Selective hybridization between the main band and the superstructure band in the Bi₂Sr₂CaCu₂O₈₊δ superconductor

Qiang Gao,[1,2] Hongtao Yan,[1,2] Jing Liu,[1,2,3] Ping Ai,[1,2] Yongqing Cai,[1,2] Cong Li,[1,2] Xiangyu Luo,[1,2] Cheng Hu,[1,2] Chunyao Song,[1,2] Jianwei Huang,[1,2] Hongtao Rong,[1,2] Yuan Huang ,[1] Qingyan Wang ,[1] Guodong Liu,[1,2] Genda Gu,[4] Fengfeng Zhang,[5] Feng Yang,[5] Shenjin Zhang,[5] Qinjun Peng,[5] Zuyan Xu,[5] Lin Zhao,[1] Tao Xiang,[1,2,3,6] and X. J. Zhou [1,2,3,6,*]

[1]*Beijing National Laboratory for Condensed Matter Physics, Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China*
[2]*School of Physical Sciences, University of Chinese Academy of Sciences, Beijing 100049, China*
[3]*Beijing Academy of Quantum Information Sciences, Beijing 100193, China*
[4]*Condensed Matter Physics and Materials Science Department, Brookhaven National Laboratory, Upton, New York 11973, USA*
[5]*Technical Institute of Physics and Chemistry, Chinese Academy of Sciences, Beijing 100190, China*
[6]*Songshan Lake Materials Laboratory, Dongguan, Guangdong 523808, China*

High-resolution laser-based angle-resolved photoemission measurements have been carried out on Bi₂Sr₂CaCu₂O₈₊δ (Bi2212) and Bi₂Sr₂₋ₓLaₓCuO₆₊δ (Bi2201) superconductors. Unexpected hybridization between the main band and the superstructure band in Bi2212 is clearly revealed. In the momentum space where one main Fermi surface intersects with one superstructure Fermi surface, four bands are observed instead of two. The hybridization exists in both the superconducting state and the normal state, and in Bi2212 samples with different doping levels. Such a hybridization is not observed in Bi2201. This phenomenon can be understood by considering the bilayer splitting in Bi2212, the selective hybridization of two bands with peculiar combinations, and the altered matrix element effects of the hybridized bands. These observations indicate that the origin of the superstructure bands is intrinsic to the CuO₂ planes. Therefore, understanding physical properties and superconductivity in Bi2212 should consider the complete Fermi surface topology which involves the main bands, the superstructure bands, and their interactions.

I. INTRODUCTION

High-temperature cuprate superconductors have been extensively studied for more than 30 years due to its unusually high critical temperature (T_C), anomalous normal state, and challenging mechanism of high temperature superconductivity [1–5]. Angle-resolved photoemission spectroscopy (ARPES) [3,6,7] has played a key role in studying the electronic structure of the cuprate superconductors, including the revelation of the distinct d-wave superconducting gap symmetry [8–10], the pseudogap above T_C [11–14], and many-body effects [15–26]. The majority of these significant ARPES results are obtained from Bi₂Sr₂CaCu₂O₈₊δ (Bi2212) because of the availability of high-quality single crystals and its easiness to cleave to get smooth and clean surfaces. Bi2212 is also widely used in other experimental techniques like scanning tunneling microscopy and scanning tunneling spectroscopy [27–31].

It is well-known that the bismuth-based cuprate superconductors have incommensurate modulations in their crystal structure along the b^* direction [32–36]. This modulation leads to the formation of superstructure bands of various orders in the measured ARPES results [37–42]. A related issue under debate is about the origin of the superstructure

bands whether they are extrinsic due to diffraction of the BiO layers or intrinsic coming directly from the CuO₂ planes [38,39,43]. Another feature of Bi2212 is that there are two CuO₂ planes in one structural unit separated by calcium. The interaction between these two structurally equivalent CuO₂ planes gives rise to bilayer splitting, i.e., two Fermi surface sheets, bonding and antibonding, which correspond to different doping levels with a distinct superconducting gap [44–52]. While most ARPES measurements focus on the main band, little attention has been paid to the interaction between the main and superstructure bands and between the bonding and antibonding bands. These interactions may provide important information on the origin of superstructure bands and adequate consideration of Fermi surface topology in studying the physical properties and superconductivity in Bi2212.

In this paper, we report the observation of an unexpected hybridization between the main band and the superstructure band in Bi2212 due to the interaction between the bonding and antibonding bands. Taking advantage of high-resolution laser-based ARPES measurements, we observed four bands, instead of two, in the momentum space where one main band intersects with one superstructure band. Such an unusual hybridization is present in both the normal state and the superconducting state and in Bi2212 samples with different doping levels. This hybridization phenomenon is not observed in similar ARPES measurements on Bi₂Sr₂₋ₓLaₓCuO₆₊δ (Bi2201) excluding the possibility of its spin-related origin. We propose

