

**Optical spectroscopy of complex open-4*d*-shell ions Sn<sup>7+</sup>–Sn<sup>10+</sup>**F. Torretti,<sup>1,2,\*</sup> A. Windberger,<sup>1,3</sup> A. Ryabtsev,<sup>4,5</sup> S. Dobrodey,<sup>3</sup> H. Bekker,<sup>3</sup> W. Ubachs,<sup>1,2</sup> R. Hoekstra,<sup>1,6</sup> E. V. Kahl,<sup>7</sup> J. C. Berengut,<sup>7</sup> J. R. Crespo López-Urrutia,<sup>3</sup> and O. O. Versolato<sup>1</sup><sup>1</sup>*Advanced Research Center for Nanolithography, Science Park 110, 1098 XG Amsterdam, The Netherlands*<sup>2</sup>*Department of Physics and Astronomy, and LaserLab, Vrije Universiteit, De Boelelaan 1081, 1081 HV Amsterdam, The Netherlands*<sup>3</sup>*Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany*<sup>4</sup>*Institute of Spectroscopy, Russian Academy of Sciences, Troitsk, Moscow, 108840 Russia*<sup>5</sup>*EUV Labs, Ltd., Troitsk, Moscow, 108840 Russia*<sup>6</sup>*Zernike Institute for Advanced Materials, University of Groningen, Nijenborgh 4, 9747 AG Groningen, The Netherlands*<sup>7</sup>*School of Physics, University of New South Wales, Sydney 2052, Australia*

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We analyze the complex level structure of ions with many-valence-electron open-[Kr] 4*d<sup>m</sup>* subshells ( $m = 7-4$ ) with *ab initio* calculations based on configuration-interaction many-body perturbation theory (CI+MBPT). Charge-state-resolved optical and extreme ultraviolet (EUV) spectra of Sn<sup>7+</sup>–Sn<sup>10+</sup> ions were obtained using an electron beam ion trap. Semiempirical spectral fits carried out with the orthogonal parameters technique and COWAN code calculations lead to 90 identifications of magnetic-dipole transitions and the determination of 79 energy ground-configuration levels, questioning some earlier EUV-line assignments. Our results confirm the *ab initio* predictive power of CI+MBPT calculations for these complex electronic systems.

DOI: [10.1103/PhysRevA.95.042503](https://doi.org/10.1103/PhysRevA.95.042503)**I. INTRODUCTION**

The electronic structure [Kr] 4*d<sup>m</sup>* ( $m = 7-4$ ) of the highly charged ions (HCI) Sn<sup>7+</sup>–Sn<sup>10+</sup> is extremely complicated due to the many electrons that occupy their open-4*d* subshell, and remains inaccessible to even some of the most advanced atomic theories. Furthermore, the unresolved transition arrays [1] formed by the Sn ions are particularly useful for the production of 13.5-nm-wavelength extreme ultraviolet (EUV) radiation for nanolithographic applications [2–5]. Unfortunately, experimental assessments [6–20] of spectral data are hampered by the prevalence of strong configuration-interaction contributions, and by a high density of states which approaches the quantum-chaos regime for high excitation energies [21–24]. In a recent study [16], we found evidence calling for a revision of earlier identifications [19,20] in Sn<sup>11+</sup>–Sn<sup>14+</sup> ions having 3 to 0 electrons in their 4*d* subshell, and successfully demonstrated the suitability of Fock space coupled-cluster (FSCC) calculations for systems with up to two valence electrons or holes. We now investigate other charge states relevant for the EUV production in plasmas, namely, Sn<sup>7+</sup>–Sn<sup>10+</sup>.

We focus on optical spectroscopy in this work, which can resolve the complex manifold fine-structure splittings of these ions. Therefore, the analysis of optical transitions in heavy multielectron, open-shell ions enables the most stringent tests of *ab initio* atomic-structure calculations of strongly correlated systems with non-negligible many-electron Breit contributions. For such systems, a suitable theoretical tool is a combination of configuration interaction and many-body perturbation theory (CI+MBPT). The CI+MBPT method was first developed to very accurately treat neutral thallium as a three-valence-electron atom [25]. Since then, it has been markedly successful in treating also four- [26–29] and even five-valence-electron [30] systems. However, as the number

of valence electrons increases, it becomes less accurate. A recent extension of the CI+MBPT method, used here, includes particle-hole interaction, and improves the accuracy of the calculations [31]. This makes it possible to treat systems with several vacancies which, e.g., are currently inaccessible to FSCC calculations.

We present charge-state-resolved optical and EUV spectral measurements of Sn<sup>7+</sup>–Sn<sup>10+</sup> ions trapped in an electron beam ion trap (EBIT), FLASH-EBIT [32], at the Max Planck Institute for Nuclear Physics (Max-Planck-Institut für Kernphysik, MPIK) in Heidelberg. EUV spectra were obtained simultaneously with the optical ones in order to identify the charge states and assign the optical lines to them. Then, we compare the Sn<sup>7+</sup> data to the level structure accurately determined in Ref. [6], whereby a good agreement further validates our charge-state assignments. Subsequently, we perform line and level identifications for Sn<sup>8+</sup>–Sn<sup>10+</sup> using semiempirical calculations by employing the orthogonal parameters technique [33,34] and the COWAN code [35]. The observed Ritz combinations strongly support our semiempirical spectral analysis. Analogous to our recent work [16], we compare our experimental findings to previous level-structure determinations from measurements of EUV spectra [18] and find noteworthy discrepancies. Armed with these experiment-fitted level structure, we test our state-of-the-art *ab initio* CI+MBPT calculations, and find them in very good agreement with the data. Both the important practical applications of the ions under study as well as the relative novelty of using CI+MBPT calculations for systems with such large numbers of valence electrons make our theory-experiment comparisons very valuable.

**II. EXPERIMENT**

Tin ions were produced and subsequently trapped and excited using FLASH-EBIT [16,32]. In this device, the electron beam is compressed to a diameter of approximately 50 μm by

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the 6-T magnetic field generated by a pair of superconducting coils in Helmholtz configuration. Tin atoms were brought to the trapping region by injecting a tenuous molecular beam of tetra-*i*-propyltin ( $C_{12}H_{28}Sn$ ), which dissociated while crossing the electron beam. Tin HCl were subsequently produced through electron impact ionization, while tuning the electron beam acceleration potential allowed the selection of the desired charge states. The heavier tin HCl were trapped longitudinally by the trapping potential created using a set of drift tubes and radially by the electron beam space-charge potential, while the lighter elements in the compound (C, H) escaped from the trap. Electron collisions populate levels close to the corresponding ionization continua and profusely induce fluorescence which was recorded by two instruments: a flat-field grating spectrometer and a Czerny-Turner spectrometer for EUV and optical emissions, respectively.

In the EUV spectrometer [36], light emitted by the trapped ion cloud is diffracted by a 1200-lines/mm flat-field, grazing-incidence grating [37] and recorded with a Peltier-cooled charge-coupled device (CCD) sensor. Calibration was performed using resolved bright lines of Sn in the 12–17 nm range, for which the wavelengths were known from Ref. [18], yielding a root-mean-square deviation of the calibration function residuals of 0.03 nm. Typical observed linewidths are in the order of 0.04 nm, giving an experimental resolving power  $\lambda/\delta\lambda$  of approximately 300 in the region near 13.5 nm.

In order to measure optical spectra, FLASH-EBIT is equipped with two *in vacuo* and two in-air lenses imaging the ion cloud onto the entrance slit of a 320-mm focal-length Czerny-Turner spectrometer equipped with a 300-lines/mm grating. For wavelength calibration, Ne-Ar and Hg spectral lamps were used, depending on the spectral region. They exhibited an instrument-dominated linewidth of approximately 1 nm at full-width at half-maximum (FWHM) around 400 nm. This setup, despite its relatively low resolving power compared to typical work of the MPIK group, is very convenient for quickly covering the whole optical range in these cases where no data were available.

A typical acquisition cycle consisted of a short calibration of the optical spectrometer, and a series of 30-min-long simultaneous exposures of both the optical and EUV spectrometers. After each acquisition, the electron beam acceleration potential was increased by 10 V, stepping from 137 to 477 V at a constant beam current of 10 mA. This low current gives rise to a modest space-charge potential correction of approximately 25 V [16,36]. The chosen range of the acceleration potential enabled the production of charge states from  $Sn^{7+}$  up to at least  $Sn^{14+}$  [16]. After each energy scan, the grating was rotated to measure an adjacent wavelength range while keeping a certain overlap. Next, the acceleration potential was stepped through its entire range again. This procedure was performed thrice, encompassing the full accessible wavelength range from 260 to 780 nm. Gaussian fits were carried out to determine the centroid positions of the recorded lines. Associated error bars of approximately 0.4 nm are dominated by the calibration uncertainty [16]. Intensities are taken from the Gaussian fits and corrected for the grating efficiency.

### III. THEORY

Two calculation methods are presented in this work. First, we present dedicated *ab initio* CI+MBPT calculations performed with the AMBIT code, and benchmark them by comparison with our experimental data. Second, in order to identify the measured transitions and the associated energy levels, we utilize semiempirical calculations using orthogonal energy scaling parameters which can be tuned to fit the spectral data. We also use the COWAN code results on weighted transition rates  $gA$  to predict line strengths and branching ratios.

#### A. CI+MBPT

The detailed electronic structure of  $Sn^{7+}$ – $Sn^{10+}$  was calculated using the AMBIT code which combines configuration interaction and many-body perturbation theory (CI+MBPT). Full details of this method have been presented previously [26,30,31]. Here, we explain some of the physics and details relevant to the current calculations of tin ions. A more formal discussion, including mathematical details, may be found in Ref. [26]. Atomic units ( $\hbar = m_e = e = 1$ ) are used in this section.

We start with a Dirac-Fock (relativistic Hartree-Fock) calculation in the  $V^N$  approximation. In this approximation, all  $N$  electrons of the tin ion are included in the self-consistency procedure, creating a Dirac-Fock potential and electron orbitals that are optimized for the  $[Kr] 4d^m$  ground-state configuration. This is particularly important for this study because between  $m = 4$  and 7, the  $4d$  orbitals pass through the half-filled subshell ( $4d^5$ ), in which the exchange contribution is maximal. We will use  $Sn^{9+}$  ( $m = 5$ ) as a working example in the following.

A large orbital basis is formed by diagonalizing a set of B splines [38–40] over the Dirac-Fock operator

$$\hat{h}_{DF} = c \boldsymbol{\alpha} \cdot \mathbf{p} + (\beta - 1)m_e c^2 - \frac{Z}{r} + V^N(r). \quad (1)$$

The resulting basis is then ordered by energy. The lowest few valence orbitals in each wave function are close to their “spectroscopic” counterparts, while the higher-energy orbitals, so-called pseudostates, include large contributions from the continuum.

We now form a set of many-body configurations for the CI method. The CI basis includes all configurations formed by allowing single and double excitations from the  $4d^5$  ground-state configuration up to  $8spdf$  orbitals (i.e., including  $5s$ – $8s$ ,  $5p$ – $8p$ ,  $4d$ – $8d$ , and  $4f$ – $8f$  orbitals). The configurations included in CI are defined to be within a subspace here denoted  $P$ ; all others are within its complementary subspace  $Q$ . For each configuration, a complete set of projections is generated, specifying the total angular momentum and projection of each electron in the configuration. These projections are diagonalized over the  $\hat{J}^2$  operator to obtain configuration-state functions (CSFs). The CSFs are diagonal in total angular momentum, projection, and relativistic configuration, and they form the CI basis which we denote  $|I\rangle$ . All CSFs corresponding to configurations in the subspace  $P$  are included in CI.

The many-electron wave function  $\psi$  is expressed as a linear combination of CSFs from the subspace  $P$ ,

$$\psi = \sum_{I \in P} C_I |I\rangle,$$

where the  $C_I$  are obtained from the matrix eigenvalue problem. The Hamiltonian for the CI problem is

$$\hat{H} = E_{\text{core}} + \sum_i \hat{h}_{\text{CI}} + \sum_{i < j} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}, \quad (2)$$

where the indices  $i$  and  $j$  run over the valence electrons only. Note that the one-body operator  $\hat{h}_{\text{CI}}$  is not equal to the Dirac-Fock operator:  $\hat{h}_{\text{CI}}$  has a potential term  $V^{N_{\text{core}}}$  due to the core electrons only. Therefore, the basis orbitals are not eigenvalues of the one-body CI operator, which must then be included explicitly.

