Optimal usage of quantum random access memory in quantum machine learning

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(Received 17 September 2018; published 16 January 2019)

By considering an unreliable oracle in a query-based model of quantum learning, we present a tradeoff relation between the oracle's reliability and the reusability of the quantum state of the input data. The tradeoff relation manifests as the fundamental upper bound on the reusability. This limitation on the reusability would increase the quantum access to the input data, i.e., the usage of quantum random access memory (qRAM), repeating the preparation of a superposition of large (or big) input data on the query failure. However, it is found that a learner can obtain a correct answer even from an unreliable oracle without any additional usage of qRAM; i.e., the complexity of the qRAM query does not increase even with an unreliable oracle. This is enabled by repeatedly cycling the quantum state of the input data to the upper bound on the reusability.

DOI: 10.1103/PhysRevA.99.012326

I. INTRODUCTION

Quantum machine learning (QML) is a rapidly growing research field. A primary issue in QML is to prove a quantum learning speedup over the classical counterparts [1–3]. A recent proposal of the quantum support vector machine (QSVM) [4], providing an exponential speedup in a classification, can be considered as a paradigmatic achievement. Currently, the QSVM (and other variant QML proposals [4–8]) provides a standardized approach to achieve the quantum speedup—to use the quantum state provided that a set of input data is superposed so that useful quantum algorithms, e.g., so-called HHL (named after the inventors, Harrow, Hassidim, and Lloyd [9]), can be utilized as the kernel.

However, unclear aspects still exist in QML. In particular, whether the quantum advantage remains significant even when the cost to access a large size of input data is considered needs to be clarified, i.e., whether classical input data can be transformed to a quantum superposition [10,11]. In theory, at least, quantum random access memory (qRAM) can accomplish the aforementioned task [12,13], even though its realization is far from trivial [14]. Subsequently, a question arises as to whether it is possible to reduce the qRAM query by reusing the quantum state of the input data that has been initialized once. The reuse of quantum data is limited because the information extraction causes the state disturbance (or equivalently, based on the no-cloning theorem [15]), contrary to the classical machine learning that has no limitation in reusing the data [16]. However, an original state can be retrieved using weak measurements with nonzero probability [17,18]. Hence, it is important to explore whether the reuse or recycle of the quantum state of the input data is possible, the quantum limit on the reusability, and whether it offers any advantage in QML.

The oracle's reliability also affects the learning performance significantly [19]. The effects by the unreliable oracle with a missed answer or an evasive answer (e.g., "I do not know") have been studied and shown to be tolerable in querybased models of classical learning [20,21]. In particular, the learner can be polynomially dominated by a failure query rate of less than 1/2 [20,21]. Such results were also drawn in QML [3,19]. Furthermore, it was claimed that some quantum advantages are achievable with noisy oracles [22,23]. However, the effects of the oracle's reliability on the complexity of the qRAM query have not been studied in QML, even though recent speedups of QML hinge crucially on the low qRAM queries.

Herein, by casting a query-based model of quantum learning with an unreliable oracle, we explore the fundamental limit on the reusability of the quantum state of the input data quantitatively. In particular, we present a tradeoff relation between the oracle's reliability and the reusability of the quantum state of the input data. The tradeoff relation indicates that the more reliable is the oracle, the lower is the reusability. It also manifests the fundamental upper bound on the reusability for given oracle reliability. Such a limited reusability would impose the additional usage of qRAM, thus repeating the quantum access to the input data with query failure. However, it is found that the learner can, in principle, arrive at the correct answer with a single run of qRAM [24], repeatedly cycling the quantum state of the input data to the upper bound of the reusability. This result implies that, if the traveling cost of the input data is neglected, an incomplete-oracle learner has the same complexity of qRAM query as that of a completeoracle learning.