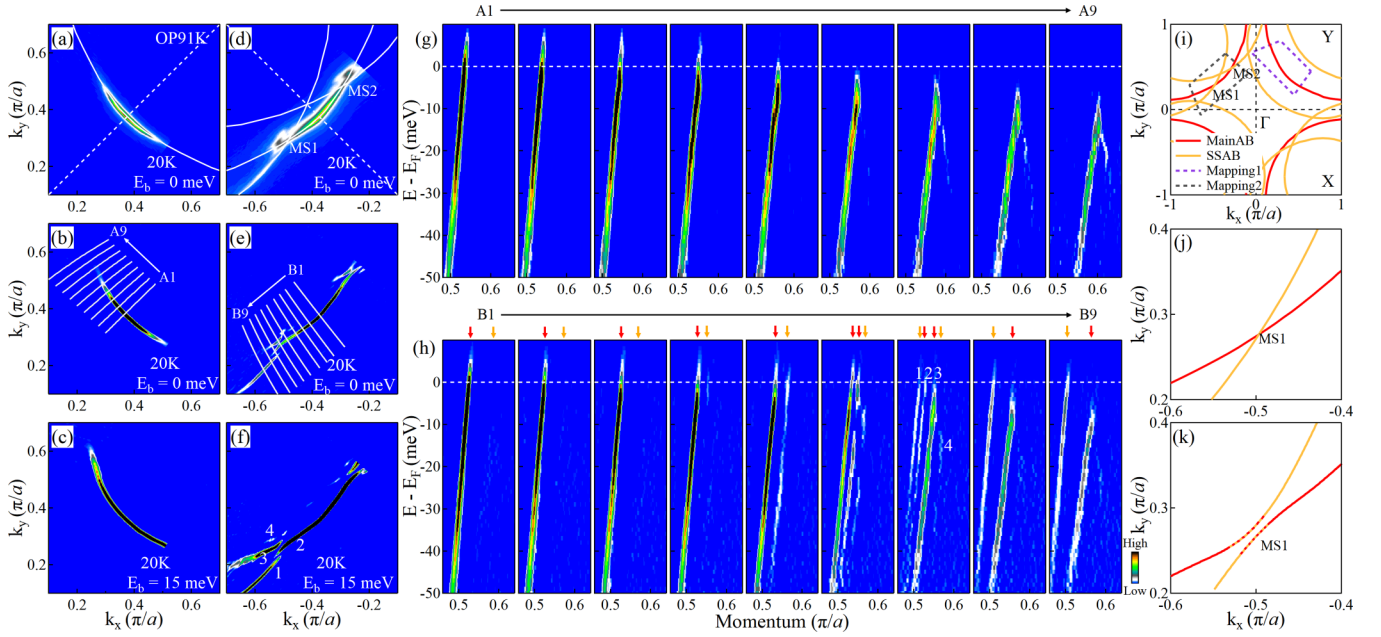*Corresponding author: XJZhou@iphy.ac.cn

FIG. 1. Fermi surface and band structure of Bi2212 (OP91K) measured at a temperature of 20 K by laser-based ARPES. (a) Fermi surface mapping in the first quadrant of the Brillouin zone. It is obtained by integrating the spectral weight over an energy window of $[-2, +2]$ meV with respect to the Fermi level. (b) The second derivative of the original image (a) with respect to momentum. The second derivative image helps enhance the contrast to reveal some detailed structures. (c) The second derivative image of the constant energy contour obtained by integrating the spectral weight over an energy window of $[-3, +3]$ meV with respect to a binding energy of 15 meV. (d–f) The same as panels (a)–(c) but measured in the second quadrant of the Brillouin zone. (g) Band structure along the different momentum cuts A1–A9 in the first quadrant. The location of the momentum cuts is shown by white lines in panel (b). The images are second derivative of the original band structure with respect to momentum. (h) Band structure along the different momentum cuts B1–B9 in the second quadrant. The location of the momentum cuts is shown by white lines in panel (e). The images are second derivative of the original band structure with respect to momentum. The red and yellow arrows mark the locations of the observed bands. (i) Schematic Fermi surface of Bi2212 containing the main antibonding Fermi surface (MainAB, red lines) and its first-order superstructure replicas (SSAB, yellow lines). The dashed purple box represents the mapping area of panel (a), and the dashed gray box represents the mapping area of panel (d). (j) The enlarged momentum area where the main band and the superstructure band intersect at the MS1 point. (k) Same as panel (j) by assuming hybridization of two bands.

this unusual band hybridization in Bi2212 can be understood by considering the bilayer splitting effect. We find that in this case only the bonding and antibonding bands can hybridize and the resultant hybridized bands exhibit distinct matrix element effects. These observations provide key insights on the origin of the superstructure band that is intrinsic to the CuO_2 planes. They also provide a new overall picture of the Fermi surface topology of Bi2212 that needs to be taken into account in studying its electronic structure and physical properties.

II. EXPERIMENT

ARPES measurements were carried out on a vacuum ultraviolet laser-based ARPES system [53,54]. The photon energy is 6.994 eV with a bandwidth of 0.26 meV. The energy resolution of the electron energy analyzer (Scienta DA30L) was set at 1 meV, giving rise to an overall energy resolution of 1.0 meV. The angular resolution was ∼0.3°, corresponding to a momentum resolution of 0.004 Å^{-1}. The Fermi level is referenced by measuring on the Fermi edge of a clean polycrystalline gold that is electrically connected to the sample. High-quality optimally doped $Bi_2Sr_{2-x}La_xCuO_{6+\delta}$ ($x = 0.4$, $T_C = 32$ K, denoted as OP32K hereafter) [55] and $Bi_2Sr_2CaCu_2O_{8+\delta}$ ($T_C = 91$ K, denoted as OP91K hereafter)

single crystals were grown by the floating zone method. Overdoped Bi2212 sample ($T_C = 78$ K, denoted as OD78K hereafter) was prepared by annealing the optimally doped sample in a high-pressure oxygen atmosphere [56]. Underdoped Bi2212 sample ($T_C = 64$ K, denoted as UD64K hereafter) was prepared by annealing the optimally doped sample in vacuum. The T_C was measured using a Quantum Design SQUID magnetometer. All samples were cleaved *in situ* and measured in vacuum with a base pressure better than 3×10^{-11} mbar.

III. RESULTS AND DISCUSSIONS

Figure 1 shows the Fermi surface and the corresponding band structure of the Bi2212 OP91K sample. Two momentum spaces are covered, one in the first quadrant and the other in the second quadrant, as indicated in Fig. 1(i) by the dashed box areas. The Fermi surface measured in the first quadrant [Figs. 1(a) and 1(b)] consists of only one sheet, as also shown by the measured band structures in Fig. 1(g). This Fermi surface corresponds to the antibonding sheet when measured by using 6.994 eV laser [23]; the bonding sheet is fully suppressed at this measurement condition. In the second quadrant, the main Fermi surface is expected to intersect with two first-order superstructure replicas at the MS1 and MS2
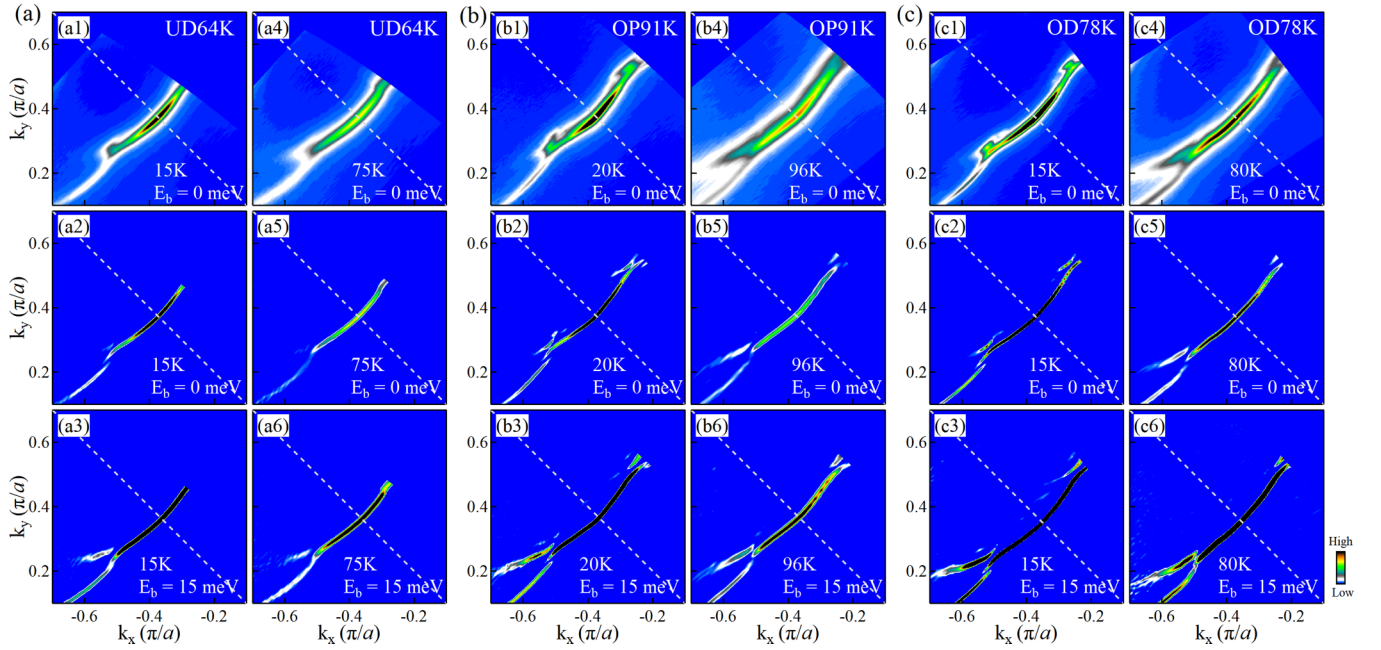
FIG. 2. Doping and temperature dependencies of Fermi surface and band hybridization in Bi2212. (a) Fermi surface and constant energy contour of underdoped Bi2212 with $T_C = 64$ K (UD64K) measured in the superconducting state at 15 K (a1–a3) and the normal state at 75 K (a4–a6). Panels (a2), (a3), (a5), and (a6) are obtained by the second derivative of the original data with respect to momentum. Panels (a1), (a2), (a4), and (a5) represent the measured Fermi surface, while panels (a3) and (a6) are constant energy contours at a binding energy of 15 meV. (b) Same as panel (a) but the Fermi surface and constant energy contour are measured on an optimally doped Bi2212 with $T_C = 91$ K (OP91K) in the superconducting state (20 K) and the normal state (96 K). (c) Same as panel (a) but the Fermi surface and constant energy contour are measured on an overdoped Bi2212 with $T_C = 78$ K (OD78K) in the superconducting state (15 K) and the normal state (80 K).

