

First-Principles Study of Monolayer *penta*-CoS₂ as a Promising Anode Material for Li/Na-ion Batteries

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Using first-principles calculations, we investigate the properties of the CoS₂ monolayer pentagonal structure as a possible anode material for Li- and Na-ion batteries. The geometrical optimization reveals that the metal atom prefers to be adsorbed on the hollow site, and that the diffusion barrier of Li and Na can be as low as 0.22 eV, which would allow a relatively fast diffusion on the surface. Moreover, our calculations demonstrate that the CoS₂ monolayer has a theoretical specific capacity of 653.31 and 326.77 mA h g⁻¹ for Li and Na, respectively, which in the case of Li is larger than the capacity of other two-dimensional materials used as an anode material. Overall, the CoS₂ monolayer in this structure shows excellent electrochemical properties, making it a promising candidate for use as the anode material in metal-ion batteries.

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I. INTRODUCTION

The environmental pollution and the increasing energy demand have pushed researchers to replace dwindling supplies of fossil fuels by developing high-performance energy-storage devices [1]. Therefore it is urgent to popularize and explore renewable energy sources with high performances in term of high-energy density, high-power density, and long cycling life. Rechargeable ion batteries are alleged to have an enormous potential, given their attractive performances of high energy density, rational working voltage, and good cyclability [1–3]. So far, a variety of systems have been explored, including lithium-ion batteries (LIBs), sodium-ion batteries (SIBs), etc. Although the performance of these batteries have largely improved in the past years, their further development is still limited [4]. Recently, considerable effort has been made to improve their performance by developing electrode materials [4,5]. The required electrode materials must not only have high electrical and thermal conductivities, but also need to have a sufficiently high strength, hardness, and heat resistance [6].

At present, considerable effort has been made to develop cathode materials with large interstitial space. Recently

and since the discovery of graphene [7], two-dimensional (2D) materials have emerged in various research fields with a wide range of applications from electronic devices [8–11], including energy-storage technologies such as batteries [4–6,12–16]. This is due to the excellent characteristics of their surfaces, which permits a fast ion diffusion and large ionic insertion channels, compared to bulk materials and other dimensional nanomaterials [17]. Numerous 2D materials have been reported as potential electrodes for energy-storage applications: for example, transition metal dichalcogenides (TMDCs) [4], transition metal carbides and nitrides (MXenes) [18–20], and transition metal oxides [21,22]. Among these, transition-metal sulfides are known to be potential candidates for anode materials due to their high theoretical capacity, versatile material types, sufficient abundance, cost effectiveness, and robust structures [4,23]. For instance, the potential of VS₂ monolayers has been demonstrated by first-principles calculations as the electrode material for Li-, K-, Mg-, and Al-ion batteries with theoretical capacities of 1397, 1863, 466, and 78 mA h g⁻¹, respectively, and a diffusion barrier of less than 0.4 eV [18]. Also, XS₂(X = Mo, Nb, Zr, and Ti) are found to have a suitable open-circuit voltage range as a Na anode material (0.49–0.95 V versus Na⁺/Na), with theoretical capacities of 260–339 mA h g⁻¹. These materials also showed Na ion migration barriers as low as 0.22 (TiS₂)

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and 0.07 eV (NbS_2) [24]. Moreover, Veenu *et al.* [25] performed first-principles calculations of Li adsorption on a Mo_2CS_2 monolayer and predicted a diffusion barrier of 0.24 eV. They also showed that Mo_2CS_2 has a large storage capacity (410 mA h g^{-1}), larger than some available electrode materials such as graphite, MoS_2 , TiO_2 , Ti_3C_2 , and Mn_2CF_2 . Also tin (Sn)-based sulfides SnS and SnS_2 materials are considered as promising anode materials for LIBs and NIBs due to their high theoretical capacities [(1137 mA h g^{-1} for LIBs and 1022 mA h g^{-1} for NIBs) and (1231 mA h g^{-1} for LIBs and 1136 mA h g^{-1} for NIBs), respectively] [26].