II. INCOMPLETE-QUERY LEARNING

Typically, machine learning is often formulated as an identification of a function c (referred to as a "concept" in the language of machine learning); it maps the input data

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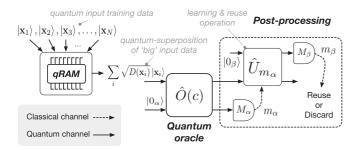


FIG. 1. A schematic picture of the query-based QML model. First, the quantum random access memory (qRAM) initializes a superposition of a large size, say N, of input data. Then, the quantum oracle $\hat{O}(c)$ is employed. For a general scenario, $\hat{O}(c)$ is assumed to be unreliable, yielding an incorrect answer with a nonzero probability. The last postprocessing block (dashed box) is responsible for the learning and the reuse of the superposed state initialized by qRAM (see the main text).

 $\mathbf{x} \in \{0, 1\}^n$ (in arbitrary *n*-bit strings) to the target $c(\mathbf{x}) \in \{0, 1\}$, i.e., a task of classification [25]. In contrast to classical machine learning, QML employs a set of quantum training data, i.e., $|\mathbf{x}\rangle$ and $|c(\mathbf{x})\rangle$. Hence, we design a query-based QML model, as shown in Fig. 1.

Our model is roughly composed of three blocks. The first block is the initialization of a quantum superposition of the input data $|\mathbf{x}_i\rangle$ (i = 1, 2, ..., N). At least in theory, this task can be accomplished by casting the qRAM [12]; more specifically, qRAM allows for the data to be read (or to be written) from arbitrary *i*th memory cells [26] and creates the superposition of all *N* input data, denoted hereinafter as

$$|\psi_0\rangle = \sum_{i=1}^N \sqrt{D_i} |\mathbf{x}_i\rangle,\tag{1}$$

where D_i is a probability distribution of memory cells.

Then we consider the quantum learning oracle $\hat{O}(c)$ that is assumed to be unreliable, yielding an incorrect answer $|c(\mathbf{x}_i) \oplus 1\rangle$ with a certain probability [22]. The oracle operation is defined as

$$\begin{aligned} |\psi_{0}\rangle|0_{\alpha}\rangle &\xrightarrow{\hat{O}(c)} \sum_{i=1}^{N} \sqrt{D(\mathbf{x}_{i})} (\sqrt{\lambda_{+}}|\mathbf{x}_{i}\rangle|c(\mathbf{x}_{i})\rangle \\ &+ \sqrt{\lambda_{-}}|\mathbf{x}_{i}\rangle|c(\mathbf{x}_{i})\oplus1\rangle), \end{aligned}$$
(2)

where

$$\lambda_{\pm} = \frac{1 \pm \mathcal{L}}{2} \tag{3}$$

is the qubit state of the oracle-answer register. Here, we define the *oracle's reliability* with the factor $\mathcal{L} \in [0, 1]$; for example, $\hat{O}(c)$ is perfectly reliable when $\mathcal{L} = 1$, but is less reliable when $\mathcal{L} < 1$. For the case when $\mathcal{L} = 0$, the oracle $\hat{O}(c)$ produces a completely random answer, yielding no information. We clarify that the queries to qRAM and $\hat{O}(c)$ are distinct; the qRAM query is engaged as the process of initializing a superposition [as in Eq. (1)] of the input data, while the oracle $\hat{O}(c)$ is queried about the legitimate learning output for the (superposed) inputs. The last block is for postprocessing, i.e., learning and the reuse of $|\psi_0\rangle$ in Eq. (1). Prior to those processes, the oracle's answer, or equivalently the learning information, should be identified within this block. Thus, a projection measurement M_{α} is assumed first to yield (the information of) the oracle's answer, followed by sequential operations denoted by $\hat{U}_{m_{\alpha}}$. The information for learning can be extracted from the measurement outcome m_{α} . After the measurement, m_{α} is delivered to and utilized by $\hat{U}_{m_{\alpha}}$ for the learning and/or reuse process. To recycle $|\psi_0\rangle$, the operation $\hat{U}_{m_{\alpha}}$ is manipulated according to the outcome m_{α} and is applied to the output state of $\hat{O}(c)$. Another projection measurement M_{β} is performed after $\hat{U}_{m_{\alpha}}$. Here, we define the *reusability*, denoted by \mathcal{R} , in terms of the overall probability of attaining $|\psi_0\rangle$ after the measurement M_{β} .

