

# Powerful Terahertz Emission from a $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ Mesa Operating Above 77 K

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Mesa-shaped structures of the high critical temperature (high- $T_c$ ) superconductor  $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$  contain stacked intrinsic Josephson junctions. As such, they are a promising source of coherent radiation in the “terahertz gap” range, spanning from approximately 0.3 to 2.0 THz. Technological applications of these devices become far more practical if they can be operated at a cryogenic bath temperature of 77 K or higher. Previous works have reported emission from this type of device at high terahertz power levels at lower operating temperatures, 40–60 K, while at  $T_{\text{bath}} \geq 77$  K observed power levels have generally been low. Here we report generation of 130  $\mu\text{W}$  of coherent power at 0.456 THz from a mesa of  $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$  doped with 0.16 holes per Cu atom, at a bath temperature of 77.4 K. We find that the device radiates terahertz power when clearly identifiable cavity modes are excited, and that the frequency and bias voltage corresponding to each of these modes is almost independent of temperature. This is consistent with these modes having terahertz-frequency electric fields with very little dependence on vertical position within the mesa. We also find that the terahertz power radiated from any given mode decreases monotonically as the mesa temperature is increased. On the other hand, the low-frequency modes become inaccessible at low temperatures due to retrapping of the intrinsic Josephson junctions, and the maximum radiation power for the emitting mode is typically achieved at the temperature at which the retrapping voltage reaches the resonance voltage for this mode.

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## I. INTRODUCTION

Stacked “intrinsic” Josephson junctions in the high- $T_c$  superconductor  $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$  (Bi-2212) offer the possibility of a tunable and highly compact source of continuous-wave coherent terahertz radiation in a part of the electromagnetic spectrum that is not served well by existing technologies. Such a source would be useful for a wide range of applications in the fields of telecommunications and medical imaging, as well as in security and defense [1–4]. Operation of a mesa-shaped (Bi-2212) stack containing approximately 700 series-connected intrinsic Josephson junctions (IJJs) was demonstrated by Ozyuzer

*et al.* [5] in 2007, achieving an estimated 0.5  $\mu\text{W}$  at 0.49 THz. Subsequent development of this type of device has resulted in terahertz emission power levels as high as 0.61 mW at 0.51 THz [6], operation temperatures as high as 86.0 K [7], and emission frequencies covering the entire frequency range from 0.3 up to 2.4 THz [8,9]. In series-biased IJJ devices with a voltage bias of  $V_{\text{IJJ}}$  per junction, coherent terahertz emission occurs when the Josephson frequency  $f_{\text{Josephson}} = 2eV_{\text{IJJ}}/h$  matches a geometric resonant mode of the mesa-shaped stack, as has been studied in detail in many previous works (see, for instance, Refs. [5–20] and references therein). Here,  $e$  is the electronic charge, and  $h$  is Planck’s constant.

In order to employ high- $T_c$  terahertz sources in field-portable applications, it is highly desirable for them to be able to operate at cryogenic bath temperatures of 77.4

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K or more, while still generating technologically useful amounts of coherent terahertz power. This would permit the Bi-2212 stacks to be cooled by liquid nitrogen boiling at atmospheric pressure, or by very compact Stirling-cycle microcryocoolers using helium gas as the working fluid [21]. Bi<sub>2</sub>Sr<sub>2</sub>CaCu<sub>2</sub>O<sub>8+δ</sub> doped “optimally” with  $p = 0.16$  holes per Cu atom has a maximum possible  $T_c$  of around 95 K, depending on cation stoichiometry [22,23]. Nonetheless, achieving high terahertz emission power at bath temperatures exceeding 77.4 K is far from trivial. Previous works report emission power levels of the order a few microwatts at these temperatures [7,14,24–27]. By comparison, it is generally estimated that approximately 1 mW of time-averaged terahertz power is needed for real-time imaging applications [1,2]. Here we demonstrate that through careful design of the geometry of the Bi-2212 stack and the chip on which it stands, it is possible to achieve 130 μW of coherent terahertz power at a bath temperature of 77.4 K, which represents a significant step towards this goal. (The total emitted terahertz power levels reported in this work are calculated when integrated over  $2\pi$  sr, taking measured angle-dependence of the power density into account.) A general three-dimensional overview of one of the mesas and sketch of its vertical cross section are provided in Figs. 1(a) and 1(b), respectively, while an optical micrograph of the terahertz sources we employ in this work is shown in Fig. 1(c).

The frequencies of the geometric resonances of the stack are determined by the in-plane mode velocity  $c_n$ , which depends on the vertical periodicity of the mode and the boundary conditions at the top and bottom. Following the notation of Ref. [14], a junction stack with thickness  $t$  that contains a large number of junctions ( $N = t/D \gg 1$  with  $D$  being the interlayer periodicity), and assuming boundary conditions that can be expressed as nodes or antinodes, we have a discrete set of  $c$ -axis wave vectors  $k_z = n\pi/2t$  with  $n = 0, 1, 2, \dots$ , giving the mode velocities [28]

$$c_n^2 \approx \frac{c_0^2}{\epsilon_{rc}[1 + (n\pi\lambda_{ab}/2t)^2]} \quad (1)$$

for mode indices  $n \ll N/2$ . Here,  $c_0$  is the speed of light in free space,  $\epsilon_{rc}$  is the relative out-of-plane electric permittivity, and  $\lambda_{ab}$  is the magnetic flux penetration depth in the direction of the CuO<sub>2</sub> planes. Permitted values for  $n$  depend upon the boundary conditions at the top and bottom of the mesa, as discussed in Refs. [14,28]. If the stacked IJJ device has the shape of a rectangular mesa, it supports Fabry-Perot type resonances that approximately correspond to integer numbers of half-periods of the  $c$ -axis electric field across the width,  $w = L_x$ , and length,  $l = L_y$ , of the mesa [14,29], such that the in-plane electromagnetic wavelengths are given as  $\Lambda_{x,y} \approx 2L_{x,y}/m_{x,y}$ . The frequency

of the  $(n, m_x, m_y)$  cavity mode is, therefore,

$$f_{n, m_x, m_y} \approx \frac{c_n}{2} \sqrt{\left(\frac{m_x}{w}\right)^2 + \left(\frac{m_y}{l}\right)^2}. \quad (2)$$

Typically, the strongest emission is seen from rectangular IJJ stacks, radiating in the  $(n, 1, 0)$  cavity mode with the smallest allowed index  $n$ .

Numerical simulations of stacked IJJs show that the probability of phase synchronization of the junctions increases with the amount of energy stored in the cavity [28]. Consequently, a higher quality factor of the collective electromagnetic excitation mode of the stack is more desirable for the purposes of enhancing terahertz emission power. This is given by

$$\frac{1}{Q_{\text{total}}} = \frac{1}{Q_{\text{IJ}}} + \frac{1}{Q_{\text{radiation}}} + \frac{1}{Q_{\text{leakage}}}. \quad (3)$$

At temperatures above approximately 75 K, the largest contribution to the damping  $1/Q_{\text{total}}$  is from the quality factor of the Josephson oscillations themselves. At the Josephson plasma frequency  $\omega_p$ , this is given by

$$Q_{\text{IJ}} = RC\omega_p = RC\sqrt{\frac{2\pi I_c}{C\Phi_0}}, \quad (4)$$

where  $R$  is the resistance of each IJJ in the stack,  $C$  is its shunt capacitance,  $I_c$  is its critical current, and  $\Phi_0$  is the flux quantum (see Stewart [30]). Equivalently, we can express Eq. (4) in terms of the bulk  $c$ -axis properties of the extremely anisotropic Bi-2212 superconductor,

$$Q_{\text{IJ}} = \rho_c n_c \sqrt{\frac{2\pi \epsilon_0 D J_{zc}}{\Phi_0}}, \quad (5)$$

where  $\rho_c(T)$  is the resistive-state out-of-plane resistivity, and  $n_c$  is the far-infrared refractive index for electromagnetic waves with electric field polarized along the  $c$  axis.  $\epsilon_0$  is the permittivity of free space,  $J_{zc}(T)$  is the out-of-plane critical current density, and  $D$  is the junction periodicity along the  $c$  axis.  $\rho_c(T)$  can be estimated from the slope of the downward-swept branch of the mesa’s (hysteretic) current-voltage characteristic, as described in [31], while  $J_{zc}(T)$  can be estimated from the same data sets. The relative out-of-plane permittivity  $\epsilon_{rc} \equiv n_c^2$  is usually estimated to be about 16 for this material [32]. Using data for our Bi-2212 crystals doped with  $p = 0.16$  obtained in Ref. [31] we can calculate  $Q_{\text{IJ}}(T)$ . Our results are plotted in Fig. 2, showing that for  $T \geq 77$  K, the quality factor of the Josephson oscillations is comparatively modest, due to quasiparticle damping in the current-biased IJJs. Significant energy losses due to resistive damping of the Josephson oscillations in the mesa itself are therefore unavoidable