Cobalt is a popular metal within material science: cobalt-based materials have attracted much attention in many fields like energy storage [27], especially for Li and Na batteries [5,6]. Among them, cobalt sulfides have emerged as promising semiconductor materials and possess a wide range of applications like supercapacitors, dye-sensitized solar cells, and lithium- and sodium-ion batteries [28]. Cobalt sulfides with many phases and different stoichiometry coefficients, including CoS , CoS_2 , Co_4S_3 , and Co_8S_9 have attracted increasing attention due to their electrochemical performance (high first-cycle efficiency and a high theoretical capacity) [29–31]. Motivated by the excellent properties of CoS_2 -based electrode (high capacity, lower over potential, and better cycling stability) compared with other sulfide materials [32], and the advantages of this allotrope pentagonal structure, which is found to be mainly responsible for the high Li/Na ion-storage capacity and the fast diffusivity [33], we study, by density functional theory (DFT), the adsorption, diffusion behavior, and electrochemical performance of metal ion adsorbed in CoS_2 -pentagonal Cairo-tiled structures.

II. COMPUTATIONAL DETAILS AND METHODS

Our calculations are performed using density functional theory as implemented in the Quantum Espresso package [34]. We choose the projector-augmented-wave (PAW) method for the basis set [35]. The exchange-correlation functional is treated by the Perdew-Burke-Ernzerhof (PBE) [36] form of the generalized gradient approximation (GGA). The energy cutoff of the plane waves is set to 550 eV and the vacuum spacing between two layers is set to 20 Å in order to avoid unwanted interactions between periodic images of the layers. The Brillouin zone (BZ) is integrated using a $12 \times 12 \times 1$ Γ -centered Monkhorst-Pack k -mesh method during the relaxation [37]. The convergence criterion is set to 10^{-6} eV/cell in energy and 10^{-3} eV/Å in forces. To treat the d orbitals of the Co atoms, we adopt the GGA+ U approach formulated by Cococcioni *et al.* [38]. The effective U parameter U_{eff} is set to 3.32 eV, consistently with a previous work [39]. For the calculation of the electronic structure, the Brillouin zone is sampled with a k -point grid

TABLE I. Calculated structural parameters (a and b in Å), magnetic moment of Co atom m in Bohr magneton, band gap (eV) and relative energy ΔE in meV/unit cell of the CoS_2 monolayer for the nonmagnetic (NM), and the different magnetic configurations (FM, AFM).

Configuration	a	b	m	ΔE	E_g
NM	5.381	5.382	00	678	Metal
FM	5.337	5.431	1.107	34.00	$1.19^{\text{a}}, 1.96^{\text{b}}$
AFM	5.337	5.432	1.192	00	1.06

^aSpin-up component.

^bSpin-down component.

of $20 \times 20 \times 1$. The climbing image nudged-elastic-band (CI NEB) method [40] is used to calculate the potential energy-diffusion pathway and the minimum diffusion energy barrier of the Li/Na ions on the CoS_2 monolayer. Also a Bader charge analysis is carried out to study the charge distribution and transfer [41].

III. RESULTS AND DISCUSSION

A. Structure and electronic properties of CoS_2 monolayers

The CoS_2 monolayer displays a buckled pentagonal structure with the space group $P2_1/c$, where each pentagon is formed by two Co atoms and three S atoms. Each Co atom adopts a planar tetracoordination with four S atoms, forming a pentagonal ring network known as the Cairo pentagonal tiling, as shown in Fig. 1. For the structural optimization of the CoS_2 monolayer, both spin and non-spin polarized solutions are considered, and the results are given in Table I. It is found that the AFM structure has the lowest total energy with in-plane lattice constants ($a = 5.337 \text{ \AA}$, $b = 5.432 \text{ \AA}$), which agree well with the parameters reported by Lei *et al.* [39]. In addition, our calculated localized magnetic moments for each Co ion ($1.19 \mu_B$) and band-gap energy (1.06 eV) also coincide with earlier theoretical results given by Lei *et al.* [39]. The computed band structure and projected density of states (PDOS) of a CoS_2 monolayer are shown in Fig. 2. It can be seen that the CoS_2 monolayer is a semiconductor with an indirect band gap (X - M) of 1.06 eV.