It is worth noting that our model is general and equivalent to the conventional query-based model of learning [27], by which the best speedup is polynomial [28,29]. However, employing such a model is sufficient to derive a quantitative relation between the oracle's reliability and the reusability of $|\psi_0\rangle$.

III. TRADEOFF RELATION BETWEEN ORACLE RELIABILITY AND INPUT REUSABILITY

We herein present a tradeoff relation between the reliability \mathcal{L} of the oracle $\hat{O}(c)$ and the reusability \mathcal{R} of the quantum state $|\psi_0\rangle$ of the input data. For convenience in calculations, we rewrite Eq. (1) as the following form:

$$|\psi_0\rangle = \sum_{i=1}^N \sqrt{D(\mathbf{x}_i)} |\mathbf{x}_i\rangle = \sum_{\tau=0,1} \sqrt{\xi_\tau} |X_\tau\rangle, \qquad (4)$$

where

$$|X_{\tau}\rangle = \sum_{\mathbf{x}_i \in X_{\tau}} \sqrt{\frac{D(\mathbf{x}_i)}{\xi_{\tau}}} |\mathbf{x}_i\rangle \text{ and } \xi_{\tau} = \sum_{\mathbf{x}_i \in X_{\tau}} D(\mathbf{x}_i).$$
(5)

Here, $X_{\tau} \subset {\mathbf{x}_i : i = 1, ..., N}$ denotes a set of \mathbf{x}_i , satisfying $c(\mathbf{x}_i) = \tau$ and $\cup_{\tau=0,1} X_{\tau} = {\mathbf{x}_i : i = 1, ..., N}$.

Subsequently, we introduce a set of Kraus operators $\hat{A}_{m_{\alpha}}$ $(m_{\alpha} = 0, 1)$, defined by the combination of $\hat{O}(c)$ and M_{α} . By adopting a fixed form of $\hat{O}(c)$ as

$$\hat{O}(c) = \sum_{\tau=0,1} |X_{\tau}\rangle \langle X_{\tau}| (\sqrt{\lambda_{+}} | c = \tau\rangle \langle 0 | + \sqrt{\lambda_{-}} | c = \tau \oplus 1\rangle \langle 0 |),$$
(6)

we can characterize $\hat{A}_{m_{\alpha}}$ such that

$$\hat{A}_{0} = \sqrt{\lambda_{+}} |X_{0}\rangle \langle X_{0}| + \sqrt{\lambda_{-}} |X_{1}\rangle \langle X_{1}|,$$
$$\hat{A}_{1} = \sqrt{\lambda_{-}} |X_{0}\rangle \langle X_{0}| + \sqrt{\lambda_{+}} |X_{1}\rangle \langle X_{1}|,$$
(7)

with the eigenvalues $\sqrt{\lambda_{\pm}}$. The process of extracting the learning information is subsequently expressed as follows:

$$\hat{A}_{m_{\alpha}}|\psi_{0}\rangle = \sqrt{P_{m_{\alpha}}}|\varphi_{m_{\alpha}}\rangle(m_{\alpha}=0,1), \qquad (8)$$

where $P_{m_{\alpha}}$ is given as

$$P_0 = \frac{1 + (\xi_0 - \xi_1)\mathcal{L}}{2}$$
 and $P_1 = \frac{1 - (\xi_0 - \xi_1)\mathcal{L}}{2}$. (9)

Here, $|\varphi_{m_{\alpha}}\rangle$ is the postmeasurement state as a general term. In our scenario, $P_{m_{\alpha}}$ is the probability of obtaining m_{α} from M_{α} and $|\varphi_{m_{\alpha}}\rangle$ is the remaining state after M_{α} , as shown in a later section.