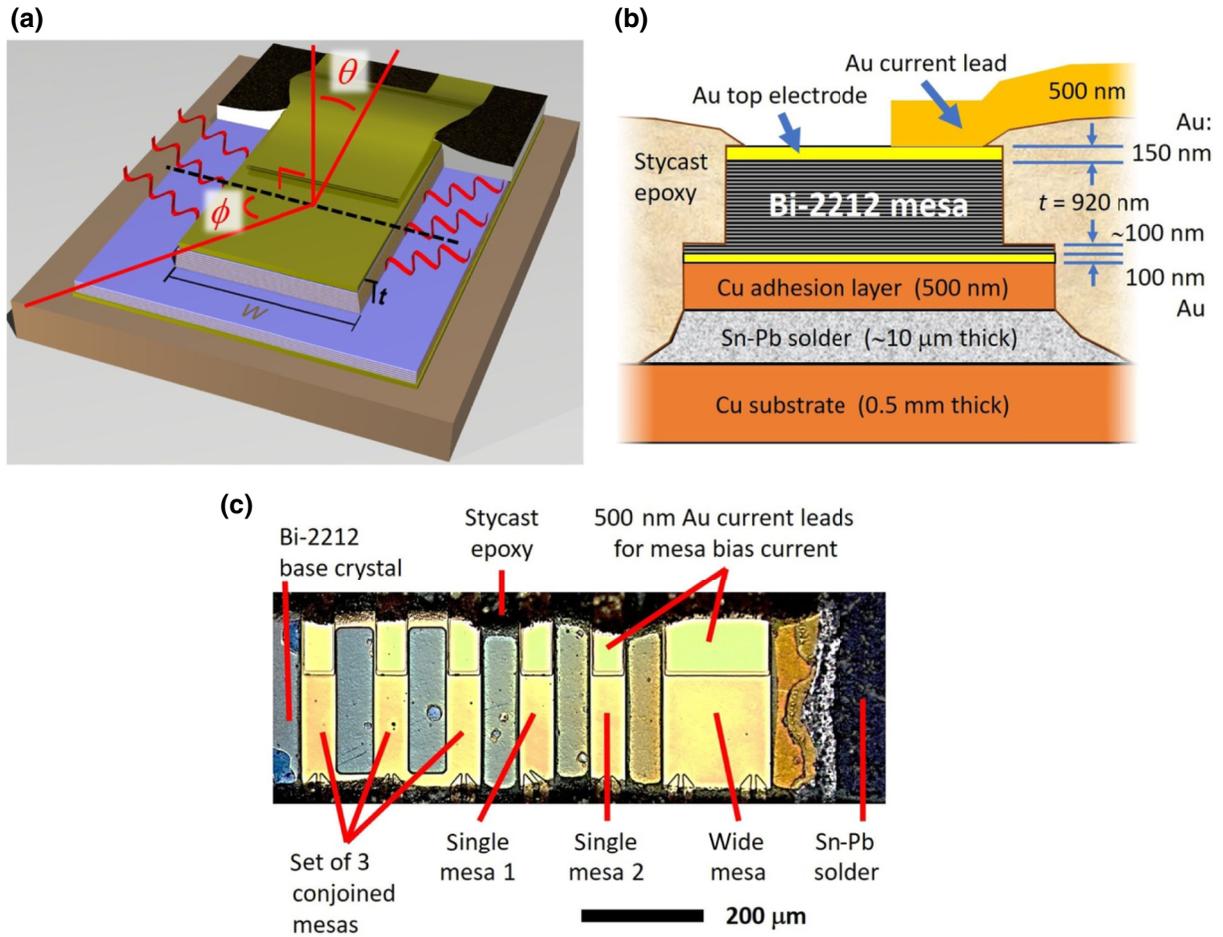


FIG. 1. (a) Schematic showing the geometry of a typical Bi-2212 mesa designed for coherent terahertz emission, after Ref. [28], with the IJJ stack having width  $w$  and total thickness  $t$ . For clarity, vertical scale is shown greatly exaggerated with respect to lateral scale. Red wavy lines are a schematic indication of the direction of greatest far-field radiated terahertz intensity, and the indicated wavelength is not drawn to scale. Actual wavelength of the emitted radiation is approximately  $600 \mu\text{m}$ , which is twice the mesa length. Polar angle  $\theta$  and azimuthal angle  $\phi$  as defined in this work are indicated. Current is supplied to the stacked IJJs via the gold top electrode and sunk via the base crystal. Any dc power dissipated in the device is also heat sunk via the base crystal. (b) Schematic showing vertical cross section of the devices studied in this work, including the metal thin films deposited on the top and bottom surfaces of the Bi-2212 crystal. For clarity, vertical distances are shown greatly exaggerated and not to scale. (c) Optical micrograph of the four devices studied in this work, with important parts of microchip labeled. Note that the Bi-2212 base crystal is so thin that it is almost transparent, and the metal adhesion layers underneath can be seen. The gold tracks at the top end of the mesas are for supplying bias current to the mesas themselves, while those at the bottom end are current leads for future installation of NiCr thin-film heaters, intended for studying the effects of localized heating of the mesas. Note that the metal tracks deposited on the epoxy appear slightly out of focus in this image, since the epoxy surface lies a few microns above the plane of the top of the mesas.

at these bath temperatures. However, when the  $Q$  factor of the Josephson oscillations is high, leakage of additional terahertz energy into the base crystal becomes relevant, whether into the Bi-2212 base crystal material directly below the mesa, or around its perimeter. In principle, these losses can be minimized by reducing the thickness of the base crystal, and/or by reducing the perimeter-to-area ratio of the mesa, thus reducing the ratio of the volume of Bi-2212 into which terahertz power leakage occurs to the amount of energy stored in the resonant mode of the stack.

The present work explores these approaches as a strategy for improving terahertz emission power for  $T \geq 77.4 \text{ K}$ .

## II. EXPERIMENT

High-quality Bi-2212 single crystals are grown using the traveling solvent floating zone method. These are annealed in flowing zero-grade air at  $600^\circ\text{C}$  to achieve optimal oxygen doping ( $p = 0.16$  holes per Cu atom), and quenched by tipping them into liquid nitrogen, resulting in a

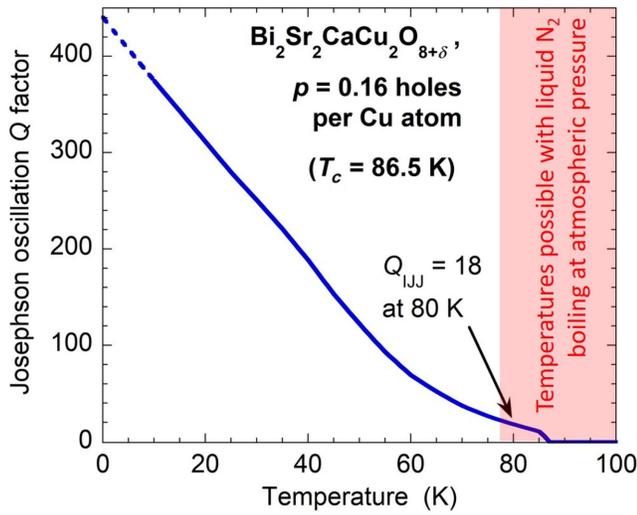


FIG. 2.  $Q_{\text{IJJ}}(T)$  as calculated using Eq. (5) for optimally doped Bi-2212, with  $p = 0.16$  holes per Cu. The experimental data sets obtained from Ref. [31] only extended as low as 10 K, so data below this temperature are a linear extrapolation (shown dotted).

superconducting  $T_c$  of 86.5 K. The crystals (typically cut to dimensions of approximately  $2 \times 2 \times 0.02 \text{ mm}^3$ ) are then mounted on  $10 \times 10 \times 0.5 \text{ mm}^3$  copper substrates using a metal adhesion coating and tin-lead solder. After the Bi-2212 crystal is cleaved to expose a fresh and planar surface, mesa-shaped terahertz sources are fabricated using optical lithography (using a direct-write laser pattern generator) and argon-ion etching. The widths of the mesas studied in this experiment vary as described in the following, but all mesas are  $300 \text{ }\mu\text{m}$  long and  $924 \text{ nm}$  high, containing close to 596 IJJs.