B. Li/Na adsorption on CoS_2

In order to examine the Li/Na-adsorption capability of the CoS_2 monolayer, which is the prerequisite for using it as an electrode material in Li/Na batteries, we investigate the different lithium and sodium adsorption sites on CoS_2 . Three possible adsorption sites are examined, as shown in Fig. 1: (1) the top of a Co atom (site T1), (2) the top of a S atom (site T2), and the hollow sites above the center of the

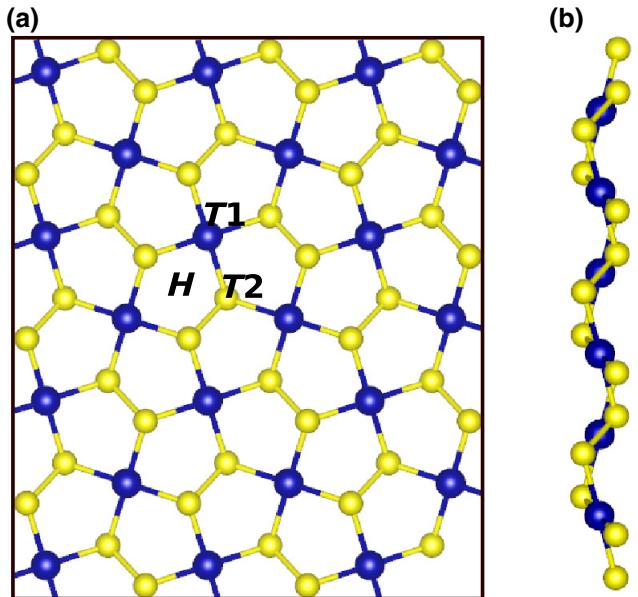


FIG. 1. (a) Top and (b) side views of the optimized $3 \times 3 \times 1$ CoS_2 monolayer. H , $T1$, and $T2$ represent three possible adsorption sites for the adatoms. The cobalt and sulfur atoms are shown with blue and yellow spheres, respectively.

pentagon (site H). We perform a full structural optimization for all the adsorption sites under both spin and nonspin polarized conditions, and a $2 \times 2 \times 1$ supercell of CoS_2 is used in our calculations. The most favorable adsorption site is identified by calculating the adsorption energy using the following equation:

$$E_{\text{ad}} = E(\text{MCoS}_2) - E(\text{CoS}_2) - E(M), \quad (1)$$

where $E(\text{MCoS}_2)$, $E(\text{CoS}_2)$, and $E(M)$ are the total energies per atom of CoS_2 with the M adsorbate, pristine CoS_2 , and the bulk metal atom, respectively.

According to Eq. (1), a negative adsorption energy represents an exothermic reaction and a strong attractive

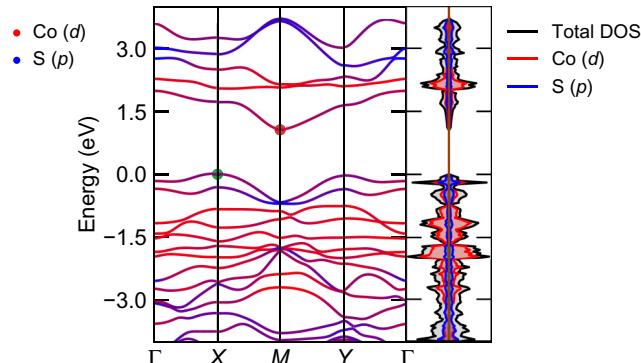


FIG. 2. Orbital-resolved band structure and PDOS of a CoS_2 monolayer.