points, as indicated in Fig. 1(i). However, the measured Fermi surface in the second quadrant [Figs. 1(d) and 1(e)] shows rather complex topology near the crossing areas, MS1 and MS2, of the main and superstructure Fermi surface sheets. First, both the main Fermi surface and the superstructure Fermi surface get broken and become disconnected near the MS1 crossing point. Three branches of Fermi surface can be clearly seen near the crossing point, although the main Fermi surface and the superstructure Fermi surface still consist of only one band away from the crossing point. In the constant energy contour at a binding energy of 15 meV [Fig. 1(f)], these features show up more clearly by reducing the effect of the superconducting gap, and four branches of contours can be identified near the MS1 crossing point. Here branch 1 comes from the superstructure band, and branches 2 and 3 come from the main band broken at the crossing point. But the weak branch, branch 4, appears unexpectedly and cannot be attributed to either the main band or the superstructure band.

When the momentum cuts move from the nodal direction to the antinodal direction [B1 to B9 cuts in Fig. 1(e)], the measured band structure [Fig. 1(h)] changes from two bands for the B1–B5 cuts, to three bands for the B6 cut, to four bands for the B7 cut, and finally back to two bands for the cuts B8 and B9, as indicated by the arrows on the top of Fig. 1(h). The band structure evolution in Fig. 1(h) agrees with the measured Fermi surface in Fig. 1(e) and the constant energy contour in Fig. 1(f). In particular, the simultaneous observation of four bands in the cut B7 corresponds to the four branches of contours in Fig. 1(f). We note that band 4 in the cut B7 panel of Fig. 1(h) gives rise to branch 4

in Fig. 1(f); it is not due to the superconductivity-induced back bending of band 3 because it is parallel to band 3 [57]. The bands caused by superconductivity-induced back bending exhibit different slopes from the original band, as seen in the panels for the cuts A7–A9 in Fig. 1(g). The two bands observed away from the crossing point, for the cuts B1-B5, B8, and B9, are well understood as from one main Fermi surface and one superstructure Fermi surface as depicted in Fig. 1(d). However, the simultaneous observation of four branches of Fermi surface sheets [Fig. 1(f)] and four bands [Fig. 1(h)] B7 near the crossing point MS1 is rather different from the expected pictures where only two bands can be observed; either there is no hybridization [Fig. 1(j)] or there is hybridization [Fig. 1(k)] between them.

To check the effect of superconductivity on this unusual band hybridization and its doping dependence, we measured Bi2212 samples with different doping levels in both the normal state and the superconducting state. Figure 2 shows the Fermi surface mappings and the constant energy contours at a binding energy of 15 meV for the underdoped UD64K [Fig. 2(a)], the optimally doped OP91K [Fig. 2(b)], and the overdoped OD78K [Fig. 2(c)] samples both above and below their corresponding superconducting transition temperatures. The measured data in the normal state show behaviors quite similar to that measured in the superconducting state; i.e., the Fermi surface topology near the crossing point is characterized by a broken Fermi surface and the appearance of multiple branches. Some fine structures, in particular the appearance of branch 4, become less clear in the normal state data possibly due to thermal broadening at relatively high temperature. This
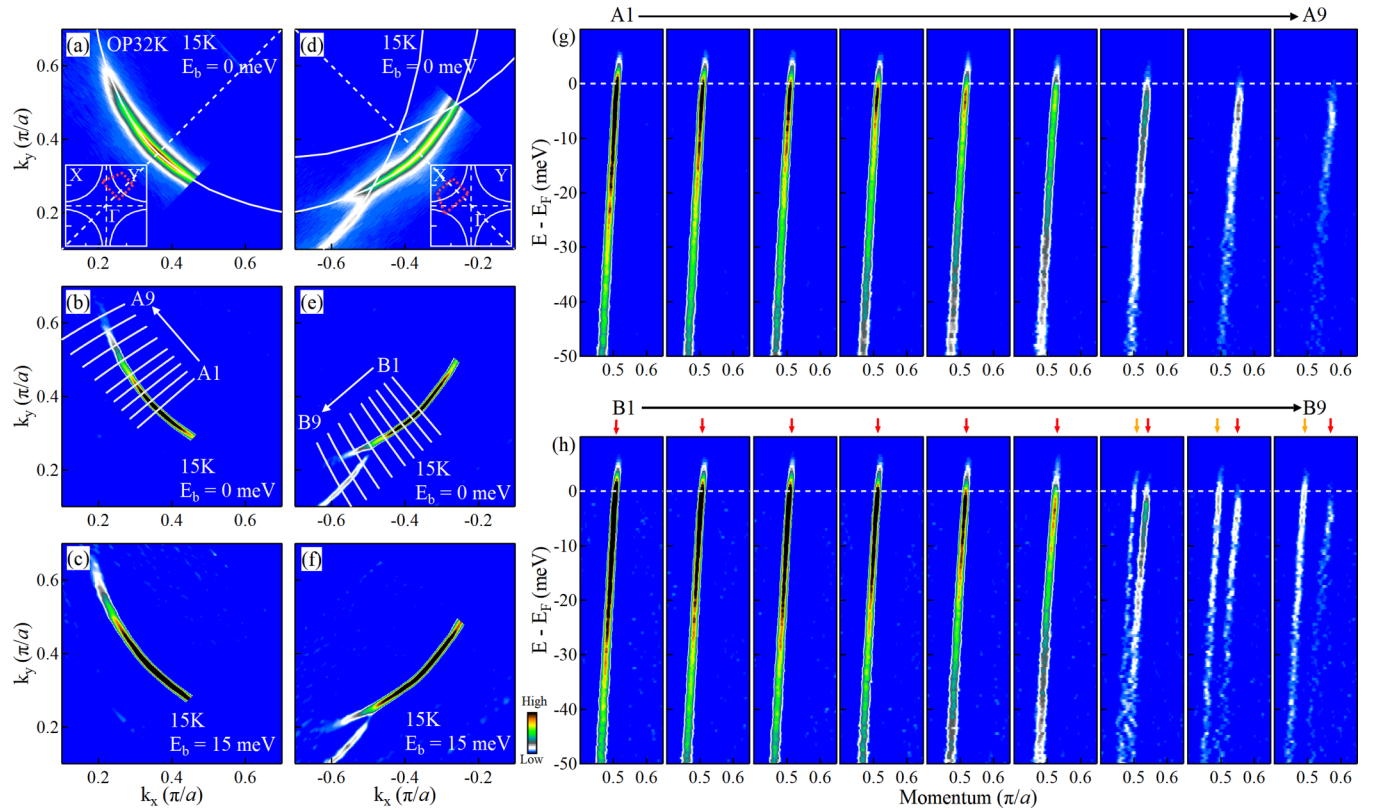
FIG. 3. Fermi surface and band structure of Bi2201 (OP32K) measured at a temperature of 20 K by laser-based ARPES. (a) Fermi surface mapping in the first quadrant of the Brillouin zone. It is obtained by integrating the spectral weight over an energy window of $[-2, +2]$ meV with respect to the Fermi level. The corresponding momentum area is marked as the dashed red box in the inset. (b) The second derivative of the original data corresponding to panel (a) with respect to momentum. (c) Constant energy contour at a binding energy of 15 meV. The image is obtained by the second derivative of the original data with respect to momentum. (d) Fermi surface mapping in the second quadrant of the Brillouin zone at 15 K. The corresponding momentum area is marked as the dashed red box in the inset. (e) The second derivative of the original data corresponding to panel (d) with respect to momentum. (f) Constant energy contour at a binding energy of 15 meV. The image is obtained by the second derivative of the original data with respect to momentum. Panels (g) and (h) show the band structure measured at 15 K. The location of the corresponding momentum cuts is shown as white lines in panels (b) and (e) labeled with A1–A9 and B1–B9. The images are obtained by the second derivative of the original data with respect to momentum. The red and yellow arrows in panel (h) mark the locations of the observed bands.