During the mesa fabrication process, the crystal thickness is controlled such that the base crystal on which the resulting devices stand is only approximately  $100 \text{ nm}$  thick. This combination of minimal base crystal thickness plus a highly thermally conductive mounting arrangement is essential to ensure the highest possible thermal conductance between the base of the electrically biased IJJ stack and the cryogenic bath. Since the stack is approximately 9 times thicker than the base crystal, the thermal resistance between the devices studied in this work and the cryogenic bath is dominated by the thermal resistance of the stack itself. At the bath temperatures of interest, the thermal conductivity of Bi-2212 along the crystalline  $c$  axis ( $\kappa_c$ ) is approximately  $0.3 \text{ W m}^{-1} \text{ K}^{-1}$  [33]. For the  $300 \times 180 \times 0.92 \text{ }\mu\text{m}^3$  device studied here, this implies a thermal conductance between the base plane of the mesa and the thermal bath of approximately  $0.162 \text{ W/K}$ , since the thermal resistance of the metal adhesion layer, the solder, and the copper substrate can all be neglected. The epoxies that have been used to mount Bi-2212 crystals onto substrates in terahertz generation experiments typically

have a thermal conductivity of around  $0.12 \text{ W m}^{-1} \text{ K}^{-1}$  at typical mesa operation temperatures [34,35], while the layer of epoxy between the substrate and the bottom of the Bi-2212 single crystal is usually of the order of  $20 \text{ }\mu\text{m}$  thick [34]. As a result, in most terahertz experiments employing epoxy-mounted Bi-2212 crystals, the thermal conductance from the device to the cryogenic bath is typically much lower than in the present work. One alternative approach to thermal anchoring developed recently by Tsujimoto *et al.* [36] is to mount the mesa on a sapphire substrate (which is both highly thermally conductive and transparent to terahertz radiation) using an unfilled epoxy such as Stycast 1266. The entire mesa is then encapsulated in a filled—but electrically insulating—epoxy with significant thermal conductance, such as Stycast 2850. This configuration allows the mesa to be heat sunk from the top and the sidewalls, as well as via the substrate, in addition to protecting the mesa from contaminants such as water vapor. However, the resulting thermal conductance to the bath is nonetheless much lower than in the present work, since the thermal conductance of Stycast 2850—even at room temperature—is about  $1.36 \text{ W m}^{-1} \text{ K}^{-1}$  [37], as compared with  $50 \text{ W m}^{-1} \text{ K}^{-1}$  for the 63% Sn, 37% Pb solder employed here [38]. At cryogenic temperatures, the difference in thermal conductivity between the two materials is likely to be even greater. Another important consequence of a thin base crystal is modification of the effective height  $t$  entering the mode velocity in Eq. (1). Contrary to the standard situation of a very thick base crystal, where the parameter  $t$  is equal to the mesa height, we believe that in our case this effective height is given by the distance between the mesa top and the bottom of the crystal.

Three types of device are constructed on the same crystal: two mesas with dimensions  $300 \times 60 \times 0.92 \text{ }\mu\text{m}^3$  designed to operate in the  $(n, 1, 0)$  mode; a mesa with dimensions  $300 \times 180 \times 0.92 \text{ }\mu\text{m}^3$  designed to operate in the  $(n, 3, 0)$  mode; and a device consisting of an array of three conjoined mesas, each with dimensions  $300 \times 60 \times 0.92 \text{ }\mu\text{m}^3$  and with  $60\text{-}\mu\text{m}$  spacing between nearest neighbors. This latter device is intended to achieve enhanced emission power by acting effectively as three  $300 \times 60 \times 0.92 \text{ }\mu\text{m}^3$  stacks, each operating in its  $(n, 1, 0)$  mode and electromagnetically coupled together [see Fig. 1(c)]. All four devices are designed to emit in the frequency range from approximately  $0.4$  to  $0.7 \text{ THz}$ , while operating at a range of bath temperatures from  $70 \text{ K}$  and above. For efficient terahertz radiation coupling to free space, it is preferable for the vertically oriented electric fields at opposite edges of the mesa to oscillate in antiphase with respect to each other, meaning that the width of the device (or of its synchronized emitting parts) must ideally be an odd number of half-wavelengths in Bi-2212. For instance, making the width of the mesa an even number of half-wavelengths (such as by making its dimensions  $300 \times 120 \times 0.92 \text{ }\mu\text{m}^3$ ) would be expected to be

comparatively inefficient at emitting terahertz radiation with a half-wavelength of 60  $\mu\text{m}$  in Bi-2212 (see Ref. [14]).

Devices are mounted in a closed-cycle cryostat with high-density polyethylene optical windows, while the emitted terahertz power is monitored using a calibrated bolometer employing a niobium transition-edge sensor in closed-loop power feedback mode, supplied by QMC Instruments Limited. The terahertz beam is chopped at 233.3 Hz, and the detector signal is measured using a lock-in amplifier. The power responsivity of the bolometer is determined at 0.599 THz by comparing its response with that of a calibrated pyroelectric detector employing identical terahertz collection geometry. The QMC bolometer collects terahertz radiation using a circular Winston cone with diameter of 5 mm, resulting in a collection angle of  $2.18 \times 10^{-4}$  sr in our measurement configuration. It is found to have an optical terahertz responsivity of 9350 V/W (before signal preamplification) for radiation incident upon the bolometer's (circular) outer window that falls within the field of view of its collection optics, independent of terahertz polarization orientation. We also measure terahertz emission power levels from the devices studied in this work using a liquid-helium-cooled silicon bolometer supplied by Infrared Laboratories, also employing a circular Winston cone. This bolometer is also calibrated in the same manner using the same pyroelectric detector, and the implied terahertz power levels are found to agree within uncertainties (of approximately  $\pm 5\%$ ) when measured with either of these bolometer systems. These power level uncertainties are primarily limited by small uncertainties associated with the precise dimensions and alignment of the terahertz collection geometries. The Infrared Laboratories bolometer is found to have an optical terahertz responsivity of 4970 V/W (before signal preamplification) for radiation incident upon the bolometer's outer window that falls within the field of view of its collection optics. Further details of our calibration procedures can be found in the Supplemental Material of Ref. [6].

The mesas are current-biased, with the soldered connection between the base crystal and the copper substrate acting as a return current lead. The voltage across the current-driven mesa is measured by using an adjacent unbiased mesa as a negative voltage tap. (Owing to the extreme anisotropy of Bi-2212, this mesa is effectively at the same potential as the base of the driven mesa.) At each bath temperature studied, the bias current is swept up until all of the IJJs are driven into the resistive state, and then swept back down continuously at 40  $\mu\text{A/s}$  to measure the bias-dependence of the terahertz emission.

An example of such a current-voltage sweep for the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa at a bath temperature of 77.4 K (just above the boiling point of liquid nitrogen at sea level) is plotted in Fig. 3(a), together with the

voltage-dependence of the emitted terahertz power. The angle-dependence of the emitted terahertz radiation is measured for the devices studied in this work, and is found to be very similar to that reported by Minami *et al.* [39]. Power levels plotted in Figs. 3, 4, 5, and 7 are obtained by normalizing the power level measured at  $(\theta = 45^\circ, \phi = 0^\circ)$  to  $2\pi$  sr, incorporating an appropriate angle-dependence (see Fig. 6). Figures 3(b) and 3(c), respectively, show details of the terahertz emission and  $I$ - $V$  characteristics in the region in which the mesa emits maximum terahertz radiation at 77.4 K. The dashed black line in Fig. 3(c) is a straight line fitted through the origin, allowing us to estimate the current drawn by the mesa in the absence of terahertz resonance modes. [We note that in the absence of energy loss to any geometric or bosonic resonant modes, the current-voltage characteristics are weakly nonlinear with only small upward curvature,  $d^2I(V)/dV^2 \geq 0$ , meaning that we can estimate the small-bias resistance of the mesa by identifying a straight line through the origin that is the tangent to  $I(V)$ .] As was done in Ref. [28], we may then estimate the excess current drawn by the mesa due to terahertz resonances being excited, and thus the power delivered to the terahertz mode. At peak emission intensity at 77.4 K, we estimate 6.8 mA of excess current and 4.0 mW of excess power. Since we measure 130  $\mu\text{W}$  of emitted terahertz power under these conditions, this implies that approximately 3.3% of the power delivered to the terahertz mode is being converted to far-field radiation. Since the base crystal in this experiment is only 100 nm thick and the temperature is quite close to  $T_c$ , it is likely that dissipation under these emission conditions is dominated by quasiparticle losses inside the mesa itself. Equivalent calculations performed in Ref. [28] estimated that emitted terahertz radiation was 2.5% of the power delivered to the cavity resonance in that experiment. However, the mesa in Ref. [28] stood on a Bi-2212 base crystal approximately 20  $\mu\text{m}$  thick, was underdoped ( $T_c = 76.3$  K), and was operated at 20–30 K, so such a level of agreement in the emission efficiency of these two devices is actually quite surprising. Similar dramatic deviations from a linear  $I$ - $V$  characteristic at terahertz emission resonances were previously reported by Kitamura *et al.* [26]. Cattaneo *et al.* have very recently reported a terahertz emission power efficiency of 12% of the input dc power, as obtained from a small mesa microfabricated upon a Bi-2212 whisker, emitting at 4.0 THz [40].