Because the size of the CI matrix grows rapidly with the inclusion of additional orbitals, we must account for these configurations using many-body perturbation theory. The matrix-eigenvalue equation for the combined CI+MBPT method in second-order of perturbation theory is

$$\sum_{J \in P} \left( H_{IJ} + \sum_{M \in Q} \frac{\langle I | \hat{H} | M \rangle \langle M | \hat{H} | J \rangle}{E - E_M} \right) C_J = E C_I, \quad (3)$$

where the CSFs  $|M\rangle$  belong to configurations outside of the subspace  $P$ . They are, in fact, in the subspace  $Q$ .

Because of the extremely large number of CSFs in the subspace  $Q$ , it is prohibitively expensive computationally to modify all matrix elements  $H_{IJ}$  directly. Instead, the CI+MBPT method includes Eq. (3) by modifying the radial integrals of the one- and two-body matrix elements [25]. The Slater-Condon rules for calculating matrix elements of Slater determinants ensure that this is equivalent to Eq. (3), except for the energy denominator (for a detailed discussion beyond the scope of this work, see [25,26,41]). Because in this work  $\hat{h}_{\text{DF}} \neq \hat{h}_{\text{CI}}$ , so-called “subtraction diagrams” must be included with terms proportional to  $\hat{h}_{\text{CI}} - \hat{h}_{\text{DF}}$ . These diagrams can become very large when there are many valence electrons since  $V^{N_{\text{core}}} - V^N$  is large, but there is cancellation between some of the largest subtraction diagrams and the three-body MBPT operator [30]. For this reason, it is important to include three-body operators when calculating these tin ions. An alternative is to calculate the orbitals in the  $V^{N-m}$  approximation (equal to  $V^{N_{\text{core}}}$ ) as suggested in Ref. [42]; however, in this case the orbitals are much further from “spectroscopic” and the CI basis must be made considerably larger to correct them. In this work, all one-, two-, and three-body second-order diagrams are included.

Until recently, only core-valence correlations were taken into account using MBPT. These correlations incorporate the effects of configurations  $|M\rangle$ , which include an excitation from the  $N_{\text{core}}$  electrons. It was shown in Ref. [31] that valence-valence correlations could also be included in the same manner. Thus, in this work, valence-valence correlations with excited orbitals up to  $30spdfg$  are included; this incorporates the effect of configurations that have one or two pseudo-orbitals above  $8spdf$ , but have no core excitations. Furthermore, the valence-valence subtraction diagrams presented in Ref. [31] are also

TABLE I. Energy levels of the  $\text{Sn}^{9+} 4d^5$  configuration (in  $\text{cm}^{-1}$ ) calculated by AMBiT CI+MBPT code. The first column gives the approximate  $LS$ -term of the calculated energy levels. The CI values give the energy as calculated using only configuration-interaction, while the  $\Sigma^{\text{core}}$ ,  $\Sigma^{\text{val}}$ , Breit, and QED are the successive corrections to the CI energy by including core-valence MBPT, valence-valence MBPT, Breit, and QED contributions, respectively. The total energy including all corrections is also presented, as are the available experimentally determined values and the differences  $\Delta E$  (Expt minus Total) (see main text).

Level	Energy ( $\text{cm}^{-1}$ )						Expt	$\Delta E$
	CI	$\Sigma^{\text{core}}$	$\Sigma^{\text{val}}$	Breit	QED	Total		
$6S_{5/2}$	0	0	0	0	0	0	0	0
$4G_{5/2}$	39469	-4203	-2141	284	-28	33381	33784	403
$4G_{7/2}$	42840	-4593	-1833	11	-2	36421	36874	453
$4G_{11/2}$	43706	-4676	-1756	-132	3	37145	37535	390
$4G_{9/2}$	44212	-4606	-1734	-120	7	37759	38170	411
$4P_{5/2}$	43692	-3649	-2067	60	-4	38032	38315	283
$4P_{3/2}$	44398	-3174	-2316	138	-12	39035	39190	155
$4P_{1/2}$	47021	-2711	-2281	32	-1	42060		
$4D_{7/2}$	51789	-4612	-2351	-98	8	44737	44915	178
$4D_{5/2}$	55276	-3752	-2521	-106	10	48907		
$4D_{1/2}$	55286	-3812	-2310	-190	22	48996		
$4D_{3/2}$	56627	-3340	-2319	-241	25	50753		
$2I_{11/2}$	62330	-7093	-2270	-110	5	52863	53692	829
$2I_{13/2}$	65768	-7344	-2186	-318	18	55937	56792	855
$4F_{7/2}$	66988	-5849	-3102	60	-9	58088	58487	399
$2D_{5/2}$	65152	-3732	-2808	-150	18	58479	58756	277
$4F_{3/2}$	65795	-4189	-3004	-17	4	58588	58891	303
$4F_{9/2}$	67897	-5777	-3104	-33	-2	58981	59417	436
$4F_{5/2}$	71298	-4933	-2901	-193	20	63291	63643	352
$2H_{9/2}$	74999	-5532	-3005	-207	17	66273	66824	551
$2G_{7/2}$	75308	-4572	-3146	-292	27	67325	67698	373
$2D_{3/2}$	76386	-4767	-3007	-325	34	68321		
$2F_{7/2}$	80012	-6786	-3048	-351	32	69859	70199	340
$2F_{5/2}$	81165	-5771	-3713	-170	19	71529	71806	277
$2H_{11/2}$	82714	-5812	-2746	-527	45	73674	74311	637
$2F_{7/2}$	85363	-6347	-3651	-323	31	75073	75470	397
$2G_{9/2}$	85188	-6283	-3135	-465	42	75347	75795	448
$2F_{5/2}$	90363	-7289	-4328	-145	16	78616	78700	84
$2S_{1/2}$	87288	-5338	-2910	-647	64	78457		
$2D_{3/2}$	99595	-6555	-4503	-149	18	88405	88649	244
$2D_{5/2}$	102913	-6472	-4465	-373	37	91640	91927	287
$2G_{9/2}$	111086	-8615	-4736	-273	23	97485	98217	732
$2G_{7/2}$	112328	-8403	-4689	-341	32	98927	99649	722

included. They vanished in that work because  $\hat{h}_{\text{DF}}$  was the same as  $\hat{h}_{\text{CI}}$ , but play a role in the present context.

Finally, Breit and Lamb shift corrections are taken into account. The latter include the vacuum polarization (Uehling) [43] and self-energy [44] corrections in the radiative potential formulation of Flambaum and Ginges [45]. Because both of these effects arise from the electron density near the nucleus, they have a fairly constant ratio for all the levels we calculated.

The ion  $\text{Sn}^{9+}$  has a half-filled  $4d$  shell, and for it the results are presented broken down into different contributions (Table I). The MBPT corrections are separated into

TABLE II. Mean differences and standard deviation between our CI+MBPT calculations and experiment for measured transitions in different Sn ions (all this work).

Ion	Configuration	No. of lines	$\Delta E_{\text{th-expt}}$ (eV)
Sn <sup>7+</sup>	4d <sup>7</sup>	8	-0.004 ± 0.013
Sn <sup>8+</sup>	4d <sup>6</sup>	24	-0.005 ± 0.023
Sn <sup>9+</sup>	4d <sup>5</sup>	30	-0.010 ± 0.026
Sn <sup>10+</sup>	4d <sup>4</sup>	28	-0.008 ± 0.034

core-valence contributions  $\Sigma^{\text{core}}$  (which correspond to un-freezing of the 4sp3d core), and valence-valence contributions  $\Sigma^{\text{val}}$  (introduced in Ref. [31]), which account for configurations that include orbitals above 8spdf. The column marked QED shows the vacuum polarization and self-energy corrections.

The Sn<sup>9+</sup> and Sn<sup>10+</sup> ions were treated with CI+MBPT calculations using only electron excitations (the approach of Refs. [25,26]). However, as the number of valence electrons increases, this electron-only approach becomes inaccurate due to very large contributions from the subtraction diagrams. To avoid this inaccuracy, the particle-hole CI+MBPT calculations are instead used for the Sn<sup>7+</sup> and Sn<sup>8+</sup> ions. This approach, described in Ref. [31], places the Fermi level above the 4d shell and treats the 4d<sup>m</sup> ground-state configuration as a corresponding number of valence holes in an otherwise filled shell. That is, the one-body CI operator includes the potential due to a completely filled 4d shell,  $V^{N_{\text{core}}+10}$ . Our complete CI+MBPT results for Sn<sup>7+</sup> and Sn<sup>8+</sup> in this particle-hole framework are presented in Tables II–VI together with the results for Sn<sup>9+</sup> and Sn<sup>10+</sup>.

TABLE III. Energy levels of the ground configuration 4d<sup>7</sup> for Sn<sup>7+</sup> (all in cm<sup>-1</sup>). The level energies  $E_{\text{expt}}$  were experimentally determined in Ref. [6] and are provided along with their approximate LS-term.  $E_{\text{CI+MBPT}}$  are energy levels calculated by the AMBiT code. The difference between the two data sets  $\Delta E_{\text{CI+MBPT}} = E_{\text{expt}} - E_{\text{CI+MBPT}}$  are presented in the last column.

Level	Term	$E_{\text{expt}}$	$E_{\text{CI+MBPT}}$	$\Delta E_{\text{CI+MBPT}}$
0	<sup>4</sup> F <sub>9/2</sub>	0	0	0
1	<sup>4</sup> F <sub>7/2</sub>	6986	6944	42
2	<sup>4</sup> F <sub>5/2</sub>	10341	10318	23
3	<sup>4</sup> F <sub>3/2</sub>	12153	12137	16
4	<sup>4</sup> P <sub>3/2</sub>	18280	18126	154
5	<sup>4</sup> P <sub>5/2</sub>	20373	20123	250
6	<sup>2</sup> G <sub>9/2</sub>	22636	22523	113
7	<sup>4</sup> P <sub>1/2</sub>	23946	23698	248
8	<sup>2</sup> G <sub>7/2</sub>	29001	28924	77
9	<sup>2</sup> H <sub>11/2</sub>	30312	30047	265
10	<sup>2</sup> P <sub>3/2</sub>	30657	30487	170
11	<sup>2</sup> D <sub>5/2(3)</sub>	33670	33762	-92
12	<sup>2</sup> P <sub>1/2</sub>	35458	35329	129
13	<sup>2</sup> H <sub>9/2</sub>	37751	37486	265
14	<sup>2</sup> D <sub>3/2(3)</sub>	44177	44051	126
15	<sup>2</sup> F <sub>5/2</sub>	45452	45083	369
16	<sup>2</sup> F <sub>7/2</sub>	49476	49087	389
17	<sup>2</sup> D <sub>3/2(1)</sub>	73321	71994	1327
18	<sup>2</sup> D <sub>5/2(1)</sub>	75377	75089	288

## B. Orthogonal energy parameters

Line and level identifications in the Sn<sup>8+</sup>–Sn<sup>10+</sup> ions were performed using the *ab initio* MCDF (multiconfiguration Dirac-Fock) code [46], followed by semiempirical calculations based on the orthogonal energy scaling parameters methods for the predictions of the energy levels. The orthogonal parameters method [34,47] has several advantages in comparison with the more usual Slater-Condon approximation as, for instance, used in the COWAN code [35]. First, the energy parameters are maximally independent, facilitating the fitting of the radial integrals of the interactions to experimental energy levels. Second, it is possible to include an additional number of small interactions, such as two-particle magnetic and three- and four-particle electrostatic parameters. These qualities of the orthogonal parameters method, in general, improve the agreement between calculated and measured energy levels when sufficient experimental data are available to fit its parameters. This method has been shown to be apt even for complex electronic configurations, where configuration-interaction plays a relevant role [33]. For instance, it has been applied successfully in the identification of 4d<sup>4</sup>–4d<sup>3</sup> 5p transitions in Pd<sup>6+</sup> [48] reducing the standard deviation of the fits up to nine times compared to the COWAN code.

