Rh_xSi₂ by applying a critical field of 26 T. A pulsed-field neutron diffraction study has shown that this magnetic order corresponds to a commensurate up-up-down ferrimagnetic structure with propagation vector $\mathbf{q}_{\text{IN}} = (\frac{1}{3}, 0, 0)$

[47,49]. Since the high-field magnetization of P and Rh substituted samples in the paramagnetic regime have very similar characteristics [62], and the magnetic order at higher P and Rh concentrations is identical, we speculate that the field-induced magnetic phase in URu₂Si_{2-x}P_x is also the one described by \mathbf{q}_{IN} .

In summary, it is not straightforward to identify the mechanism of the large moment antiferromagnetic state discovered in URu₂Si_{2-x}P_x discovered at $x \gtrsim 0.27$, since the consequences of ionic substitution are necessarily entangled. However, such a large catalog of tuning studies is now available [5] that it becomes possible to recognize the key effects by comparison. In the present case, we thus arrive at a simplified picture in which the rise in chemical potential forms local magnetic moments (associated with the destruction of HO), before correlations are increasingly enhanced by chemical pressure, eventually stabilizing long-range magnetic order. The chemical potential is likely decisive in selecting a propagation vector that is distinct from the large moment phase of pure URu₂Si₂ under pressure.

VI. CONCLUSION

Using state-of-the-art neutron time-of-flight diffractometers, we were able to determine the magnetic structure in a minute single crystal of URu₂Si_{2-x}P_x ($x > 0.27$). Our measurements indicate c -axis collinear antiferromagnetic order with a propagation vector $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ and an ordered moment of $\sim 2.1\text{--}2.6 \mu_B$. This highlights that the phase diagrams of the two substitution series URu₂Si_{2-x}P_x [59–62] and URu_{2-x}Rh_xSi₂ [45–49] are nearly identical with respect to the observed sequence of ground states. By comparison of various tuning studies, we infer that the localization of $5f$ electrons as well as the selection of the $\mathbf{q}_m = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ order parameter is a common consequence of the increased chemical potential, whereas enhanced exchange interactions are induced by chemical pressure, in which the increase in interlayer spacing may play a special role.

In spite of these parallels, we note that the (Si,P) and (Ru,Rh) substitutions must act differently in terms of the spin-orbit coupling, lattice strain, and local crystal electric field, and may alter different aspects of the Fermi surface. Detailed investigations of the underlying electronic structures and local degrees of freedom would be in place to shed more light on these differences.

It is also important to keep in mind that the observed behavior near quantum phase transitions between different correlated phases is known to be highly sensitive to the ionic disorder introduced by chemical substitution. It has been shown to impact both electrical transport properties [89] as well as the nature of the quantum phase transition itself [90–92]. In turn, the impact of disorder remains another important open question when comparing these two substitution series.

Neutron diffraction or NMR measurements of URu₂Si_{2-x}P_x with $x \gtrsim 0.26$ as a function of magnetic field would be of great interest to confirm whether the high-field induced magnetic order in this system [62] is indeed the same as that found in the (Ru,Rh) series [49].

Finally, it is worth highlighting that extensive work on URu₂Si₂ has demonstrated that the various forms of magnetic order that emerge by destabilizing the HO phase are representative of magnetism in the extended family of UT₂Si₂ compounds (*T*: transition metal). This poses the fascinating question of whether HO may be stabilized in these related materials as well.

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