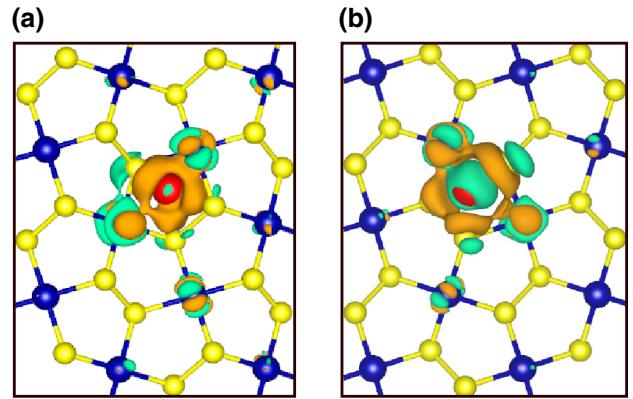


FIG. 3. Charge-density difference of (a) Li atom adsorption and (b) Na atom adsorption on the CoS_2 monolayer with an isovalue of $0.005 \text{ e}/\text{\AA}^3$. The orange and light green isosurfaces represent the electron accumulation and depletion, respectively.

interaction occurs between the adsorbate and the substrate, which indicates that this system is potentially a suitable cathode material for Li/Na-ion batteries. Our results suggest that all the above metal atoms prefer to be adsorbed above the hollow site, and the corresponding adsorption energies are -211 and -198 meV for Li and Na, respectively.

Moreover, to understand the adsorption behavior of Li/Na ions on the monolayer, the three-dimensional charge-density difference of metal/ CoS_2 is shown in Fig. 3, which is obtained as the difference between the valence charge density before and after bonding using the following equation:

$$\Delta\rho = \rho_{\text{CoS}_2\text{Li/Na}} - \rho_{\text{CoS}_2} - \rho_{\text{Li/Na}}, \quad (2)$$

where $\rho_{\text{CoS}_2\text{Li/Na}}$ and ρ_{CoS_2} denote the total electron densities of the relaxed CoS_2 monolayer with and without metal atoms, respectively, and $\rho_{\text{Li/Na}}$ is the total electron density of the isolated metal atom. In Fig. 3, the orange and light green areas represent the electron accumulation and depletion zones, respectively, and it is found that the metal atoms behave as charge donors. Then a Bader analysis is performed to quantify the charge transfer by using a $250 \times 250 \times 400$ grid for the electron density. It is found that approximately $0.93e$ and $0.89e$ are transferred to the substrate from the adatoms Li and Na, respectively, implying that the Li interacts more strongly with CoS_2 than Na. These findings are in accordance with the value of the adsorption energy.

To further understand the interaction between the metal atom and the CoS_2 monolayer, we calculate the electronic structure of the resultant compound. Figure 4 shows the computed PDOS for these metal-adsorbed systems. It is seen that the CoS_2 monolayer becomes metallic after the adsorption. This transition from a semiconductor to a metal

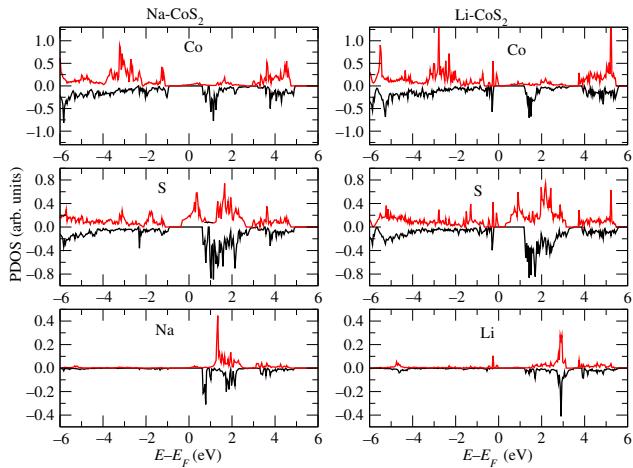


FIG. 4. PDOS when the Li and Na atoms are adsorbed on the CoS_2 monolayer.

is due to the charge transfer from the adatoms to CoS_2 . This metallic behavior ensures a good electrical conductivity, which is a desired condition for a battery electrode candidate. The same phenomenon has also been observed in other 2D semiconducting electrode materials such as GaN [42], PC [43], and MX ($M = \text{Ge}, \text{Sn}; X = \text{S}, \text{Se}$) [44].