Prediction of the energy levels in the 4d<sup>6</sup>–4d<sup>4</sup> configuration in the Sn<sup>8+</sup>–Sn<sup>10+</sup> spectra was performed by interpolation of the energy parameters between Sn<sup>6+</sup> (4d<sup>8</sup>) [49,50], Sn<sup>7+</sup> (4d<sup>7</sup>) [6], Sn<sup>11+</sup> (4d<sup>3</sup>) [16], and Sn<sup>12+</sup> (4d<sup>2</sup>) [16]. These spectra were recalculated in the framework of the orthogonal parameters to determine the scaling parameters needed for the interpolation (see also Sec. IV C).

The transition probabilities of the magnetic dipole (*M*1) transitions were calculated using the COWAN code. In first approximation, the COWAN code energy levels were fitted to the energy levels predicted with the orthogonal parameters method to determine the level wave functions. The transition probabilities estimated with these wave functions were then used for the spectrum analyses and for the identification of the spectral lines. The energy levels established from the identified lines were optimized using Kramida’s code LOPT [51]. The line uncertainty relevant for the optimization was taken to be 0.4 nm corresponding to 60 cm<sup>-1</sup> at 260 nm and 10 cm<sup>-1</sup> at 600 nm. The uncertainty was increased only for doubly classified (i.e., lines that can be ambiguously assigned to two transitions), blended, or masked lines. Final transition probabilities were obtained after the fitting of the COWAN code to the experimentally established levels. The details of the identifications are given in the following section.

## IV. RESULTS

In the following, we present optical and EUV spectra of tin ions in the resolved charge states Sn<sup>7+</sup>–Sn<sup>10+</sup> (see Figs. 1–3). We interpret the data using orthogonal parameters and semiempirical COWAN code calculations, delivering the most complete data set available to date for the ground configurations of Sn<sup>8+</sup>–Sn<sup>10+</sup> (with the semiempirical results providing data also on level energies that were not directly probed experimentally). A detailed comparison of the thus obtained lines and levels with the CI+MBPT calculations is

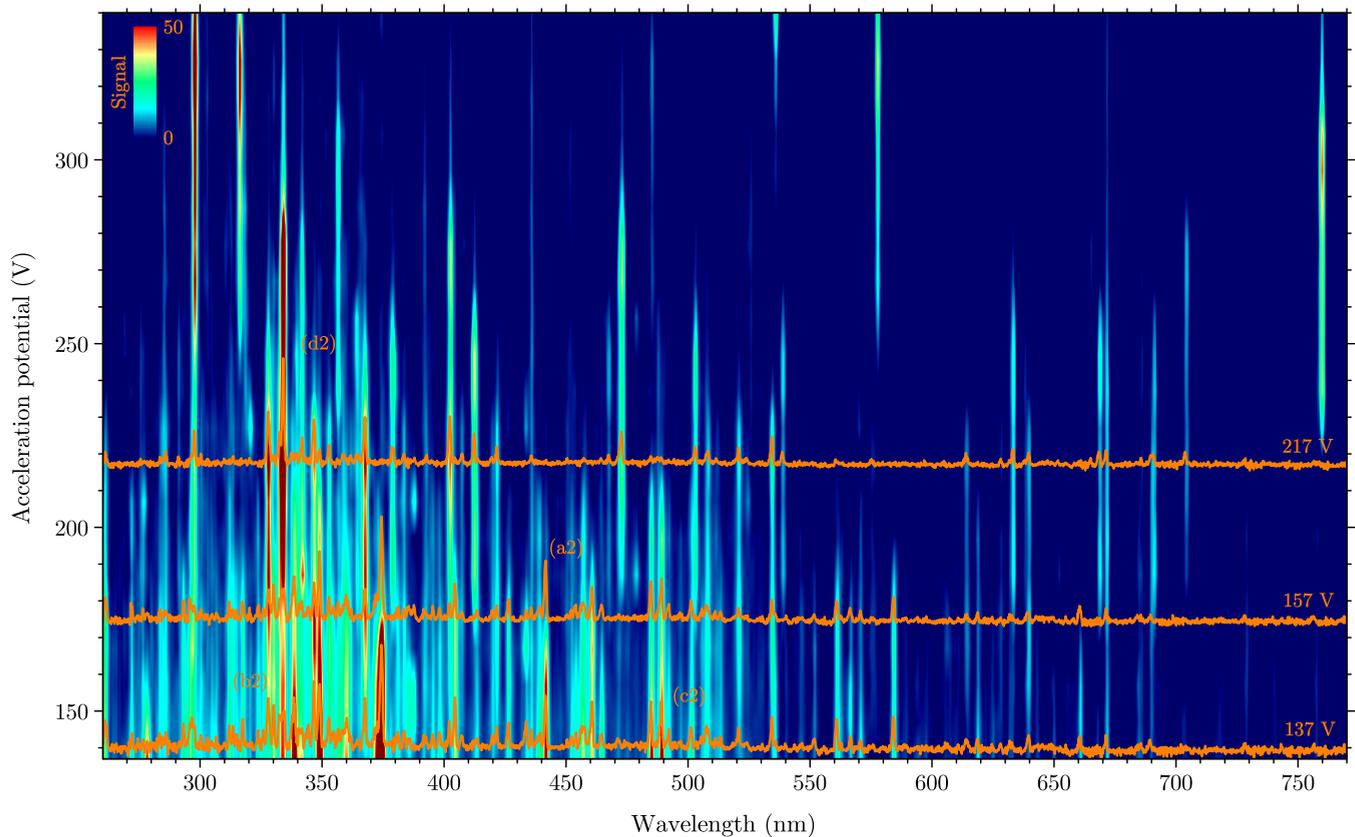


FIG. 1. Spectral map of Sn ions in the optical regime obtained by interpolating discrete spectra acquired at different electron beam energies (uncorrected for space-charge effects). The inset color map represents the fluorescence signal strength scale in arbitrary units. The projections highlight spectra at three acceleration potentials at which the fluorescence of a certain charge state is highest. The lines labeled with (a2), (b2), (c2), and (d2) are shown in more detail in Fig. 3, alongside with the features recorded in the EUV (Fig. 2) to assign the charge state.

presented. Furthermore, we perform a comparison with energy levels available from existing data obtained in the EUV regime. We discuss first the charge-state identification and second the line identifications. All results are summarized in Tables II–VI. The result of line and level identifications is presented in the form of Grotrian diagrams in Fig. 4.

### A. Charge-state identification

In an EBIT, charge-state identification can be performed by evaluating the intensities of groups of spectral lines belonging to the same charge state as a function of the electron beam acceleration potential [16,36,52]. The doubly peaked structure of the  $\text{Sn}^{10+}$  fluorescence curve (cf. Fig. 3), and its premature onset, has been previously observed in the optical domain for the charge states  $\text{Sn}^{11+}$ – $\text{Sn}^{14+}$  [16]. This phenomenon was interpreted as being caused by the existence of strongly populated high- $J$  metastable states. They act as stepping stones for reaching the next charge state at an energy below the corresponding ionization threshold, which is derived from the ground-state binding energy. Moreover, as in Ref. [16], we observe that the onset of charge breeding of  $\text{Sn}^{8+}$  and  $\text{Sn}^{9+}$  takes place well before the respective ionization potentials are reached. Once again, this is a signature of the presence of metastable states.

In this work, the charge-state identification was somewhat hampered due to the low-energy onset of the charge states

$\text{Sn}^{7+}$ ,  $\text{Sn}^{8+}$ , and  $\text{Sn}^{9+}$ , which could not be clearly discerned in the optical data. Therefore, we relied on simultaneously obtained charge-state-resolved EUV spectra to assign optical spectra to their respective charge states. The EUV spectra were compared to previously observed clusters of lines [17,18], as shown in Fig. 2. These lines stem from transitions to the ground configurations  $[\text{Kr}] 4d^m$  ( $m = 7-4$ ) from the  $4p^6 4d^{m-1} 4f + 4p^5 4d^{m+1}$  excited electronic configurations. Four main features have been identified in the EUV spectra that could reliably be attributed to the charge states of interest. Therefore, tracking these features as a function of the electron beam acceleration potential and comparing them to their counterparts in the optical enabled the charge-state assignment, as shown in Fig. 3.

### B. Line and energy level identification

Line wavelengths and intensities for each charge state were determined at the acceleration potential that maximized their yield by fitting Gaussians to the data. The intensities, given by the curve areas, were corrected for the grating spectral efficiency. Their uncertainties arise from statistics, as well as chromatic aberrations of the optical system and its finite aperture [16]. They are minimized when comparing close-lying wavelengths. The strong magnetic field of the EBIT could potentially influence the observed intensities as

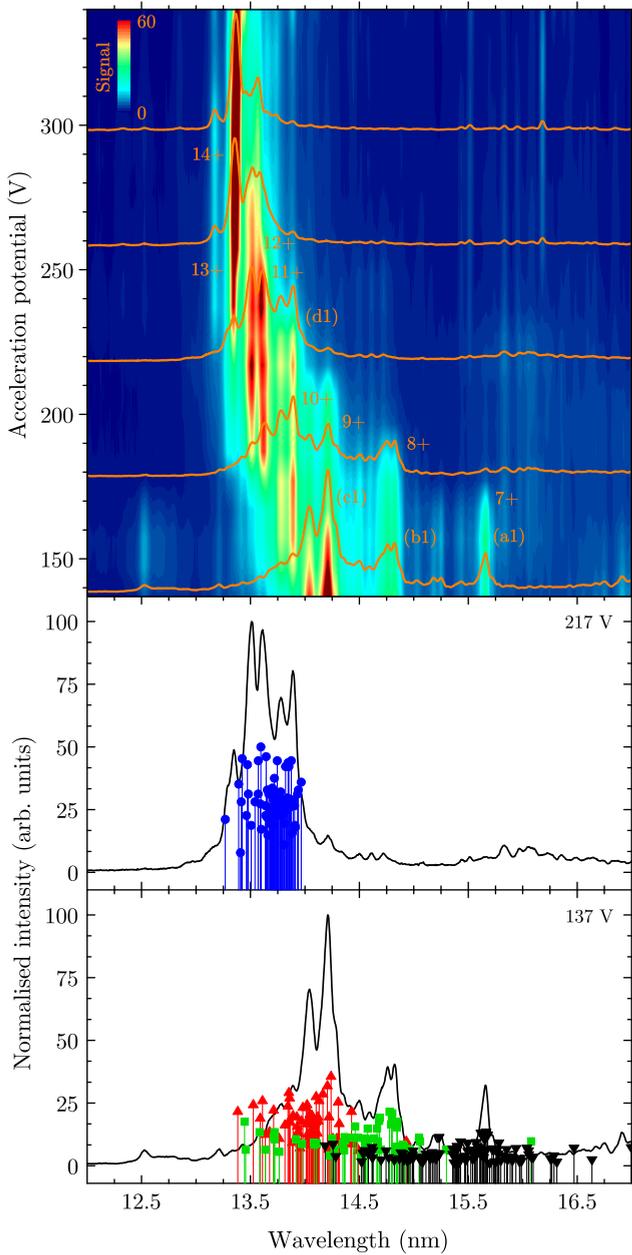


FIG. 2. (Upper) Spectral map of Sn ions in the extreme ultraviolet obtained by interpolating discrete spectra acquired at 10-V acceleration potential steps starting from 137 V (uncorrected for space charge). The inset color map indicates the fluorescence strength scale in arbitrary units. The overlaid spectra are individually scaled for visibility and spaced by approximately 40 V, highlighting the onset of various spectral features. Labels (a1), (b1), (c1), and (d1) indicate features shown in detail in Fig. 3, where they are used for charge-state identification. (Lower) Spectra obtained at acceleration potentials 137 and 217 V, individually normalized to 100 (arb. units). Scatter points represent transitions previously observed in a vacuum spark discharge [17,18], arbitrarily scaled for visibility and to facilitate comparison (black inverted triangles: Sn<sup>7+</sup>; green squares: Sn<sup>8+</sup>; red triangles: Sn<sup>9+</sup>; blue circles: Sn<sup>10+</sup>).

well. First, this is because emission is polarized. However, since no polarizing optical element was employed, such possible effects are minimal and limited to small differences in the

overall detection efficiency of the two polarizations. Second, anisotropic excitation of magnetic sublevels by the directed electron beam leads indeed to a differential population of electronic states, but the large number of cascades from highly excited levels feeding the upper states of the optical transitions under study leads to an almost complete depolarization of the emission lines. These effects have been studied in detail in the case of x-ray transitions with far less complex cascading networks populating their respective upper states [55–58], and should lead in the present cases to almost isotropic emission.