As the bias current is swept through the level corresponding to maximum terahertz emission, the mesa voltage undergoes a jump of  $-35$  mV in total at constant current, followed by a second jump of  $+10$  mV as the current is swept below the level associated with maximum emission. [The jumps are labeled by the arrows and red asterisks in Figs. 3(b) and 3(c).] The  $-35$  mV jump is highly reproducible, as shown in the Appendix, and does not show the current hysteresis that would be expected if

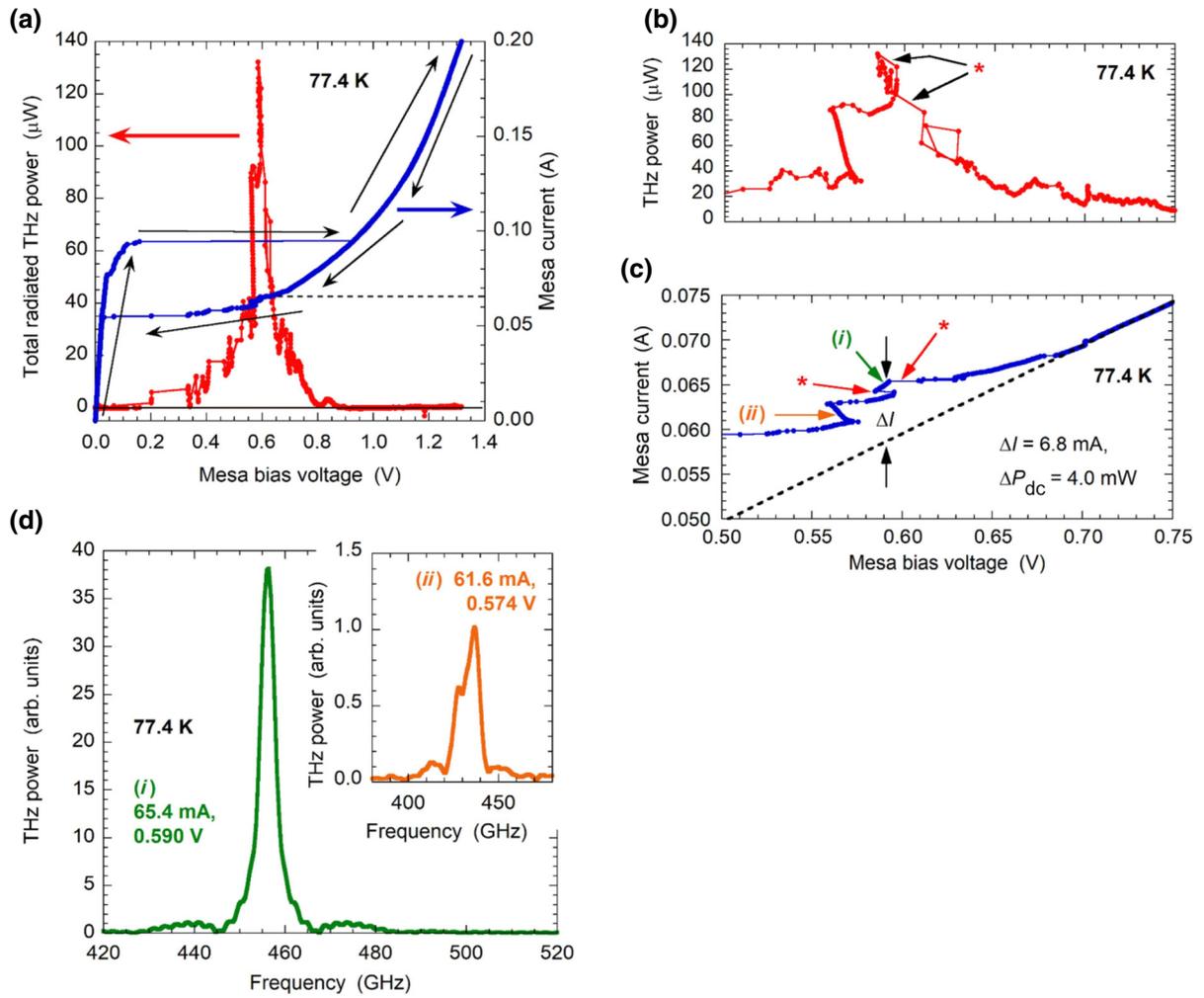


FIG. 3.  $p = 0.16$ ,  $N = 596$ ,  $300 \times 180 \times 0.92 \mu\text{m}^3$  ( $l \times w \times t$ ). (a) Mesa bias current for  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa and total integrated radiated terahertz power, plotted against mesa bias voltage, at a bath temperature of 77.4 K. Black arrows indicate the direction in which current is swept. (b) Detail of the terahertz power emission plotted in (a), close to the emission mode. (c) Detail of the  $I$ - $V$  characteristic plotted in (a), over the same range of voltages as in (b). The dashed black line is a straight line fit through the origin assuming a constant device resistance, allowing us to estimate the excess current and power drawn by the resonant terahertz mode. The voltage steps labeled with the arrows and red asterisks are discussed in the text. (d) Main image shows FTIR spectrum for terahertz emission from the same device, collected close to maximum emission power at this temperature, labeled as point (i) in (c). Bias current of 65.4 mA is indicated by the dotted horizontal line in (a) above. As-measured FWHM linewidth is limited by the 2.25 GHz resolution of the spectrometer. Inset shows emission spectrum measured at 61.6 mA on the reentrant section of the current-voltage curve at the same temperature, labeled as point (ii) in (c). The emission spectrum is much broader and multimodal, while emission power is much weaker than in the main plot. (Please note the different scales for the power and frequency axes in the inset.)

it were associated with retrapping of one or more IJJs in the mesa itself, or in the base crystal immediately below it. Since this jump cannot be explained via Josephson retrapping, we suggest that these features are most probably due to a discontinuous change in the power drawn by the cavity resonance while the mesa is being biased at constant current. Such voltage jumps are most probably due to sudden changes in the fraction of synchronized junctions in the stack, and were predicted by our earlier numerical Lawrence-Doniach simulations in Ref. [28] (see, for instance, Fig. 5 of Ref. [28]). The +10 mV jump occurs

at a highly reproducible current (see the Appendix), but the size of the jump shows a certain amount of statistical variation, implying that the collapse of the highly synchronized state associated with this jump may not be complete, and that after the jump occurs, some of the IJJs remain synchronized and generating terahertz radiation.

At lower bias currents, the IJJs retrap over a range of currents that vary by a factor of between 10% and 20%, depending upon the cryogenic bath temperature. Retrapping of Bi-2212 IJJs is a statistical process that for small-size stacks has previously been thoroughly

studied and found to agree well with theoretical expectations [41]. In larger stacks, such as those studied here, retrapping is found to occur at higher current densities than expected, and the reasons for this have not been

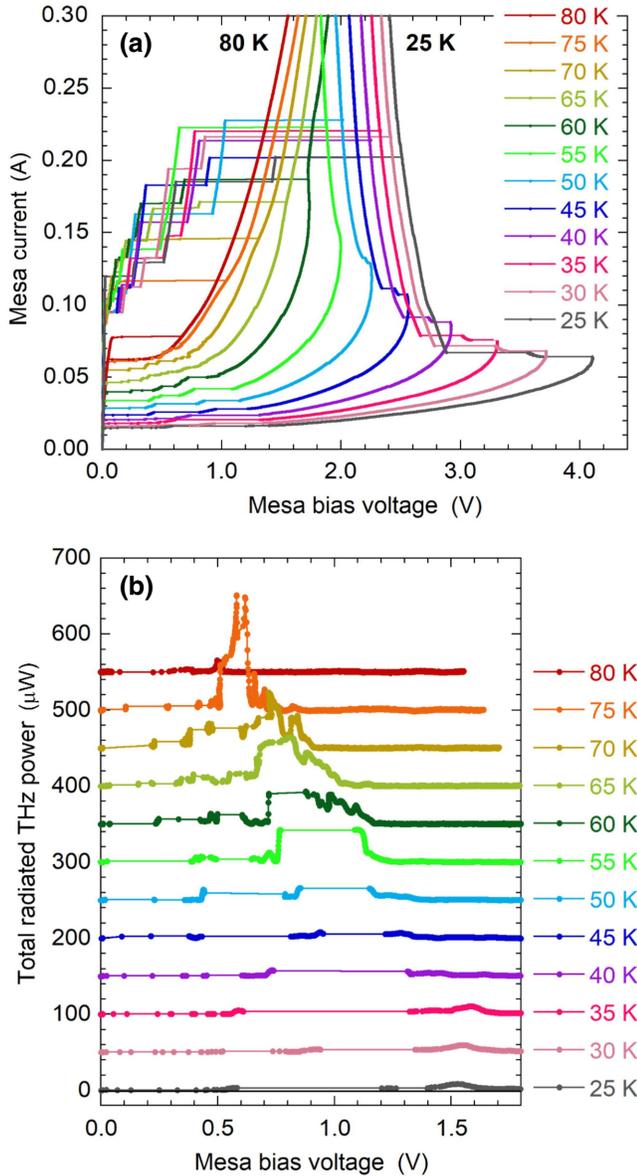


FIG. 4.  $p = 0.16$ ,  $N = 596$ ,  $300 \times 180 \times 0.92 \mu\text{m}^3$  ( $l \times w \times t$ ). (a) Current-voltage characteristics for  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa at a range of cryogenic bath temperatures, showing hysteretic switching of stacked IJJs. For this device, terahertz emission always occurs on the downward sweep, and only in the low-current bias regime (i.e., not in the back-bending regions). (b) Total integrated radiated terahertz power plotted against voltage for the same device. For clarity, successive curves are offset by 50  $\mu$ W, and only data recorded on the current downsweep are shown. Note also that the voltage scale differs from that shown in (a). (No terahertz emission occurs at mesa bias voltages exceeding 1.8 V for this device.)

