# Plasma cleaning of the SLAC Linac Coherent Light Source II high energy verification cryomodule cavities

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Plasma cleaning is a technique that can be applied in superconducting radio-frequency cavities *in situ* in cryomodules to decrease their level of field emission (FE). We developed the technique for Linac Coherent Light Source II (LCLS-II) cavities and we present in this paper the full development and application of plasma processing to the LCLS-II High Energy verification cryomodule (vCM). We validated our plasma processing procedure on the vCM, fully processing four out of eight cavities of this CM, demonstrating that cavity performance was preserved in terms of both accelerating field and quality factor. Applying plasma processing to this clean, record breaking cryomodule also showed that no contaminants were introduced in the string, maintaining the vCM FE-free up to the maximum field reached by each cavity. We also found that plasma processing eliminates multipacting (MP)-induced quenches that are frequently observed within the MP band field range. This suggests that plasma processing could be employed *in situ* in CMs to mitigate both FE and MP, significantly decreasing the testing time of cryomodules, the linac commissioning time and cost, and increasing the accelerator reliability.

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#### I. INTRODUCTION

The term plasma cleaning refers to a process in which impurities are removed from a surface via a mixture of inert and reactive gases in the form of plasma. This technique, also called plasma processing, was developed for superconducting radio-frequency (SRF) cavities some years ago by Oak Ridge National Laboratory (ORNL). The developers demonstrated that a plasma composed of a mixture of neon and a small percentage of oxygen reduced hydrocarbon-related field emission (FE) in the spallation neutron source cavities [1–4]. Starting from this successful experience, plasma processing studies are being conducted at multiple laboratories for different accelerating structures [5–12].

Plasma processing for 1.3 GHz SRF cavities was developed at Fermi National Accelerator Laboratory (FNAL) [5–7], in collaboration with SLAC National Accelerator Laboratory and ORNL with the purpose of mitigating FE [13,14] in the 1.3 GHz Linac Coherent Light Source II (LCLS-II) [15] and LCLS-II High Energy (LCLS-II-HE) [16] cavities.

As discussed in Berrutti *et al.* [5], a new method of plasma ignition was developed for LCLS-II 1.3 GHz cavities. This innovative method utilizes the higher order modes (HOM) and HOM couplers (Fig. 1) to ignite and move the glow discharge inside the cavity rf volume. A systematic study of plasma processing applied to LCLS-II 1.3 GHz nitrogen-doped [17] cavities was carried out at FNAL, demonstrating that plasma cleaning can successfully mitigate hydrocarbon-related FE without affecting the high Q-factors and quench fields characteristic of N-doped cavities [7].

In this paper, the development and first application of plasma processing to a full size 1.3 GHz cryomodule (CM) are presented. The cryomodule under study is the LCLS-II-HE verification cryomodule (vCM). Before plasma processing, the vCM underwent extensive tests, and the results are reported in Posen *et al.* [18]. The LCLS-II-HE project

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FIG. 1. Photo of TESLA shape nine-cell 1.3 GHz cavity, with highlighted ports for the higher order modes (HOM) couplers, the fundamental power coupler (FPC), and the pickup (PU) coupler.

requires 1.3 GHz nine-cell niobium cavities operating at quality factor  $Q_0 > 2.7 \times 10^{10}$  and accelerating gradient  $E_{\rm acc} \approx 21$  MV/m. To reach these specifications, a new N-doping treatment was developed, called "2/0" since cavities are treated with nitrogen for only 2 min and subsequently cooled down in the furnace [19]. The vCM showed record performance in terms of both quality factor and accelerating gradient, exceeding the project specification. Particularly relevant for plasma processing, only one cavity (CAV5) exhibited detectable radiation (0.6 mR/h), but the FE source was processed during later testing, leaving the vCM FE-free. Nevertheless, the vCM represented a unique opportunity to scale the 1.3 GHz plasma processing technique from a single TESLA-style [20] cavity to an entire cryomodule. For this reason, we plasma processed four out of eight cavities of the vCM and we compared the rf cavity performance before and after plasma processing. The details of the development of the technique from the single cavity system to the cryomodule are presented, together with the risk analysis and mitigation strategy that was conducted. We also summarize the plasma processing plan, the details of the experimental system, and the rf test results. Of particular interest is the multipacting (MP) quench analysis that was conducted before and after plasma processing. The four cavities that underwent plasma cleaning did not show any MP quenches, demonstrating that this treatment can fully eliminate MP in cavities in cryomodules. Therefore, performing plasma processing after a cryomodule has been assembled may significantly decrease its overall testing time and the subsequent accelerator commissioning time, thus reducing the overall project cost. In addition, cryomodules with MP-free cavities would offer greater stability during the accelerator operation, increasing its reliability.

## **II. RISK ANALYSIS AND MITIGATION**

All major risks were identified before deploying plasma processing in a fully assembled and tested cryomodule. The main risks that were identified are the following: (i) potential damage to the HOM cables, connectors, and feedthroughs in the module due to the rf power passing through them during plasma ignition; (ii) the potential ignition of the plasma in the fundamental power coupler (FPC) and consequent damage of the FPC; and (iii) potential pressure instabilities throughout the cryomodule.