Listed in Tables IV and V are their centers and integrated intensities obtained by Gaussian fits. A direct comparison to the CI+MBPT calculations is also displayed. Deviations from the experimental data are quantified by the mean difference and standard deviation between theory and experiment for all measured transitions (see Table II). We find very good agreement with experiment for all Sn ions studied. In the following, the results per charge state are discussed in detail.

### 1. Spectrum of the Sn<sup>7+</sup> ion

All levels of the  $4d^7$  configuration in Sn<sup>7+</sup> are known from the analysis of the  $4d^7-4d^65p$  transitions in the EUV region [6], with estimated uncertainties of less than  $12\text{ cm}^{-1}$ . The position of the  $M1$  optical transitions can be accurately obtained from the energy differences of these levels. Weighted transition rates  $gA$  for these  $M1$  transitions were calculated by the COWAN code to facilitate the comparison, shown in Table IV, of the eight lines measured in this work with the energy levels in Ref. [6]. Most of the transitions that are predicted from the available structure [6] have a relatively small calculated  $gA$  value, and as such are not observable in our experiments. Three of the predicted stronger transitions (here taking  $gA > 35\text{ s}^{-1}$ ), at 372.1 nm ( $^2F_{5/2}-^2D_{3/2}$ ), nearby 372.6 nm ( $^2G_{9/2}-^2F_{7/2}$ ), and at 488.4 nm wavelength ( $^2G_{7/2}-^2F_{7/2}$ ), were not reliably identified. In all three instances, this can be explained by line blending and by masking of such transitions by stronger emissions of the other charge states in the trap. The differences between our experimental wavelengths and wavelengths predicted from Ref. [6] are well within mutual uncertainties, which are dominated by the 0.4-nm uncertainty in our spectrometer calibration. Branching ratios cannot straightforwardly be used for comparison purposes, as the relevant observed transition sets (between levels 1-8/2-8, and 15-18/16-18, see Table IV and Fig. 4) are affected by blends with neighboring lines. We do not experimentally reinvestigate the  $4d^7$  configuration in Sn<sup>7+</sup> level structure because of the limited number of lines here well resolved and the high-accuracy and detailed results available from Ref. [6]. The good agreement between the present data and previous experimental observations serves as further validation of our charge-state identifications.

### 2. Spectrum of the Sn<sup>8+</sup> ion

The list of the identified transitions between levels within the  $4d^6$  configuration is presented in Table V. In total, 22 spectral lines were uniquely identified. Of these, we found 17 levels connected to the ground  $^5D_4$  level. Identification of 9 levels is supported by observation of 11 Ritz combinations. The level energies, optimized by LOPT [51], are presented in

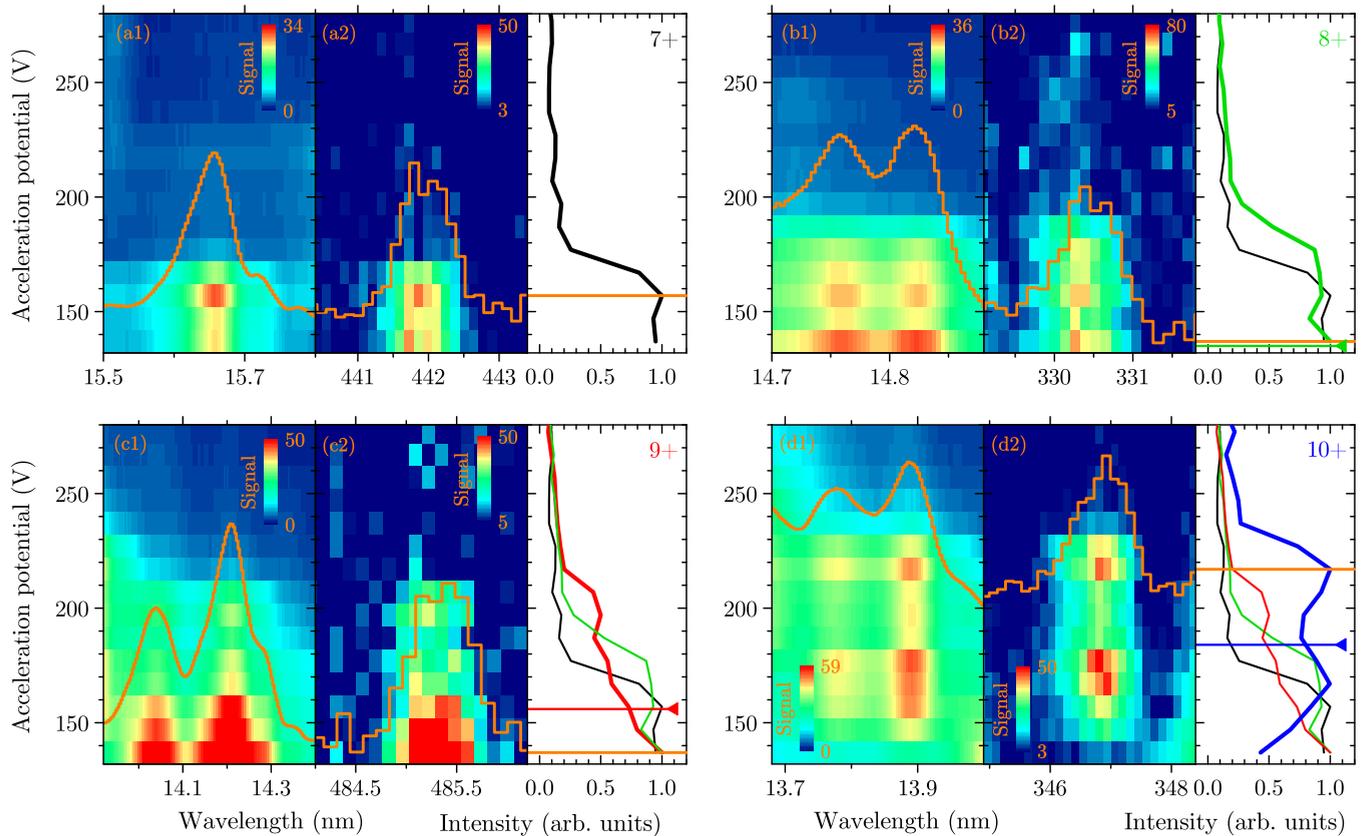


FIG. 3. Enlarged view of selected features from Figs. 1 and 2. Independent color map scales for the fluorescence signal strength are given in arbitrary units. The fluorescence curves are determined by the averaged projections of all line intensities onto the acceleration potential axis of all the lines belonging to the same charge state, normalized to 1 at their respective maximum. Arrows indicate theoretical ionization potentials ( $\text{Sn}^{6+}$ : 113 eV,  $\text{Sn}^{7+}$ : 135 eV,  $\text{Sn}^{8+}$ : 156 eV, and  $\text{Sn}^{9+}$ : 184 eV [53,54]). Vertical axes show acceleration potentials (not corrected for the space-charge contribution).

Table VI with their respective uncertainties. The fitting of the orthogonal parameters was performed with these optimized levels. The resulting optimized sets of parameters are given in Table VII. The 354.8-, 360.2-, and 381.2-nm lines are isolated lines, i.e., the upper and lower levels of these lines are not involved in any other transition. Therefore, the lower levels of these isolated transitions were placed at the position as calculated with the orthogonal parameters method, with an estimated uncertainty of  $16 \text{ cm}^{-1}$  (one-standard-deviation value of the orthogonal parameters fit to the experimental values). These levels were not used in the parameter-fitting procedure.

Most levels can be uniquely designated by the largest contributor in the  $LS$ -coupling decomposition of their wave functions. For example, the level labeled  $^5D_3$  in Table VI is composed of 97%  $^5D_3$ , 2%  $^3F_3(2)$ , and 1%  $^3D_3$ . Here, the number in brackets serves to distinguish between different  $LSJ$  wave functions designated by the same  $LSJ$  values, supplementing a sequential index as defined by Nielson and Koster [59]. In this intermediate coupling regime, the choice of such a designating  $LS$ -term does not imply this term also provides the largest, or any at all, contribution to the transition magnetic dipole matrix elements. Indeed, some transitions occur that appear to violate  $M1$  selection rules based on the chosen  $LS$ -term denomination, but nonetheless satisfy them

since they arise from the admixture of other appropriate terms in the  $LS$  decomposition of both the upper- and lower-state wave functions.

Two exceptions to the naming convention explained in the previous paragraph are the  $^1S_0(4)$  and  $^3F_4(2)$  levels in the  $\text{Sn}^{8+}$  ground configuration, which we uniquely designate by the second-largest component of the wave-function decomposition.

There are seven branched upper levels [numbered 6, 9, 11, 20, 21, 27, and 30 in Table VI (cf. Fig. 4)]. Most of the associated COWAN-calculated branching ratios are in reasonable agreement with the experimental data, except for lines affected by blends (such as in the ratios 9-21/11-21 and 0-9/1-9) or for short-wavelength transitions below 300 nm (featuring in the ratio 0-11/1-11 at 293.3 nm), the intensities of which are affected by a significant drop of detection efficiency. This reduction could not be assessed in our experimental setup.

The levels of  $\text{Sn}^{8+}$  found from EUV measurements on vacuum sparks [18] belong to four disjointed groups, where levels within a single group are connected to each other by measured lines but no transitions were identified connecting the different groups. The uncertainty of the level energies within each of the groups was estimated at  $10 \text{ cm}^{-1}$ , but between the groups as several  $100 \text{ cm}^{-1}$ . For this reason in Table VI, which contains comparison of the energies of