We also define an operator $\hat{R}^{(m_{\alpha})}$ for the reuse process as the combination of $\hat{U}_{m_{\alpha}}$ and M_{β} . The reuse process is subsequently expressed as

$$\hat{R}^{(m_{\alpha})}\hat{A}_{m_{\alpha}}|\psi_{0}\rangle = \sqrt{\eta^{(m_{\alpha})}}|\psi_{0}\rangle, \qquad (10)$$

where $\eta^{(m_{\alpha})}$ is a nonzero complex number. Because $\hat{1} - \hat{R}^{(m_{\alpha})\dagger}\hat{R}^{(m_{\alpha})}$ is positive semidefinite,

$$\sup_{|\chi\rangle} \langle \chi | \hat{R}^{(m_{\alpha})\dagger} \hat{R}^{(m_{\alpha})} | \chi \rangle \leqslant 1$$
(11)

for arbitrary normalized states $|\chi\rangle$. Meanwhile, for the state $|\varphi_{m_{\alpha}}\rangle = \frac{\hat{A}_{m_{\alpha}}|\psi_{0}\rangle}{\sqrt{P_{m_{\alpha}}}}$ [17],

$$\sup_{|\chi\rangle} \langle \chi | \hat{R}^{(m_{\alpha})\dagger} \hat{R}^{(m_{\alpha})} | \chi \rangle$$

$$\geqslant \sup_{|\varphi_{m_{\alpha}}} \langle \varphi_{m_{\alpha}} | \hat{R}^{(m_{\alpha})\dagger} \hat{R}^{(m_{\alpha})} | \varphi_{m_{\alpha}} \rangle$$

$$= \sup_{|\psi_{0}\rangle} \frac{\langle \psi_{0} | \hat{A}^{\dagger}_{m_{\alpha}} \hat{R}^{(m_{\alpha})\dagger} \hat{R}^{(m_{\alpha})} \hat{A}_{m_{\alpha}} | \psi_{0} \rangle}{P_{m_{\alpha}}}$$

$$= \inf_{|\psi_{0}\rangle} \frac{\eta^{(m_{\alpha})}}{P_{m_{\alpha}}}, \qquad (12)$$

and by Eq. (11), we obtain $\eta^{(m_{\alpha})} \leq \inf_{|\psi_0\rangle} P_{m_{\alpha}}$. Subsequently, from Eq. (9), we can verify that $\inf_{|\psi_0\rangle} P_{m_{\alpha}} = \lambda_-$ when $\xi_0 - \xi_1 = -1$ for $m_{\alpha} = 0$ and when $\xi_0 - \xi_1 = 1$ for $m_{\alpha} = 1$. It is worth noting that generally the initial state can be written with an arbitrary orthonormal basis and coefficients according to the choice of M_{α} . Thus, for all $m_{\alpha} = 0$, 1, we can obtain $\eta^{(m_{\alpha})} \leq \lambda_-$. Then, the probability of attaining the reusable $|\psi_0\rangle$, for an $m_{\alpha} \in \{0, 1\}$, is bounded as

$$\left|\langle\psi_0|\hat{R}^{(m_{\alpha})}|\varphi_{m_{\alpha}}\rangle\right|^2 = \frac{\eta^{(m_{\alpha})}}{P_{m_{\alpha}}} \leqslant \frac{\lambda_-}{P_{m_{\alpha}}}.$$
 (13)

We can finally obtain the overall success probability of the reuse, i.e., the reusability, as

$$\mathcal{R} = \sum_{m_{\alpha}=0,1} P_{m_{\alpha}} \left| \langle \psi_0 | \hat{R}^{(m_{\alpha})} | \varphi_{m_{\alpha}} \rangle \right|^2 \leqslant 1 - \mathcal{L}.$$
(14)

This clearly shows that \mathcal{R} is inversely correlated with and limited by \mathcal{L} , i.e., that of a tight tradeoff relation between the reusability and the oracle's reliability. Note that our proof is valid for arbitrary $\hat{U}_{m_{\alpha}}$ and M_{β} . This result is in agreement

with the theorem made in the information-theoretic perspectives [18].