will not affect the result that unusual band hybridization exists both above and below T_C, and therefore it is not caused by superconductivity. The three samples with different doping levels show quite similar behaviors near the crossing area of the main band and the superstructure band. Here some fine structures in the measured data of the underdoped UD64K sample [Fig. 2(a)] are less clear than those in the optimally doped and overdoped samples, which may be related to photoemission spectrum broadening by strong electron scattering in the underdoped samples [58–61]. The key features of the unusual band hybridization, broken Fermi surfaces and multiple branches, remain in the underdoped sample. Therefore, the unusual band hybridization between the main and the superstructure bands is a general phenomenon in Bi2212 with different doping levels.

Since superstructure modulation is present in all the bismuth-based superconductors [32–35], it is natural to ask whether similar band hybridization phenomenon can also occur in other systems besides Bi2212. To this end, we performed similar ARPES measurements on Bi2201. The measured Fermi surface in two quadrants and the corresponding

band structure are shown in Fig. 3. In the first quadrant, a single Fermi surface sheet is clearly observed [Fig. 3(a)]. The corresponding band structure consists of a single band for the momentum cuts from nodal to antinodal regions [Fig. 3(g)]. In the second quadrant, both the main Fermi surface and the superstructure replica are clearly observed. The main Fermi surface is continuous and there is no additional Fermi surface branch in the crossing area of the main and superstructure bands [Fig. 3(d)]. There appears to be only the main band and the superstructure band in the measured band structure for the momentum cuts from the nodal to the antinodal directions and no sign of any additional bands is observed [Fig. 3(h)]. These results can be well understood by the picture in Fig. 1(j) where the main band and the superstructure band cross without hybridization. The unusual band hybridization observed in Bi2212 does not occur in Bi2201, although in both cases only one main Fermi surface is observed.

Now we come to discuss the possible origin of the peculiar band hybridization discovered in Bi2212. When two bands hybridize with each other, they usually still give rise to two avoided bands as shown in Fig. 1(k). When we first observed
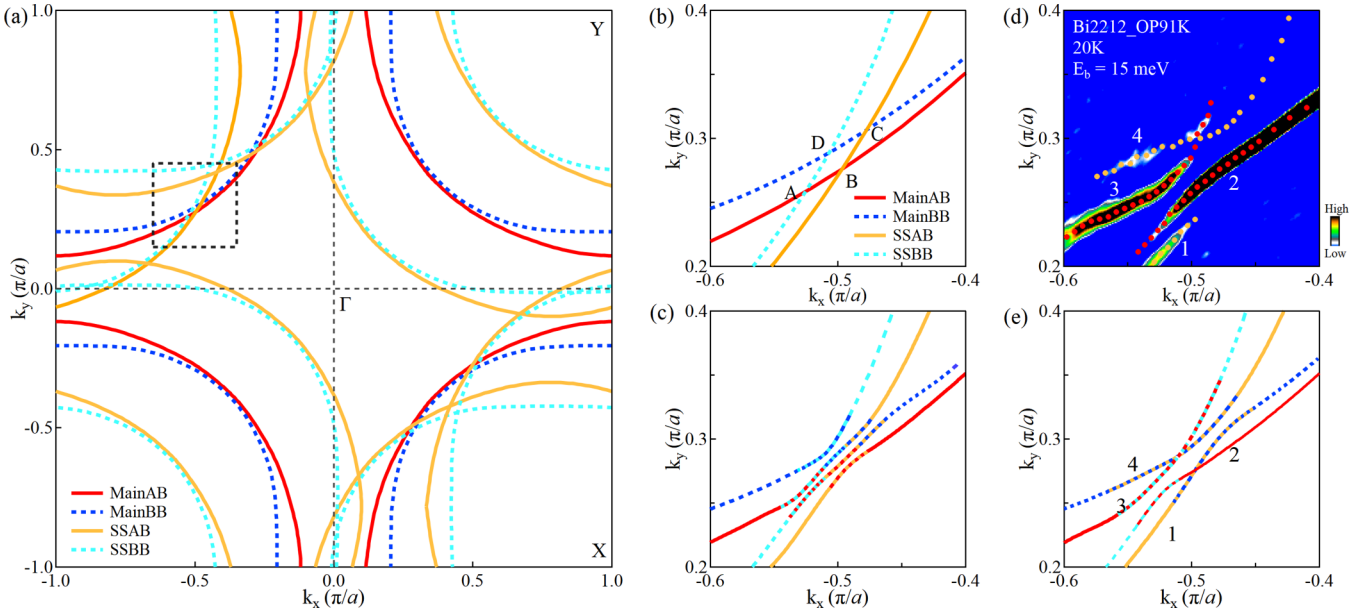
FIG. 4. Schematic illustration of band hybridization and its comparison with the measured results in Bi2212. (a) Schematic Fermi surface of Bi2212 containing the main Fermi surface with bilayer splitting: Antibonding sheet (MainAB, red lines) and bonding sheet (MainBB, dashed blue lines); and first-order superstructure replicas: Antibonding sheet (SSAB, yellow lines) and bonding sheet (SSBB, dashed green lines). The dashed gray box marks the momentum area enlarged in panels (b)–(e). (b) The enlarged momentum area marked in panel (a). The main bonding and antibonding Fermi surface sheets intersect with the superstructure bonding and antibonding replicas that give rise to four crossing points labeled as A, B, C, and D. (c) Same as panel (b) but considering band hybridization at all four crossing points. (d) The measured constant energy contour at a binding energy of 15 meV at the crossing point MS1 as shown in Fig. 1(f). The observed four branches of sheets are marked. (e) Same as panel (b) but considering the band hybridization only at A and C crossing points.