conclusively established. Once the bias current is sufficiently low for retrapping to begin, the shape of the  $I$ - $V$  curve is different for each realization of the measurement. Since self-heating in our devices at these bias currents is comparatively low (as we discuss in the following) the per-IJJ electrical resistance changes little as a function of the number of junctions switched to the resistive state, implying that the spread of retrapping voltages is only about 10% to 20%, even as the total stack voltage varies substantially. As shown in Figs. 3–5, terahertz emission can occur even when only a subset of the stack is still in the resistive state. Our results show that this can be observed for as few as 200 IJJs remaining in the resistive state.

An example of a resulting terahertz emission spectrum measured using an FTIR spectrometer is plotted in Fig. 3(d), with greatest intensity at 0.456 THz. The linewidth shown in Fig. 3(d) is limited by the 2.25 GHz resolution of our spectrometer, and the true linewidth is likely to be even narrower than this [42]. The linewidth would be determined by the quality factor in the situation when the cavity is excited by unsynchronized IJJs with a broad frequency distribution. This scenario is clearly not realized in our system, since the as-measured linewidth would imply a quality factor of at least  $Q = f / \Delta f \approx 114$ , which is much narrower than  $Q \approx 22$  as calculated using Eq. (5) for Bi-2212 doped with  $p = 0.16$  at  $T = 77.4$  K, as can be seen from Fig. 2. This result is an experimental demonstration of the established theoretical prediction that for emission modes of stacked IJJs, the linewidth of the terahertz emission is determined *not* by the quality factor of the Josephson oscillations, but instead by the level of thermal noise present in the stack (which adds to the junction bias voltage provided by the power supply), plus—in some cases—the failure of all of the IJJs in the stack to phase-lock to the same frequency 100% of the time. This phenomenon of partial synchronization was simulated and analyzed using a Lawrence-Doniach model for a nonuniform IJJ stack with a trapezoidal cross section including thermal noise by some of the authors in Ref. [28].

Figure 4 shows the evolution of the current-voltage and terahertz emission characteristics for the same  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa, as the bath temperature is changed in 5 K steps. As usual, the increasing current drives IJJs in the mesa into dynamic resistive states until all junctions are switched and the full dynamic state is reached. This state is preserved with decreasing current until the IJJs start to switch back to static states (retrap). The retrapping current increases with decreasing temperature. The nonmonotonic behavior of  $I$ - $V$  curves at higher currents observed at temperatures below 60 K is caused by heating effects. This is usual behavior for mesas fabricated on top of Bi-2212 crystals [5–18]. The maximum voltage ranging from 1.5 V at 60 K to 4.1 V at 25 K, however, is

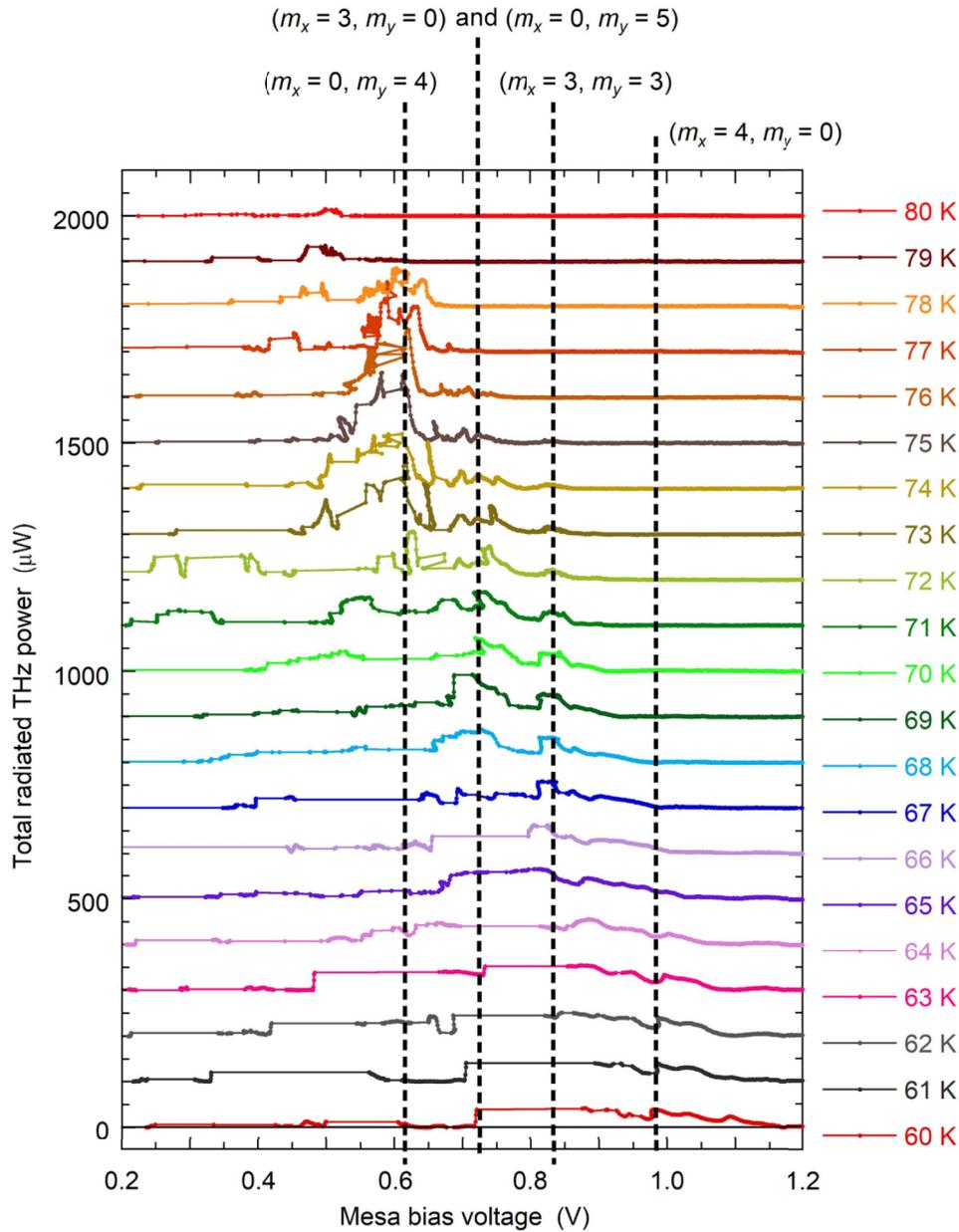


FIG. 5.  $p = 0.16$ ,  $N = 596$ ,  $300 \times 180 \times 0.92 \mu\text{m}^3$  ( $l \times w \times t$ ). Total integrated radiated terahertz power for  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa plotted against voltage. These data are for the same device as shown in Fig. 4, but at more closely spaced temperatures. For clarity, successive curves are offset by  $100 \mu\text{W}$ . Vertical dashed lines show the four strongest modes of terahertz emission, labeled with their most probable indices for  $(m_x, m_y)$ , assuming  $n = 0$ .

unusually high due to the efficient heat management in our structure.

The voltage jumps seen at bias voltages in excess of 2 V for  $T \leq 45$  K are not due to any kind of Josephson switching, but are instead due to the hysteretic collapse of localized current filaments in the mesa, so called “hot spots.” These filaments can occur due to the fact that for underdoped or optimally doped Bi-2212 (i.e., for  $p \leq 0.16$ ),  $d\sigma_c/dT > 0$  for all  $T \leq T_c$ , which can give rise to localized thermal runaway for certain combinations of bath

temperature and biasing conditions. This phenomenon was studied and analyzed in detail by Wang *et al.* [43,44], and by some of the authors, in Ref. [45].