To verify the first point, the HOM cables were tested extensively by subjecting them to different levels of input power, from 10 to 125 W. These tests were carried out offline, on a test bench setup with a single nine-cell cavity, and the cables were not placed in a vacuum environment. The heat and temperature distributions along the input and transmitted power cables were monitored through temperature sensors installed on the cable connectors and at the cable center point. Another set of temperature sensors was installed on the cavity HOM feedthroughs in which there are fragile components, such as the ceramic window. An image of the temperature sensors installed on a cable end and HOM feedthrough is shown in Fig. 2. The HOM1 cable was used as the input line, and the HOM2 cable was used as the transmitted line, following the same configuration used for plasma processing in the CM, where the HOM1 cable is connected to the HOM coupler on the FPC side, and the HOM2 is connected to the HOM coupler on the pickup side. The rf characteristics of the cables were measured before and after each test (i.e., after each input power from 10 to 125 W), and both the reflected  $(S_{11})$  and transmitted  $(S_{21})$  signals were recorded. The measurements of both  $S_{11}$ and  $S_{21}$  did not change after any of the tests, therefore, the cable performance was not affected by the rf power. The frequency chosen for these tests was 1829 MHz. This frequency was chosen to be close to the mode used to ignite the plasma in LCLS-II cavities and to maximize the amount of reflected power within an MHz range from the cavity rf mode. This ensured exceeding the reflected power and the heat dissipated in the cables under regular plasma conditions: during cavity processing, power is sent to the resonator at its mode frequency, thus minimizing reflected power. The most significant cable test was conducted at 100 W of input power for 30 min: 100 W indeed corresponds to the power required for plasma ignition. That level of power is needed only for a few seconds at ignition, and 10-20 W of rf power is sufficient to sustain the



FIG. 2. Temperature sensors installed on the cable end (left,  $90^{\circ}$  connector) and cavity HOM feedthroughs (right).



FIG. 3. Power versus time during cable testing for LCLS-II plasma processing: input power of 100 W for 30 min. The input and reflected power,  $P_{\rm IN}$ ,  $P_{\rm R}$ , respectively, follow the left y axis, while the right y axis refers to the transmitted power  $P_{\rm T}$ .

plasma afterward. Details regarding input, reflected, and transmitted power during plasma cleaning are discussed in Sec. IV. The input power versus time is shown in Fig. 3 and the temperature profile is presented in Fig. 4. The maximum cable temperature increase, of



FIG. 4. Temperature of HOM1 and HOM2 cables and feedthroughs versus time during cable testing for plasma processing: input power of 100 W for 30 min. HOM1 denotes the cable connected to the HOM FPC side and used as an input line, HOM2 denotes the cable connected to the HOM pickup side and used as the transmitted line. For both HOM1 and HOM2 cables, the 90° connector (shown in Fig. 2, left) is the side connected to the HOM coupler, while the straight connector is the opposite end of the cable. The HOM-FPC and HOM-PU copper connectors denote the sensors placed on the HOM antenna feedthroughs, shown in Fig. 2, right.

approximately 40 K, does not damage or change the rf performance of the cables. One can observe that the central point of the HOM1 input cable is the hottest spot, followed by the 90° connector (connected to the HOM FPC antenna), the HOM-FPC copper connector (i.e., the feedthrough of the input HOM), and finally the straight connector (connected to the power amplifier). As expected, the difference in cable temperature between the input (HOM1) and the transmitted (HOM2) power sides is quite large: the HOM2 cable temperature increase is well below 5 K.

To verify whether plasma ignition in the FPC represented a risk during plasma processing in a cryomodule, an experiment was carried out in a nine-cell cavity on a test bench: a variable coupler was installed on the cavity, with an identical antenna tip design and coupling to the cavity as the production FPC for cryomodule operation. The cavity was then plasma processed following the standard procedure. The results of this test were encouraging because plasma ignition in the coupler volume was not observed. In addition, after the test, both the input HOM antenna and the FPC were optically inspected and found to be free of any kind of damage or discoloration.

Last, to minimize the pressure fluctuations along the beamline, the vacuum portion of the cryomodule, a new vacuum cart was assembled and tested on a single cavity prior to plasma processing of the vCM.

#### III. EXPERIMENTAL SETUP AND PLASMA PROCESSING PROCEDURE

Plasma cleaning of the vCM took place at the Cryomodule Test Facility (CMTF) at FNAL after the completion of extensive rf tests, the results of which are discussed in Posen *et al.* [18] and the subsequent warm-up to room temperature. The vCM was rf tested again after plasma cleaning to assess the cavity performance changes after plasma processing.

Plasma cleaning was carried out in four out of the eight cavities of the vCM. This enabled the ability to (i) compare the performance of cavities subjected (and not subjected) to plasma cleaning and to (ii) perform cleaning within the time constraint imposed by the project. To monitor the cavity's surface temperature during processing, the four cavities with additional temperature sensors installed inside the helium vessel were chosen for plasma processing, namely, CAV1, CAV4, CAV5, and CAV8 (the number indicates the position along the cavity string). These four cavities were instrumented with four temperature sensors placed on the outside wall of the cavity on cells no. 1 and no. 9 at the top and bottom positions. The same cavities were also instrumented with temperature sensors on the HOM connectors, offering the advantage of being able to monitor the temperature of both higher order mode cables. In addition, CAV1 and CAV5 each have two more temperature sensors placed on the outside of the helium vessel.