C. Ion diffusion pathway and energy barrier

The value of the energy barrier represents a useful parameter to evaluate the charge and discharge rates of a rechargeable battery. The calculated diffusion pathway from one stable site to the neighboring stable site is obtained by taking nine images through the diffusion pathways on a $3 \times 3 \times 1$ supercell of the CoS_2 monolayer. The structures of all images are relaxed with a force convergence criterion set up to $10^{-3} \text{ eV}/\text{\AA}$.

The energy barrier E_b is calculated as the energy difference between the transition state (E_{TS}) and the initial state (E_{IS}): $E_b = E_{\text{TS}} - E_{\text{IS}}$. In our case, two possible diffusion pathways are considered between two neighboring adsorption sites, denoted as P1($\text{H} \leftrightarrow \text{H}$) and P2($\text{H} \leftrightarrow \text{T2} \leftrightarrow \text{H}$). The corresponding diffusion energy profiles along the optimized pathway for the two metal atoms (Li, Na) are shown in Fig. 5. The nonsymmetrical energy barrier is due to numerical inaccuracies. For the two systems, P1 is the most favorable diffusion pathway, which possesses the lowest diffusion barrier, that is 0.15 eV for Na and 0.21 eV for Li. These barrier values are comparable to that of VS_2 (0.22 eV) [45] and B_2S [46], and lower than typical 2D anode materials, such as graphene (0.31 eV) [47], MoC_2 (0.266 eV) [48], borophene (0.60 eV) [49], MoN_2 (0.78 eV) [50], Mo_2CS_2 (0.24 eV) and other functionalized MXenes $\text{Ti}_3\text{C}_2\text{F}_2$ (0.36 eV), $\text{Ti}_3\text{C}_2(\text{OH})_2$ (0.90 eV) [25].

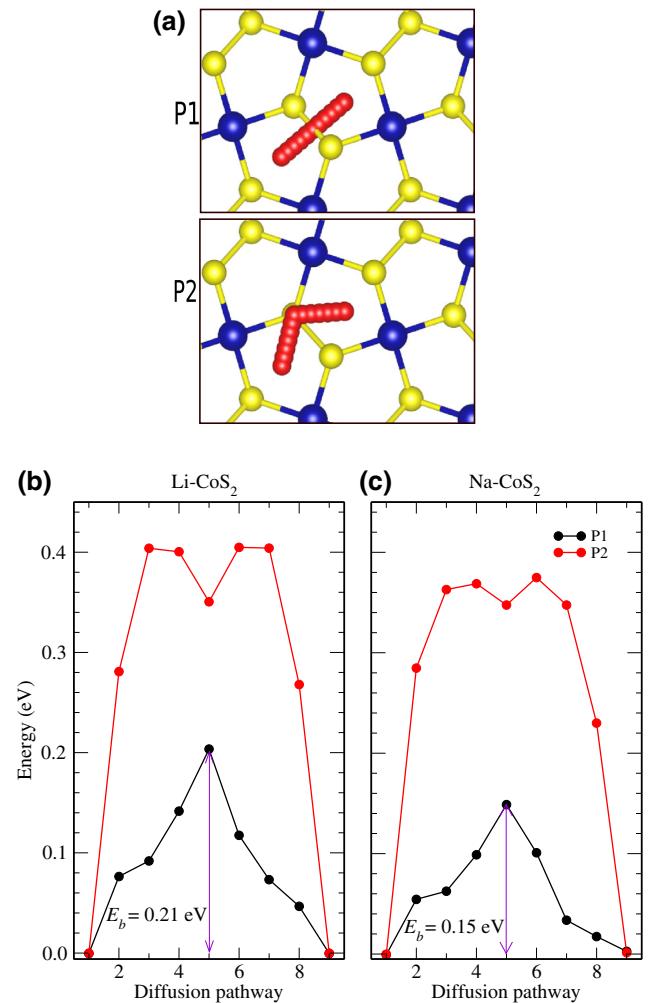


FIG. 5. (a) Top view of the two possible migration paths of Na and Li on the CoS_2 monolayer and (b) the corresponding diffusion barrier profiles.