two hybridized bands split into four bands, the first possibility coming to mind was whether this was a spin-related phenomenon. Special momentum-locked spin texture has been reported in the Bi2212 superconductor [62]. If spin-up and spin-down bands split during the band hybridization between the main band and the superstructure band, it will naturally produce four bands as we have observed. We believe this scenario is less likely. First, it remains to be seen whether the observed spin texture in Bi2212 can cause the splitting of the spin-up and spin-down bands. Second, if we assume the spin texture observed in Bi2212 is general in cuprate superconductors, the absence of the unusual band hybridization in Bi2201 is not consistent with the spin-related scenario.

The observation of the unusual band hybridization in Bi2212 and its absence in Bi2201 indicates that it is a Bi2212-specific phenomenon. Compared with Bi2201, the unique features of Bi2212 lie in the presence of two CuO_2 planes in one structural unit that gives rise to bilayer splitting [44–52]. Although there is only one main Fermi surface that is observed during our measurements, there are actually two Fermi surface sheets, bonding and antibonding, that are present. This prompted us to examine whether this unusual hybridization could be related to the bilayer splitting in Bi2212. Figure 4(a) shows the Fermi surface topology of Bi2212 by considering both the bilayer splitting and the superstructure bands. In the same momentum area as covered by our experiment, two main Fermi surface sheets are expected to cross with two superstructure Fermi surface sheets in the second quadrant, giving rise to four crossing points (A, B, C, and D) as shown in Fig. 4(b). If there is no band hybridization at the four crossing

points, it is obviously not consistent with the measured result shown in Fig. 4(d). When we consider band hybridization on all the four crossing points in Fig. 4(c), it does not agree with the measured result in Fig. 4(d) either. A consistent picture emerges when we consider the following.

(i) Only the bonding band and the antibonding band can hybridize. This applies to the crossing point A where the main antibonding band (MainAB) intersects with the superstructure bonding band (SSBB) and the crossing point C where the main bonding band (MainBB) intersects with the superstructure antibonding band (SSAB).

(ii) The bonding band does not hybridize with the bonding band like the crossing point D. The antibonding band does not hybridize with the antibonding band like the crossing point B.

(iii) Under our present ARPES measurement condition, only the antibonding band is observed while the bonding band is fully suppressed.

The hybridized bands show up because they are mixed with both bonding and antibonding characters. The resultant Fermi surface picture [Fig. 4(e)] agrees well with the observed results [Fig. 4(d)]. In particular, the simultaneous observation of four bands in the crossing area can be well understood in this picture.

The present work provides a unique opportunity to look into the interaction between the bonding and the antibonding bands in Bi2212. Our results indicate that only the bonding and the antibonding bands can hybridize with each other. The underlying mechanism of this selective hybridization is interesting and asks for further theoretical understanding. This band interaction rule can also explain why there is no band

hybridization between the main and the superstructure bands in Bi2201.

The origin of superstructure bands in the bismuth-based superconductors in which the superstructure modulation is present has been under debate [38,39,43]. In one scenario, the superstructure bands that are formed when the photoelectrons from the CuO_2 plane are diffracted in passing through the BiO layer with superstructure modulation are extrinsic. In the other scenario, those that come from the CuO_2 planes directly are intrinsic. Our present work provides strong evidence in favor of their intrinsic nature. If the superstructure band originates from photoelectron diffraction, it is impossible for it to hybridize with the main band. Both the main band and the superstructure band, as well as their hybridization, come directly from the CuO_2 plane. This is consistent with the observation of superstructure modulation in CuO_2 planes [36,43,63]. We note that a similar situation occurs in graphene where several new replicas of Dirac cones are observed which are attributed to the underlying weak modulation of the ionic potential by the substrate experienced by the electrons in the graphene [64]. The new features are intrinsic and important in understanding the related transport properties [64]. Likewise, the intrinsic nature of the main bands, the superstructure bands, and their hybridization provides a complete Fermi surface topology for Bi2212 that should be considered in understanding its electron structure, physical properties, and high-temperature superconductivity.

IV. CONCLUSIONS

In summary, by taking high-resolution laser-based ARPES measurements on Bi2212, we have discovered an unexpected band hybridization between the main band and the superstructure band. Such a hybridization is observed in Bi2212 with different dopings and in both the normal and the superconducting states. It is not present in Bi2201. This phenomenon can be understood by considering the bilayer splitting and the selective interaction of bonding and antibonding bands in Bi2212. Our results provide strong evidence to support the intrinsic nature of the superstructure bands. They also provide a complete Fermi surface topology of Bi2212 that needs to be considered in theoretical and experimental understanding of its physical properties.

[1] C. C. Tsuei and J. R. Kirtley, Pairing symmetry in cuprate superconductors, Rev. Mod. Phys. **72**, 969 (2000).

[2] T. Timusk and B. Statt, The pseudogap in high-temperature superconductors: An experimental survey, Rep. Prog. Phys. **62**, 61 (1999).

[3] A. Damascelli, Z. Hussain, and Z. X. Shen, Angle-resolved photoemission studies of the cuprate superconductors, Rev. Mod. Phys. **75**, 473 (2003).

[4] P. A. Lee, N. Nagaosa, and X. G. Wen, Doping a Mott insulator: Physics of high-temperature superconductivity, Rev. Mod. Phys. **78**, 17 (2006).

[5] B. Keimer, S. A. Kivelson, M. R. Norman, S. Uchida, and J. Zaanen, From quantum matter to high-temperature superconductivity in copper oxides, Nature (London) **518**, 179 (2015).

[6] J. C. Campuzano, M. R. Norman, and M. Randeria, in *The Physics of Superconductors*, edited by K. H. Bennemann and J. B. Ketterson (Springer, New York, 2004), Vol. 2.

[7] X. J. Zhou, T. Cuk, T. Devereaux, N. Nagaosa, and Z. X. Shen, in *Handbook of High-Temperature Superconductivity: Theory and Experiment*, edited by J. R. Schrieffer (Springer, New York, 2007).

[8] Z. X. Shen, D. S. Dessau, B. O. Wells, D. M. King, W. E. Spicer, A. J. Arko, D. Marshall, L. W. Lombardo, A. Kapitulnik, P. Dickinson, S. Doniach, J. DiCarlo, A. G. Loeser, and C. H. Park, Anomalously Large Gap Anisotropy in the *a-b* Plane of $Bi_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. Lett. **70**, 1553 (1993).

[9] H. Ding, M. R. Norman, J. C. Campuzano, M. Randeria, A. F. Bellman, T. Yokoya, T. Takahashi, T. Mochiku, and K. Kadowaki, Angle-resolved photoemission spectroscopy study of the superconducting gap anisotropy in $Bi_2Sr_2CaCu_2O_{8+x}$, Phys. Rev. B **54**, R9678 (1996).