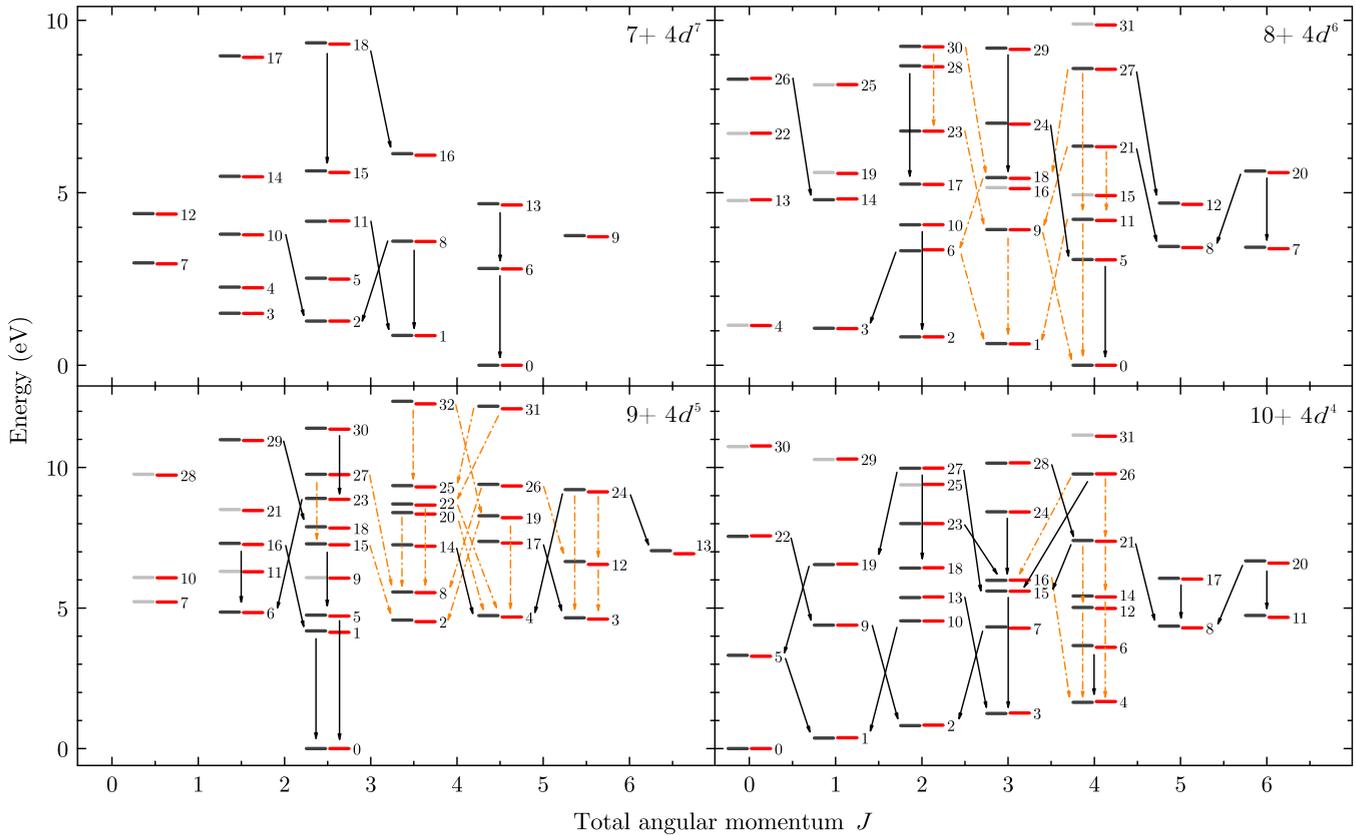


FIG. 4. Grotrian diagrams for the ions  $\text{Sn}^{7+}$ – $\text{Sn}^{10+}$ . The energy levels in dark gray (left) are determined experimentally, for  $\text{Sn}^{7+}$  they are taken from [6], whereas the  $\text{Sn}^{8+}$ – $\text{Sn}^{10+}$  energy levels are results of this work (see Table VI). Light gray levels (left) in  $\text{Sn}^{8+}$ – $\text{Sn}^{10+}$  could not be determined experimentally, and thus the values from orthogonal parameters calculations are used. Levels in red (right) are calculated with the AMBIT code. Arrows indicate experimentally observed optical transitions. Orange (dashed-dotted) arrows indicate transitions which are part of one or more Ritz combinations.

these levels ( $E_{\text{vs}}$ ) with our results ( $E_{\text{expt}}$ ), the four groups are identified by their respective uncertainties  $z_i$  ( $i = 1$ – $4$ , as given in Ref. [18]). These values should be interpreted as systematic common shifts, with respect to the ground state. Statistics of the differences  $\Delta E_{\text{vs}}$  (see Table VI) between previous identifications and the current experimental results provide a meaningful comparison between the two data sets. It is found that the two levels with shift  $z_1$  are consistent with the present experimental values, with differences  $\Delta E_{\text{vs}}$  equaling  $-36$  and  $-22 \text{ cm}^{-1}$ , the scatter in which is well within the uncertainty of  $16 \text{ cm}^{-1}$  obtained from the orthogonal parameter fitting of the experimental data. In contrast, the average value of  $z_2$  is found to be  $270 \pm 377 \text{ cm}^{-1}$ , where the latter number represents the one-standard-deviation spread in the former. This large spread, compared to the experimental uncertainties (see Table VI), indicates that the respective levels from the previous work [18] are not consistent with the present experimental values and that the classification of EUV transitions therein requires a revision. The consistency of the remaining shifts for  $\text{Sn}^{8+}$ ,  $z_3$  and  $z_4$ , cannot be ascertained from our data.

### 3. Spectrum of the $\text{Sn}^{9+}$ ion

Table V contains 28 identified lines between the levels of the  $4d^5$  configuration in  $\text{Sn}^{9+}$ . Two lines, at 296.0- and 457.0-nm

wavelength, are doubly classified. Measured intensities of branching ratios (from the 10 upper levels numbered 15, 16, 19, 20, 22, 24, 26, 27, 31, and 32 in Table VI, cf. Fig. 4) are in reasonable agreement with the COWAN calculations, except for blended or doubly classified lines (affecting the transition coupling levels 2-15, 2-19, 4-20, 12-24, and 8-27), and ultraviolet transitions (8-27, 3-24, and 4-24) due to the lower detection efficiency, similar to the identified  $\text{Sn}^{8+}$  lines. With the identified transitions, 24 levels connected to the ground  ${}^6S_{5/2}$  level were established. Their values, optimized using LOPT [51], are listed in Table VI with their respective uncertainties. The parameter fitting to the available levels resulted in an uncertainty of  $41 \text{ cm}^{-1}$  (one standard deviation of the fit). As in the case of  $\text{Sn}^{8+}$ , the majority of the levels can be uniquely designated by the largest component of the  $LS$ -coupling decomposition of the wave function. Four levels are named by the second-largest component. Only a single transition, at 399.9-nm wavelength, is found to be isolated. Thus, as before, the energy for the corresponding lower level  ${}^4F_{5/2}$  was assumed to be equal to the value obtained from the orthogonal parameters method. These two isolated levels were not used in the parameter fitting. In addition to the agreement with the calculations, 15 levels are connected by transitions composing four Ritz combination chains, thus further supporting the identifications.

From the analysis of EUV vacuum spark observations [18], 21 levels of the  $4d^5$  configuration were found as one group not connected to the ground  $^6S_{5/2}$  level in the  $\text{Sn}^{9+}$  spectrum. The uncertainty of these levels relative to the ground level was estimated as several  $100 \text{ cm}^{-1}$  common to the whole group, as indicated by the single value  $z_5$  in Table VI. As in  $\text{Sn}^{8+}$ , the comparison between vacuum spark measurements and the current results allows to obtain information on the systematic uncertainty  $z_5$ . We find that the average value of this systematic shift is  $460 \pm 422 \text{ cm}^{-1}$ , the one-standard-deviation spread of which exceeds the current uncertainty of  $41 \text{ cm}^{-1}$  by a factor 10, and therefore points to inconsistencies in the previous assignments. From the statistics of the differences  $\Delta E_{vs}$  (see Table VI), we identify two groups of levels with common deviations  $131 \pm 10$  and  $197 \pm 13 \text{ cm}^{-1}$ . These groups comprise four (level numbers 3, 12, 13, and 24) and five levels (numbered 1, 4, 5, 6, and 14), respectively. However, the identification in Ref. [18] of more than half of the levels in the  $\text{Sn}^{9+} 4d^5$  configuration presents too large deviations ( $\Delta E_{vs} > 250 \text{ cm}^{-1}$ ) from current values.

#### 4. Spectrum of the $\text{Sn}^{10+}$ ion

The list of 26 identified lines between the levels of the  $4d^4$  configuration is presented in Table V. Three of these (at 297.4-, 328.1-, and 614.1-nm wavelength) are doubly classified because they can be ambiguously assigned to two transitions. The measured intensities of the branched transitions (from upper levels 20, 21, 26, and 27 in Table VI, cf. Fig. 4) are in reasonable agreement with the COWAN code calculations except for the same two situations seen in the previous subsections: lines observed in the ultraviolet near the edge of the observable range, affecting the branching ratios related to the upper level 27; blending and double classifications which affect the transitions 15-21, 15-26, 16-26, and 18-27.

The level energies obtained from the analysis of the experimental spectra belong to two isolated groups. One group consists of 23 levels with the  $^5D_1$  level being the lowest in energy. The remaining four levels numbered 2, 7, 9, and 22 form another group. Their relative energy values are optimized using LOPT [51] and are collected in Table VI. All of the found levels can be uniquely designated by the largest component of the  $LS$ -coupling decomposition of the wave function. The position of the two groups relative to the ground  $^5D_0$  level could not be established from the present identifications. Thus, we assume the spacing between the  $^5D_1$  and  $^5D_0$  to be equal to the values obtained from the orthogonal parameters method, with a one-standard-deviation uncertainty of  $14 \text{ cm}^{-1}$  obtained from the fit. In a similar fashion, the value calculated employing the orthogonal parameters method was assigned to the lowest,  $^5D_2$  level of the smaller group. The thus determined energy levels of these groups fall well within statistical uncertainties of the calculated values (cf.  $\Delta E_{\text{orth}}$  in Table VI). However, they were not used in the fitting procedure to determine the energy parameters shown in Table VII.

The level energies thus obtained in this work can be compared to the levels established in Ref. [18] as determined from EUV spectra. The levels in that work form four isolated groups, three of which are not connected to the ground level.

The uncertainties in relative positions of these three groups were estimated to be several  $100 \text{ cm}^{-1}$  [18], parametrized by the values  $z_{6,7,8}$  in Table VI. Analogous to the cases of  $\text{Sn}^{8+}$  and  $\text{Sn}^{9+}$ , the differences  $\Delta E_{vs}$  are used to probe the agreement of our results with the previous analysis. We find average values for the shifts  $z_6 = 339 \pm 95 \text{ cm}^{-1}$  and  $z_7 = 383 \pm 43 \text{ cm}^{-1}$ , when removing a single outlier in the latter group (level number 14). These one-standard-deviation values are reasonably consistent with the experimental uncertainties. Thus, our data support the identification of five levels with shift  $z_6$  and seven levels of the  $z_7$  group. The outlier, as well as the levels with shift  $z_8$ , show much larger discrepancies implying that affected energy levels in Ref. [18] are called into question.

#### C. Orthogonal scaling parameters

The orthogonal parameters for the isonuclear sequence  $\text{Sn}^{6+}$ - $\text{Sn}^{12+}$  obtained from a fit to the experimental levels are collected in Table VII. Here, the orthogonal parameters  $O2$ ,  $O2'$ ,  $E'_a$ , and  $E'_b$  are the orthogonal counterparts of the traditional COWAN parameters  $F^2$ ,  $F^4$ ,  $\alpha$ , and  $\beta$  [16,35]. The one-electron magnetic (spin-orbit) operator  $\zeta(4d)$  and the effective three-particle electrostatic operators  $T1$  and  $T2$  are the same as in the COWAN code and ( $A_c$ - $A_0$ ) are additional two-particle magnetic parameters. The fitting was performed for the matrices of the interacting  $4d^k + 4d^{k-1}5s + 4d^{k-2}5s^2$  configurations,  $k = 8-2$ , for  $\text{Sn}^{6+}$ - $\text{Sn}^{12+}$ , respectively. The energy parameters of the unknown  $4d^{k-1}5s + 4d^{k-2}5s^2$  configurations therein were fixed at the MCDF-calculated values for the average energies and spin-orbit interactions. The corresponding electrostatic and configuration-interaction parameters (the latter ones calculated by the COWAN code) were also kept fixed for the  $4d^{k-1}5s + 4d^{k-2}5s^2$  configurations, after scaling them by 0.85 from their *ab initio* values. The average energy  $E_{\text{av}}$  is defined such that the ground level energy of the  $4d^k + 4d^{k-1}5s + 4d^{k-2}5s^2$  configurations is equal to zero. The two-particle magnetic parameters ( $A_c$ - $A_0$ ) were fixed to the MCDF-calculated values in all instances. For better stability of the fitting parameters,  $E'_b$  in  $\text{Sn}^{11+}$  and  $E'_a$  in  $\text{Sn}^{12+}$  were fixed to the extrapolated values.

