

## Crossover from multiple- to single-gap superconductivity in $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ alloys

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By using mostly the muon-spin rotation/relaxation ( $\mu\text{SR}$ ) technique, we investigate the superconductivity (SC) of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  ( $x = 0$  and  $1.6$ ) alloys, with  $T_c = 10.5$  and  $9.1$  K, respectively. At a macroscopic level, their superconductivity was studied by electrical resistivity, magnetization, and specific-heat measurements. In both compounds, the electronic specific heat and the low-temperature superfluid density data suggest a nodeless SC. The superconducting gap value and the specific heat discontinuity at  $T_c$  are larger than that expected from BCS theory in the weak-coupling regime, indicating strong-coupling superconductivity in the  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  family. In  $\text{Nb}_5\text{Ir}_3\text{O}$ , multigap SC is evidenced by the field dependence of the electronic specific heat coefficient and the superconducting Gaussian relaxation rate, as well as by the temperature dependence of the upper critical field. Pt substitution suppresses one of the gaps, and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  becomes a single-gap superconductor. By combining our extensive experimental results, we provide evidence for a multiple- to single-gap SC crossover in the  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  family.

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### I. INTRODUCTION

The  $A_5B_3$  compound family, where  $A$  is a transition or rare-earth metal and  $B$  is a (post)transition metal or a metalloid element, consists of more than 500 compounds with three distinct crystal structures. They are orthorhombic  $\text{Yb}_5\text{Sb}_3$ -type ( $Pnma$ , No. 62), tetragonal  $\text{Cr}_5\text{B}_3$ -type ( $I4/mcm$ , No. 140), and hexagonal  $\text{Mn}_5\text{Si}_3$ -type ( $P6_3/mcm$ , No. 193) structures. The latter possesses an interstitial  $2b$  site, allowing for the intercalation of light atoms, e.g., oxygen, boron, and carbon, to engineer the band topology. The ordered variant of the  $\text{Mn}_5\text{Si}_3$ -type structure is also known as the  $\text{Ti}_5\text{Ga}_4$ - or  $\text{Hf}_5\text{CuSn}_3$ -type structure.  $\text{Nb}_5\text{Ir}_3\text{O}$ , whose crystal structure is presented in the inset of Fig. 1, also belongs to the latter. Superconductivity (SC) has been reported to occur in several  $\text{Ti}_5\text{Ga}_4$ -type compounds, including  $\text{Nb}_5\text{Ir}_3\text{O}$  [1],  $(\text{Nb,Zr})_5\text{Pt}_3\text{O}$  [2,3],  $\text{Nb}_5\text{Ge}_3\text{C}_{0.3}$  [4], and  $\text{Zr}_5\text{Pt}_3\text{C}_x$  [5], with the highest superconducting transition temperature  $T_c$  reaching  $\sim 15$  K. Interestingly, specific heat and penetration depth results suggest a nodal superconducting gap in  $\text{Zr}_5\text{Pt}_3\text{C}_x$  and possibly unconventional SC [5].

The parent compound  $\text{Nb}_5\text{Ir}_3$  consists of mixed tetragonal and hexagonal phases, both showing superconducting behavior below  $T_c = 2.8$  and  $9.4$  K, respectively. The gradual intercalation of oxygen suppresses the tetragonal phase, making the hexagonal phase the dominant one. Accordingly,

the  $T_c$  value increases linearly with the interstitial oxygen content, reaching  $10.5$  K in the purely hexagonal  $\text{Nb}_5\text{Ir}_3\text{O}$  [1].  $\text{Nb}_5\text{Ir}_3$  exhibits an unusual interplay of electronegativity and SC states, whereas only SC remains in  $\text{Nb}_5\text{Ir}_3\text{O}$  [1]. Density functional theory calculations indicate  $\text{Nb}_5\text{Ir}_3$  is a multiband metal, whose density of states (DOS) at the Fermi level is dominated by the Nb  $4d$  and Ir  $5d$  orbitals [1]. The increase of  $T_c$  in  $\text{Nb}_5\text{Ir}_3\text{O}_x$  is most likely attributed to the enhanced electron-phonon coupling strength or to an increased DOS at the Fermi level (with extra Nb  $4d$  and O  $2p$  contributions) [1]. On the other hand, by applying external pressure,  $T_c$  decreases monotonically from  $10.5$  K at ambient pressure to  $9.5$  K at  $13$  GPa [1,6]. Unlike the  $\text{Nb}_5\text{Ir}_3\text{O}_x$  case, in  $\text{Zr}_5\text{Pt}_3$  the addition of oxygen reduces the  $T_c$  value from  $6.4$  to  $3.2$  K (in  $\text{Zr}_5\text{Pt}_3\text{O}_{0.6}$ ) [3]. Similarly, in  $\text{Zr}_5\text{Sb}_3$ , the addition of oxygen reduces  $T_c$ , with  $\text{Zr}_5\text{Sb}_3\text{O}$  being a normal metal down to  $1.8$  K [7].

Although the superconductivity of several  $\text{Mn}_5\text{Si}_3$  and  $\text{Ti}_5\text{Ga}_4$  types of compounds has been studied by magnetic and transport measurements and, in many cases, electronic band structure calculations are available, the microscopic nature of their superconducting phase remains largely unexplored. In  $\text{Nb}_5\text{Ir}_3\text{O}$ , the low- $T$  electronic specific heat data suggest a nodeless SC with multiple gaps [1,6]. However, in the  $\text{Nb}_5\text{Pt}_3\text{O}$  case, the electronic specific heat shows a single-exponential temperature dependence below  $T_c$ , more consistent with a single-gap superconductivity [2]. In  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ , Pt substitution increases the  $a$ -axis lattice constant while reducing the  $c$  axis. As shown in Fig. 1,  $T_c$  is almost independent of Pt content for  $x \leq 0.5$ , but it starts to decrease continuously

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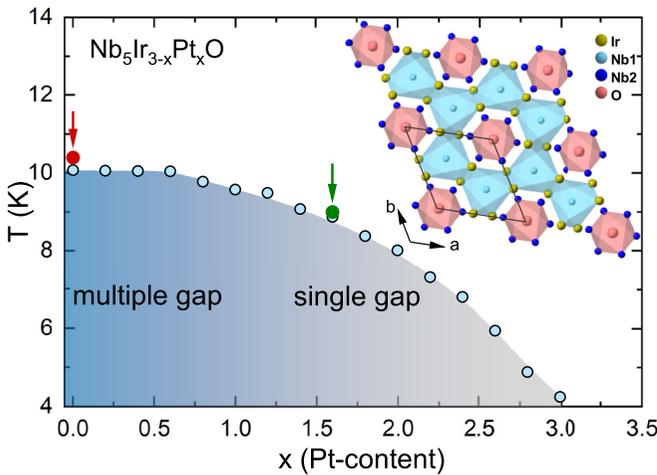


FIG. 1. The  $T$ - $x$  phase diagram of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ . The arrows indicate the two samples studied here. The inset shows the crystal structure of  $\text{Nb}_5\text{Ir}_3\text{O}$  viewed along the  $c$  axis (solid lines mark the unit cell). A crossover from multiple- to single-gap SC with increasing Pt content has been proposed. Data were adopted from Ref. [8].

for  $x > 0.5$ , reaching 4.3 K in  $\text{Nb}_5\text{Pt}_3\text{O}$  [3]. As hinted by specific heat data, one expects a crossover from multiple- to single-gap SC in  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ .

To provide further evidence of such a crossover, we initiated an extensive study of the superconducting properties of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  for  $x = 0$  and 1.6, two representative cases in the multiple- and single-gap regions (see arrows in Fig. 1), probing them at both the macroscopic and microscopic levels. Both compounds were found to be fully gapped superconductors, with  $\text{Nb}_5\text{Ir}_3\text{O}$  showing multiple gaps and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  being a single-gap superconductor, indicating a crossover at an intermediate Pt doping.