The diffusion coefficient of the metal atoms on the CoS_2 monolayer can be also evaluated on the basis of transition state theory. By neglecting the change in vibrational free energy, the diffusion coefficient (D) can be calculated by [51,52]

$$D = d^2 \omega / \exp \left(\frac{E_a}{k_B T} \right), \quad (3)$$

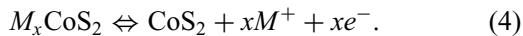
where d is the diffusion distance, E_a is the activation energy of the diffusion, k_B is Boltzmann's constant, T is the absolute temperature, and ω is the attempt frequency with ω taken as 10^{13} Hz [51,53]. According to our calculations, the diffusion coefficient at 300 K for Li and Na on the CoS_2 monolayer is 2.9×10^{-6} and $3.13 \times 10^{-5} \text{ cm}^2/\text{s}$, respectively. Because of the relatively smaller activation energy of the Na atom compared to the Li atom on CoS_2 monolayer, the diffusion is predicted to be faster

in NaCoS_2 at room temperature. Our calculated diffusion coefficients are in qualitative agreement with the one for the $\text{Na/Li-Nb}_2\text{S}_2\text{C}$ monolayer [54], Na/Li-TiSSe [51] and larger than that of graphene, C_2N and $\text{C}_2\text{N}/\text{graphene}$ [52,54].

D. Open-circuit voltage and storage capacity of Li and Na on the CoS_2 monolayer

The storage capacity and open-circuit voltage are two helpful parameters for the performance of electrode materials. We increase the number of adsorbed metal atoms on the CoS_2 monolayer in the $2 \times 2 \times 1$ supercell by constructing a series of configurations with the chemical stoichiometry $M_x\text{CoS}_2$. Here, the different numbers of metal atoms are placed on both sides of the CoS_2 monolayer. Firstly, we begin by adding the metal atom one by one on the side of the most stable site to form the first layer, when the sites are occupied on this side, we follow on the other side of the CoS_2 to form the second layer. After the complete occupation of stable sites of the CoS_2 monolayer, the adsorption will occur on the top of the next stable site to form the third and fourth layers of adsorption etc., as shown in Fig. 6(a). For each structure, the ground state is obtained after testing all possible adsorption sites into consideration.

The charge-discharge processes can be described by the following chemical reaction:



The open-circuit voltage can be obtained by computing the total energies before and after Li/Na intercalation. To obtain accurate energy values, both the lattice parameters and atomic positions are fully relaxed for the intercalated configurations. By neglecting the volume and entropy effects, the open-circuit voltage (OCV) can be estimated by [55,56]

$$\text{OCV} \simeq \frac{E(\text{CoS}_2) + xE(M) - E(M_x\text{CoS}_2)}{xe} \quad (5)$$

where $E(M_x\text{CoS}_2)$, $E(\text{CoS}_2)$, and $E(M)$ are the total energies of CoS_2 with the M adsorbate, pristine CoS_2 , and bulk metal atom, respectively.