The high values of the maximum voltages in our device allow us to observe excitation of the higher-order modes, which are usually inaccessible. Indeed, the voltage dependences of the radiation power at different temperatures presented in Fig. 4(b) demonstrate the rich peak structure. At low temperatures, we only observe the modes with high indices yielding low radiation power, because the

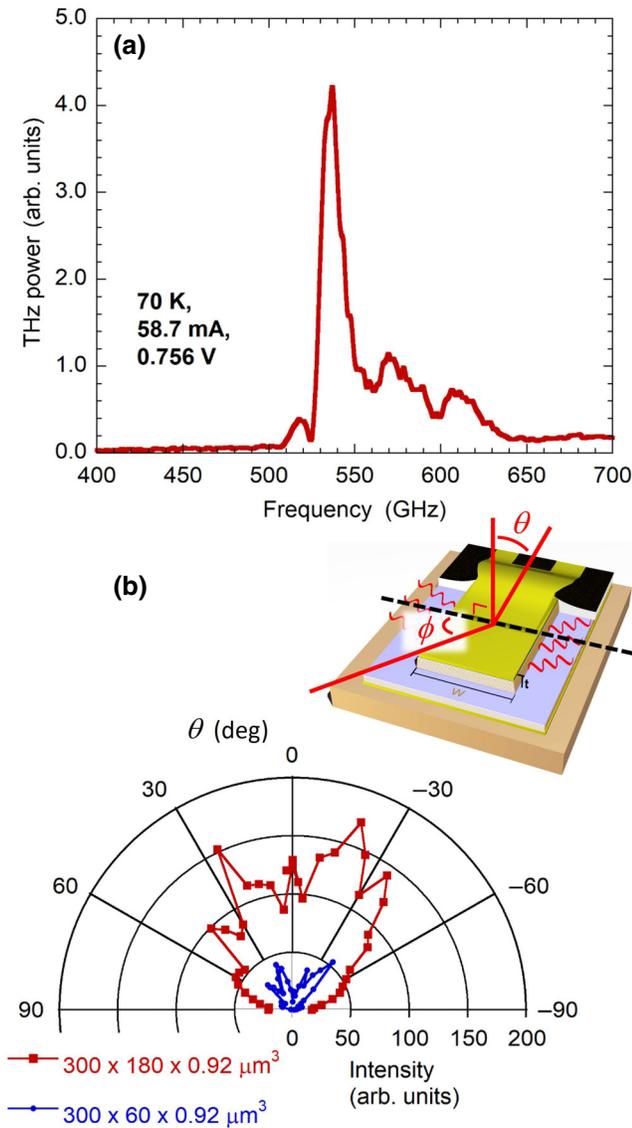


FIG. 6.  $p = 0.16$ ,  $N = 596$ ,  $300 \times 180 \times 0.92 \mu\text{m}^3$  ( $l \times w \times t$ ). (a) Terahertz emission spectrum as measured by FTIR from  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa at a cryogenic bath temperature of 70 K and a bias current of 58.7 mA. This spectrum is very broad and multi-peaked, indicating that synchronization of the junctions is quite unstable under these conditions. The most probable modes being excited are ( $n = 0$ ,  $m_x = 3$ ,  $m_y = 0$ ) and/or ( $n = 0$ ,  $m_x = 0$ ,  $m_y = 5$ ). (b) Angle dependence of the terahertz emission for this device as measured under the same operating conditions as in (a) as a function of polar angle  $\theta$ , i.e., when rotating about the long axis of the mesa. For comparison, data from one of the  $300 \times 60 \times 0.92 \mu\text{m}^3$  mesas on the same chip are plotted, also operating at 70 K, in the ( $n = 0$ ,  $m_x = 1$ ,  $m_y = 0$ ) and/or ( $n = 0$ ,  $m_x = 0$ ,  $m_y = 5$ ) mode. Diagram shows the definition of the polar angle  $\theta$  and azimuthal angle  $\phi$  as referred to in this work.

low-frequency modes are inaccessible due to retrapping. With increasing temperature, the per-IJJ voltage at which retrapping occurs decreases, and the modes with lower

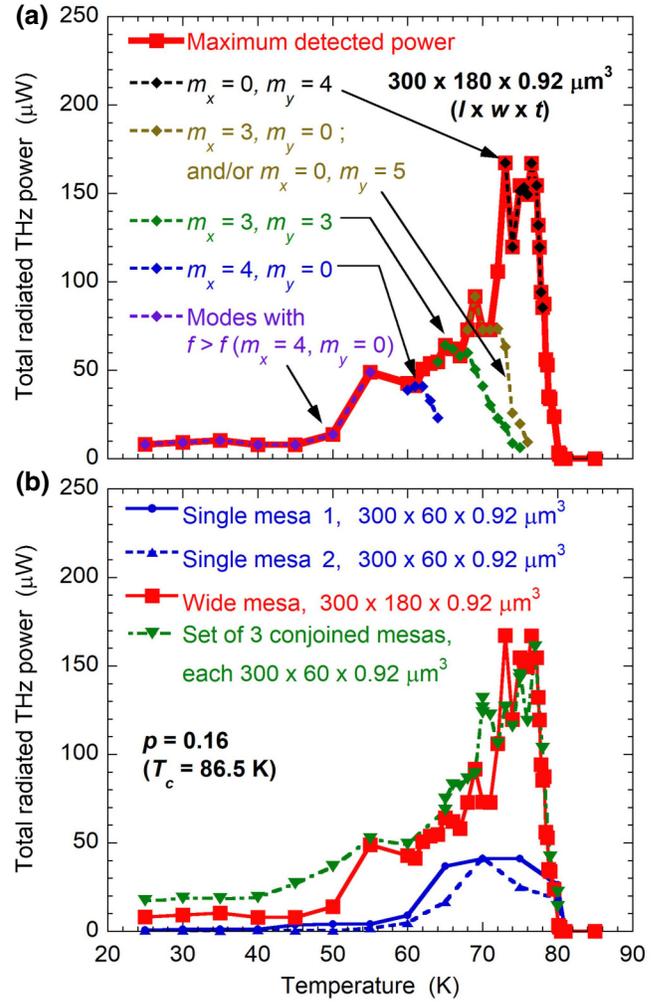


FIG. 7. (a)  $p = 0.16$ ,  $N = 596$ ,  $300 \times 180 \times 0.92 \mu\text{m}^3$  ( $l \times w \times t$ ). Maximum total radiated terahertz power that could be generated by the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa at each temperature (red curve, squares). Where one or more modes could be clearly identified at each temperature, the emission power from each of these is plotted also (diamonds). (b) Maximum total radiated terahertz power plotted against temperature, for all four devices fabricated on this chip. Blue curves (circles and triangles) show emission for the two individual  $300 \times 60 \times 0.92 \mu\text{m}^3$  mesas. Red curve (squares) shows emission for  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa, while the green dashed curve (inverted triangles) shows emission for the set of three conjoined  $300 \times 60 \times 0.92 \mu\text{m}^3$  mesas.

indices, which have higher radiation power, become accessible. As a consequence, the maximum achievable radiation power *increases* with increasing temperature, before peaking at around 77 K. Figure 5 shows the evolution of the emitted terahertz power with respect to bias voltage over a range of more closely spaced temperatures, closer to  $T_c$ . Between 60 and 80 K, the voltage dependence of the power displays a complicated structure, partly caused by hysteretic retrapping events. Nevertheless, groups of emission lines can be identified that do not shift with

temperature (marked by vertical dashed lines), while the entire pattern seems to shift to lower frequencies at higher  $T$ . A temperature-independent emission frequency for any given mode clearly implies that the vertical mode index  $n$  in Eq. (1) is zero for all observed excited modes. Similar behavior has been observed in stand-alone mesas in Refs. [46,47].

By contrast, using experimental values for  $\lambda_{ab}(T)$  for optimally doped Bi-2212 with  $p = 0.16$  measured by Anukool *et al.* [48], we would expect terahertz modes with  $n = 1$  to drop in frequency and corresponding bias voltage by approximately 30% over this temperature range. Modes with  $n > 1$  would be expected to drop in frequency even further [see Eqs. (1) and (2)]. Such temperature-dependent behavior has been observed in mesas fabricated on the top of large crystals [28]. We can therefore conclude that any terahertz modes being strongly excited in this device have  $n = 0$ , which is consistent with these modes having antinodes of the  $c$ -axis electric field at both the top of the mesa and the bottom of the Bi-2212 base crystal. The presence of these antinodes is in turn consistent with this mesa having a thick gold electrode on top, while standing on a 100-nm base crystal, which itself stands on a thick layer of gold. Since our base crystal is only approximately 100 nm thick, the terahertz modes are expected to extend from the top of the mesa to the bottom of the base crystal, with minimal variation over this height of the  $c$ -axis electric field strength associated with each mode that has  $n = 0$ . At terahertz frequencies, the skin depth in thermally evaporated Au thin films can be estimated to be about 20 nm, using experimental values for Au thin-film sheet resistance obtained from Refs. [49–51]. The 150-nm Au top electrode is sufficiently thick that it substantially exceeds the gold's skin depth, so that the presence of the 500-nm Au current leads at either end of the mesa are not expected to affect the nature of the resonant modes there. This contrasts with the situation identified for a 30-nm top electrode in Ref. [14].