## II. EXPERIMENTAL DETAILS

Polycrystalline samples of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  were prepared by the arc-melting method (the full details are reported in Ref. [8]). The magnetic susceptibility, electrical resistivity, and specific heat measurements were performed on a 7-T Quantum Design magnetic property measurement system (MPMS-7) and a 9-T physical property measurement system (PPMS-9). The muon-spin rotation/relaxation ( $\mu\text{SR}$ ) experiments were carried out at the general-purpose surface-muon (GPS) spectrometer at the Swiss muon source ( $\text{S}\mu\text{S}$ ) at Paul Scherrer Institut, Villigen, Switzerland [9]. Both transverse-field (TF) and zero-field (ZF)  $\mu\text{SR}$  measurements were performed. The  $\mu\text{SR}$  data were analyzed using the MUSRFIT software package [10].

## III. RESULTS AND DISCUSSION

### A. Magnetization

The SC of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  ( $x = 0, 1.6$ ) was first characterized by magnetic susceptibility measurements, carried out in a 5-mT field, using both field-cooled (FC) and zero-field-cooled (ZFC) protocols. As shown in Fig. 2(a), the ZFC suscepti-

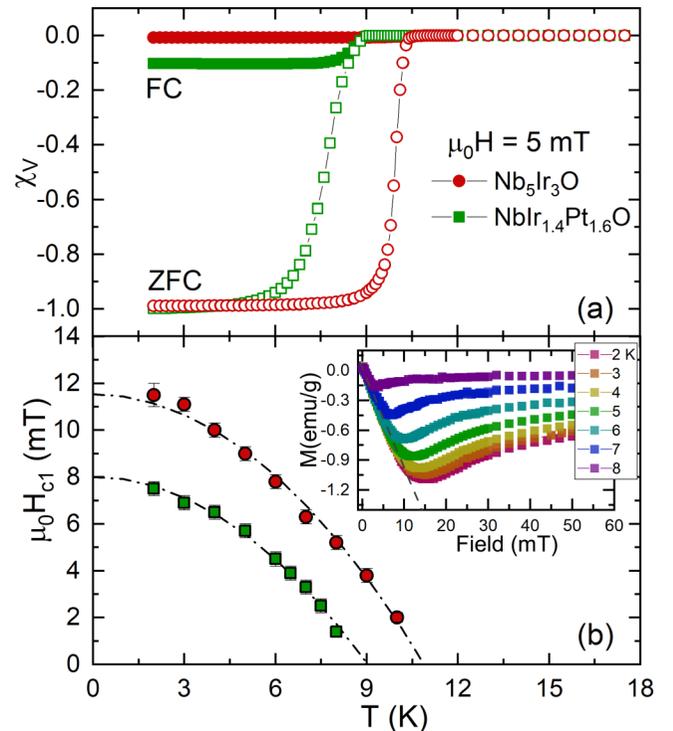


FIG. 2. (a) Temperature dependence of the magnetic susceptibility of  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ . (b) Estimated lower critical field  $\mu_0 H_{c1}$  vs temperature. The dash-dotted lines are fits to  $\mu_0 H_{c1}(T) = \mu_0 H_{c1}(0)[1 - (T/T_c)^2]$ . The inset shows the field-dependent magnetization  $M(H)$  recorded at various temperatures up to  $T_c$  for  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ . For each temperature, the lower critical field  $\mu_0 H_{c1}$  was determined as the value where  $M(H)$  starts deviating from linearity (dashed line). The magnetic susceptibilities were corrected by using the demagnetization factor obtained from the field-dependent magnetization at 2 K (base temperature).

bility, corrected to account for the demagnetization factor, indicates bulk SC below  $T_c = 10.5$  and 9.1 K for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , respectively. The well-separated ZFC and FC susceptibilities indicate strong flux-line pinning across the  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  series. To perform TF- $\mu\text{SR}$  measurements on superconductors, the applied magnetic field should exceed the lower critical field  $\mu_0 H_{c1}$ , so that the additional field-distribution broadening due to the flux-line lattice (FLL) can be quantified from the muon-spin relaxation rate. To determine  $\mu_0 H_{c1}$ , the field-dependent magnetization  $M(H)$  was measured at various temperatures up to  $T_c$ . As an example, the  $M(H)$  data for  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  are plotted in the inset of Fig. 2(b), with  $\text{Nb}_5\text{Ir}_3\text{O}$  showing similar features. The estimated  $\mu_0 H_{c1}$  values are summarized in Fig. 2(b) as a function of temperature for both samples. The dash-dotted lines represent fits to  $\mu_0 H_{c1}(T) = \mu_0 H_{c1}(0)[1 - (T/T_c)^2]$  and yield lower critical fields of 11.5(2) and 8.0(1) mT for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , respectively.

### B. Upper critical field

The upper critical field  $\mu_0 H_{c2}$  of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  was determined from the measurements of electrical resistivity  $\rho(T, H)$ , specific heat  $C(T, H)/T$ , and magnetization

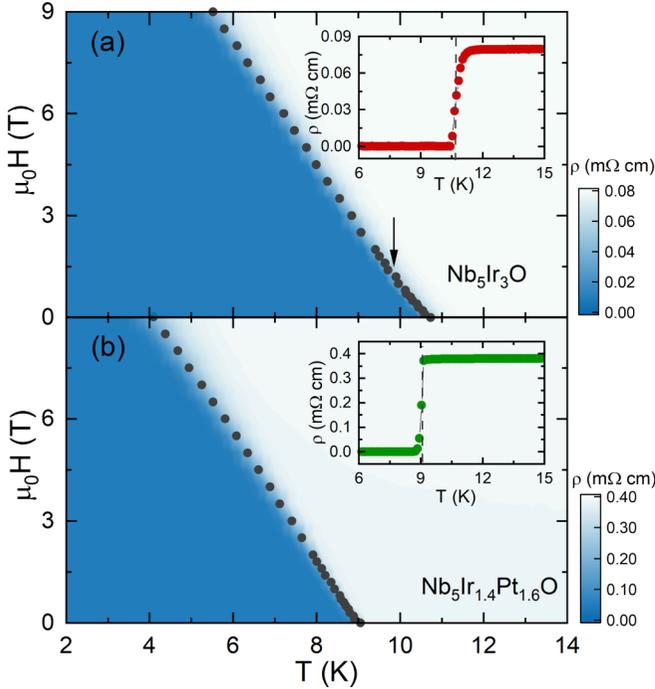


FIG. 3. Contour plots showing the electrical resistivity vs temperature (down to 2 K) and magnetic field (up to 9 T) for (a)  $\text{Nb}_5\text{Ir}_3\text{O}$  and (b)  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ . Color hues represent the absolute value of electrical resistivity. The symbols indicate the critical temperatures  $T_c$ , as determined from the middle of superconducting transition (see dashed lines in the insets). Insets: In both cases, the zero-field electrical resistivity shows a sharp transition. The arrow in (a) indicates a change in slope (see also inset in Fig. 11).

$M(T, H)$  under various applied magnetic fields up to 9 T. Here, in Figs. 3 and 4, we show  $\rho(T, H)$  and  $C(T, H)/T$  for both samples respectively. Upon applying a magnetic field, the superconducting transition, as detected from either  $\rho(T)$

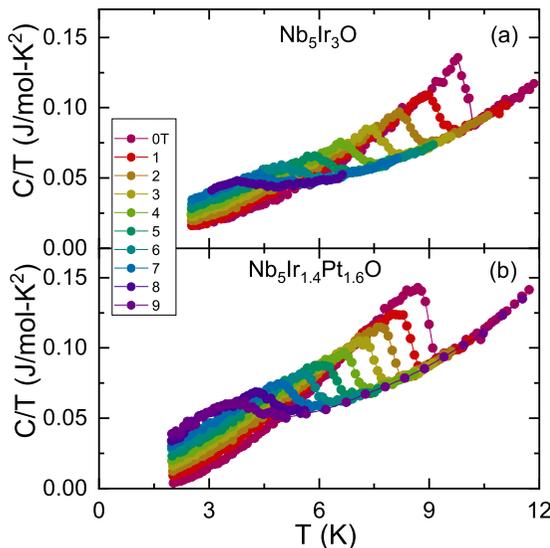


FIG. 4. Temperature dependence of the specific heat of (a)  $\text{Nb}_5\text{Ir}_3\text{O}$  and (b)  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  measured under various magnetic fields up to 9 T. The data were collected upon sample heating.