Table VII furthermore contains the ratios of the fitted parameters to the parameters obtained from MCDF calculations. Along the  $\text{Sn}^{6+}$ - $\text{Sn}^{12+}$  isonuclear sequence the orthogonal energy parameters and the scaling factors can be approximated by linear or weak quadratic dependencies as is visible from Fig. 5 for the  $O2$  and  $O2'$  parameters. However, the scaling factor for the average energy presents a discontinuity going from the  $4d^5$  to  $4d^6$  configurations. Figure 5 also shows that a similar dependence of the  $E_{\text{av}}$  scaling factors occurs in the Pd and Ag isonuclear sequences. For comparison purposes, the ions  $\text{Pd}^{4+}$ - $\text{Pd}^{8+}$  and  $\text{Ag}^{4+}$ - $\text{Ag}^{9+}$  have been analyzed using the same orthogonal parameters method here used for Sn. The data were taken from Refs. [60] ( $\text{Pd}^{4+}$ ), [61] ( $\text{Pd}^{5+}$ ), [48] ( $\text{Pd}^{6+}$ ), [62] ( $\text{Pd}^{7+}$ ), [63] ( $\text{Pd}^{8+}$ ,  $\text{Ag}^{9+}$ ), [64] ( $\text{Ag}^{4+}$ ), [65] ( $\text{Ag}^{5+}$ ), and [66] ( $\text{Ag}^{6+}$ - $\text{Ag}^{8+}$ ). The three elements strikingly exhibit the same irregularity in scaling factors of the average energies for the configuration  $4d^5$  ( $\text{Sn}^{9+}$ ,  $\text{Ag}^{6+}$ ,  $\text{Pd}^{5+}$ ), which may be related to the fact that the  $4d^5$  configuration is a half-filled shell. This physical phenomenon, resulting from the maximal exchange contribution in half-filled shells, yields

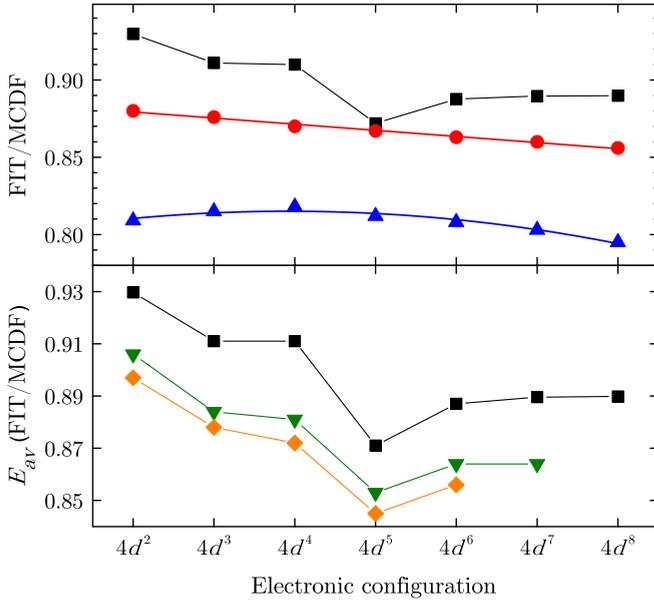


FIG. 5. (Upper) Empirical adjustments of scaling factors compared to the MCDF-calculated values: The ratios (FIT/MCDF) for the electrostatic parameters  $O2$  (red circles) and  $O2'$  (blue triangles), and for the average energy  $E_{av}$  (black squares) were obtained by fitting (FIT) to available data. Solid lines represent quadratic fits. (Lower) Ratio of the semiempirical final value (FIT) to the MCDF value for the average energy  $E_{av}$  along the isonuclear sequences of three elements: black squares  $\text{Sn}^{12+}$  to  $\text{Sn}^{6+}$ , green inverted triangles  $\text{Ag}^{9+}$  to  $\text{Ag}^{4+}$ , orange diamonds  $\text{Pd}^{8+}$  to  $\text{Pd}^{4+}$ .

for them a higher binding energy. It is also the cause for the often discussed and somewhat anomalous ground-state configurations of the chemical elements Cr ( $3d^5 4s$ ), Mn ( $3d^5 4s^2$ ), Mo ( $4d^5 5s$ ), and Tc ( $4d^5 5s^2$ ), and analogously for Eu ( $4f^7 5s^2 p^6 6s^2$ ), Gd ( $4f^7 5s^2 5p^6 5d^1 6s^2$ ), Am ( $5f^7 6s^2 6p^6 7s^2$ ), and Cm ( $5f^7 6s^2 6p^6 6d^1 7s^2$ ). The *ab initio* MCDF calculations do not accurately account for this exchange contribution, and thus the required empirical correction does not follow a continuous trend, otherwise seen

TABLE IV. Experimental vacuum wavelengths  $\lambda_{\text{expt}}$  and line intensities for  $\text{Sn}^{7+}$  within its ground electronic configuration  $[\text{Kr}] 4d^7$ . Spectra recorded at acceleration potential of 157 V, which yielded maximum fluorescence signal. Intensity integrals from Gaussian fits were corrected for the grating efficiency. Wavelengths  $\lambda_{\text{Ritz}}$  are determined from the energy levels of  $\text{Sn}^{7+}$  given in Ref. [6]. Transition probabilities  $g A_{ij, \text{COWAN}}$  are calculated with the COWAN code based on those levels. Wavelengths  $\lambda_{\text{CI+MBPT}}$  are *ab initio* results from our CI+MBPT calculations. “Transition” column shows lower and upper levels as used in Fig. 4. Approximate *LS*-terms are given in the last column. Numbers in brackets are sequential indices as defined by Nielson and Koster [59] to differentiate levels with the same *LSJ* values. Superscripts *bl* mark spectral blends.

$\lambda_{\text{expt}}$ (nm)	Intensity (arb. units)	$\lambda_{\text{Ritz}}$ (nm)	$g A_{ij, \text{COWAN}}$ ( $\text{s}^{-1}$ )	$\lambda_{\text{CI+MBPT}}$ (nm)	Transition (see Fig. 4)	Terms
333.9 <sup>bl</sup>	132	334.2	167	333.3	15-18	$^2F_{5/2}-^2D_{5/2}(1)$
374.5	271	374.8	375	372.9	1-11	$^4F_{7/2}-^2D_{5/2}(2)$
386.0	68	386.1	57	384.6	16-18	$^2F_{7/2}-^2D_{5/2}(1)$
441.9	142	441.8	262	444.0	0-6	$^4F_{9/2}-^2G_{9/2}$
454.3	12	454.2	66	455.0	1-8	$^4F_{7/2}-^2G_{7/2}$
492.2	69	492.2	44	495.8	2-10	$^4F_{5/2}-^4P_{3/2}$
536.0 <sup>bl</sup>	35	535.9	32	537.5	2-8	$^4F_{5/2}-^2G_{7/2}$
660.9	116	661.6	170	668.3	6-13	$^2G_{9/2}-^2H_{9/2}$

in the filling of the  $nd^m$  subshell in the isonuclear sequences observed in this experiment and earlier work [48,60–66] (cf. Fig. 5). In contrast to this, our CI+MBPT calculations with AMBiT include this effect from the start.

## V. CONCLUSIONS

We performed optical and EUV spectroscopy on open- $4d$ -shell ions  $\text{Sn}^{7+}$ – $\text{Sn}^{10+}$  in a charge-state-resolved manner using an electron beam ion trap and recorded 90 magnetic dipole transitions. Line and level identifications were performed using the semiempirical orthogonal parameters method and COWAN code calculations. Our measurements of transitions in the ground configuration of  $\text{Sn}^{7+}$  are in good agreement with previous measurements in the EUV [6]. The lines and level energies obtained for the  $4d^m$  ( $m = 6$ – $4$ ) configurations in  $\text{Sn}^{8+}$ – $\text{Sn}^{10+}$  present the most complete data available to date for these ground configurations, with a total of 79 energy levels experimentally determined. Analogous to our recent work on  $\text{Sn}^{11+}$ – $\text{Sn}^{14+}$  [16], we conclude that the classification of certain cataloged EUV transitions in previous work [18] needs to be revised. Furthermore, these many-valence-electron, open- $4d$ -shell ions provide an excellent testing ground for state-of-the-art CI+MBPT calculations, performed with the AMBiT code. Our *ab initio* calculations are shown to be in very good agreement with our data, validating the predictive power of this theoretical method for these until now challengingly complex electronic systems.

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## APPENDIX

In this Appendix, the tabulated values for all experimentally determined and calculated quantities are presented. Table IV collects the measured *M1* transitions of  $\text{Sn}^{7+}$ . Comparison

is made between these transitions, the transitions as inferred from the levels determined in Ref. [6], and the transitions predicted by CI+MBPT theory. Moreover, in Table III, the AMBiT-calculated energy levels of the  $4d^7$  ground configuration of  $\text{Sn}^{7+}$  are compared to the level energies determined experimentally in Ref. [6]. Table V shows the wavelengths of the lines measured in the optical domain, along with their identification and values as determined from AMBiT

calculations. The energy levels of the ground configuration of the ions  $\text{Sn}^{8+}$ – $\text{Sn}^{10+}$  are given in Table VI. Here, the level energies optimized with Kramida's LOPT [51] are shown alongside with the levels calculated with the orthogonal energy parameters method, the results of CI+MBPT calculations, and levels from previous work [18]. Finally, the orthogonal parameters used in the semiempirical calculations are collected in Table VII.

TABLE V. Vacuum wavelengths and line intensities for  $\text{Sn}^{8+}$ – $\text{Sn}^{10+}$  ions at the acceleration potential maximizing ion fluorescence. Intensities are taken from Gaussian fits and corrected for the grating efficiency. Wavelengths  $\lambda_{\text{orth}}$  are calculated from level energies from Table VI. Transition probabilities  $gA_{ij,\text{COWAN}}$  are determined with the COWAN code using the same level energies. Wavelengths  $\lambda_{\text{CI+MBPT}}$  calculated *ab initio* with the AMBiT CI+MBPT code. “Transition” refers to levels shown in Fig. 4. Configurations and approximate  $LS$ -coupling terms are given in the last two columns (the numbers in brackets are sequential indices [59] for distinction of levels with the same  $LSJ$  values). The superscript *bl* indicates line blends, and the superscript *D* marks doubly classified lines (i.e., which can be ambiguously assigned to two transitions).

Ion	$V_{\text{max}}$ (V)	$\lambda_{\text{expt}}$ (nm)	Intensity (arb. units)	$\lambda_{\text{orth}}$ (nm)	$gA_{ij,\text{COWAN}}$ ( $\text{s}^{-1}$ )	$\lambda_{\text{CI+MBPT}}$ (nm)	Transition (see Fig. 4)	Configuration	Term symbol
8+	137	283.4	18	284	104	283.1	11-27	[Kr] $4d^6$	$^3F_4(2)$ - $^3F_4(1)$
		293.3	42	293	226	295.2	0-11		$^5D_4$ - $^3F_4(2)$
		313.5	17	314	54	315.1	5-24		$^3H_4$ - $^1F_3$
		315.0	16	315	60	315.4	0-9		$^5D_4$ - $^3F_3(2)$
		317.6	45	318	192	316.6	12-27		$^3G_5$ - $^3F_4(1)$
		326.0	15	326	105	325.2	18-30		$^3D_3$ - $^3P_2(1)$
		330.4	70	330	279	331.6	18-29		$^3D_3$ - $^3F_3(1)$
		344.6	25	344	31	346.5	1-11		$^5D_3$ - $^3F_4(2)$
		354.8	33	355	102	352.3	14-26		$^3P_1(2)$ - $^3P_0(1)$
		360.2 <sup>bl</sup>	43	360	120	363.5	17-28		$^3D_2$ - $^3F_2(1)$
		374.5 <sup>bl</sup>	271	374	131	374.6	1-9		$^5D_3$ - $^3F_3(2)$
		381.2	36	381	65	381.3	2-10		$^5D_2$ - $^3F_2(2)$
		392.1	24	392	78	392.2	18-27		$^3D_3$ - $^3F_4(1)$
		404.6	60	404	239	405.7	0-5		$^5D_4$ - $^3H_4$
		426.6	54	428	186	424.0	8-21		$^3H_5$ - $^1G_4(2)$
		434.1	36	434	107	434.2	9-23		$^3F_3(2)$ - $^1D_2(2)$
		460.9	138	461	315	454.5	1-6		$^5D_3$ - $^3P_2(2)$
		505.7	38	505	107	507.3	23-30		$^1D_2(2)$ - $^3P_2(1)$
		513.2 <sup>bl</sup>	85	515	89	516.1	9-21		$^3F_3(2)$ - $^1G_4(2)$
		551.8	40	552	65	542.9	3-6		$^5D_1$ - $^3P_2(2)$
560.9	149	560	185	568.9	7-20	$^3H_6$ - $^1I_6$			
566.8	91	565	87	576.5	8-20	$^3H_5$ - $^1I_6$			
584.4 <sup>D</sup>	118	585	62	599.1	6-18	$^3P_2(2)$ - $^3D_3$			
584.4 <sup>D</sup>	118	585	291	581.0	11-21	$^3F_4(2)$ - $^1G_4(2)$			
9+	137	261.0	66	261	304	262.9	0-5	[Kr] $4d^5$	$^6S_{5/2}$ - $^4G_{5/2}$
		272.0	53	272	418	273.8	3-24		$^4G_{11/2}$ - $^2H_{11/2}$
		276.6	30	276	209	278.4	4-24		$^4G_{9/2}$ - $^2H_{11/2}$
		296.0 <sup>D</sup>	46	296	115	295.2	8-27		$^4D_{7/2}$ - $^2F_{5/2}(1)$
		296.0 <sup>D</sup>	46	296	149	299.6	0-1		$^6S_{5/2}$ - $^4G_{5/2}$
		304.8	25	303	69	306.2	19-32		$^2H_{9/2}$ - $^2G_{7/2}(1)$
		306.6 <sup>bl</sup>	34	306	16	307.7	6-23		$^4P_{3/2}$ - $^2F_{5/2}(2)$
		312.1	62	312	151	311.5	4-22		$^4G_{9/2}$ - $^2F_{7/2}(1)$
		323.9	50	324	105	326.7	8-26		$^4D_{7/2}$ - $^2G_{9/2}(2)$
		333.9 <sup>bl</sup>	70	333	132	335.0	2-19		$^4G_{7/2}$ - $^2H_{9/2}$
		338.7 <sup>bl</sup>	219	339	222	338.2	4-20		$^4G_{9/2}$ - $^2G_{7/2}(2)$
		349.1 <sup>bl</sup>	271	349	297	350.7	4-19		$^4G_{9/2}$ - $^2H_{9/2}$