The plot of the voltage and metal atoms' content are depicted in Fig. 6(b). It is seen that the OCV gradually decreases with the increase of the number of adsorbed metal atoms because of the enhanced repulsion among metal cations. The evolution of the OCV exhibits two effective plateaus corresponding to the first and second layers of sodium adsorptions. However, for Li adsorption, there are three plateaus. The positive OCV demonstrates that the accommodation of one-layer metal atoms on the CoS_2 monolayer is energetically favorable for both Na and Li metals. For the second layer, the accommodation

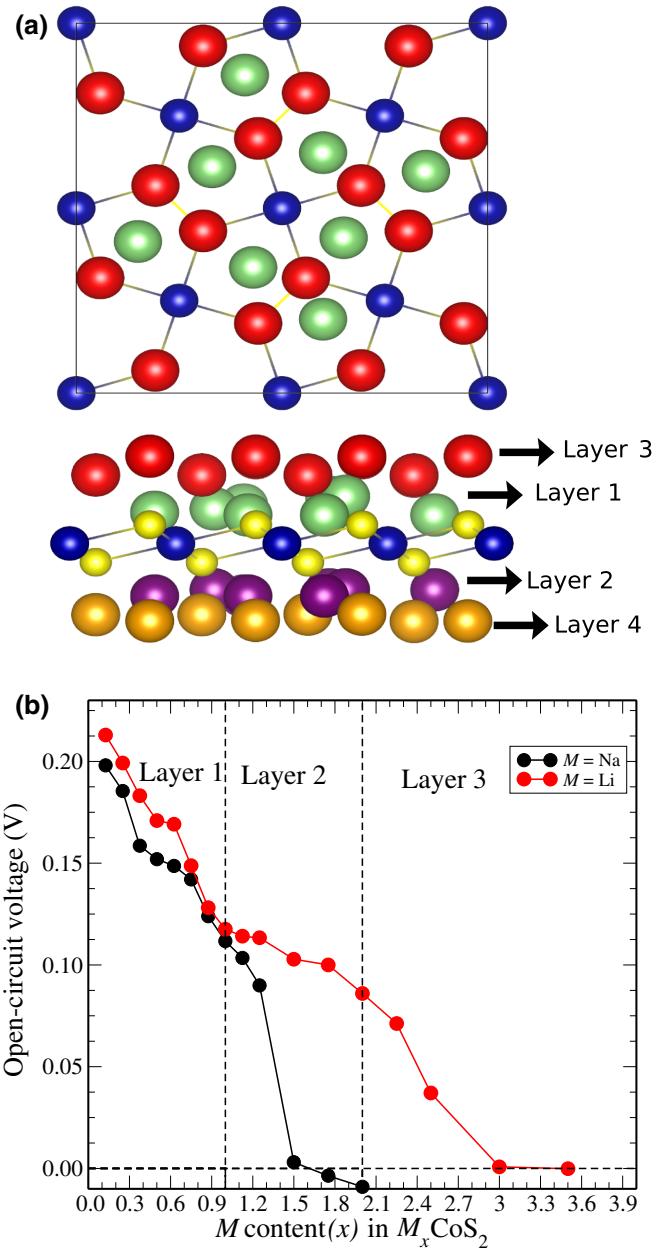


FIG. 6. (a) The top and side view configurations of the adsorbed metal layer on CoS_2 monolayer. (b) The OCV profile with metal ions adsorbed. The red, green, purple, and golden rod represent the metal.

of Li is still energetically favorable, however, the adsorption of Na atoms becomes unstable when x is larger than 1.5, owing to the strong repulsive interaction caused by the Na cations. The corresponding chemical stoichiometry are Li_3CoS_2 and $\text{Na}_{1.5}\text{CoS}_2$ with an average OCV in the range of 0.0–1.0 V. These values are comparable to some typical electrode materials such as Mo_2C , MoC_2 , MoC , Nb_2C , MoN_2 , and Ti_3C_4 [48,50], demonstrating the feasibility of a monolayer of CoS_2 as a 2D anode for rechargeable ion batteries.