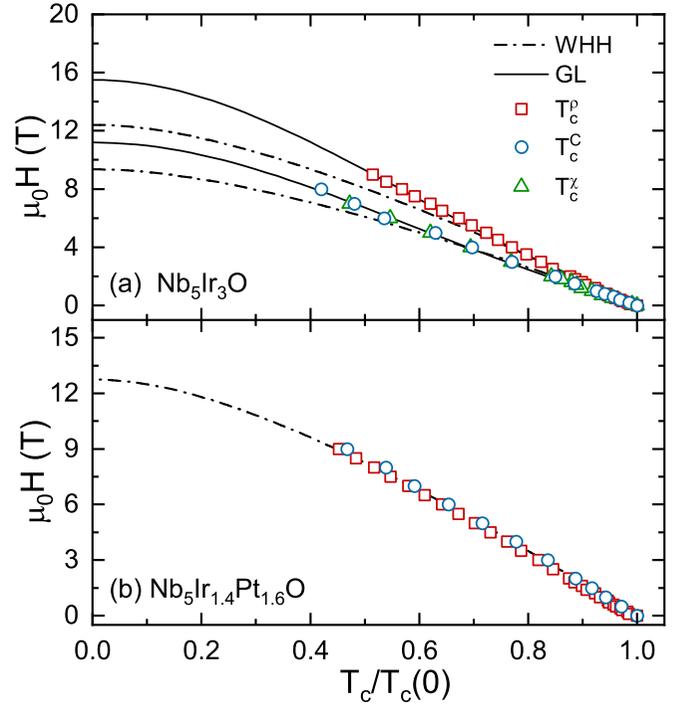


FIG. 5. Upper critical field  $\mu_0 H_{c2}$  vs reduced transition temperature  $T_c/T_c(0)$  for (a)  $\text{Nb}_5\text{Ir}_3\text{O}$  and (b)  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ . The  $T_c$  values were determined from temperature-dependent electrical resistivity  $\rho(T, H)$ , magnetization  $M(T, H)$ , and specific heat  $C(T, H)/T$ . For  $\rho(T, H)$  measurements,  $T_c$  is defined as the middle of the superconducting transition. Two different models, the GL (solid lines) and WHH models (dash-dotted lines), were used to analyze the  $\mu_0 H_{c2}(T)$  data. For the WHH model, the spin-orbit scattering was neglected.

or  $C(T)/T$  data, shifts towards lower temperatures. The insets of Fig. 3 show enlarged plots of zero-field electrical resistivity.  $\text{Nb}_5\text{Ir}_3\text{O}$  exhibits an onset of the SC at  $T_c^{\text{onset}} = 10.9$  K, and its resistivity drops to zero at  $T_c^{\text{zero}} = 10.5$  K. For  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ ,  $T_c^{\text{onset}} = 9.1$  K, and  $T_c^{\text{zero}} = 8.9$  K. The sharp superconducting transitions ( $\Delta T \sim 0.2\text{-}0.4$  K) in zero field, confirmed also by specific heat data (see Fig. 4), indicate the good quality of the samples.

The determined  $\mu_0 H_{c2}$  values as a function of the reduced temperature  $T_c/T_c(0)$  are summarized in Fig. 5. Here  $T_c(0)$  is the transition temperature in zero field. For  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , the  $\mu_0 H_{c2}$  values determined using different techniques are quite consistent. Conversely, for  $\text{Nb}_5\text{Ir}_3\text{O}$ , the data sets agree well only at low fields (below 2 T) since at higher fields the transition temperatures determined from  $\rho(T)$  data are systematically higher than those derived from  $C(T)/T$ . The surface or filamentary SC above bulk  $T_c$  might cause the different  $T_c$  values, but why it shows up only in  $\text{Nb}_5\text{Ir}_3\text{O}$  is not yet clear. The temperature dependence of  $\mu_0 H_{c2}(T)$  was analyzed by means of Ginzburg-Landau (GL) [11] and Werthamer-Helfand-Hohenberg (WHH) models [12]. For  $\text{Nb}_5\text{Ir}_3\text{O}$ , the better agreement of the GL model with the data is clearly seen in Fig. 5(a). At low fields, both the GL and WHH models reproduce very well the experimental data. However, at higher fields, the WHH model deviates significantly from the experimental data, giving underestimated values of  $\mu_0 H_{c2}^{\text{WHH}}(0) =$

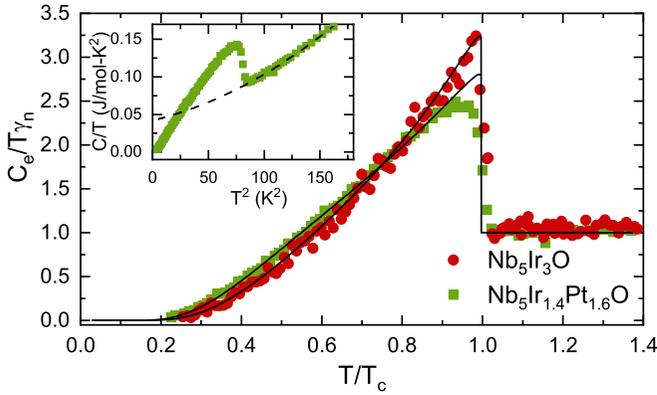


FIG. 6. Normalized electronic specific heat  $C_e/\gamma_n T$  of  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  as a function of the reduced temperature  $T/T_c$ . Inset: Specific heat  $C/T$  vs  $T^2$  for  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ ; the dashed line is a fit to  $C/T = \gamma_n + \beta T^2 + \delta T^4$  for  $T > T_c$ . The solid lines in the main panel represent the electronic specific heat calculated by considering a fully gapped  $s$ -wave model.

9.3(1) T and 12.4(1) T for  $C(T, H)/T$  and  $\rho(T, H)$ , respectively. In contrast, the GL model fits the data over the entire field range, providing  $\mu_0 H_{c2}^{\text{GL}}(0) = 11.2(1)$  T ( $C/T$ ) and 15.5(1) T ( $\rho$ ). In the case of  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ ,  $\mu_0 H_{c2}(T)$  data are reproduced very well by the WHH model, which yields  $\mu_0 H_{c2}^{\text{WHH}}(0) = 12.7(1)$  T.

The superconducting coherence length  $\xi$  can be calculated from  $\xi = \sqrt{\Phi_0/2\pi H_{c2}}$ , where  $\Phi_0 = 2.07 \times 10^{-3}$  T  $\mu\text{m}^2$  is the magnetic flux quantum. With a bulk  $\mu_0 H_{c2}(0) = 11.2(1)$  and 12.7(1) T, the calculated  $\xi(0)$  is 5.4(1) and 5.1(1) nm for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , respectively. The lower critical field  $\mu_0 H_{c1}$  is related to the magnetic penetration depth  $\lambda$  and the coherence length  $\xi$  via  $\mu_0 H_{c1} = (\Phi_0/4\pi\lambda^2)[\ln(\kappa) + 0.5]$ , where  $\kappa = \lambda/\xi$  is the GL parameter [13]. By using  $\mu_0 H_{c1} = 11.5(2)$  mT [8.0(1) mT] and  $\mu_0 H_{c2} = 11.2(1)$  T [12.7(1) T], the resulting magnetic penetration depth  $\lambda_{\text{GL}} = 249(3)$  nm [308(3) nm] for  $\text{Nb}_5\text{Ir}_3\text{O}$  ( $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ ) is comparable to 230(2) nm [314(2) nm], the experimental value evaluated from TF- $\mu$ SR data (see Sec. III D). A large GL parameter  $\kappa \sim 50$ –60 clearly indicates that  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  compounds are type-II superconductors.