TABLE V. (*Continued.*)

Ion	$V_{\max}$ (V)	$\lambda_{\text{expt}}$ (nm)	Intensity (arb. units)	$\lambda_{\text{orth}}$ (nm)	$gA_{ij,\text{COWAN}}$ ( $\text{s}^{-1}$ )	$\lambda_{\text{CI+MBPT}}$ (nm)	Transition (see Fig. 4)	Configuration	Term symbol
		356.8	41	354	57	362.0	22-31		${}^2F_{7/2}(1)-{}^2G_{9/2}(1)$
		395.5	64	394	131	398.1	8-22		${}^4D_{7/2}-{}^2F_{7/2}(1)$
		398.3	68	398	86	396.7	1-16		${}^4G_{5/2}-{}^4F_{3/2}$
		399.9	26	401	42	398.2	18-29		${}^4F_{5/2}-{}^2D_{3/2}(2)$
		413.5	43	415	133	419.2	25-32		${}^2F_{7/2}(2)-{}^2G_{7/2}(1)$
		438.9	46	439	155	442.7	8-20		${}^4D_{7/2}-{}^2G_{7/2}(2)$
		439.7	34	441	109	446.2	25-31		${}^2F_{7/2}(2)-{}^2G_{9/2}(1)$
		452.4	32	453	105	444.8	12-26		${}^2I_{11/2}-{}^2G_{9/2}(2)$
		457.0 <sup>D</sup>	58	456	99	453.4	2-15		${}^4G_{7/2}-{}^2D_{5/2}(3)$
		457.0 <sup>D</sup>	58	456	82	458.0	3-17		${}^4G_{11/2}-{}^4F_{9/2}$
		485.0 <sup>bl</sup>	358	484	472	480.5	12-24		${}^2I_{11/2}-{}^2H_{11/2}$
		489.2	414	489	562	489.1	5-15		${}^4G_{5/2}-{}^2D_{5/2}(3)$
		492.2	107	492	69	491.9	4-14		${}^4G_{9/2}-{}^4F_{7/2}$
		497.0	55	496	107	497.3	23-30		${}^2F_{5/2}(2)-{}^2D_{5/2}(2)$
		501.4	167	502	339	496.6	15-27		${}^2D_{5/2}(3)-{}^2F_{5/2}(1)$
		507.6	105	508	103	511.4	6-16		${}^4P_{3/2}-{}^4F_{3/2}$
		570.8	76	569	151	563.8	13-24		${}^2I_{13/2}-{}^2H_{11/2}$
		618.9	137	621	119	636.2	3-12		${}^4G_{11/2}-{}^2I_{11/2}$
10+	217	283.7	13	283	278	283.0	15-27	[Kr] 4d <sup>4</sup>	${}^3F_3(2)-{}^3F_2(1)$
		284.7	18	285	114	286.3	3-15		${}^5D_3-{}^3F_3(2)$
		286.0	33	286	163	287.4	4-16		${}^5D_4-{}^3D_3$
		297.4 <sup>D</sup>	15	297	110	298.8	1-10		${}^5D_1-{}^3F_2(2)$
		297.4 <sup>D</sup>	15	298	91	296.7	15-26		${}^3F_3(2)-{}^3F_4(1)$
		300.5	15	300	73	300.2	3-13		${}^5D_3-{}^3P_2(2)$
		328.1 <sup>D</sup>	112	328	137	333.1	4-14		${}^5D_4-{}^3G_4$
		328.1 <sup>D</sup>	112	328	261	327.4	16-26		${}^3D_3-{}^3F_4(1)$
		346.8	91	347	278	348.2	2-9		${}^5D_2-{}^3P_1(2)$
		349.1 <sup>bl</sup>	271	349	93	349.7	18-27		${}^3D_2-{}^3F_2(1)$
		353.0	38	353	82	359.0	2-7		${}^5D_2-{}^3G_3$
		361.9	14	361	95	363.1	19-27		${}^3D_1-{}^3F_2(1)$
		367.7	145	368	227	374.6	4-12		${}^5D_4-{}^3F_4(2)$
		383.7	28	384	84	378.0	5-19		${}^3P_0(2)-{}^3D_1$
		392.7	68	392	124	390.6	9-22		${}^3P_1(2)-{}^1S_0(2)$
		407.4	38	408	87	402.4	8-21		${}^3H_5-{}^3G_4(2)$
		421.8	26	421	139	428.3	1-5		${}^5D_1-{}^3P_0(2)$
		450.5	18	450	88	444.7	21-28		${}^1G_4(2)-{}^3F_3(1)$
		508.2	127	508	193	510.3	16-24		${}^3D_3-{}^1F_3$
		520.7	84	521	310	518.6	12-21		${}^3F_4(2)-{}^1G_4(2)$
		524.2	17	525	45	516.5	21-26		${}^1G_4(2)-{}^3F_4(1)$
		534.4	117	535	234	538.4	8-20		${}^3H_5-{}^1I_6$
		614.1 <sup>D</sup>	95	614	102	643.9	4-6		${}^5D_4-{}^3H_4$
		614.1 <sup>D</sup>	95	616	106	616.1	16-23		${}^3D_3-{}^1D_2(2)$
		628.3	69	630	116	626.8	14-21		${}^3G_4-{}^1G_4(2)$
		639.9	157	642	201	642.9	11-20		${}^3H_6-{}^1I_6$
		689.5 <sup>bl</sup>	85	690	164	697.0	15-21		${}^3F_3(2)-{}^1G_4(2)$
		728.1	39	727	103	713.1	8-17		${}^3H_5-{}^3G_5$

TABLE VI. Energy levels of the  $\text{Sn}^{8+} 4d^6$ ,  $\text{Sn}^{9+} 4d^5$ , and  $\text{Sn}^{10+} 4d^4$  configurations (in  $\text{cm}^{-1}$ ) adjusted with the LOPT algorithm [51] based on the measured transitions. Levels' labels use approximate  $LS$ -coupling terms. Numbers in brackets display sequential indices [59] to differentiate levels having the same  $LSJ$  values. Uncertainties  $x_j$  ( $j = 1-5$ ) and  $y$  are given as the one standard deviation of the orthogonal parameters fit for the respective configuration:  $x_{1,2,3} = \pm 16 \text{ cm}^{-1}$ ;  $x_4 = \pm 41 \text{ cm}^{-1}$ ;  $x_5 = y = \pm 14 \text{ cm}^{-1}$ . The dispersive energy uncertainty  $D_1$  is close to the minimum uncertainty of separation from other levels, and the energy uncertainty  $D_2$  is that relative to the ground level of the configuration (cf. [51]).  $N$  is the total number of lines connected to the level.  $E_{\text{orth}}$  values are semiempirical energy levels calculated with the orthogonal parameters in Table VII. The  $E_{\text{CI+MBPT}}$  values are *ab initio* energy levels calculated using the AMBIT CI+MBPT code. Differences between experimental and calculated values appear in columns  $\Delta E_{\text{orth}}$  ( $E_{\text{expt}} - E_{\text{orth}}$ ) and  $\Delta E_{\text{CI+MBPT}}$  ( $E_{\text{expt}} - E_{\text{CI+MBPT}}$ ). Energies determined from previous vacuum spark measurement [18] shown as  $E_{\text{vs}}$ ;  $\Delta E_{\text{vs}} = E_{\text{expt}} - E_{\text{vs}}$ , their deviations. Uncertainties in the systematic common shifts of the identified level groups  $z_i$  ( $i = 1-8$ ) [18] are of the order of several hundreds of  $\text{cm}^{-1}$  (see main text). The uncertainty of the relative level energies within each of these groups was estimated at  $10 \text{ cm}^{-1}$ .

Ion	Level	Term	$E_{\text{expt}}$	$D_1$	$D_2$	$N$	$E_{\text{orth}}$	$\Delta E_{\text{orth}}$	$E_{\text{CI+MBPT}}$	$\Delta E_{\text{CI+MBPT}}$	$E_{\text{vs}}$	$\Delta E_{\text{vs}}$
8+ $4d^6$	0	$^5D_4$	0	30	0	3	-5	5	0	0	0	0
	1	$^5D_3$	5 075	13	30	3	5 064	11	5 011	64	5 050	25
	2	$^5D_2$	$6\,634 + x_2$	0	0	1	6 634	0	6 626	8	$6\,670 + z_1$	-36
	3	$^5D_1$	8 648	13	40	1	8 636	12	8 593	55	$8\,670 + z_1$	-22
	4	$^5D_0$					9 345		9 307			
	5	$^3H_4$	24 716	24	24	1	24 726	-10	24 651	65	$24\,685 + z_2$	31
	6	$^3P_2(2)$	26 785	16	40	3	26 771	14	27 011	-226		
	7	$^3H_6$	27 592	13	43	1	27 604	-12	27 270	322	$27\,610 + z_2$	-18
	8	$^3H_5$	27 778	22	39	2	27 781	-3	27 503	275	$27\,710 + z_2$	68
	9	$^3F_3(2)$	31 740	12	30	4	31 736	4	31 709	31	$31\,747 + z_3$	-7
	10	$^3F_2(2)$	$32\,847 + x_2$	28	28	1	32 847	0	32 855	-8	$33\,028 + z_3$	-181
	11	$^3F_4(2)$	34 103	11	30	4	34 102	1	33 873	230	$34\,220 + z_2$	-117
	12	$^3G_5$	37 908	40	59	1	37 930	-22	37 616	292	$37\,950 + z_2$	-42
	13	$^1S_0(4)$					38 532		38 684			
	14	$^3P_1(2)$	$38\,694 + x_1$	0	0	1	38 694	0	38 903	-209		
	15	$^3G_4$					39 872		39 674		$39\,609 + z_2$	
	16	$^3G_3$					41 548		41 310			
	17	$^3D_2$	$42\,340 + x_3$	0	0	1	42 340	0	42 287	53	$41\,787 + z_2$	553
	18	$^3D_3$	43 879	17	40	4	43 887	-8	43 704	175		
	19	$^3D_1$					45 061		44 847			
	20	$^1I_6$	45 421	13	41	2	45 399	22	45 032	389	$45\,440 + z_2$	-19
	21	$^1G_4(2)$	51 219	11	30	3	51 217	2	51 085	134	$50\,840 + z_4$	379
	22	$^3P_0(2)$					54 202		54 250			
	23	$^1D_2(2)$	54 777	13	40	2	54 795	-18	54 742	35		
	24	$^1F_3$	56 613	41	47	1	56 586	27	56 385	228		
	25	$^3P_1(1)$					65 561		65 611			
	26	$^3P_0(1)$	$66\,874 + x_1$	32	32	1	66 875	-1	67 067	-193		
	27	$^3F_4(1)$	69 394	23	40	3	69 401	-7	69 198	196	$68\,566 + z_2$	828
	28	$^3F_2(1)$	$70\,006 + x_3$	31	31	1	70 006	0	69 800	206		
	29	$^3F_3(1)$	74 146	37	54	1	74 144	2	73 860	286	$73\,385 + z_2$	761
	30	$^3P_2(1)$	74 552	14	40	2	74 548	4	74 454	98		
	31	$^1G_4(1)$					79 767		79 565		$79\,186 + z_2$	
	32	$^1D_2(1)$					101 675		101 319		$99\,838 + z_4$	
33	$^1S_0(1)$					131 838		131 874		$130\,008 + z_4$		
9+ $4d^5$	0	$^6S_{5/2}$	0	59	0	2	-17	17	0	0	0	0
	1	$^4G_{5/2}$	33 784	61	61	2	33 748	36	33 381	403	$33\,582 + z_5$	202
	2	$^4G_{7/2}$	36 874	21	70	2	36 834	40	36 421	453	$36\,610 + z_5$	264
	3	$^4G_{11/2}$	37 535	10	90	3	37 576	-41	37 145	390	$37\,399 + z_5$	136
	4	$^4G_{9/2}$	38 170	20	80	5	38 173	-3	37 759	411	$37\,958 + z_5$	212
	5	$^4G_{5/2}$	38 315	16	59	2	38 282	33	38 032	283	$38\,110 + z_5$	205
	6	$^4P_{3/2}$	39 190	16	68	2	39 183	7	39 035	155	$39\,010 + z_5$	180
	7	$^4P_{1/2}$					42 159		42 060			
	8	$^4D_{7/2}$	44 915	15	80	4	44 958	-43	44 737	178	$44\,470 + z_5$	445
	9	$^4D_{5/2}$					49 065		48 907			
	10	$^4D_{1/2}$					49 104		48 996			
	11	$^4D_{3/2}$					50 861		50 753			
12	$^2I_{11/2}$	53 692	8	80	3	53 685	7	52 863	829	$53\,554 + z_5$	138	