The tradeoff relation in Eq. (14) manifests the fundamental limit on the reusability of the quantum state of the input data in QML. The average reusable number is given by

$$\overline{n} = \sum_{n=0}^{\infty} n \mathcal{R}^n = \frac{\mathcal{R}}{(1-\mathcal{R})^2}$$
(15)

and by Eq. (14), where we have $\overline{n} \leq \mathcal{L}^{-1}(\mathcal{L}^{-1} - 1)$; i.e., for a single run of qRAM, it is possible to continue the reuse of $|\psi_0\rangle$, on average, less than $\mathcal{L}^{-1}(\mathcal{L}^{-1} - 1)$. This implies that the higher the learning efficiency or equivalently of the oracle's reliability, the lower the reusability of the state of the input data. Such a limited reusability imposes the requirement of a higher rate of qRAM query.

IV. OPTIMAL USAGE OF QRAM

We herein demonstrate that the usage of qRAM can be optimized by cycling the state $|\psi_0\rangle$ of the input data to the fundamental bound to saturate the tradeoff relation. Hence, we consider an exemplary protocol as described below. The oracle operation is described by

$$\begin{aligned} |\psi_{0}\rangle|0_{\alpha}\rangle \xrightarrow{O(c)} &\sum_{\tau=0,1} (\sqrt{\xi_{\tau}\lambda_{+}}|X_{\tau}\rangle|c=\tau) \\ &+\sqrt{\xi_{\tau}\lambda_{-}}|X_{\tau\oplus1}\rangle|c=\tau\oplus1\rangle), \end{aligned} \tag{16}$$

with the states of the correct $|c = \tau\rangle$ and incorrect answer $|c = \tau \oplus 1\rangle$. Subsequently, a measurement M_{α} is performed, yielding the oracle's answer with outcomes $m_{\alpha} \in \{0, 1\}$. Given the measurement result $m_{\alpha} \in \{0, 1\}$, the postmeasurement states $|\varphi_{m_{\alpha}}\rangle$ defined in Eq. (8) can be obtained [30]. The processes including the oracle and the subsequent measurement, $\hat{O}(c) + M_{\alpha}$, result in a specific form of remaining state $|\varphi_{m_{\alpha}}\rangle$, such that

$$|\psi_{0}\rangle|0_{\alpha}\rangle \rightarrow \begin{cases} |\varphi_{0}\rangle = \sqrt{\frac{\xi_{0}\lambda_{+}}{P_{0}}}|X_{0}\rangle + \sqrt{\frac{\xi_{1}\lambda_{-}}{P_{0}}}|X_{1}\rangle, \\ |\varphi_{1}\rangle = \sqrt{\frac{\xi_{0}\lambda_{-}}{P_{1}}}|X_{0}\rangle + \sqrt{\frac{\xi_{1}\lambda_{+}}{P_{1}}}|X_{1}\rangle, \end{cases}$$
(17)

where P_0 and P_1 are given in Eq. (9) and denote the probabilities of getting $m_{\alpha} = 0$ and $m_{\alpha} = 1$, respectively.

Subsequently, $\hat{U}_{m_{\alpha}}$ is applied on the state $|\varphi_{m_{\alpha}}\rangle$ and an ancillary state $|0\rangle_{\beta}$. The optimal $\hat{U}_{m_{\alpha}}$ can be chosen, according to the identified m_{α} , to maximize the reusability \mathcal{R} . Note that the optimal form of $\hat{U}_{m_{\alpha}}$ is not unique; i.e., many structures with the maximum \mathcal{R} may be constructed. Here, we can select $\hat{U}_{m_{\alpha}}$ in the form of

$$\hat{U}_{m_{\alpha}} = \left(\cos\Theta\left(\hat{\sigma}_{x}^{m_{\alpha}\oplus1}\otimes\hat{\mathbb{1}}_{N}\right) + i(-1)^{m_{\alpha}\oplus1}\sin\Theta\,\hat{C}_{X_{0},X_{1}}\right)(\hat{\mathbb{1}}_{2}\otimes\hat{R}(\Theta)),\tag{18}$$