[10] M. Hashimoto, E. A. Nowadnick, R. H. He, I. M. Vishik, B. Moritz, Y. He, K. Tanaka, R. G. Moore, D. H. Lu, Y. Yoshida, M. Ishikado, T. Sasagawa, K. Fujita, S. Ishida, S. Uchida, H. Eisaki, Z. Hussain, T. P. Devereaux, and Z. X. Shen, Direct spectroscopic evidence for phase competition between the pseudogap and superconductivity in $Bi_2Sr_2CaCu_2O_{8+\delta}$, Nat. Mater. **14**, 37 (2015).

[11] D. S. Marshall, D. S. Dessau, A. G. Loeser, C.-H. Park, A. Y. Matsuura, J. N. Eckstein, I. Bozovic, P. Fournier, A. Kapitulnik, W. E. Spicer, and Z. X. Shen, Unconventional Electronic Structure Evolution with Hole Doping in $Bi_2Sr_2CaCu_2O_{8+\delta}$: Angle-Resolved Photoemission Results, Phys. Rev. Lett. **76**, 4841 (1996).

[12] A. G. Loeser, Z. X. Shen, D. S. Dessau, D. S. Marshall, C. H. Park, P. Fournier, and A. Kapitulnik, Excitation gap in the normal state of underdoped $Bi_2Sr_2CaCu_2O_{8+\delta}$, Science **273**, 325 (1996).

[13] H. Ding, T. Yokoya, J. C. Campuzano, T. Takahashi, M. Randeria, M. R. Norman, T. Mochiku, K. Kadowaki, and J. Giapintzakis, Spectroscopic evidence for a pseudogap in the normal state of underdoped high-T_C superconductors, Nature (London) **382**, 51 (1996).

[14] I. M. Vishik, Photoemission perspective on pseudogap, superconducting fluctuations, and charge order in cuprates: A review of recent progress, Rep. Prog. Phys. **81**, 062501 (2018).

[15] P. V. Bogdanov, A. Lanzara, S. A. Kellar, X. J. Zhou, E. D. Lu, W. J. Zheng, G. Gu, J. I. Shimoyama, K. Kishio, H. Ikeda, R. Yoshizaki, Z. Hussain, and Z. X. Shen, Evidence for an Energy Scale for Quasiparticle Dispersion in $Bi_2Sr_2CaCu_2O_8$, Phys. Rev. Lett. **85**, 2581 (2000).

[16] P. D. Johnson, T. Valla, A. V. Fedorov, Z. Yusof, B. O. Wells, Q. Li, A. R. Moodenbaugh, G. D. Gu, N. Koshizuka, C. Kendziora, Sha Jian, and D. G. Hinks, Doping and Temperature Dependence of the Mass Enhancement Observed in the Cuprate $Bi_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. Lett. **87**, 177007 (2001).

[17] A. Kaminski, M. Randeria, J. C. Campuzano, M. R. Norman, H. Fretwell, J. Mesot, T. Sato, T. Takahashi, and K. Kadowaki, Renormalization of Spectral Line Shape and Dispersion below T_C in $Bi_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. Lett. **86**, 1070 (2001).

[18] A. Lanzara, P. V. Bogdanov, X. J. Zhou, S. A. Kellar, D. L. Feng, E. D. Lu, T. Yoshida, H. Eisaki, A. Fujimori, K. Kishio, J. I. Shimoyama, T. Noda, S. Uchida, Z. Hussain, and Z. X. Shen, Evidence for ubiquitous strong electron-phonon coupling in high-temperature superconductors, Nature (London) **412**, 510 (2001).

[19] X. J. Zhou, T. Yoshida, A. Lanzara, P. V. Bogdanov, S. A. Kellar, K. M. Shen, W. L. Yang, F. Ronning, T. Sasagawa, T. Kakeshita, T. Noda, H. Eisaki, S. Uchida, C. T. Lin, F. Zhou, J. W. Xiong, W. X. Ti, Z. X. Zhao, A. Fujimori, Z. Hussain, and Z. X. Shen, Universal nodal fermi velocity, Nature (London) **423**, 398 (2003).

[20] A. D. Gromko, A. V. Fedorov, Y. D. Chuang, J. D. Koralek, Y. Aiura, Y. Yamaguchi, K. Oka, Y. Ando, and D. S. Dessau, Mass-renormalized electronic excitations at $(\pi, 0)$ in the superconducting state of $Bi_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. B **68**, 174520 (2003).

[21] T. K. Kim, A. A. Kordyuk, S. V. Borisenko, A. Koitzsch, M. Knupfer, H. Berger, and J. Fink, Doping Dependence of the Mass Enhancement in (Pb, Bi)$_2$Sr$_2$CaCu$_2$O$_8$ at the Antinodal Point in the Superconducting and Normal States, Phys. Rev. Lett. **91**, 167002 (2003).

[22] T. Cuk, F. Baumberger, D. H. Lu, N. Ingle, X. J. Zhou, H. Eisaki, N. Kaneko, Z. Hussain, T. P. Devereaux, N. Nagaosa, and Z. X. Shen, Coupling of the B_{1g} Phonon to the Antinodal Electronic States of $Bi_2Sr_2Ca_{0.92}Y_{0.08}Cu_2O_{8+\delta}$, Phys. Rev. Lett. **93**, 117003 (2004).

[23] W. T. Zhang, G. D. Liu, L. Zhao, H. Y. Liu, J. Q. Meng, X. L. Dong, W. Lu, J. S. Wen, Z. J. Xu, G. D. Gu, T. Sasagawa, G. L. Wang, Y. Zhu, H. B. Zhang, Y. Zhou, X. Y. Wang, Z. X. Zhao, C. T. Chen, Z. Y. Xu, and X. J. Zhou, Identification of a New Form of Electron Coupling in the $Bi_2Sr_2CaCu_2O_8$ Superconductor by Laser-Based Angle-Resolved Photoemission Spectroscopy, Phys. Rev. Lett. **100**, 107002 (2008).

[24] J. F. He, W. T. Zhang, J. M. Bok, D. X. Mou, L. Zhao, Y. Y. Peng, S. L. He, G. D. Liu, X. L. Dong, J. Zhang *et al.*, Coexistence of Two Sharp-Mode Couplings and Their Unusual Momentum Dependence in the Superconducting State of $Bi_2Sr_2CaCu_2O_{8+\delta}$ Revealed by Laser-Based Angle-Resolved Photoemission, Phys. Rev. Lett. **111**, 107005 (2013).

[25] J. M. Bok, J. J. Bae, H. Y. Choi, C. M. Varma, W. T. Zhang, J. F. He, Y. X. Zhang, L. Yu, and X. J. Zhou, Quantitative determination of pairing interactions for high-temperature superconductivity in cuprates, Sci. Adv. **2**, e1501329 (2016).

[26] Y. He, M. Hashimoto, D. Song, S. D. Chen, J. He, I. M. Vishik, B. Moritz, D. H. Lee, N. Nagaosa, J. Zaanen, T. P. Devereaux, Y. Yoshida, H. Eisaki, D. H. Lu, and Z. X. Shen, Rapid change of superconductivity and electron-phonon coupling through critical doping in Bi-2212, Science **362**, 62 (2018).