During plasma cleaning, a mixture of neon and oxygen in gaseous form is injected into the cavity string, and a constant flow is maintained from CAV1 to CAV8. The glow discharge is ignited in one cell at a time following the direction of the gas flow; the cavities were also processed following the direction of the gas flow from CAV1 to CAV8. As explained in Giaccone *et al.* [7], two identical passes of plasma cleaning are applied to each cavity. Due to uncertainty regarding the time available to apply plasma processing to the vCM, CAV1 subsequently received the two runs of plasma processing, meaning that as soon as the first pass was completed, the second started. CAV4, 5, and 8 instead received the first plasma pass one after the other, then the second run restarted from CAV4 to 8.

In addition to the cavity and HOM cable temperature, the pressure of the gas mixture, the partial pressure of oxygen and the plasma byproducts (C, CO, CO<sub>2</sub>, H<sub>2</sub>O), the power levels (input, reflected, and transmitted), and the cavity spectrum were monitored.

# **IV. DATA ANALYSIS AND DISCUSSION**

The vCM cavity string was maintained under a high vacuum during and after the connections to the plasma processing gas and vacuum carts. While at room temperature the insulating vacuum in the CM, around the cavity string, was released bringing the CM to atmospheric pressure. Before starting the cavity plasma cleaning, the gas flow was established through the cavity string and the pressure was slowly increased up to the working point ( $\approx$ 134 mTorr). Once the pressure stabilized, the pressure drop across the two ends of the vCM was measured to be  $\approx$ 24 mTorr. The relative concentrations of neon and oxygen in the gas mixture were monitored with the



FIG. 5. Experimental setup used to apply plasma processing to the LCLS-II-HE vCM. The gas station is connected to one end of the CM (a) and the vacuum and RGA station to the other end (b). The rf rack and computers are located in between the gas and vacuum stations (c).

RGA and adjusted as necessary by operating the gas cart leak valves controlling the mixture of the pure neon and the 20% oxygen balance neon cylinders. Once the  $O_2$  partial pressure stabilized to the desired level (between 1.5% and 2%), the glow discharge was ignited inside the cavity rf volume.

The RGA data collected during the first run of plasma processing on CAV1 are shown in Fig. 6. Once the glow discharge ignites inside the cavity, the carbon-related signals (CO, CO<sub>2</sub>, C) increase as shown in Fig. 6. The subsequent peak corresponds to the transfer of the plasma and the tuning of its density in cell no. 1. Concurrently, a small decrease in the O<sub>2</sub> level can also be observed, confirming that the reactive oxygen ions present in the glow discharge are creating volatile byproducts with the hydrocarbons knocked off the cavity surface. However, as shown in Fig. 6, not all cells showed a significant change in the RGA signals after the plasma was transferred. Only in this first plasma cleaning run was the processing conducted from cell no. 1 to 9, while in all subsequent cleaning runs, it was conducted from cell no. 9 to 1, following the direction of the gas flow. Figure 5 shows the experimental setup used to apply plasma processing to the vCM.

At the end of each plasma cleaning session (usually corresponding to the end of the day), the gas flow was stopped, and the cavity string was evacuated using the plasma vacuum cart. Then, the cavity string was isolated by manually closing the two right angle valves connecting the string to the gas and vacuum carts. During these



FIG. 6. Plot of RGA signals measured on CAV1, during the first day of plasma processing from cell no. 1 to no. 5 and cell no. 9. As the gas flow is started and the pressure in the vCM increases, there is an initial peak in the oxygen signals (at 16 and 32 amu), which then decreases and stabilizes to  $\approx 1.6\%$  of the neon content. The partial pressures measured by the RGA differ from the total pressure in the cavity string since the RGA is positioned on a bypass line near a turbo pump on the vacuum cart to allow the RGA filament to sample the gas composition. A leak valve is used to regulate the gas flow through the RGA.



FIG. 7. Pressure measured at the two ends of the vCM during initial neon-oxygen mixture injection and during the flow maintained for the entire duration of plasma processing. The pressure drop across the string is approximately 24 mTorr. The pressure values are scaled to account for the gauge sensitivity to neon. These data were acquired during plasma processing of CAV4.

downtimes, the pressure in the cavity string was monitored using beamline pressure gauges.

Figures 7 and 8 show a complete set of data collected during the first plasma cleaning run on CAV4. This cavity is used as an example to explain the data monitored during plasma processing and their typical trends and values. Figure 7 shows the pressure data acquired by the gauges at the two ends of the cryomodule at the gas injection and vacuum sides. This graph shows the pressure increase at both ends as the flow of gas is initiated into the beamline, the pressure stabilization phase, and the pressure drop due to the beamline evacuation at the end of the processing. The pressure difference between the gas injection side and the vacuum side indicates the pressure drop between the two ends.