The multiple terahertz emission power peaks shown in Fig. 5 imply that modes with different  $(n = 0, m_x, m_y)$  are being excited at different cryogenic bath temperatures. The values of the emission frequencies for the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa can be interpreted assuming a far-infrared refractive index of 4.24 for optimally doped Bi-2212 with  $p = 0.16$ , and resulting frequencies of important resonant modes of this device are listed in Table I. This interpretation suggests that the strongest emission achieved at 77.4 K is due to the  $(n = 0, m_x = 0, m_y = 4)$  mode, resulting in 130  $\mu\text{W}$  of total radiated power at 0.456 THz. Even though this mode has even  $m_y$ , it is likely to radiate terahertz power more strongly than might naïvely be expected, since the 500-nm Au current lead almost completely covers one end of the mesa, meaning that the terahertz radiation from the ends of the mesa would not be expected to cancel itself out in the far-field limit. Since the responsivity of our terahertz detector setup is

independent of polarization angle, we expect that the differences observed in mode intensities are most probably due to some modes radiating more efficiently in the far-field limit than others. It is also likely that different mode indices result in different angular emission patterns, resulting in different collection efficiencies for any given choice of  $\theta$  and  $\phi$ . An important future experiment will be to compare angular dependence of the emission for different modes. An additional important future experiment would be to collect and focus as much of the emitted terahertz power as possible using a photoconductive antenna, as was performed by An *et al.* [52], allowing a more direct comparison for the terahertz power emitted by different modes.

At lower bath temperatures, emission is also seen from higher frequency modes. The strongest of these is observed most clearly at 70 K, and is most likely to be the  $(n = 0, m_x = 3, m_y = 0)$  mode, which is approximately degenerate with the  $(n = 0, m_x = 0, m_y = 5)$  mode, due to the dimensions of this device. Terahertz emission from these two degenerate modes is found to consistently result in broad, multi-peaked spectra. It is likely that the near-degeneracy of these two modes results in unstable synchronization characteristics of the IJJs, suggesting that this is a characteristic that should be avoided if narrow emission linewidth is desirable, as is usually the case. Figure 6 shows the emission spectrum from this mode, together with the angle dependence of its emission intensity. A relevant future experiment will be to study the emission modes from this device with much higher frequency resolution using heterodyne mixing techniques, as was performed in Refs. [39,42].

A further strong emission peak seen at approximately 0.84 V over a wide range of temperatures (see Fig. 5) is most likely to correspond to the  $(n = 0, m_x = 3, m_y = 3)$  mode, while a broad terahertz emission peak at approximately 1.00 V is likely to correspond to the  $(n = 0, m_x = 4, m_y = 0)$  mode. At voltages above 1.00 V, emission is seen over a broad range of voltages, which is likely to be the result of a large number of modes whose different values of  $(n = 0, m_x, m_y)$  result in closely spaced mode energies.

The total radiated terahertz power for each of these modes of the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa is plotted as a function of temperature in Fig. 7(a). The major limiting factor preventing excitation of low-frequency modes at low temperatures is retrapping of the IJJs when the current drops below a certain value. With increasing temperature, a decrease of the voltage at which retrapping occurs allows the excitation of modes with lower frequencies, excited at lower per-junction bias voltages. At the same time, emission from higher-frequency terahertz modes becomes suppressed more quickly as the bath temperature is increased. This is most probably because higher-frequency modes are excited at higher bias voltage and current leading to

TABLE I. Resonant frequencies for modes of the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa that are anticipated to radiate significant amounts of terahertz power.<sup>a</sup>

$n$	$m_x$	$m_y$	$f_{n, m_x, m_y}$ (THz)	Comment
0	0	3	0.354	Not observed due to retrapping
0	2	0	0.393	Not observed due to retrapping
0	0	4	0.471	Strong emission seen with narrow linewidth from 74 to 78 K
0	0	5	0.589	These two degenerate modes give strong emission, but with a
0	3	0	0.589	broad and multipeaked emission spectrum, from 68 to 75 K
0	3	3	0.687	Emission observed with broad and multipeaked emission spectrum, from 65 to 75 K
0	4	0	0.786	Emission observed with comparatively weak intensity but narrow linewidth, from 60 to 64 K

<sup>a</sup>Frequencies are calculated using Eqs. (1) and (2) assuming mode index  $n=0$ , and a far-infrared refractive index of 4.24. Figure 7 shows the temperature dependences of the emission power for several modes.

higher self-heating. As a result, IJJ synchronization for these modes is suppressed at lower bath temperatures as compared with lower-frequency modes.

Figure 7(b) shows the behavior of the maximum power achievable at each bath temperature studied for all four devices investigated. In all four cases, emitted terahertz power increases as temperature is increased, peaking between 70 and 75 K. We wish to emphasize, however, that such behavior is caused by the emergence of the modes with low indices characterized by higher emission power, while the power emitted from any given mode monotonically decreases with increasing temperature, as shown in Fig. 7(a). Maximum emission power decreases rapidly as the bath temperature is increased above 77.0 K. The two mesas with dimensions  $300 \times 60 \times 0.92 \mu\text{m}^3$  show comparatively consistent emission power characteristics, while the wider mesa with dimensions  $300 \times 180 \times 0.92 \mu\text{m}^3$  shows significantly enhanced emission above 70 K relative to the  $300 \times 60 \times 0.92 \mu\text{m}^3$  mesas. Making the mesa wider is likely to have two effects. Firstly, it increases the total length of the mesa sidewalls, enhancing the terahertz power coupling to free space. Secondly, tripling the mesa width decreases the mesa's perimeter-to-area ratio, thus reducing the importance of energy losses into the surrounding base crystal as described previously. Both of these effects are likely to improve emission performance, especially at bath temperatures close to  $T_c$ . The greatest achievable power from this device is just over  $160 \mu\text{W}$ , obtained at a bath temperature of 76.5 K. At bath temperatures achievable with liquid nitrogen at atmospheric pressure (77.4 K), over  $130 \mu\text{W}$  is generated from the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa. At a bath temperature of 55 K, this mesa is able to emit at much higher bias voltages and emission frequencies, namely generating in excess of  $40 \mu\text{W}$  at approximately 0.910 THz.

The device consisting of three conjoined mesas shows very similar emission intensity as a function of temperature, at the cost of having a total area footprint some 40% larger than that of the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa. The fact that this set of conjoined mesas shows a roughly fourfold increase in maximum power output

relative to the  $300 \times 60 \times 0.92 \mu\text{m}^3$  mesas implies that only two out of the three conjoined mesas synchronize effectively, and that the third mesa does not significantly contribute to the terahertz emission power. All the mesas studied show back-bending of their  $I$ - $V$  curves at sufficiently high current bias. However, due to the very strong heat sinking in these devices, the bias voltages at which this occurs are always much larger than the Josephson voltage for any resonant mode that could be efficiently excited. Consequently, no terahertz emission is detected in the back-bending regime, in contrast to what is commonly observed for Bi-2212 mesas whose base crystals are mounted on epoxy.

### III. DISCUSSION AND CONCLUSIONS

By constructing  $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$  mesas doped with  $p=0.16$ , and with greatly enhanced heat sinking allowing us to design larger devices with a correspondingly lower perimeter-to-area ratio, we obtain  $130 \mu\text{W}$  of coherent continuous-wave radiation at 0.456 THz and a cryogenic bath temperature of 77.4 K. This represents an improvement of over an order of magnitude beyond previously reported results for this type of device operating at  $T_{\text{bath}} \geq 77$  K, as reported in Refs. [7, 14, 17–20, 24]. At 76.5 K, just below the boiling point of liquid nitrogen, the same device emits over  $160 \mu\text{W}$ . We attribute these advances to the fact that our samples are far more strongly heat sunk than those previously reported. This is achieved by using a copper substrate instead of sapphire, and mounting by using tin-lead solder instead of electrically conductive epoxy, as well as by using a very thin base Bi-2212 crystal of approximately only 100 nm in thickness. This allows the bath temperature to be brought much closer to  $T_c$  before self-heating begins to significantly reduce the superfluid density, and thus in turn, reduces the emitted THz power.