[27] S. H. Pan, J. P. O'Neal, R. L. Badzey, C. Chamon, H. Ding, J. R. Engelbrecht, Z. Wang, H. Eisaki, S. Uchida, A. K. Gupta, K. W. Ng, E. W. Hudson, K. M. Lang, and J. C. Davis, Microscopic electronic inhomogeneity in the high-T_C superconductor $Bi_2Sr_2CaCu_2O_{8+x}$, Nature (London) **413**, 282 (2001).

[28] J. E. Hoffman, E. W. Hudson, K. M. Lang, V. Madhavan, H. Eisaki, S. Uchida, and J. C. Davis, A four unit cell periodic pattern of quasi-particle states surrounding vortex cores in $Bi_2Sr_2CaCu_2O_{8+\delta}$, Science **295**, 466 (2002).

[29] Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, An intrinsic bond-centered electronic glass with unidirectional domains in underdoped cuprates, Science **315**, 1380 (2007).

[30] M. H. Hamidian, S. D. Edkins, S. H. Joo, A. Kostin, H. Eisaki, S. Uchida, M. J. Lawler, E. A. Kim, A. P. Mackenzie, K. Fujita, J. Lee, and J. C. S. Davis, Detection of a Cooper-pair density wave in $Bi_2Sr_2CaCu_2O_{8+x}$, Nature (London) **532**, 343 (2016).

[31] Ø Fischer, M. Kugler, I. M. Aprile, C. Berthod, and C. Renner, Scanning tunneling spectroscopy of high-temperature superconductors, Rev. Mod. Phys. **79**, 353 (2007).

[32] Y. Gao, P. Lee, P. Coppens, M. A. Subramanian, and A. W. Sleight, The incommensurate modulation of the 2212 Bi-Sr-Ca-Cu-O superconductor, Science **241**, 954 (1988).

[33] R. L. Withers, J. G. Thompson, L. R. Wallenberg, J. D. FitzGerald, J. S. Anderson, and B. G. Hyde, A transmission electron microscope and group theoretical study of the new Bi-based high-T_C superconductors and some closely related Aurivillius phases, J. Phys. C **21**, 6067 (1988).

[34] O. Eibl, Displacive modulation and chemical composition of (Bi, Pb)$_2$Sr$_2$Ca$_{n-1}$Cu$_n$O$_{2n+4}$ ($n = 2, 3$) high-T_C superconductors, Phys. C (Amsterdam, Neth.) **175**, 419 (1991).

[35] H. Heinrich, G. Kostorz, B. Heeb, and L. J. Gauckler, Modelling the atomic displacements in $Bi_2Sr_2Ca_{n-1}Cu_nO_x$ superconductors, Phys. C (Amsterdam, Neth.) **224**, 133 (1994).

[36] A. Bianconi, M. Lusignoli, N. L. Saini, P. Bordet, A. Kvick, and P. G. Radaelli, Stripe structure of the CuO_2 plane in $Bi_2Sr_2Ca_{n-1}Cu_nO_{8+y}$ by anomalous x-ray diffraction, Phys. Rev. B **54**, 4310 (1996).

[37] P. Aebi, J. Osterwalder, P. Schwaller, L. Schlapbach, M. Shimoda, T. Mochiku, and K. Kadowaki, Complete Fermi Surface Mapping of $Bi_2Sr_2CaCu_2O_{8+x}$(001): Coexistence of Short Range Antiferromagnetic Correlations and Metallicity in the Same Phase, Phys. Rev. Lett. **72**, 2757 (1994).

[38] J. Osterwalder, P. Aebi, P. Schwaller, L. Schlapbach, M. Shimoda, T. Mochiku, and K. Kadowaki, Angle-resolved photoemission experiments on $Bi_2Sr_2CaCu_2O_{8+\delta}$ (001): Effects of the incommensurate lattice modulation, Appl. Phys. A **60**, 247 (1995).

[39] H. Ding, A. F. Bellman, J. C. Campuzano, M. Randeria, M. R. Norman, T. Yokoya, T. Takahashi, H. K. Yoshida, T. Mochiku, K. Kadowaki, G. Jennings, and G. P. Brivio, Electronic Excitations in $Bi_2Sr_2CaCu_2O_8$: Fermi Surface, Dispersion, and Absence of Bilayer Splitting, Phys. Rev. Lett. **76**, 1533 (1996).

[40] J. Mesot, M. R. Norman, H. Ding, and J. C. Campuzano, Hot Spots on the Fermi Surface of $Bi_2Sr_2CaCu_2O_{8+\delta}$: Stripes versus Superstructure, Phys. Rev. Lett. **82**, 2618 (1999).

[41] N. L. Saini, A. Bianconi, A. Lanzara, J. Avila, M. C. Asensio, S. Tajima, G. D. Gu, and N. Koshizuka, Hot Spots on the Fermi Surface of $Bi_2Sr_2CaCu_2O_{8+\delta}$: Stripes versus Superstructure–Reply, Phys. Rev. Lett. **82**, 2619 (1999).

[42] H. M. Fretwell, A. Kaminski, J. Mesot, J. C. Campuzano, M. R. Norman, M. Randeria, T. Sato, R. Gatt, T. Takahashi, and K. Kadowaki, Fermi Surface of $Bi_2Sr_2CaCu_2O_8$, Phys. Rev. Lett. **84**, 4449 (2000).

[43] J. Liu, L. Zhao, Q. Gao, P. Ai, L. Zhang, T. Xie, J. W. Huang, Y. Ding, C. Hu, H. T. Yan *et al.*, Evolution of incommensurate superstructure and electronic structure with Pb substitution in $(Bi_{2-x}Pb_x)Sr_2CaCu_2O_{8+\delta}$ superconductors, Chin. Phys. B **28**, 077403 (2019).

[44] P. V. Bogdanov, A. Lanzara, X. J. Zhou, S. A. Kellar, D. L. Feng, E. D. Liu, H. Eisaki, J. I. Shimoyama, K. Kishio, Z. Hussain, and Z. X. Shen, Photoemission study of Pb doped $Bi_2Sr_2CaCu_2O_8$: A Fermi surface picture, Phys. Rev. B **64**, 180505(R) (2001).

[45] D. L. Feng, N. P. Armitage, D. H. Lu, A. Damascelli, J. P. Hu, P. Bogdanov, A. Lanzara, F. Ronning, K. M. Shen, H. Eisaki, C. Kim, Z. X. Shen, J.-I. Shimoyama, and K. Kishio, Bilayer Splitting in the Electronic Structure of Heavily Overdoped $Bi_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. Lett. **86**, 5550 (2001).

[46] Y. D. Chuang, A. D. Gromko, A. Fedorov, Y. Aiura, K. Oka, Y. Ando, H. Eisaki, S. I. Uchida, and D. S. Dessau, Doubling of the Bands in Overdoped $Bi_2Sr_2CaCu_2O_{8+x}$: Evidence for *c*-Axis Bilayer Coupling, Phys. Rev. Lett. **87**, 117002 (2001).