where $\hat{C}_{X_0,X_1} = \hat{\mathbb{1}}_2 \otimes |X_0\rangle \langle X_0| + \hat{\sigma}_x \otimes |X_1\rangle \langle X_1|$, and $\hat{R}(\Theta) = |X_0\rangle \langle X_0| + e^{i\Theta} |X_1\rangle \langle X_1|$. Here, $\hat{\mathbb{1}}_d$ is the identity of the *d*-dimensional Hilbert space. For each case of m_α , the state $|\varphi_{m_\alpha}\rangle$ undergoes the transformation with \hat{U}_{m_α} as

$$\begin{array}{l} 0_{\beta} |\varphi_{0}\rangle \xrightarrow{\hat{U}_{m_{\alpha}=0}} -i\sqrt{\frac{\xi_{0}\lambda_{+}}{P_{0}}\sin\Theta|0\rangle|X_{0}\rangle} + \sqrt{Q_{0}}|1\rangle \left(\sqrt{\frac{\xi_{0}\lambda_{+}}{P_{0}Q_{0}}\cos\Theta|X_{0}\rangle} + \sqrt{\frac{\xi_{1}\lambda_{-}}{P_{0}Q_{0}}}|X_{1}\rangle\right), \\ 0_{\beta} |\varphi_{1}\rangle \xrightarrow{\hat{U}_{m_{\alpha}=1}} \sqrt{Q_{1}}|0\rangle \left(\sqrt{\frac{\xi_{0}\lambda_{-}}{P_{1}Q_{1}}}|X_{0}\rangle + \sqrt{\frac{\xi_{1}\lambda_{+}}{P_{1}Q_{1}}}\cos\Theta|X_{1}\rangle\right) + i\sqrt{\frac{\xi_{1}\lambda_{+}}{P_{1}}}\sin\Theta|1\rangle|X_{1}\rangle, \end{array}$$

$$(19)$$

where

$$Q_0 = \frac{\xi_0 \lambda_+ \cos^2 \Theta + \xi_1 \lambda_-}{P_0}, \quad Q_1 = \frac{\xi_0 \lambda_- + \xi_1 \lambda_+ \cos^2 \Theta}{P_1}.$$
(20)

Because Θ can be written in terms of \mathcal{L} , the optimal $\hat{U}_{m_{\alpha}}$ is determined depending on a given oracle's reliability \mathcal{L} . Here, if we set $\Theta = \arccos \sqrt{\frac{\lambda_{-}}{\lambda_{+}}}$, then Q_j becomes $\frac{\lambda_{-}}{P_j}$ (j = 0, 1) and the transformations in Eq. (19) are rewritten as

$$\begin{aligned} |0_{\beta}\rangle|\varphi_{0}\rangle \xrightarrow{\hat{U}_{m_{\alpha}=0}} -i\sqrt{\frac{\xi_{0}\mathcal{L}}{P_{0}}}|0\rangle|X_{0}\rangle + \sqrt{\frac{\lambda_{-}}{P_{0}}}|1\rangle\underbrace{\left(\sum_{\tau=0,1}\sqrt{\xi_{\tau}}|X_{\tau}\rangle\right)}_{\text{reusable state }|\psi_{0}\rangle} \\ |0_{\beta}\rangle|\varphi_{1}\rangle \xrightarrow{\hat{U}_{m_{\alpha}=1}} i\sqrt{\frac{\xi_{1}\mathcal{L}}{P_{1}}}|1\rangle|X_{1}\rangle + \sqrt{\frac{\lambda_{-}}{P_{1}}}|0\rangle\underbrace{\left(\sum_{\tau=0,1}\sqrt{\xi_{\tau}}|X_{\tau}\rangle\right)}_{\text{reusable state }|\psi_{0}\rangle} \end{aligned}$$

$$(21)$$

After a secondary measurement M_{β} is performed on the first mode of Eq. (21), the probabilities of the cases when the results are consistent ($m_{\beta} = m_{\alpha}$) and inconsistent ($m_{\beta} \neq m_{\alpha}$) are obtained, respectively, as

$$Q_{m_{\beta}=m_{\alpha}} = \frac{\xi_{m_{\alpha}}\mathcal{L}}{P_{m_{\alpha}}}, \quad Q_{m_{\beta}\neq m_{\alpha}} = \frac{\lambda_{-}}{P_{m_{\alpha}}}.$$
 (22)