Figure 8 shows different signals recorded during plasma processing, and it also illustrates a rare case in which plasma ignition occurred at the HOM. The ignition of the plasma at the coupler, which we refer to as HOM ignition, rather than in the cavity volume can be recognized from multiple features in the data: the RGA always shows an immediate and sharp peak in CO and CO<sub>2</sub>; the reflected power would normally increase when the plasma is correctly ignited in the cavity volume, while in the case of HOM ignition, it does not show an abrupt change; in the case of HOM ignition, the transmitted power decreases drastically, and the  $S_{21}$  measurement on the vector network analyzer (VNA) fails to measure the cavity spectrum and only noise is observed. Figure 8(a) shows the partial pressure as a function of the time of some of the RGA signals measured on CAV4. Only the signals relevant for plasma processing are plotted: oxygen, neon, carbonrelated peaks, and water. The initial peak in oxygen is always observed by the RGA as soon as the gas flow is established in the cavity string. This peak quickly decreases, and the oxygen concentration reaches equilibrium. The first peak in CO,  $CO_2$  is due to the plasma being accidentally ignited at the HOM coupler rather than inside the cavity volume. This event lasted only a few seconds [5 s according to the power levels shown in Fig. 8(b)]. As reported in the next section, the results of the rf test after plasma cleaning confirmed that this event exerted no negative effect on the cavity performance. However, it is best practice to avoid HOM ignition, as prolonged plasma ignition in the coupler could cause the formation of an arc between the antenna and the cavity and could potentially cause sputtering of particles into the cavity. For this reason, after the first HOM ignition, we collected additional  $S_{21}$ ,  $S_{11}$ , and  $S_{22}$  measurements on CAV4 with the VNA to optimize the plasma ignition technique for this particular cavity by matching the rf frequency to the peak frequency measured through  $S_{11}$  rather than  $S_{21}$ . The additional measurements provided a better understanding of the coupling of the HOM1 and HOM2 antennas to the 2D-1 mode (second dipole pass-band, first resonant mode) used for plasma ignition and allowed us to successfully ignite the plasma inside the cavity volume, as shown by the small, broad peak in CO, CO<sub>2</sub> in the RGA plot [approximately 11:00 in Fig. 8(a)]. No other cell in this cavity showed peaks in the carbon-related signals when the plasma was ignited.

The input, reflected, and transmitted power,  $P_{\rm IN}$ ,  $P_{\rm R}$ ,  $P_{\rm T}$ , respectively, measured during the first run of plasma processing on CAV4 are shown in Fig. 8(b). The first peak at approximately 9:30 corresponds to HOM ignition, which was reached at  $P_{\rm IN} = 94.2$  W,  $P_{\rm R} = 19.1$ W, and  $P_{\rm T} =$ 1.3 mW and lasted  $\approx$ 5s. Just before 11:00, the plasma was ignited inside the cavity volume in cell no. 5, requiring  $P_{\rm IN} = 67.3$  W,  $P_{\rm R} = 14.9$  W, and  $P_{\rm T} = 1.4$  mW. Once ignited using mode 2D-1, mode 1D-5 (first dipole passband, fifth resonant mode) is quickly added using the second rf signal generator, and the first high power mode is turned off. At this point, the plasma is maintained using only  $P_{\rm IN} = 6.5$  W. The following cluster of peaks at approximately 11:00 corresponds to the plasma being transferred from cell no. 5 to cell no. 9 and tuned to the desired plasma density. After almost an hour, just before 12:00, the plasma is detuned, meaning that its density is reduced, and it is moved to the following cell, where its density is once again maximized. This process is repeated for all the cells. Every time the plasma is moved from one cell to the other, the detection program is used to confirm the cell of ignition. Each cell is processed for approximately 1 h and typically  $P_{\rm IN} \approx 10$  W,  $P_{\rm R} \approx 5$  W, and  $P_{\rm T} \approx 0.15$  mW. When a combination of two modes is used to transfer the plasma or



FIG. 8. Experimental data collected during plasma processing of CAV4. (a) contains a plot of the RGA signals measured during plasma processing of cells no. 9 to no. 4. The peak at approximately 9:30 measured for CO and CO<sub>2</sub> indicates the accidental ignition of the plasma in the HOM coupler. The following peak at approximately 11:00 is measured when the plasma is successfully ignited inside the cavity. (b) shows the input and reflected power measured at the CAV4 HOM1 coupler and transmitted by the HOM2 during plasma processing. Note the different scales on the *y* axis between the top and bottom plots. Additionally, in this case, the 9:30 peak corresponds to accidental HOM ignition, while the next peak corresponds to the ignition of the plasma in the cavity, followed by a rapid sequence of smaller intensity peaks that identify the plasma being moved from cell no. 5 to no. 9 and then tuned in cell no. 9. After approximately 1 h, the plasma density is decreased, and the glow discharge is moved to the next cell. This procedure is repeated for each cell. The temperature increase measured on the HOM1 and HOM2 connectors of all four plasma-processed cavities is shown in (c). In the period plotted here, the glow discharge was ignited in CAV4, as confirmed by the temperature increase registered by its HOM1 and HOM2 sensors. (d) shows the temperature increase measured on the outside of cells no. 1 and no. 9. The four sensors are placed inside the He vessel at the top and bottom of the end cells.

to tune its density, the total input power can temporarily reach  $P_{\rm IN} \approx 20$  W.