TABLE VI. (*Continued.*)

Ion	Level	Term	$E_{\text{expt}}$	$D_1$	$D_2$	$N$	$E_{\text{orth}}$	$\Delta E_{\text{orth}}$	$E_{\text{CI+MBPT}}$	$\Delta E_{\text{CI+MBPT}}$	$E_{\text{vs}}$	$\Delta E_{\text{vs}}$
	13	$^2I_{13/2}$	56 792	12	83	1	56 765	27	55 937	855	$56\,660 + z_5$	132
	14	$^4F_{7/2}$	58 487	16	77	1	58 491	-4	58 088	399	$58\,300 + z_5$	187
	15	$^2D_{5/2}(3)$	58 756	14	60	3	58 721	35	58 479	277	$58\,370 + z_5$	386
	16	$^4F_{3/2}$	58 891	25	66	2	58 848	43	58 588	303		
	17	$^4F_{9/2}$	59 417	28	87	1	59 469	-52	58 981	436	$58\,850 + z_5$	567
	18	$^4F_{5/2}$	$63\,643 + x_4$	0	0	1	63 643	0	63 291	352		
	19	$^2H_{9/2}$	66 824	22	70	3	66 846	-22	66 273	551	$66\,427 + z_5$	397
	20	$^2G_{7/2}(2)$	67 698	18	80	2	67 687	11	67 325	373	$66\,975 + z_5$	723
	21	$^2D_{3/2}(3)$					68 584		68 321			
	22	$^2F_{7/2}(1)$	70 199	18	80	3	70 228	-29	69 859	340	$70\,185 + z_5$	14
	23	$^2F_{5/2}(2)$	71 806	43	80	2	71 837	-31	71 529	277		
	24	$^2H_{11/2}$	74 311	16	80	4	74 338	-27	73 674	637	$74\,195 + z_5$	116
	25	$^2F_{7/2}(2)$	75 470	16	80	2	75 423	47	75 073	397	$74\,385 + z_5$	1 085
	26	$^2G_{9/2}(2)$	75 795	18	80	2	75 816	-21	75 347	448	$75\,345 + z_5$	450
	27	$^2F_{5/2}(1)$	78 700	16	60	2	78 654	46	78 616	84		
	28	$^2S_{1/2}$					78 719		78 457			
	29	$^2D_{3/2}(2)$	$88\,649 + x_4$	25	25	1	88 702	-53	88 405	244		
	30	$^2D_{5/2}(2)$	91 927	16	81	1	91 976	-49	91 640	287	$90\,911 + z_5$	1 016
	31	$^2G_{9/2}(1)$	98 217	18	80	2	98 228	-11	97 485	732	$96\,800 + z_5$	1 417
	32	$^2G_{7/2}(1)$	99 649	21	80	2	99 568	81	98 927	722	$98\,277 + z_5$	1 372
	33	$^2P_{3/2}$					114 830		114 351			
	34	$^2P_{1/2}$					117 607		117 122			
	35	$^2D_{5/2}(1)$					128 906		128 281			
	36	$^2D_{3/2}(1)$					130 802		130 180			
10+ 4d <sup>4</sup>	0	$^5D_0$	0 + y	0	0	0	0	0	0	0	0	0
	1	$^5D_1$	3 043	22	0	2	3 043	0	3 141	-98	3 035	8
	2	$^5D_2$	$6\,590 + x_5$	0	0	2	6 590	0	6 717	-127	6 545	45
	3	$^5D_3$	10 073	49	84	2	10 054	19	10 213	-140	10 005	68
	4	$^5D_4$	13 300	24	70	4	13 315	-15	13 516	-216	13 280	20
	5	$^3P_0(2)$	26 752	27	23	2	26 750	2	26 490	262		
	6	$^3H_4$	29 584	15	75	1	29 589	-5	29 046	538	$29\,380 + z_6$	204
	7	$^3G_3$	$34\,918 + x_5$	32	32	1	34 899	19	34 573	345	$34\,630 + z_6$	288
	8	$^3H_5$	35 147	24	73	3	35 143	4	34 639	508	$34\,814 + z_7$	333
	9	$^3P_1(2)$	$35\,425 + x_5$	33	33	2	35 429	-4	35 438	-13	$35\,048 + z_6$	377
	10	$^3F_2(2)$	36 669	61	61	1	36 666	3	36 613	56	$36\,297 + z_6$	372
	11	$^3H_6$	38 232	10	70	1	38 226	6	37 656	576	$37\,890 + z_7$	342
	12	$^3F_4(2)$	40 490	13	70	2	40 475	15	40 208	282	$40\,130 + z_7$	360
	13	$^3P_2(2)$	43 351	44	95	1	43 377	-26	43 530	-179	$42\,898 + z_6$	453
	14	$^3G_4$	43 777	10	70	2	43 765	12	43 539	238	$43\,710 + z_7$	67
	15	$^3F_3(2)$	45 197	8	70	3	45 196	1	45 146	51	$44\,766 + z_7$	431
	16	$^3D_3$	48 279	35	80	4	48 263	16	48 310	-31	$47\,850 + z_7$	429
	17	$^3G_5$	48 881	8	73	1	48 893	-12	48 663	218	$48\,480 + z_7$	401
	18	$^3D_2$	51 801	30	60	1	51 808	-7	51 885	-84		
	19	$^3D_1$	52 814	31	35	2	52 806	8	52 945	-131		
	20	$^1I_6$	53 860	14	74	2	53 867	-7	53 211	649	$53\,475 + z_7$	385
	21	$^1G_4(2)$	59 693	6	70	6	59 684	9	59 493	200		
	22	$^1S_0(2)$	$60\,890 + x_5$	26	42	1	60 870	20	61 041	-151		
	23	$^1D_2(2)$	64 563	15	80	1	64 549	14	64 542	21		
	24	$^1F_3$	67 957	16	80	1	67 958	-1	67 907	50	$66\,757 + z_8$	1 200
	25	$^3P_2(1)$					75 662		75 823			
	26	$^3F_4(1)$	78 771	14	70	3	78 777	-6	78 854	-83		
	27	$^3F_2(1)$	80 446	50	47	3	80 445	1	80 484	-38		
	28	$^3F_3(1)$	81 891	20	71	1	81 881	10	81 982	-91	$80\,207 + z_8$	1 684
	29	$^3P_1(1)$					82 941		83 107			
	30	$^3P_0(1)$					86 664		86 851			
	31	$^1G_4(1)$					89 965		89 627			
	32	$^1D_2(1)$					112 401		112 544			
	33	$^1S_0(1)$					144 549		145 002			

TABLE VII. Orthogonal energy parameters (all in  $\text{cm}^{-1}$ ) obtained by fitting experimental energy levels (FIT) and ratios of FIT to MCDF code (FIT/MCDF) calculations. Experimental energy levels taken from Refs. [49,50] for  $\text{Sn}^{6+}$ , Ref. [6] for  $\text{Sn}^{7+}$ , this work for  $\text{Sn}^{8+}$ – $\text{Sn}^{10+}$ , and Ref. [16] for  $\text{Sn}^{11+}$ ,  $\text{Sn}^{12+}$ . Two-particle magnetic parameters ( $A_c$ – $A_0$ ) were fixed to *ab initio* MCDF calculations, not fitted, and thus not listed here. Effective Coulomb-interaction operators  $E'_a$ ,  $E'_b$ , and effective three-particle electrostatic operators  $T1$  and  $T2$  are fit parameters for the given number of  $d$  electrons. For the fit parameters  $E'_b$  in  $\text{Sn}^{11+}$  and  $E'_a$  in  $\text{Sn}^{12+}$  were fixed (denoted by the superscript letter  $f$ ) on extrapolated values (see main text). Fits were performed for the interacting  $4d^k + 4d^{k-1}5s + 4d^{k-2}5s^2$  configurations,  $k = 8$ – $2$ , for  $\text{Sn}^{6+}$ – $\text{Sn}^{12+}$ , respectively. Energy parameters for the unknown  $4d^{k-1}5s + 4d^{k-2}5s^2$  configurations (not listed) were fixed at *ab initio* values for the average energies and spin-orbit interactions, but scaled by 0.85 for electrostatic and configuration-interaction parameters. Average energy  $E_{\text{av}}$  is defined such that the ground level energy of the  $4d^k + 4d^{k-1}5s + 4d^{k-2}5s^2$  configurations is equal to zero.  $\sigma$  is the root mean square of the fit uncertainty in  $\text{cm}^{-1}$  for the calculated configuration. “n/a” indicates nonapplicable parameters.

Parameter	$\text{Sn}^{6+}$		$\text{Sn}^{7+}$		$\text{Sn}^{8+}$		$\text{Sn}^{9+}$	
	FIT	FIT/MCDF	FIT	FIT/MCDF	FIT	FIT/MCDF	FIT	FIT/MCDF
$E_{\text{av}}$	16 279(4)	0.890	29 983(4)	0.890	42 688(4)	0.888	64 568(10)	0.872
$O2$	9 649(6)	0.856	9 979(5)	0.860	10 288(5)	0.863	10 592(9)	0.867
$O2'$	6 100(6)	0.795	6 326(4)	0.803	6 526(7)	0.808	6 702(11)	0.812
$E'_a$	243(6)		247(3)		255(3)		256(6)	
$E'_b$	22(5)		34(3)		37(5)		66(10)	
$\zeta(4d)$	3 688(4)	1.024	3 899(3)	1.021	4 119(3)	1.020	4 334(8)	1.018
$T1$	n/a	n/a	−5.2(0.2)		−5.6(0.1)		−5.9(0.3)	
$T2$	n/a	n/a	n/a	n/a	0.26(0.16)		−0.05(0.75)	
$\sigma$	10		14		16		41	
Parameter	$\text{Sn}^{10+}$		$\text{Sn}^{11+}$		$\text{Sn}^{12+}$			
	FIT	FIT/MCDF	FIT	FIT/MCDF	FIT	FIT/MCDF		
$E_{\text{av}}$	50 180(3)	0.91	37 553(9)	0.911	24 425(14)	0.93		
$O2$	10 878(3)	0.87	11 200(16)	0.876	11 480(43)	0.88		
$O2'$	6 896(6)	0.818	7 007(23)	0.815	7 102(43)	0.81		
$E'_a$	275(2)		248(15)		260 <sup>f</sup>			
$E'_b$	44(3)		50 <sup>f</sup>		n/a	n/a		
$\zeta(4d)$	4563(3)	1.017	4783(6)	1.013	5038(11)	1.016		
$T1$	−6.6(0.1)		−9.72(1.0)		n/a	n/a		
$T2$	0.37(0.16)		n/a	n/a	n/a	n/a		
$\sigma$	14		20		32			

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