### C. Zero-field specific heat

The zero-field specific heat data in the low- $T$  region (Fig. 6) show a sharp jump at  $T_c$ , again indicating a bulk superconducting transition and a good sample quality. As shown in the inset, the normal-state specific heat data of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  were fitted to  $C/T = \gamma_n + \beta T^2 + \delta T^4$ , where  $\gamma_n$  is the normal-state electronic specific heat coefficient, while the two other terms account for the phonon contribution to the specific heat. The derived values are  $\gamma_n = 37(5)$  mJ/mol K<sup>2</sup>,  $\beta = 0.32(9)$  mJ/mol K<sup>4</sup>, and  $\delta = 1.7(3)$   $\mu$ J/mol K<sup>6</sup> for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\gamma_n = 42(6)$  mJ/mol K<sup>2</sup>,  $\beta = 0.36(13)$  mJ/mol K<sup>4</sup>, and  $\delta = 2.6(5)$   $\mu$ J/mol K<sup>6</sup> for  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ . Such large  $\gamma$  values suggest a relatively large effective electron mass (see Table I) and also strong electronic correlations in  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ . The Debye temperature can be estimated

using  $\Theta_D = (12\pi^4 Rn/5\beta)^{1/3}$ , where  $R = 8.314$  J/mol K is the molar gas constant and  $n = 9$  is the number of atoms per formula unit, giving  $\Theta_D = 380(7)$  and 365(8) K for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ . The density of states at the Fermi level  $N(\epsilon_F)$  is evaluated to be  $N(\epsilon_F) = 3\gamma_n/(\pi^2 k_B^2) = 16(2)$  states/eV f.u. ( $\text{Nb}_5\text{Ir}_3\text{O}$ ) and 18(2) states/eV f.u. ( $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ ) [14], where  $k_B$  is the Boltzmann constant. Both values are comparable to those from the electronic band structure calculations [1]. The electron-phonon coupling constant  $\lambda_{\text{ep}}$ , a measure of the attractive interaction between electrons due to phonons, was estimated using the semiempirical McMillan formula [15]:

$$\lambda_{\text{ep}} = \frac{1.04 + \mu^* \ln(\Theta_D/1.45 T_c)}{(1 - 0.62 \mu^*) \ln(\Theta_D/1.45 T_c) - 1.04}. \quad (1)$$

The Coulomb pseudopotential  $\mu^*$  is material specific, typically lying in the  $0.1 \leq \mu^* \leq 0.15$  range. Here  $\mu^*$  was fixed to 0.13, a typical value for metallic samples [15]. From Eq. (1) we obtain  $\lambda_{\text{ep}} = 0.8(2)$  ( $\text{Nb}_5\text{Ir}_3\text{O}$ ) and 0.7(3) ( $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ ), which classifies both of them as relatively strongly coupled superconductors. This is consistent with previous results [1,6] and compatible with other strongly coupled superconductors, e.g.,  $\text{Ba}_{1-x}\text{K}_x\text{BiO}_3$  ( $\lambda_{\text{ep}} \sim 1$ ) and  $\text{W}_3\text{Al}_2\text{C}$  ( $\lambda_{\text{ep}} \sim 0.78$ ) [16–18].

The electronic specific heat is obtained by subtracting the phonon contribution from the total specific heat. The main panel in Fig. 6 shows the temperature dependence of  $C_e/\gamma_n T$ , from which one can evaluate the discontinuity at  $T_c$  to be  $\Delta C/\gamma T_c = 2.24(7)$  for  $\text{Nb}_5\text{Ir}_3\text{O}$  and 1.50(5) for  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , both larger than the weak-coupling BCS value of 1.43. The temperature evolution of the SC-related contribution to the entropy can be calculated from the BCS expression [19]:

$$S(T) = -\frac{6\gamma_n}{\pi^2 k_B} \int_0^\infty [f \ln f + (1-f) \ln(1-f)] d\epsilon, \quad (2)$$

where  $f = (1 + e^{E/k_B T})^{-1}$  is the Fermi distribution and  $E(\epsilon) = \sqrt{\epsilon^2 + \Delta^2(T)}$  is the excitation energy of quasiparticles, with  $\epsilon$  being the electron energy measured relative to the chemical potential (Fermi energy) [19,20]. Here  $\Delta(T) = \Delta_0 \tanh\{1.82[1.018(T_c/T - 1)]^{0.51}\}$  [21], where  $\Delta_0$  being the zero-temperature superconducting gap value. The electronic specific heat in the superconducting state can be calculated from  $C_e = T \frac{dS}{dT}$ . The solid lines in Fig. 6 are fits to the  $s$ -wave model with a single gap  $\Delta_0 = 1.89(2)$  and 1.53(1) meV for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , respectively. Both gap values are consistent with the TF- $\mu$ SR results (see Fig. 8 below) and are larger than the standard BCS value ( $\Delta_0 \sim 1.73 k_B T_c \sim 1.4$  meV) in the weak-coupling limit. This indicates strongly coupled superconducting pairs in  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ , consistent with the large  $\lambda_{\text{ep}}$  values. In previous studies, two superconducting gaps were required to describe the zero-field electronic specific heat at  $T < 1/3 T_c$  [1,6]. Note, however, that the extra gap value ( $< 0.02$  meV) was significantly smaller than the large gap and that it accounts for less than 10% of the total weight [1,6]. As such, it can easily be influenced by disorder or oxygen content. In the case of our samples, a single gap could reproduce the electronic specific heat data very well over the whole temperature range. However, the field-dependent electronic specific heat coefficient, the TF- $\mu$ SR

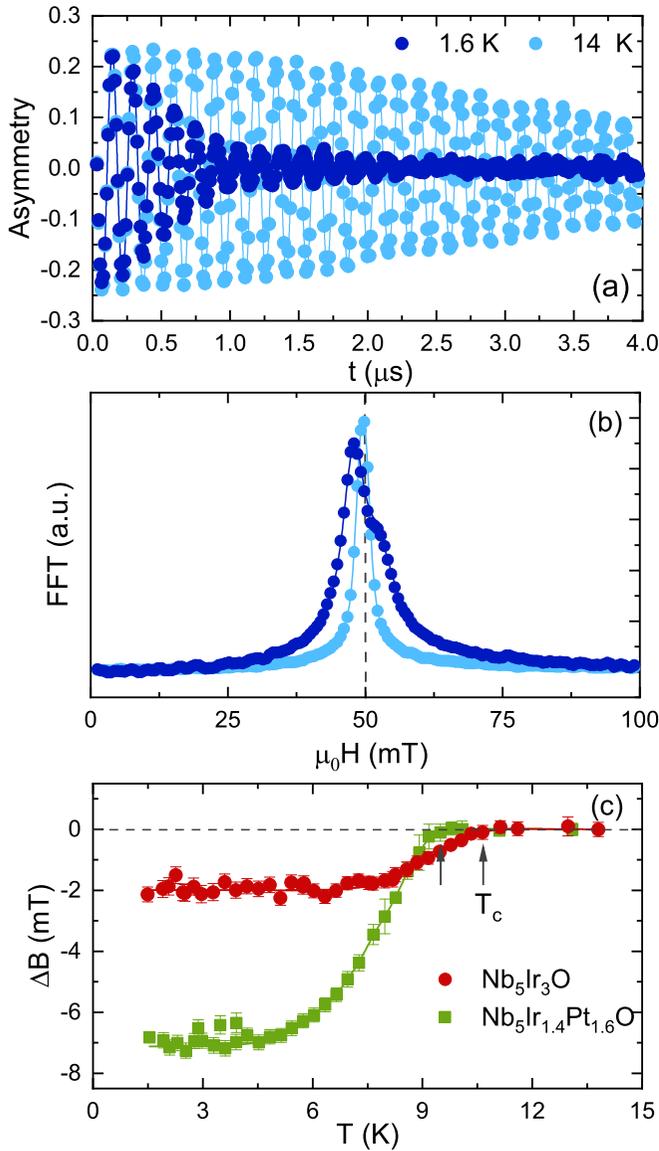


FIG. 7. (a) TF- $\mu$ SR time spectra collected at temperatures below (1.6 K) and above (14 K)  $T_c$  in an applied field of 50 mT for Nb<sub>5</sub>Ir<sub>3</sub>O. (b) The fast Fourier transform of the TF- $\mu$ SR time spectra. Note that Nb<sub>5</sub>Ir<sub>1.4</sub>Pt<sub>1.6</sub>O exhibits similar spectra. Solid lines in (a) and (b) are fits to Eq. (3) using two Gaussian relaxations. The dashed vertical line indicates the applied magnetic field, showing a clear diamagnetic shift. (c) Diamagnetic shift ( $\Delta B = \langle B \rangle - B_{\text{appl.}}$ ) vs temperature for Nb<sub>5</sub>Ir<sub>3</sub>O and Nb<sub>5</sub>Ir<sub>1.4</sub>Pt<sub>1.6</sub>O.

relaxation rates, and the temperature-dependent upper critical fields all suggest multigap features in Nb<sub>5</sub>Ir<sub>3</sub>O, while Nb<sub>5</sub>Ir<sub>1.4</sub>Pt<sub>1.6</sub>O is a single-gap superconductor (see details in Sec. III E).