Previous calculations and numerical simulations by one of the authors have studied the phenomenon of terahertz power leakage from Bi-2212 IJJ stacks into the base crystal [53]. These studies show that for Bi-2212 with  $p=0.16$  operating at temperatures above 40 K, the base crystal

becomes effectively transparent to terahertz radiation if it is less than approximately  $1 \mu\text{m}$  thick. Consequently, terahertz power losses in the base crystal become negligible, and a noble metal layer underneath the base crystal acts as an effective mirror for terahertz radiation. For devices with very thin base crystals mounted on noble metal (such as the ones studied here), coupling of terahertz radiation from the IJJ stack into free space will therefore depend only upon the electromagnetic characteristics of the stack itself, and not significantly upon the lateral dimensions of the base crystal or its precise thickness. In addition to the results reported in the present work, we have also observed 40–50  $\mu\text{W}$  of total emitted terahertz power from five other Bi-2212 mesa devices fabricated upon two other chips, both of which have base crystals that are approximately 100 nm thick and soldered to copper substrates [54]. In all of these devices, the lateral IJJ stack dimensions were  $60 \times 300 \mu\text{m}^2$ , and stack heights ranged from 0.75  $\mu\text{m}$  to 0.95  $\mu\text{m}$ .

It is likely that substantial further improvements in the terahertz power level could be achieved by either increasing the number of stacked junctions, increasing the length of the device, or both. It is already well established that the power emitted from the stack into an unmatched load scales as the square of the number of stacked junctions oscillating in phase [5]. Theoretical considerations predict that the emitted terahertz power in the  $(n, 1, 0)$  mode scales as roughly  $L_y^2$  when  $L_y$  is less than the wavelength in free space of the emitted terahertz radiation (approximately 600  $\mu\text{m}$  for the devices studied here), and as  $L_y$  thereafter [55].

For a mesa with only modest self-heating, and which is strongly heat sunk at its base to the cryogenic bath, the power dissipation density within the mesa can be approximated as uniform, allowing the heat diffusion equation to be approximately solved analytically. In this limit, the temperature within the mesa as a function of height  $z$  above the base of the mesa is given by the following:

$$T(z) = \frac{P}{\kappa_c} \left( tz - \frac{z^2}{2} \right) + \frac{Pwl}{G} + T_{\text{bath}}, \quad (6)$$

where  $P$  is the dc power density inside the mesa,  $\kappa_c$  is the  $c$ -axis thermal conductivity,  $w$  and  $l$  are the mesa's width and length,  $t$  is its  $c$ -axis thickness, and  $G$  is the thermal conductance connecting the base of the mesa to the bath. When the  $300 \times 180 \times 0.92 \mu\text{m}^3$  device is operating at its peak terahertz power output (at 0.456 THz) for  $T_{\text{bath}} = 77.4 \text{ K}$ , it dissipates approximately 39 mW. Equation (6) predicts that under such conditions, the bottom of the mesa ( $z = 0$ ) self-heats to 78.0 K, and the top of the mesa ( $z = t$ ) self-heats to 78.7 K. If the thickness of the device were increased to 1.3  $\mu\text{m}$ , it would increase the temperature at the top of the mesa only to 79.4 K (under equivalent biasing conditions, such that the voltage per IJJ

is kept the same). While it is perhaps premature to speculate, one might therefore envisage a mesa of dimensions  $1000 \times 180 \times 1.3 \mu\text{m}^3$ , which might be expected to generate approximately 2.0 mW of terahertz power at roughly 0.456 THz.

By employing a crystal with a  $T_c$  of 91.5 K, Sun *et al.* were able to generate approximately 100 nW of radiated terahertz power at  $T_{\text{bath}} = 86.0 \text{ K}$  [7]. These results suggest that the same terahertz power output levels from the type of device demonstrated in the present work could be achieved at still higher temperatures by using a  $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$  crystal with different Bi : Sr stoichiometry and therefore a higher  $T_c^{\text{max}}$  than 86.5 K (as employed here) thus increasing the amount of available thermal headroom.

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## APPENDIX

Figure 8 illustrates the reproducibility of the phenomena plotted in Fig. 3. Figure 8(a) shows the reproducibility of the voltage jumps that we provisionally attribute to the formation and collapse of a synchronized state in the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa at a cryogenic bath temperature of 77.4 K. In Fig. 8(a), we plot current-voltage measurements recorded while sweeping the current downwards or upwards through these jumps. Ten sets of data are plotted sweeping in each direction. The  $-35 \text{ mV}$  jump (as defined when sweeping the current downwards) that occurs at a bias current of around 65.7 mA is highly reproducible in size, occurs over a range of  $\pm 100 \mu\text{A}$ , and shows very little hysteresis compared to the jumps associated with Josephson switching in this device. The  $+10 \text{ mV}$  jump occurring at a bias current of around 63.5 mA occurs at a highly reproducible current, but the size of the voltage increase shows some variation from one measurement to the next, and this voltage increase is hysteretic—in other words, it cannot be reversed without sweeping the bias current above 66 mA again. (For clarity, these up-sweep curves are not shown in the plot.) This behavior was predicted by

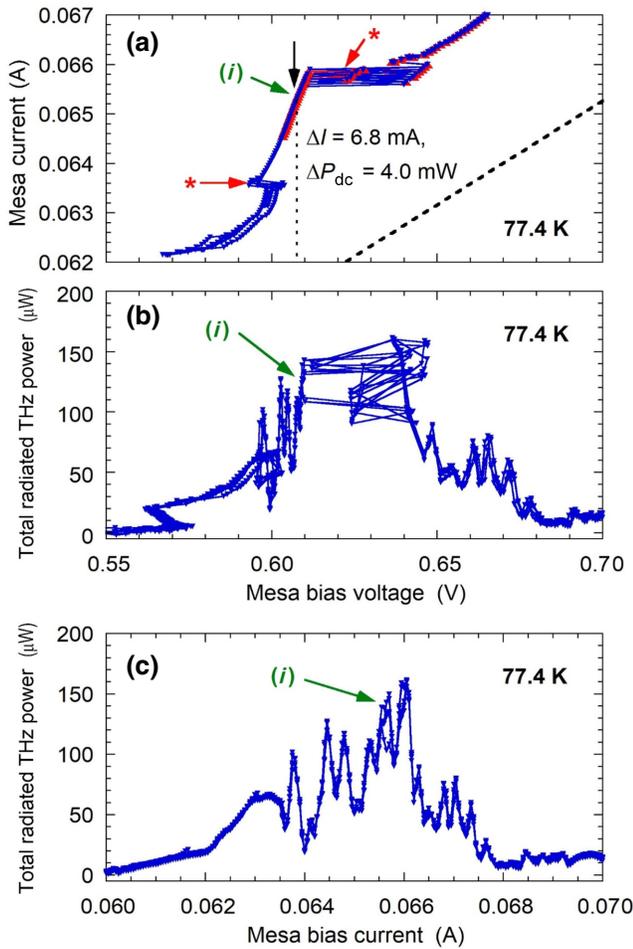


FIG. 8.  $p = 0.16$ ,  $N = 596$ ,  $300 \times 180 \times 0.92 \mu\text{m}^3$  ( $l \times w \times t$ ). (a) Plot of the unusual voltage jumps (labeled with red asterisks) seen for the  $300 \times 180 \times 0.92 \mu\text{m}^3$  mesa at a cryogenic bath temperature of 77.4 K. Blue inverted triangles are current-voltage data collected while sweeping the current downwards, while red triangles are current up-sweep data. Ten sets of data are plotted sweeping in each direction. Green label ( $i$ ) indicates the current-voltage point at which the FTIR spectrum shown in Fig. 4(d) is measured. Diagonal dashed black line indicates the tangent to the current-voltage curve, corresponding to the estimated low-bias resistance of the mesa. Vertical dotted line indicates the estimated current-voltage anomaly associated with this terahertz mode. (b) Total integrated terahertz emission power measured while sweeping the current downwards, plotted as a function of bias voltage over the same range as for (a). Data for five downward current sweeps are shown. (c) Terahertz emission power measured while sweeping the current downwards, plotted as a function of bias current over approximately the equivalent biasing range. Owing to the excellent reproducibility of our results with respect to bias current, the five data sets plotted are almost indistinguishable.

our numerical simulations performed in Ref. [28]. Since Josephson retrapping cannot account for a *positive* voltage jump, we provisionally attribute this +10 mV jump to the partial collapse of the synchronized state. Collapse of this

state may not be complete, as terahertz emission is still detected at bias currents immediately below this voltage jump.

Figures 8(b) and 8(c) show the excellent reproducibility of the terahertz emission power behavior for the same range of voltages and bias currents. Five data sets are plotted, all collected for continuous downward current sweeps. There is some variation between data sets when plotted as a function of voltage, due to the voltage jumps occurring at slightly different currents for each realization of the measurement. However, there is almost no discernable variation between these data sets when the data are plotted as a function of bias current.

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