[47] S. V. Borisenko, A. A. Kordyuk, T. K. Kim, S. Legner, K. A. Nenkov, M. Knupfer, M. S. Golden, J. Fink, H. Berger, and R. Follath, Superconducting gap in the presence of bilayer splitting in underdoped $(Pb, Bi)_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. B **66**, 140509(R) (2002).

[48] Y. H. Su, J. Chang, H. T. Lu, H. G. Luo, and T. Xiang, Effect of bilayer coupling on tunneling conductance of double-layer high-T_C cuprates, Phys. Rev. B **68**, 212501 (2003).

[49] S. V. Borisenko, A. A. Kordyuk, S. Legner, T. K. Kim, M. Knupfer, C. M. Schneider, J. Fink, M. S. Golden, M. Sing, R. Claessen, A. Yaresko, H. Berger, C. Grazioli, and S. Turchini, Circular dichroism and bilayer splitting in the normal state of underdoped $(Pb, Bi)_2Sr_2(Ca_{1-x}Y_x)Cu_2O_{8+\delta}$ and overdoped $(Pb, Bi)_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. B **69**, 224509 (2004).

[50] Y. D. Chuang, A. D. Gromko, A. V. Fedorov, Y. Aiura, K. Oka, Y. Ando, M. Lindroos, R. S. Markiewicz, A. Bansil, and D. S. Dessau, Bilayer splitting and coherence effects in optimal and underdoped $Bi_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. B **69**, 094515 (2004).

[51] A. A. Kordyuk, S. V. Borisenko, A. N. Yaresko, S. L. Drechsler, H. Rosner, T. K. Kim, A. Koitzsch, K. A. Nenkov, M. Knupfer, J. Fink, R. Follath, H. Berger, B. Keimer, S. Ono, and Y. Ando, Evidence for CuO conducting band splitting in the nodal direction of $Bi_2Sr_2CaCu_2O_{8+\delta}$, Phys. Rev. B **70**, 214525 (2004).

[52] P. Ai, Q. Gao, J. Liu, Y. X. Zhang, C. Li, J. W. Huang, C. Y. Song, H. T. Yan, L. Zhao, G. D. Liu, G. D. Gu, F. F. Zhang, F. Yang, Q. J. Peng, Z. Y. Xu, and X. J. Zhou, Distinct superconducting gap on two bilayer-split fermi surface sheets in $Bi_2Sr_2CaCu_2O_{8+\delta}$ superconductor, Chin. Phys. Lett. **36**, 067402 (2019).

[53] G. D. Liu, G. L. Wang, Y. Zhu, H. B. Zhang, G. C. Zhang, X. Y. Wang, Y. Zhou, W. T. Zhang, H. Y. Liu, L. Zhao, J. Q. Meng, X. L. Dong, C. T. Chen, Z. Y. Xu, and X. J. Zhou, Development of a vacuum ultraviolet laser-based angle-resolved photoemission system with a superhigh energy resolution better than 1 meV, Rev. Sci. Instrum. **79**, 023105 (2018).

[54] X. J. Zhou, S. L. He, G. D. Liu, L. Zhao, L. Yu, and W. T. Zhang, New developments in laser-based photoemission spectroscopy and its scientific applications: A key issues review, Rep. Prog. Phys. **81**, 062101 (2018).

[55] J. Q. Meng, G. D. Liu, W. T. Zhang, L. Zhao, H. Y. Liu, W. Lu, X. L. Dong, and X. J. Zhou, Growth, characterization and physical properties of high-quality large single crystals of $Bi_2(Sr_{2-x}La_x)CuO_{6+\delta}$ high-temperature superconductors, Supercond. Sci. Technol. **22**, 045010 (2009).

[56] Y. X. Zhang, L. Zhao, G. D. Gu, X. J. Zhou, A reproducible approach of preparing high-quality overdoped $Bi_2Sr_2CaCu_2O_{8+\delta}$ single crystals by oxygen annealing and quenching method, Chin. Phys. Lett. **33**, 067403 (2016).

[57] S. Kunisada, S. Adachi, S. Sakai, N. Sasaki, M. Nokayama, S. Akebi, K. Kuroda, T. Sasagawa, T. Watanabe, S. Shin, and T. Kondo, Observation of Bogoliubov Band Hybridization in the Optimally Doped Trilayer $Bi_2Sr_2Ca_2Cu_3O_{10+\delta}$, Phys. Rev. Lett. **119**, 217001 (2017).

[58] D. L. Feng, D. H. Lu, K. M. Shen, C. Kim, H. Eisaki, A. Damascelli, R. Yoshizaki, J. I. Shimoyama, K. Kishio, G. D. Gu, S. Oh, A. Andrus, J. O'Donnell, J. N. Eckstein, and Z. X. Shen, Signature of superfluid density in the single-particle excitation spectrum of $Bi_2Sr_2CaCu_2O_{8+\delta}$, Science **289**, 277 (2000).

[59] X. J. Zhou, T. Yoshida, D. H. Lee, W. L. Yang, V. Brouet, F. Zhou, W. X. Ti, J. W. Xiong, Z. X. Zhao, T. Sasagawa, T. Kakeshita, H. Eisaki, S. Uchida, A. Fujimori, Z. Hussain, and Z. X. Shen, Dichotomy between Nodal and Antinodal Quasiparticles in Underdoped $(La_{2-x}Sr_x)CuO_4$ Superconductors, Phys. Rev. Lett. **92**, 187001 (2004).

[60] K. M. Shen, F. Ronning, D. H. Lu, F. Baumberger, N. J. C. Ingle, W. S. Lee, W. Meevasana, Y. Kohsaka, M. Azuma, M. Takano, H. Takagi, and Z. X. Shen, Nodal quasiparticles and antinodal charge ordering in $Ca_{2-x}Na_xCuO_2Cl_2$, Science **307**, 901 (2005).

[61] K. Tanaka, W. S. Lee, D. H. Lu, A. Fujimori, T. Fujii, Risdiana, I. Terasaki, D. J. Scalapino, T. P. Devereaux, Z. Hussain, and Z. X. Shen, Distinct Fermi-momentum-dependent energy gaps in deeply underdoped Bi2212, Science **314**, 1910 (2006).

[62] K. Gotlieb, C. Y. Lin, M. Serbyn, W. T. Zhang, C. L. Smallwood, C. Jozwiak, H. Eisaki, Z. Hussain, A. Vishwanath, and A. Lanzara, Revealing hidden spin-momentum locking in a high-temperature cuprate superconductor, Science **362**, 1271 (2018).

[63] C. Guo, H. F. Tian, H. X. Yang, B. Zhang, K. Sun, X. Sun, Y. Y. Peng, X. J. Zhou, and J. Q. Li, Direct visualization of soliton stripes in the CuO2 plane and oxygen interstitials in $Bi_2(Sr_{2-x}La_x)CuO_{6+\delta}$ superconductors, Phys. Rev. Mater. **1**, 064802 (2017).

[64] L. Huang, Y. Wu, M. T. Hershberger, D. X. Mou, B. Schrunk, M. C. Tringides, M. Hupalo, and A. Kaminski, The effects of Moire lattice on the electronic properties of graphene, Phys. Rev. B **96**, 035411 (2017).