Figures 8(c) and 8(d) show the temperature increase measured for the HOM cable connectors and on the cavity external surface (on cell no. 1 and 9) during CAV4 plasma processing. Also, in this case, HOM1 indicates the line connected to the HOM-FPC side (used as the input line during plasma processing) and HOM2 indicates the line connected to the HOM-PU side (used as the transmission line during plasma cleaning). In Fig. 8(c), the CAV4 HOM1 sensor shows an increase in temperature corresponding first to the accidental ignition of the plasma in the HOM ( $\Delta T = 1.4$  K) and later to the ignition of the plasma in the cavity. A new spike in temperature is registered when two modes are superimposed to increase (or decrease) the plasma density at the beginning (or end) of



FIG. 9. Plots of the  $S_{21}$  (green),  $S_{11}$  (red), and  $S_{22}$  (black) parameters for some of the dipole modes used for plasma processing. The parameters were measured on CAV1-4-5-8 from top to bottom using a network analyzer, the *y* axis is in the unit of dB. (a), (d), (g), and (j) contain the 2D-1 peaks, (b), (e), (h), and (k) show the 1D-5 peaks, and (c), (f), (i), and (l) show the 1D-1, 1D-2, and 1D-3 peaks. The plots present an example of the different couplings to the HOM and the dipole splitting characteristics of each cavity. Exact measurements of the  $S_{21}$ ,  $S_{11}$ , and  $S_{22}$  are critical, particularly for mode 2D-1, to correctly identify the rf frequency to use during plasma ignition in order to minimize the probability of causing plasma ignition at the HOM coupler.

each cell processing and to move it to a new cell. As expected, the HOM2 cable connector exhibits a moderate temperature increase with no particular trend connected to the position of the plasma in the cavity. A plot of the temperature measured on the CAV4 external surface of cells no. 1 and 9 is shown in Fig. 8(d). The four sensors are placed inside the helium vessel and at the top and bottom of cells no. 1 and 9. When the plasma is ignited in cell no. 9, the temperature sensors on this cell measure a maximum  $\Delta T \approx 1$  K. As the plasma is moved farther away from cell no. 9, the  $\Delta T$  measured on these sensors decreases, while the temperature measured on cell no. 1 gradually increases up to  $\Delta T \approx 1.5$  K. Although clearly dependent on plasma processing, cable and cavity heating is not a source of concern, since the total cable heating was  $\Delta T_{HOM1} \approx 2$  K,  $\Delta T_{\rm HOM2} \approx 0.5$  K, and the maximum temperature increase on the cavity surface was  $\Delta T \approx 1.5$  K.

The plots discussed in Fig. 8 all report data collected during the first run of plasma processing applied to CAV4. All four plasma-processed cavities showed an initial peak in carbon-related signals on the RGA when the plasma was ignited for the first time inside the cavity and moved to cell no. 9. Only CAV1 showed additional peaks in cell no. 1, 4, and 5 and a small peak in cell no. 9 during the second pass of plasma cleaning. In all cases, the carbon-related peaks quickly decreased back to the background level. The plasma was accidentally ignited in the HOMs of CAV4 and CAV8. In both cases, the HOM ignition lasted only a few seconds and only occurred once. In both cavities, during the second attempt, the glow discharge was successfully ignited inside the cavity rf volume by carefully tuning the rf driving frequency to match the  $S_{11}$  2D-1 frequency. The plasma processing method for LCLS-II cavities uses higher order modes, which confers the advantage of offering good coupling even at room temperature, making it possible to ignite the plasma with a few 10 W and then maintain it with only 5-10 W. On the other hand, contrary to the fundamental passband, the HOMs are not tuned. The HOM couplers are only tuned so that their notch filters reject the fundamental mode. This lack of tuning of higher order modes implies that the frequencies of each higher order mode, the splittings in dipole modes, and the couplings to the HOM antennas vary from cavity to cavity. In the four cavities processed in the vCM, different configurations of degenerate dipole splittings were found, as shown in Fig. 9. Measurements of the  $S_{11}$  and  $S_{22}$ parameters for the HOMs allowed us to identify the exact frequency to send in input to each cavity to successfully ignite the plasma in the central cell using mode 2D-1. This, along with the work conducted during the R&D phase, demonstrated that the ignition of the plasma through the HOM is a robust method that guarantees successful ignition and control of the glow discharge inside the cavity rf volume.

Once the processing of the four cavities was completed, the gas cart was isolated from the string, which was evacuated up to  $\approx 1 \times 10^{-6}$  Torr using the plasma vacuum cart, and then the CMTF ion pumps were turned on. Both gas and vacuum carts were later disconnected from the vCM using a portable class-100 cleanroom.

### V. COMPARISON OF THE vCM PERFORMANCE BEFORE AND AFTER PLASMA PROCESSING

After plasma processing and evacuation, the vCM was cooled down to be rf tested again. The purpose of this test was to verify that plasma processing did not cause any deterioration in cavity performance and to understand if it had an impact on MP. The plan was to remeasure the quality factor, the maximum and the usable gradient for each cavity, the presence of x rays and dark current, and the stability of operations at 20.8 MV/m. The time necessary to process MP was also investigated.