Then, it is inferred—observing Eq. (19)—that the correct query output $|X_{\tau}\rangle$ can be extracted with the probability $Q_{m_{\beta}=m_{\alpha}}$, unless $\mathcal{L} = 0$. In other words, we can confirm that the oracle's answer obtained in M_{α} is correct if it is consistent with the outcome of M_{β} , i.e., $m_{\alpha} = m_{\beta}$. The overall probability of attaining $|X_{\tau}\rangle$ is subsequently given as $\sum_{m_{\alpha}} Q_{m_{\beta}=m_{\alpha}} P_{m_{\alpha}} = 1 - \mathcal{R} = \mathcal{L}$, satisfying the tradeoff relation [31]. Meanwhile, for the case of inconsistent results, i.e., $m_{\beta} \neq m_{\alpha}$, one can recover the state $|\psi_0\rangle$ of the input data, that is, conclusively reusable. It is worth noting that the probability $Q_{m_{\alpha}\neq m_{\beta}}$ obtained in Eq. (22) is optimal, as described in Eq. (13). Subsequently, the reusability can be calculated as $\mathcal{R} = \sum_{m_{\alpha}=0,1} P_{m_{\alpha}} Q_{m_{\beta}\neq m_{\alpha}} = \sum_{m_{\alpha}=0,1} \lambda_{-} =$ $1 - \mathcal{L}$, saturating the tradeoff relation in Eq. (14). Therefore, in principle, $|\psi_0\rangle$ is allowed to be cycled until the correct output $|X_{\tau}\rangle$ is extracted, by achieving the fundamental bound of the reusability. This indeed provides us an optimal process of query (in principle) without any additional qRAM queries caused by the incomplete oracle.

V. REMARKS

In summary, we have derived a tight tradeoff relation between the reliability of the oracle and the reusability of the quantum state of the input data. It manifests the fundamental limit on the possibility of reusing a state, initialized as a superposition of the input data for a single run of qRAM. The derived tradeoff relation indicated that the more reliable the oracle, the lower was the reusability. This would impose the additional usage of qRAM with the query failure. However, even with the limited reusability, the overall query process could be optimized by cycling the state initialized once. In particular, the optimized process was shown to saturate the fundamental upper bound of the reusability limited by the tradeoff relation. Remarkably, it was shown that the learner could, in principle, arrive at the correct answer without any additional qRAM queries caused by the incomplete oracle; for example, when the oracle produces incorrect answers, the quantum state of the input data could be recovered with postprocessing to be used again for query. Such a process could be repeated until the correct answer is extracted. Thus, the complexity of the qRAM query would not increase even with an unreliable oracle. This result will be crucial, since the low usage of qRAM is highly desirable in QML. We believe that our work will provide a fundamental and practical insight on the QML.

ACKNOWLEDGMENTS

We are grateful to Jinhyoung Lee for the fruitful discussions. J.B. and A.D. are grateful to Marcin Wieśniak, Wiesław Laskowski, and Marcin Pawłowski. J.B. would like to thank Junghee Ryu and Nana Liu for the discussions. This research was implemented as a research project on quantum machine learning (Grant No. 2018-104) by the ETRI affiliated research institute. J.B. acknowledges the support of the R&D Convergence program of NST (National Research Council of Science and Technology) of the Republic of Korea (Grant No. CAP-18-08-KRISS).