#### D. Transverse-field $\mu$ SR

To investigate the superconducting properties of Nb<sub>5</sub>Ir<sub>3-x</sub>Pt<sub>x</sub>O at a microscopic level, TF- $\mu$ SR measurements were systematically carried out, covering both the normal and superconducting states. To track the additional field-distribution broadening due to the FLL in the mixed

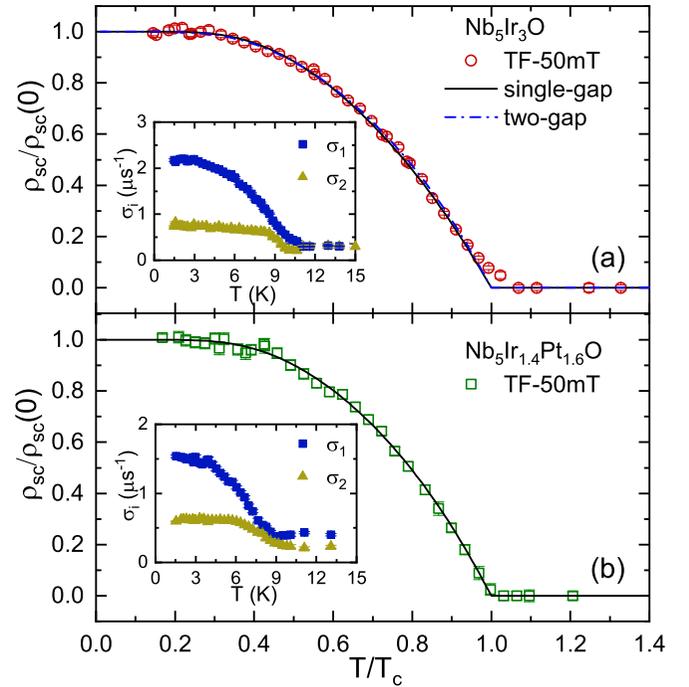


FIG. 8. Superfluid density vs reduced temperature  $T/T_c$ , as determined from TF- $\mu$ SR measurements in an applied magnetic field of 50 mT for (a) Nb<sub>5</sub>Ir<sub>3</sub>O and (b) Nb<sub>5</sub>Ir<sub>1.4</sub>Pt<sub>1.6</sub>O. The insets show the temperature dependence of the muon-spin relaxation rate  $\sigma_i(T)$ . The lines through the data represent fits to a fully gapped  $s$ -wave model with either a single gap (solid) or two gaps (dash-dotted). The goodness of fit is  $\chi_r^2 = 3.8$  (single-gap model) and 3.2 (two-gap model), respectively.

state, a magnetic field of 50 mT was applied in the normal state, before cooling the sample below  $T_c$ . The TF- $\mu$ SR time spectra were collected at various temperatures upon warming after the FC protocol. Figure 7(a) shows two representative TF- $\mu$ SR spectra of Nb<sub>5</sub>Ir<sub>3</sub>O collected at 1.6 K (i.e., below  $T_c$ ) and 14 K (above  $T_c$ ). In the normal state, the spectra show a relatively weak damping, reflecting a uniform field distribution. The enhanced depolarization rate in the superconducting state is attributed to the inhomogeneous field distribution due to the FLL, causing an additional field broadening in the mixed state. Such broadening is clearly visible in Fig. 7(b), where the fast-Fourier-transform (FFT) spectra of the corresponding TF- $\mu$ SR data in Fig. 7(a) are shown. To account for the asymmetric field distribution in the superconducting state, the  $\mu$ SR spectra were modeled by the following expression:

$$A_{\text{TF}}(t) = \sum_{i=1}^n A_i \cos(\gamma_{\mu} B_i t + \phi) e^{-\sigma_i^2 t^2 / 2} + A_{\text{bg}} \cos(\gamma_{\mu} B_{\text{bg}} t + \phi). \quad (3)$$

Here  $A_i$  (98%) and  $A_{\text{bg}}$  (2%) represent the initial muon-spin asymmetries for muons implanted in the sample and sample holder, respectively, with the latter not undergoing any depolarization.  $B_i$  and  $B_{\text{bg}}$  are the local fields sensed by implanted muons in the sample and sample holder,  $\gamma_{\mu} = 2\pi \times 135.53$  MHz/T is the muon gyromagnetic ratio,  $\phi$  is the

shared initial phase, and  $\sigma_i$  is the Gaussian relaxation rate of the  $i$ th component. More than one oscillation is required to describe the TF- $\mu$ SR spectra of Nb<sub>5</sub>Ir<sub>3-x</sub>Pt<sub>x</sub>O samples. As illustrated in Fig. 7(b), at 1.6 K, two broad peaks can clearly be seen, below and above the applied magnetic field (50 mT). The solid lines in Figs. 7(a) and 7(b) represent fits to Eq. (3) with  $n = 2$ . Below  $T_c$ , a diamagnetic field shift appears in both samples [see Fig. 7(c)]. The relaxation rate is temperature independent and small above  $T_c$ , but below  $T_c$  it starts to increase due to the onset of the FLL and the increased superfluid density (see insets in Fig. 8).

In the case of multicomponent oscillations, the first term in Eq. (3) describes the field distribution as a sum of  $n$  Gaussian relaxations (here  $n = 2$ ) [22]:

$$P(B) = \gamma_\mu \sum_{i=1}^2 \frac{A_i}{\sigma_i} \exp\left[-\frac{\gamma_\mu^2 (B - B_i)^2}{2\sigma_i^2}\right]. \quad (4)$$

The first and second moments of the field distribution can be calculated by

$$\begin{aligned} \langle B \rangle &= \sum_{i=1}^2 \frac{A_i B_i}{A_{\text{tot}}}, \\ \langle B^2 \rangle &= \frac{\sigma_{\text{eff}}^2}{\gamma_\mu^2} = \sum_{i=1}^2 \frac{A_i}{A_{\text{tot}}} \left[ \frac{\sigma_i^2}{\gamma_\mu^2} + (B_i - \langle B \rangle)^2 \right], \end{aligned} \quad (5)$$

where  $A_{\text{tot}} = A_1 + A_2$ . The superconducting Gaussian relaxation rate  $\sigma_{\text{sc}}$  can be extracted by subtracting the nuclear contribution according to  $\sigma_{\text{sc}} = \sqrt{\sigma_{\text{eff}}^2 - \sigma_n^2}$ , where  $\sigma_n$  is the nuclear relaxation rate, assumed to be temperature independent in such a narrow temperature range (see also ZF- $\mu$ SR below).

Since the upper critical fields of Nb<sub>5</sub>Ir<sub>3-x</sub>Pt<sub>x</sub>O (see Fig. 5) are significantly higher than the transverse field used in the  $\mu$ SR measurements (here 50 mT), the magnetic penetration depth  $\lambda(T)$  and the superfluid density  $\rho_{\text{sc}}(T)$  [ $\propto \lambda^{-2}(T)$ ] can be obtained from  $\sigma_{\text{sc}}(T)$  according to [13,23]

$$\frac{\sigma_{\text{sc}}^2(T)}{\gamma_\mu^2} = 0.00371 \frac{\Phi_0^2}{\lambda^4(T)}. \quad (6)$$

The derived superfluid density normalized to the zero-temperature values is shown in the main panels of Fig. 8. The superfluid density is almost constant at temperatures below  $T_c/3$ , indicating fully gapped SC in Nb<sub>5</sub>Ir<sub>3-x</sub>Pt<sub>x</sub>O, in good agreement with the specific heat results (see Fig. 6).