	Before plasma processing				After plasma processing			
Cavity	Maximum E <sub>acc</sub> (MV/m)	Usable E <sub>acc</sub> (MV/m)	$Q_0$ at 21 MV/m × 10 <sup>1</sup>	MP <sup>0</sup> quenches	Maximum E <sub>acc</sub> (MV/m)	Usable E <sub>acc</sub> (MV/m)	$Q_0$ at 21 MV/m × 10 <sup>10</sup>	MP quenches
1	23.4	22.9	3.0	Yes	23.8	23.3	3.4	No
2	24.8	24.3	3.0	Yes	25.2	24.7	3.2	Yes
3	25.4	24.9	2.6	Yes	26.0	26.0	3.4	Yes
4	26.0	26.0	3.2	Yes	26.0	26.0	3.2	No
5	25.3	24.8	2.9	Yes	25.5	25.0	2.8	No
6	26.0	25.5	3.4	Yes	26.0	26.0	3.2	Yes
7	25.7	25.2	3.4	Yes	25.9	25.4	3.3	Yes
8	24.4	23.9	2.7	Yes	24.7	24.2	2.6	No
Average	25.1	24.7	3.0		25.3	25.1	3.1	
Total	209	205			210	208		

TABLE I. Comparison of vCM performance measured before and after plasma processing. The four cavities to which plasma processing was applied are highlighted in bold.

The results of the vCM prior to plasma processing are reported in Table I, along with the comparison with the performance measured after plasma cleaning. For all eight cavities, it was verified that the performance was preserved in terms of both quality factor and accelerating gradient. This demonstrates that plasma processing did not introduce any observably detrimental contamination or particulates inside the cavity string, maintaining the vCM FE-free. After plasma processing, the CM test confirmed the total module voltage of 210 MV versus the 173 MV required by the LCLS-II-HE specification. The vCM average quality factor still exceeded the specification ( $Q_0 = 2.7 \times 10^{10}$ at 20.8 MV/m), with an average of  $Q_0 = 3.1 \times 10^{10}$ . As explained in Posen et al. [18], the administrative limit for the gradient, in case the ultimate quench field of the cavity was not reached, was set to 26.0 MV/m while the maximum gradient defines the field level at which a cavity would quench consistently without allowing any further field increase. The usable gradient instead is defined as the maximum field at which (i) the cavity can operate for more than 1 h without quenching, (ii) the radiation remains below 50 mR/h, and (iii) the cavity is 0.5 MV/m below its ultimate quench field limit.

As shown in Figs. (10) and (11) from Posen *et al.* [18], during the first vCM rf test, all eight cavities suffered from MP induced quenches during both the power rise to 21 MV/m (Fig. 10 from Ref. [18]) or the long duration operation at a high gradient (Fig. 11 from Ref. [18]). As highlighted by the last column in Table I, after plasma cleaning, the four cavities that were not processed still showed MP quenches, while the four processed cavities did not experience any MP-related quench.

In Table II, the numbers of quenches attributed to MP for each cavity, as measured during three different cooldowns, are reported. The first two cooldowns were both measured during the initial test of the vCM. No data of MP quenches

TABLE II. Number of MP quenches measured for each cavity in two different cooldowns prior to plasma processing and the one cooldown after plasma cleaning. The four cavities to which plasma processing was applied are highlighted in bold.

	Ν		
	Befor	After plasma processing	
Cavity	First cooldown	Second cooldown	
1	not available	157	0
2	135	106	205
3	41	44	53
4	68	3	0
5	10	16	0
6	46	7	69
7	68	33	82
8	128	108	0

for CAV1 relative to the first cooldown are available, as initially there was a loose connection in the coupler that was later secured when the cryomodule was warmed to room temperature. The "After Plasma Processing" column shows the MP quenches after plasma cleaning of the vCM cavities. For all three rf tests, the data includes the MP quenches encountered during the initial power rise to 16 MV/m, the power rise to the maximum accelerating gradient (often accompanied by MP processing as necessary), the measurement of the usable gradient followed by additional MP processing, the stability test at 21 MV/m (usually 1 h or more), and additional MP processing as necessary. The MP quenches that took place during pulsed rf processing (applied to CAV3 and CAV5 during the first cooldown and to CAV1 and CAV6 during the second cooldown) are not counted by our algorithm, causing underestimation in the number of quenches estimated for the two preplasma cooldowns.

To identify and quantify the MP quenches, the data concerning the evolution of the cavity gradient and radiation level as a function of time were analyzed for each cavity. An algorithm was used to identify all the quenches, seen as a sudden drop in the gradient signal, above 16 MV/m and below the maximum accelerating field, such as due to MP whenever a radiation signal above the background was detected within the 2 min around the quench event. For this purpose, all the radiation sensors were considered, not only the sensors closest to the cavity under examination.