- M. Schuld, I. Sinayskiy, and F. Petruccione, Contemp. Phys. 56, 172 (2015).
- [2] J. Biamonte, P. Wittek, N. Pancotti, P. Rebentrost, N. Wiebe, and S. Lloyd, Nature (London) 549, 195 (2017).
- [3] C. Ciliberto, M. Herbster, A. D. Ialongo, M. Pontil, A. Rocchetto, S. Severini, and L. Wossnig, Proc. R. Soc. A 474, 20170551 (2018).
- [4] P. Rebentrost, M. Mohseni, and S. Lloyd, Phys. Rev. Lett. 113, 130503 (2014).
- [5] S. Lloyd, M. Mohseni, and P. Rebentrost, Nat. Phys. 10, 631 (2014).

- [6] M. Schuld, I. Sinayskiy, and F. Petruccione, Phys. Rev. A 94, 022342 (2016).
- [7] I. Kerenidis and A. Prakash, arXiv:1603.08675.
- [8] N. Liu and P. Rebentrost, Phys. Rev. A 97, 042315 (2018).
- [9] A. W. Harrow, A. Hassidim, and S. Lloyd, Phys. Rev. Lett. 103, 150502 (2009).
- [10] S. Aaronson, Nat. Phys. 11, 291 (2015).
- [11] Z. Zhao, V. Dunjko, J. K. Fitzsimons, P. Rebentrost, and J. F. Fitzsimons, arXiv:1804.00281.
- [12] V. Giovannetti, S. Lloyd, and L. Maccone, Phys. Rev. Lett. 100, 160501 (2008).

- [13] V. Giovannetti, S. Lloyd, and L. Maccone, Phys. Rev. A 78, 052310 (2008).
- [14] S. Arunachalam, V. Gheorghiu, T. Jochym-O'Connor, M. Mosca, and P. V. Srinivasan, New J. Phys. 17, 123010 (2015).
- [15] W. K. Wootters and W. H. Zurek, Nature (London) 299, 802 (1982).
- [16] B. Custers and H. Uršič, Int. Data Privacy Law 6, 4 (2016).
- [17] M. Koashi and M. Ueda, Phys. Rev. Lett. 82, 2598 (1999).
- [18] Y. W. Cheong and S.-W. Lee, Phys. Rev. Lett. 109, 150402 (2012).
- [19] K. Iwama, R. Raymond, and S. Yamashita, in *Quantum Computation and Information* (Springer, Berlin, 2006), pp. 19–42.
- [20] D. Angluin and D. K. Slonim, Mach. Learn. 14, 7 (1994).
- [21] H. U. Simon, Theory Comput. Syst. 37, 77 (2004).
- [22] A. W. Cross, G. Smith, and J. A. Smolin, Phys. Rev. A 92, 012327 (2015).
- [23] D. Ristè, M. P. da Silva, C. A. Ryan, A. W. Cross, A. D. Córcoles, J. A. Smolin, J. M. Gambetta, J. M. Chow, and B. R. Johnson, npj Quantum Inf. 3, 16 (2017).
- [24] Throughout the work, it is assumed that the complexity of the qRAM query is primarily related to the (quantum) access to the

input data to initialize a quantum superposition of the input data [10,11].

- [25] P. Langley, *Elements of Machine Learning* (Morgan Kaufmann, San Francisco, CA, 1995).
- [26] Here, the memory arrays are either classical or quantum depending on the accessing type of qRAM.
- [27] V. Lyubashevsky, in Approximation, Randomization and Combinatorial Optimization: Algorithms and Techniques (Springer, Berlin, 2005), pp. 378–389.
- [28] R. A. Servedio and S. J. Gortler, SIAM J. Comput. 33, 1067 (2004).
- [29] R. Kothari, in 31st International Symposium on Theoretical Aspects of Computer Science (STACS 2014), Leibniz International Proceedings in Informatics (LIPIcs) Vol. 25 (Schloss Dagstuhl–Leibniz-Zentrum fuer Informatik, Dagstuhl, Germany, 2014), pp. 482–493.
- [30] M. Ueda, N. Imoto, and H. Nagaoka, Phys. Rev. A 53, 3808 (1996).
- [31] This situation together with our proof in Sec. III is suggestive of the interrelated principles in the probabilistic error correction, i.e., the upper bound of the probability of correcting the error imposed by the loss of information due to the error [17].