For more quantitative insight into the SC of Nb<sub>5</sub>Ir<sub>3-x</sub>Pt<sub>x</sub>O, the derived superfluid density  $\rho_{\text{sc}}(T)$  was further analyzed by using a fully gapped  $s$ -wave model:

$$\frac{\lambda^{-2}(T)}{\lambda_0^{-2}} = \frac{\rho_{\text{sc}}(T)}{\rho_{\text{sc}}(0)} = 1 + 2 \int_{\Delta(T)}^{\infty} \frac{E}{\sqrt{E^2 - \Delta^2(T)}} \frac{\partial f}{\partial E} dE. \quad (7)$$

Here  $f$  and  $\Delta(T)$  are the Fermi and superconducting gap functions (see details in Sec. III C). The solid black lines in Fig. 8 are fits to the above model with a single gap, which yield zero-temperature gap values  $\Delta_0 = 1.79(3)$  and  $1.67(2)$  meV and magnetic penetration depths  $\lambda_0 = 230(2)$  and  $314(2)$  nm for Nb<sub>5</sub>Ir<sub>3</sub>O and Nb<sub>5</sub>Ir<sub>1.4</sub>Pt<sub>1.6</sub>O, respectively. The estimated BCS coherence length  $\xi_0$  is larger than the electronic mean

TABLE I. Normal and superconducting state properties of Nb<sub>5</sub>Ir<sub>3</sub>O and Nb<sub>5</sub>Ir<sub>1.4</sub>Pt<sub>1.6</sub>O, as determined from electrical resistivity, magnetic susceptibility, specific heat, and  $\mu$ SR measurements. The London penetration depth  $\lambda_L$ , the effective mass  $m^*$ , carrier density  $n_s$ , BCS coherence length  $\xi_0$ , electronic mean free path  $l_e$ , Fermi velocity  $v_F$ , and effective Fermi temperature  $T_F$  were estimated following equations in Ref. [41].

Property	Units	Nb <sub>5</sub> Ir <sub>3</sub> O	Nb <sub>5</sub> Ir <sub>1.4</sub> Pt <sub>1.6</sub> O
$T_c^a$	K	10.5(1)	9.1(1)
$\mu_0 H_{c1}$	mT	11.5(2)	8.0(1)
$\mu_0 H_{c1}^{\mu\text{SR}}$	mT	13.3(1)	7.7(1)
$\mu_0 H_{c2}^p$	T	15.5(1)	12.7(1)
$\mu_0 H_{c2}^c$	T	11.2(1)	12.7(1)
$\gamma_n$	mJ/mol K <sup>2</sup>	37(5)	42(6)
$\Theta_D^C$	K	379(30)	368(30)
$\lambda_{\text{ep}}$		0.8(2)	0.73(25)
$N(\epsilon_F)$	states/eV f.u.	16(2)	18(2)
$\Delta_0^{\mu\text{SR}}$ (clean)	meV	1.79(3) <sup>b</sup>	1.67(2)
$\Delta_0^{\mu\text{SR}}$ (dirty)	meV	1.61(3)	1.51(3)
$\Delta_0^c$	meV	1.89(2)	1.53(1)
$\Delta C/\gamma_n T_c$		2.24(7)	1.50(5)
$\lambda_0^{\mu\text{SR}}$	nm	230(2)	314(2)
$\lambda_{\text{GL}}$	nm	249(3)	308(3)
$\lambda_L$	nm	73(5)	81(9)
$\xi(0)$	nm	5.4(1)	5.1(1)
$\kappa$		50(3)	60(3)
$m^*$	$m_e$	8.0(8)	9.2(4)
$n_s$	$10^{28} \text{ m}^{-3}$	4.2(5)	3.9(7)
$\xi_0$	nm	18(1)	16(2)
$l_e$	nm	2.1(1)	1.2(1)
$\xi_0/l_e$		9(1)	14(2)
$v_F$	$10^5 \text{ ms}^{-1}$	1.5(1)	1.3(1)
$T_F$	$10^4 \text{ K}$	1.27(5)	1.1(1)

<sup>a</sup>Similar values were determined via electrical resistivity, magnetic susceptibility, and specific heat measurements.

<sup>b</sup>The two-gap model provides an averaged gap value of 1.84(3) meV.

free path  $l_e$  (see Table I), implying the samples are in the dirty limit. Therefore, the temperature-dependent superfluid density was also analyzed using a dirty-limit model. In this case, in the BCS approximation, the temperature dependence of the superfluid density is given by  $\rho_{\text{sc}}(T) = \frac{\Delta(T)}{\Delta_0} \tanh[\frac{\Delta(T)}{2k_B T}]$  [19], which yields gap values of 1.61(3) and 1.51(3) meV, slightly smaller than the clean-limit value, but still larger than the weak-coupling BCS value. In the Nb<sub>5</sub>Ir<sub>3</sub>O case, to compare the zero-field electronic specific heat results [1,6] with those from  $\mu$ SR, the superfluid density was also analyzed using a two-gap model [21,24,25]. As shown by the blue dash-dotted line in Fig. 8 (a), the two-gap model shows slightly better agreement with the  $\rho_{\text{sc}}(T)$  data, as confirmed by the smaller  $\chi_r^2$  value. The derived gap values are  $\Delta_0^s = 1.34(3)$  meV (small) and  $\Delta_0^l = 1.97(3)$  meV (large), with a weight  $w = 0.2$  for the small gap. Such a small weight makes the multigap features barely visible in the  $\rho_{\text{sc}}(T)$  or  $C_e/T$  data. Although the  $\rho_{\text{sc}}(T)$  data can be well described by an  $s$ -wave model with either one or two gaps, the latter is more consistent with the field-dependent electronic specific heat coefficient (Fig. 9),

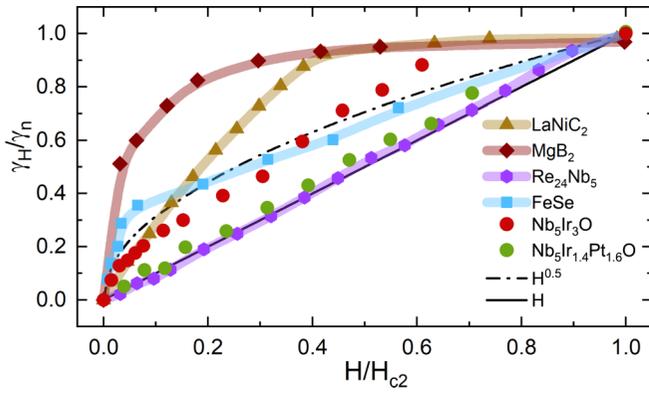


FIG. 9. The normalized specific heat coefficient  $\gamma_H/\gamma_n$  vs the reduced magnetic field  $H/H_{c2}(0)$  for  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ .  $\gamma_H$  is obtained as the linear extrapolation of  $C/T$  vs  $T^2$  (in the superconducting phase) to zero temperature. The solid line indicates a linear dependence, as predicted for a single-gap  $s$ -wave gap structure; the dash-dotted line represents the dependence expected for an anisotropic gap or a gap with nodes, e.g., the  $d$  wave. The data for the reference samples are adopted from Refs. [26–29].

the superconducting Gaussian relaxation rate (Fig. 10), the upper critical field (Fig. 11), and the electronic band structure calculations [1].

### E. Evidence of multigap superconductivity in $\text{Nb}_5\text{Ir}_3\text{O}$

The multigap nature of  $\text{Nb}_5\text{Ir}_3\text{O}$  superconductivity can be further inferred from the field-dependent electronic specific heat coefficient  $\gamma_H(H)$ . The normalized values  $\gamma_H/\gamma_n$  vs the reduced magnetic field  $H/H_{c2}(0)$  are shown in Fig. 9 (here  $\gamma_n$  is the zero-field normal-state value). For  $\text{Nb}_5\text{Ir}_3\text{O}$ , due to its multigap nature, it is difficult to describe the field dependence with a simple formula.  $\gamma_H(H)$  clearly deviates from the linear field dependence (solid line) expected for single-gap fully gapped superconductors [30] or from the square-root dependence  $\sqrt{H}$  (dash-dotted line) expected for nodal superconductors [31,32]. Different from the case of  $\text{Nb}_5\text{Ir}_3\text{O}$  but similar to  $\text{Re}_{24}\text{Nb}_5$  [29],  $\gamma_H(H)$  of  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  is practically linear in field, more consistent with a single gap SC.  $\text{Nb}_5\text{Ir}_3\text{O}$  instead exhibits features similar to other well-studied multigap superconductors, e.g.,  $\text{LaNiC}_2$  [26],  $\text{FeSe}$  [27], and  $\text{MgB}_2$  [28], although the slopes of  $\gamma_H(H)$  close to zero field are different, reflecting the different magnitudes and weights of the smaller gap.