As explained in Posen *et al.* [18], CAV5 was affected by FE during the first two cooldowns here under consideration. The electron activity caused by the FE may have contributed to CAV5 MP-induced quenches during the first two cooldowns. At the end of the unit test, the FE source was processed, leaving the CAV5 (as the rest of the vCM) FE-free. After the unit test, CAV5 still showed a few MP quenches. A higher radiation threshold was used in the cases when CAV5 was on and exhibiting FE while other cavities were powered on at the same time.

One limitation of this methodology, in general, is due to the fact that in case multiple quenches occur close in time, it is not possible to distinguish the radiation produced by each quench event.

Figure 10 shows the number of MP quenches measured for each cavity as a function of the integrated time during which the cavity was tested for the three cooldowns under study. Additionally, from the graph in Fig. 10(c), it is possible to notice that the four plasma-cleaned cavities did not experience any MP-related quench during the rf test, while the four nonprocessed cavities still showed MP quenches.

This encouraging result demonstrated that although the plasma cleaning methodology was developed and optimized to target FE, it is also effective in eliminating MP of cavities in cryomodules. Indeed, it is known that both FE



FIG. 10. The three plots show the number of MP quenches found for each cavity as a function of the integrated time that the cavity spent powered on. Plot (a) shows these results for the first cooldown, (b) for the second cooldown measured before plasma processing, and (c) for the after plasma processing cooldown.

and MP can be enhanced by the presence of adsorbates on the cavity surface [14]. For example, hydrocarbons can reduce the Nb work functions, decreasing the FE threshold, while the presence of water can increase the Nb secondary emission yield, increasing MP. Plasma processing can successfully decrease the level of different types of adsorbates, mitigating both FE and MP *in situ* at the same time. Previous studies conducted at FNAL during the R&D phase used visible carbon-based contamination to study the reach of the plasma inside the cavity cell and showed that the removal is maximal at the iris but still present and nonnegligible at the equator. More fundamental studies will follow to better understand how to optimize plasma processing for MP mitigation.

#### **VI. CONCLUSIONS**

A plasma cleaning methodology was fully developed for the Linac Coherent Light Source II (LCLS-II) cavities. We presented its first application to a full cryomodule system, i.e., to the LCLS-II High Energy (HE) verification cryomodule.

During the developmental phase, we reviewed and mitigated possible risks deriving from the application of plasma processing to the cryomodule. This phase was crucial to confirm that the procedure could be safely applied to the cryomodule without the risk of compromising its components.

After connecting all the hardware needed for plasma cleaning, four out of eight cavities of the vCM were fully processed. The cryomodule was subsequently cooled down and retested to verify the effect of the processing on the cavity performance. With this test, we demonstrated that after plasma cleaning, (i) CM integrity and cavity performance are preserved, (ii) no field emitters were introduced by the processing (the vCM is still FE-free), and (iii) plasma processing eliminates MP: plasma processed cavities did not show any sign of MP, while the other four cavities were again affected by a series of MP quenches, also observed during the previous testing. Through this test, the plasma

cleaning procedure for 1.3 GHz CMs was fully validated, showing that plasma processing has the potential to not only reduce FE but also fully eliminate the MP of cavities in CMs.

More fundamental studies will be needed to better understand how to optimize the plasma processing procedure for *in situ* MP mitigation. In addition, the plasma cleaning procedure may be further optimized to be easily implemented during the LCLS-II-HE cryomodule production phase. This could significantly reduce cryomodule testing and commissioning time and increase reliability during machine operation.

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[1] M. Doleans, R. Afanador, J. Ball, W. Blokland, M. Crofford, B. Degraff, D. Douglas, B. Hannah, M. Howell, S. Kim *et al.*, Plasma processing R&D for the SNS superconducting linac rf cavities, in *Proceedings of SRF2013, Paris, France, 2013* (JACoW, Geneva, Switzerland, 2013), TUP057, https://accelconf.web.cern.ch/SRF2013/papers/tup057.pdf.