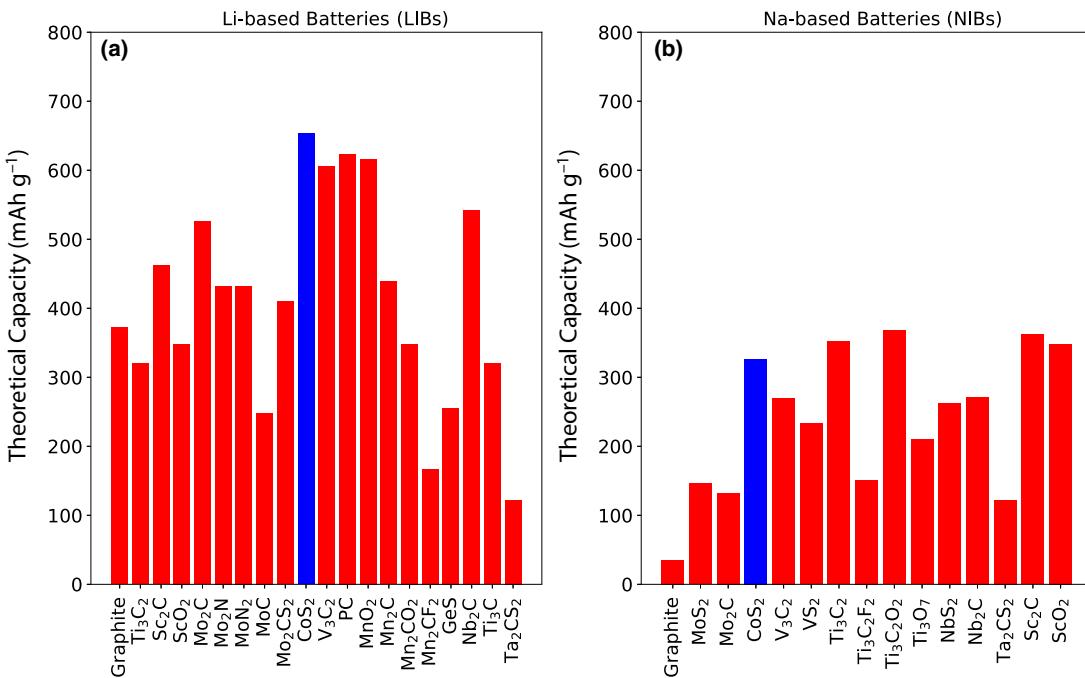


FIG. 7. Comparison of the theoretical capacity of Li/NaCoS₂(blue bar) with some selected 2D material-based electrodes for SIBs and LIBs(red bar). The values are taken from Refs. [4,5,21,25,43,45,47,50,51,57–61].

The specific capacity can be estimated as

$$C = \frac{xF}{m}. \quad (6)$$

Here m is the molecular weight of CoS₂ (123.065 g mol⁻¹), F is the Faraday constant (26.8 A h mol⁻¹), and x is the highest concentration of M ion in M_x CoS₂. The obtained maximum capacity of Li and Na are 653.31 and 326.77 mA h g⁻¹, respectively. For comparison the Li and Na storage capacities for several 2D-based structures are illustrated in Fig. 7: our calculated value for LiCoS₂ is larger than those of some other predicted 2D electrode materials, while for NaCoS₂, our calculated value is larger than most of the other known compounds but lower than the one of Sc₂C, Ti₃C₂, and Ti₃C₂O₂.

E. Conclusion

In summary, using density functional theory, we study the structural, electronic, and electrochemical properties of the CoS₂ monolayer in the pentagonal structure. We find that the metal atoms (Li or Na) prefer to absorb on the hollow site, and that this adsorption is accompanied by a large charge transfer to the CoS₂ monolayer, which becomes metallic, ensuring a good electronic conductivity. Then using the CI NEB method, the value of the migration energy barriers on the CoS₂ monolayer is found to be suitable for its use as the anode material in LIBs and NIBs. Furthermore, the specific storage capacity is found to be 653.31 and 326.77 mA h g⁻¹ for Li and Na, respectively,

which is in the case of Li, much larger than that of other 2D materials used as anodes. Overall, our results suggest that this allotrope of CoS₂ is a very promising compound to be used as the anode in Li-ion batteries. We hope that our work will stimulate further theoretical and experimental researches on this compound.

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