To get further insight into the multigap SC of  $\text{Nb}_5\text{Ir}_3\text{O}$ , TF- $\mu$ SR measurements were performed in different magnetic fields up to 780 mT at base temperature (1.6 K) in both samples. Figure 10(a) shows the TF- $\mu$ SR spectra of  $\text{Nb}_5\text{Ir}_3\text{O}$ , collected at 20 and 200 mT, with the spectra in other applied fields and in  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  showing similar features. The spectra were analyzed using the same model as described in Eq. (3), and the resulting superconducting Gaussian relaxation rates  $\sigma_{sc}$  versus the applied magnetic field are summarized in Fig. 10(b) ( $\text{Nb}_5\text{Ir}_3\text{O}$ ) and Fig. 10(c) ( $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ ).  $\sigma_{sc}(H)$  was analyzed using both a single- and a two-band model. In the latter case, each band is characterized by its own superconducting coherence length [i.e.,  $\xi_1(0)$  and  $\xi_2(0)$ ] and

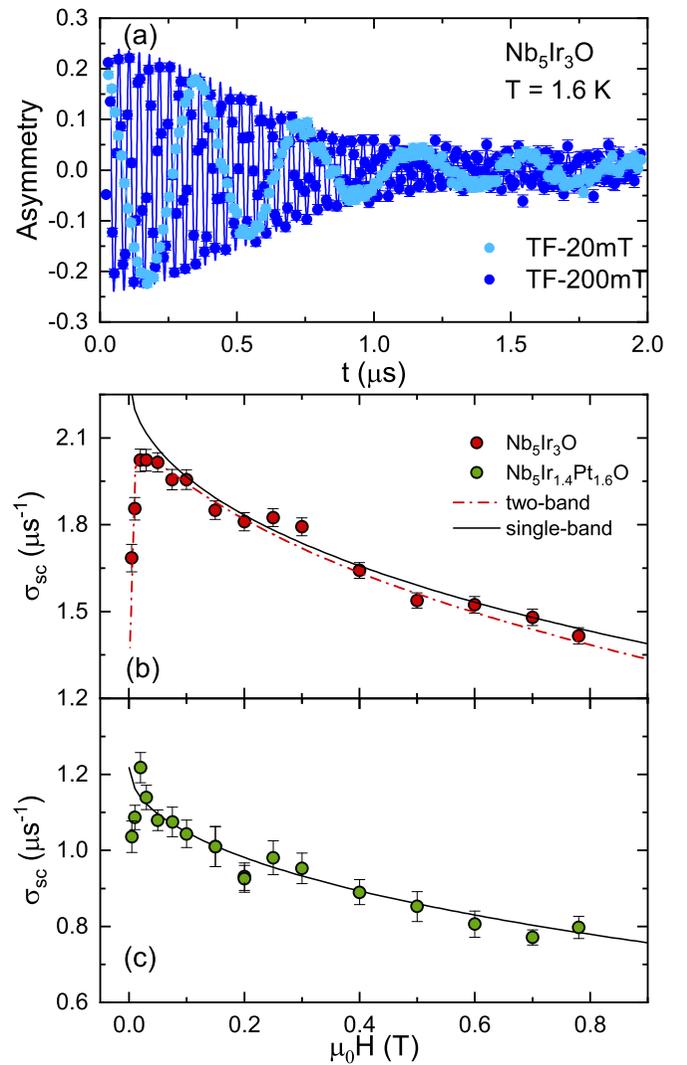


FIG. 10. (a) TF- $\mu$ SR time spectra for  $\text{Nb}_5\text{Ir}_3\text{O}$  measured at  $T = 1.5$  K (superconducting state) in a field of 20 and 200 mT.  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , too, shows a similar behavior. Field-dependent superconducting Gaussian relaxation rate  $\sigma_{sc}(H)$  for (b)  $\text{Nb}_5\text{Ir}_3\text{O}$  and (c)  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ . The dash-dotted and solid lines represent fits to two-band and single-band models, respectively. The goodness of fit is  $\chi_r^2 = 1.17$  (two-band model) and 2.65 (single-band model) for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\chi_r^2 = 1.15$  for  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  (single-band model). The data for  $H < H_{c1}$  are excluded to evaluate the  $\chi_r^2$  value.

a weight  $w$  accounting for the contribution of the second band [ $\xi_2(0)$ ] to the total superfluid density, similar to the two-gap model in Fig. 8(a). The details of the single- and two-band models can be found in Refs. [33,34]. By fixing  $w = 0.2$  and  $\xi_1(0) = 5.4$  nm, as estimated from the analysis of  $\rho_{sc}(T)$  (Fig. 8) and of the upper critical field (Fig. 5), we obtain the dash-dotted line fit in Fig. 10(b), which provides  $\lambda_0 = 222(3)$  nm and  $\xi_2(0) = 14(1)$  nm. The derived  $\lambda_0$  is consistent with the value (230 nm) estimated from the analysis of  $\rho_{sc}(T)$  in Fig. 8(a). The upper critical field of  $1.7(2)$  T calculated from the coherence length of the second band  $\xi_2(0)$  is also in good agreement with the field values where  $\gamma_H(H)$  and  $\mu_0 H_{c2}(T)$  change their slope, as shown by the arrows in Fig. 3(a) and the inset in Fig. 11. While the single-band model

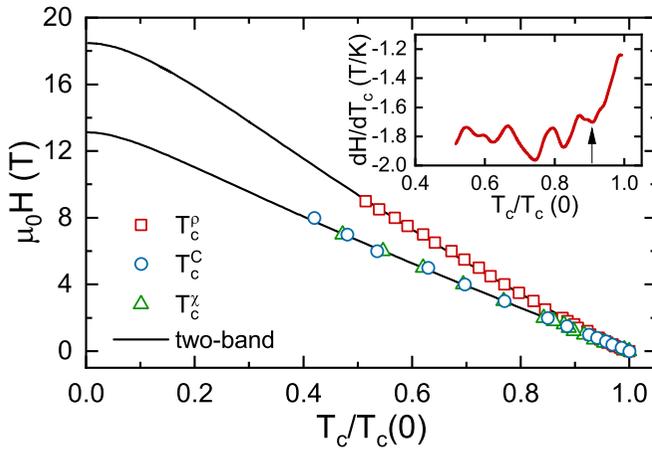


FIG. 11. Upper critical field  $\mu_0 H_{c2}$  vs reduced transition temperature  $T_c/T_c(0)$  for  $\text{Nb}_5\text{Ir}_3\text{O}$  [the same as in Fig. 5(a)]. The solid lines are fits using a two-band model. Inset: Temperature derivative of the upper critical field, as determined from the electrical resistivity.

reproduces very well the  $\sigma_{sc}(H)$  data for  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}$  [see Fig 10(c)], it is less satisfactory for  $\text{Nb}_5\text{Ir}_3\text{O}$ , where a two-band model is required to obtain a similarly small  $\chi_r^2$  value [see Fig 10(b)].

A positive curvature of the upper critical field near  $T_c$  is considered a typical feature of multiband superconductors, e.g.,  $\text{MgB}_2$  [35,36]. It reflects the gradual suppression of the small superconducting gap with increasing magnetic field. Indeed, in  $\text{Nb}_5\text{Ir}_3\text{O}$ ,  $\mu_0 H_{c2}(T)$  exhibits a clear kink close to 1.5 T [see the arrow in Fig. 3(a)], which coincides with the field value which suppresses the small superconducting gap. This is reflected also in the derivative of  $\mu_0 H_{c2}(T)$  with respect to temperature (see the arrow in the inset of Fig. 11). Also  $\gamma_H(H)$  changes its slope near this critical field (close to 0.1 in Fig. 9). Since, in  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$ , the small gap is already suppressed by Pt substitution, it exhibits a linear field dependence of  $\gamma_H(H)$  (Fig. 9), consistent with a fully gapped SC with a single gap. As shown in Fig. 11, the upper critical field of  $\text{Nb}_5\text{Ir}_3\text{O}$  was also analyzed by a two-band model [37], from which we estimate the upper critical field values  $\mu_0 H_{c2}(0) = 13.1(1)$  [from  $C(T, H)$  or  $M(T, H)$ ] and  $18.5(1)$  T [from  $\rho(T, H)$ ], both consistent with the GL model in Fig. 5(a).