- [2] M. Doleans, Ignition and monitoring technique for plasma processing of multicell superconducting radio-frequency cavities, J. Appl. Phys. **120**, 243301 (2016).
- [3] M. Doleans, D. Vandygriff, S. Stewart, S. Gold, T. Neustadt, W. Strong, C. McMahan, S.-W. Lee, P. Tyagi, D. Vandygriff *et al.*, Plasma processing to improve the performance of the SNS superconducting linac, in *Proceedings of the 28th Linear Accelerator Conference, LINAC-2016, East Lansing, MI, USA* (JACoW, Geneva, Switzerland, 2016),WE2A03, https://accelconf.web.cern.ch/linac2016/papers/we2a03.pdf.
- [4] M. Doleans, P. Tyagi, R. Afanador, C. McMahan, J. Ball, D. Barnhart, W. Blokland, M. Crofford, B. Degraff, S. Gold *et al.*, In-situ plasma processing to increase the accelerating gradients of superconducting radio-frequency cavities, Nucl. Instrum. Methods Phys. Res., Sect. A **812**, 50 (2016).
- [5] P. Berrutti, B. Giaccone, M. Martinello, A. Grassellino, T. Khabiboulline, M. Doleans, S. Kim, D. Gonnella, G. Lanza, and M. Ross, Plasma ignition and detection for in-situ cleaning of 1.3 GHz 9-cell cavities, J. Appl. Phys. **126**, 023302 (2019).
- [6] P. Berrutti, M. Doleans, D. Gonnella, A. Grassellino, T. Khabiboulline, S. Kim, G. Lanza, M. Martinello, M. Ross, and K. Tippey, Update on plasma processing R&D for LCLS-II, in *Proceedings of the 9th International Particle Accelerator Conference, IPAC-2018, Vancouver, BC, Canada* (JACoW, Geneva, Switzerland, 2018), 10.18429/JACoW-IPAC2018-WEPMK012.
- [7] B. Giaccone, M. Martinello, P. Berrutti, O. Melnychuk, D. A. Sergatskov, A. Grassellino, D. Gonnella, M. Ross, M. Doleans, and J. F. Zasadzinski, Field emission mitigation studies in the SLAC Linac Coherent Light Source II superconducting rf cavities via *in situ* plasma processing, Phys. Rev. Accel. Beams 24, 022002 (2021).
- [8] T. Powers, N. Brock, and T. Ganey, In situ plasma processing of SRF cavities at JLAB, in *Proceedings of 20th International Conference on RF Superconductivity, SRF'21,* 2021 (JACoW, Geneva, Switzerland, 2021), TUPTEV004, https://srf2021.vrws.de/papers/tuptev004.pdf.
- [9] C. Zhang, W. Chang, K. Elliott, W. Hartung, S. H. Kim, J. T. Popielarski, K. Saito, and T. Xu, Development of in-situ plasma cleaning for the FRIB SRF linac, in *Proceedings of 20th International Conference on RF Superconductivity, SRF'21, 2021* (JACoW, Geneva, Switzerland, 2021), WEPTEV011, https://srf2021.vrws .de/papers/weptev011.pdf.
- [10] A. Wu, L. Yang, C. Hu, C. Li, S. Huang, Y. Li, Q. Chu, P. Xiong, H. Guo, W. Yue, and Y. Hu, In-situ plasma cleaning

to decrease the field emission effect of half-wave superconducting radio-frequency cavities, Nucl. Instrum. Methods Phys. Res., Sect. A **905**, 61 (2018).

- [11] A. Wu, S. Huang, C. Li, Q. Chu, H. Guo, P. Xiong, Y. Song, F. Pan, T. Tan, W. Yue, and S. Zhang, The cryostat results of carbon contamination and plasma cleaning for the field emission on the SRF cavity, in *Proceedings of the 19th International Conference on RF Superconductivity, SRF'19, Dresden, Germany, 2019* (JACoW, Geneva, Switzerland, 2019), pp. 1038–1040, 10.18429/JACoW-SRF2019-THP064.
- [12] S. Huang, Q. Chu, Y. He, C. Li, A. Wu, and S. Zhang, The effect of helium processing and plasma cleaning for low beta HWR cavity, in *Proceedings of the 19th International Conference on RF Superconductivity, SRF'19, Dresden, Germany, 2019* (JACOW, Geneva, Switzerland, 2019), pp. 1228–1230, 10.18429/JACoW-SRF2019-FRCAB6.
- [13] R. H. Fowler and L. Nordheim, Electron emission in intense electric fields, Proc. R. Soc. A 119, 173 (1928).
- [14] H. Padamsee, J. Knobloch, and T. Hays, *RF Superconductivity for Accelerators* (John Wiley & Sons, New York, 2009).
- [15] J. Galayda, LCLS-II Final Design Report, Report No. LCLSII-1.1-DR-0251-R0, 2015.
- [16] R. W. Schoenlein, LCLS-II High Energy (LCLS-II-HE): A transformative x-ray laser for science, SLAC National Accelerator Laboratory, Menlo Park, CA, Technical Report No. SLAC-R-1143, 2016, 10.2172/1634206.
- [17] A. Grassellino, A. Romanenko, D. Sergatskov, O. Melnychuk, Y. Trenikhina, A. Crawford, A. Rowe, M. Wong, T. Khabiboulline, and F. Barkov, Nitrogen and argon doping of niobium for superconducting radio frequency cavities: A pathway to highly efficient accelerating structures, Supercond. Sci. Technol. 26, 102001 (2013).
- [18] S. Posen *et al.*, High gradient performance and quench behavior of a verification cryomodule for a high energy continuous wave linear accelerator, Phys. Rev. Accel. Beams **25**, 042001 (2022).
- [19] D. Gonnella, S. Aderhold, D. Bafia, A. Burrill, M. Checchin, M. Ge, A. Grassellino, G. Hays, M. Liepe, M. Martinello *et al.*, The LCLS-II HE High Q and Gradient R&D Program, in *Proceedings of International Conference on RF Superconductivity 19, SRF'19* (JACoW, Geneva, Switzerland, 2019), 10.18429/JACoW-SRF2019-MOP045.
- [20] B. Aune, R. Bandelmann, D. Bloess, B. Bonin, A. Bosotti, M. Champion, C. Crawford, G. Deppe, B. Dwersteg, D. Edwards *et al.*, Superconducting TESLA cavities, Phys. Rev. ST Accel. Beams 3, 092001 (2000).