### F. Uemura plot

According to the ratio of  $T_c$  to the effective Fermi temperature  $T_F$ , the different classes of superconductors can be classified following the so-called Uemura plot [39]. As seen in Fig. 12, conventional BCS superconductors exhibit  $T_c/T_F < 10^{-3}$ , here exemplified by the elemental Sn, Al, and Zn superconductors. In contrast, as indicated by the shaded region, several types of unconventional superconductors, including heavy fermions, organic superconductors, iron pnictides, and cuprates, all lie within a  $10^{-2} < T_c/T_F < 10^{-1}$  band. Between these two categories are located several multigap superconductors, e.g.,  $\text{LaNiC}_2$ ,  $\text{NbSe}_2$ , and  $\text{MgB}_2$ . By using the superconducting parameters obtained from the measurements presented here (see details in Table I), the calculated  $T_c/T_F$

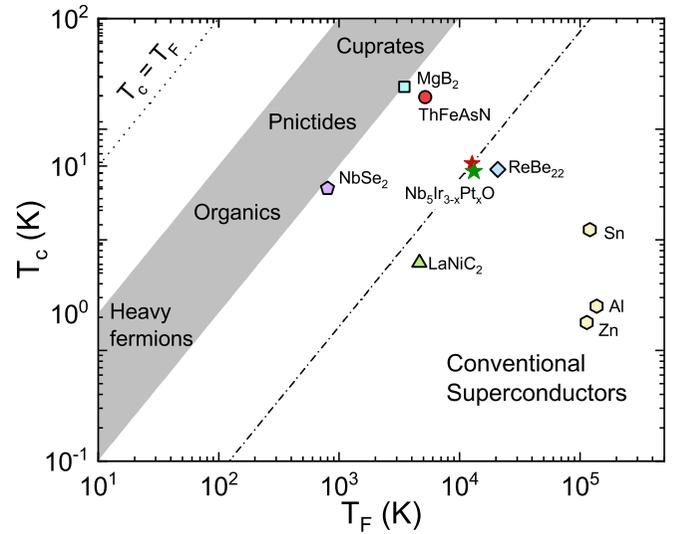


FIG. 12. Uemura plot showing  $T_c$  against the effective Fermi temperature  $T_F$  for various superconductors. The shaded region, with  $1/100 < T_c/T_F < 1/10$ , indicates the band of unconventional superconductors, such as heavy fermions, organic superconductors, iron pnictides, and cuprates. The dotted line corresponds to  $T_c = T_F$ , while the dash-dotted line indicates  $T_c/T_F = 8.2 \times 10^{-4}$  for  $\text{Nb}_5\text{Ir}_3\text{O}$ . The data for the reference samples were adopted from Refs. [25,29,38–40].

values for  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  are  $\sim 8.2\text{--}8.4 \times 10^{-4}$  (see star symbols in Fig. 12). Although there is no evidence for them to be classified as unconventional superconductors, the  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  family is clearly far off the conventional superconductors, and it shows ratios similar to other multigap superconductors, such as  $\text{LaNiC}_2$  and  $\text{ReBe}_{22}$  (both located near the dash-dotted line).

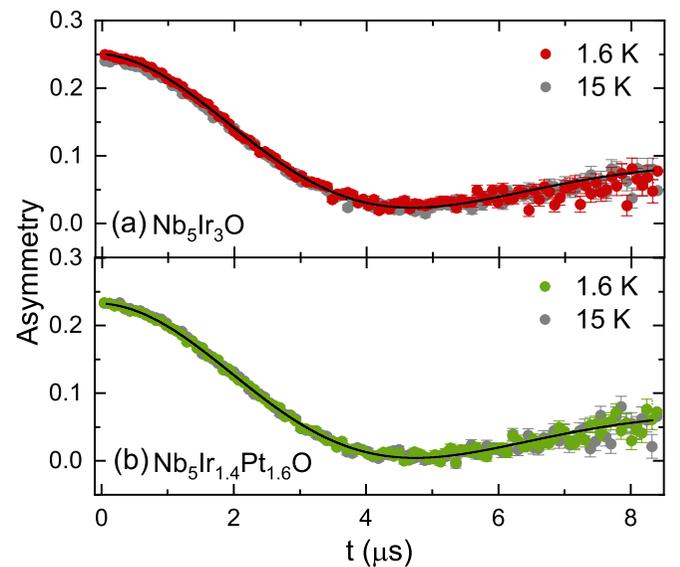


FIG. 13. ZF- $\mu\text{SR}$  spectra for (a)  $\text{Nb}_5\text{Ir}_3\text{O}$  and (b)  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  in the superconducting state (1.6 K) and the normal state (15 K). Solid lines are fits to the equation described in the text. None of the data sets shows clear changes with temperature.

### G. Zero-field $\mu$ SR

ZF- $\mu$ SR measurements were also performed in the normal and superconducting states to search for possible magnetism or time-reversal symmetry breaking in the superconducting state of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ . Representative ZF- $\mu$ SR spectra collected above (15 K) and below (1.6 K)  $T_c$  for  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  are shown in Fig. 13. Neither coherent oscillations nor fast damping could be identified at either temperature, implying the lack of any magnetic order or fluctuations in  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ . Therefore, in the absence of an applied magnetic field, the weak muon-spin relaxation is mainly determined by the randomly oriented nuclear moments, which can be modeled by a Gaussian Kubo-Toyabe relaxation function  $G_{\text{KT}} = [\frac{1}{3} + \frac{2}{3}(1 - \sigma_{\text{ZF}}^2 t^2) e^{-\frac{\sigma_{\text{ZF}}^2 t^2}{2}}]$  [42,43]. The solid lines in Fig. 13 represent fits to the data by considering an additional zero-field Lorentzian relaxation  $\Lambda$ , i.e.,  $A_{\text{ZF}} = A_s G_{\text{KT}} e^{-\Lambda t} + A_{\text{bg}}$ . Here  $A_s$  and  $A_{\text{bg}}$  are the same as in the TF- $\mu$ SR case [see Eq. (3)]. The strong Gaussian relaxation rates reflect the large nuclear moments in  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ , mostly determined by the Nb nuclear moments. In both the normal and the superconducting states, the relaxations are almost identical, as demonstrated by the practically overlapping ZF- $\mu$ SR spectra above and below  $T_c$ . This lack of evidence for an additional  $\mu$ SR relaxation below  $T_c$  excludes a possible time-reversal symmetry breaking in the superconducting state of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ .

### IV. CONCLUSION

To summarize, we investigated the superconducting properties of  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  (for  $x = 0$  and 1.6) by means of electrical resistivity, magnetization, specific heat, and  $\mu$ SR measurements.  $\text{Nb}_5\text{Ir}_3\text{O}$  and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  exhibit bulk SC with  $T_c$  of 10.5 and 9.1 K, respectively. The temperature dependence of the zero-field electronic specific heat and superfluid density reveal a nodeless SC in  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$ , well described by an isotropic  $s$ -wave model.  $\text{Nb}_5\text{Ir}_3\text{O}$ , instead, turns out to be a multigap superconductor, as demonstrated by the field dependence of the electronic specific heat coefficient, the superconducting Gaussian relaxation, and the temperature dependence of its upper critical field. Upon Pt substitution, the small superconducting gap is suppressed, and  $\text{Nb}_5\text{Ir}_{1.4}\text{Pt}_{1.6}\text{O}$  shows typical features of single-gap SC, hence indicating a crossover from multiple- to single-gap SC in the  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  family. Finally, the lack of spontaneous magnetic fields below  $T_c$  indicates that the time-reversal symmetry is preserved in  $\text{Nb}_5\text{Ir}_{3-x}\text{Pt}_x\text{O}$  superconductors.

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