

## Admissible memory kernels for random unitary qubit evolution

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We analyze the random unitary evolution of a qubit within the memory kernel approach. We provide sufficient conditions which guarantee that the corresponding memory kernel generates physically legitimate quantum evolution. Interestingly, we are able to recover several well-known examples and to generate new classes of nontrivial qubit evolution. Surprisingly, it turns out that a class of quantum evolutions with a memory kernel generated by our approach gives rise to the vanishing of a non-Markovianity measure based on the distinguishability of quantum states.

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### I. INTRODUCTION

The dynamics of open quantum systems plays an important role in the analysis of various phenomena like dissipation, decoherence, and dephasing [1,2]. The usual approach to the dynamics of an open quantum system consists of applying the Born-Markov approximation [1], which leads to a local master equation for the Markovian semigroup

$$\dot{\rho}_t = L[\rho_t], \quad (1)$$

where  $\rho_t$  is the density matrix of the investigated system and  $L$  is the time-independent generator of the dynamical semigroup defined as follows:

$$L[\rho] = -i[H, \rho] + \frac{1}{2} \sum_{\alpha} ([V_{\alpha}, \rho V_{\alpha}^{\dagger}] + [V_{\alpha} \rho, V_{\alpha}^{\dagger}]). \quad (2)$$

Here  $H$  denotes the effective system Hamiltonian, and  $V_{\alpha}$  represents noise operators [3,4]. We call Eq. (2) the Gorini-Kossakowski-Sudarshan-Lindblad (GKSL) form. The solution of Eq. (1) defines the Markovian semigroup

$$\rho_t = \Lambda_t[\rho] = e^{tL} \rho, \quad (3)$$

where  $\rho$  is an initial state. The dynamical map  $\Lambda_t = e^{tL}$  is completely positive and trace preserving (CPTP) [1,3–5]. The Born-Markov approximation assumes weak interaction and a separation of time scales between the system and its environment. Such an approach works perfectly well for many quantum optical systems [6–8]. When the above assumption is no longer valid the description based on Eq. (1) is not satisfactory. Recent technological progress and modern laboratory techniques call for a more refined approach which takes into account memory effects completely neglected in the description based on Markovian semigroups. In recent years we observed intense research activity in the field of non-Markovian quantum evolution (see the recent review [9], a collection of articles [10], and a recent comparative analysis [11]).

There are basically two approaches which generalize the standard Markovian master equation (1): a time-local approach replaces  $L$  by a time-dependent generator  $L_t$ . Interestingly, if for all  $t$  the time-dependent generator has the standard GKSL form (8), then  $\Lambda_t = \mathcal{T} \exp(\int_0^t L_u du)$  defines the so-called divisible dynamical map [12,13] which is often considered as the generalization of Markovianity (see [14] for a generalization of the notion of divisibility). The second approach is based on

the nonlocal Nakajima-Zwanzig equation [15] (see also [16]),

$$\dot{\rho}_t = \int_0^t K_{t-\tau} \rho_{\tau} d\tau, \quad (4)$$

in which quantum memory effects are taken into account through the introduction of a memory kernel  $K_t$ . This means that the rate of change of the state  $\rho_t$  at time  $t$  depends on its history (starting at  $t = 0$ ). The Markovian master equation (1) is reobtained when  $K_t = 2\delta(t)L$ . The time-dependent kernel is usually referred to as the generator of the non-Markovian master equation. Equation (4) applies to a variety of situations (see, e.g., [17]). Because of the convolution structure of Eq. (4) the time-local approach is often called a time-convolutionless approach [1,18,19]. The structure and the properties of Eq. (4) were carefully analyzed in [20–29]. In particular the generalization of the Markovian evolution to the so-called semi-Markov was investigated within the memory kernel approach by Budini [21] and Breuer and Vacchini [23] (see also discussion in [28]).

In the present article we study random unitary evolution of a qubit within the memory kernel approach. In particular we address the following problem: What is the structure of the corresponding memory kernel  $K_t$  which leads to the legitimate CPTP dynamical map  $\Lambda_t$ ? The article has the following structure: In Sec. II we recall basic facts about random unitary evolutions and in Sec. III we formulate a sufficient condition for  $K_t$  to guarantee legitimate physical evolutions. In Sec. IV we examine the issue of Markovianity. Surprisingly, it turns out that a subclass of quantum evolutions with memory kernel generated by our approach gives rise to the vanishing of a non-Markovianity measure based on the distinguishability of quantum states [30]. Section V illustrates our approach with several examples. Final conclusions are collected in Sec. VI.

### II. RANDOM UNITARY QUBIT EVOLUTION

A quantum channel  $\mathcal{E} : \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{H})$  is called random unitary [31] if its Kraus representation is given by

$$\mathcal{E}[X] = \sum_k p_k U_k X U_k^{\dagger}, \quad (5)$$

where  $U_k$  is a collection of unitary operators and  $\{p_k\}$  stands for a probability distribution. In this article we consider a random

unitary dynamical map  $\Lambda_t$  defined by

$$\Lambda_t[\rho] = \sum_{\alpha=0}^3 p_\alpha(t) \sigma_\alpha \rho \sigma_\alpha, \quad (6)$$

where  $\sigma_\alpha$  are Pauli matrices with  $\sigma_0 = \mathbb{I}_2$  [32]. The initial condition  $\Lambda_{t=0} = \mathbb{I}$  implies  $p_\alpha(0) = \delta_{\alpha 0}$ . Recently a time-local description based on the following master equation was analyzed [33,34]:

$$\dot{\Lambda}_t = L_t \Lambda_t, \quad (7)$$

where  $L_t$  is a time-local generator defined by

$$L_t[\rho] = \sum_{k=1}^3 \gamma_k(t) (\sigma_k \rho \sigma_k - \rho), \quad (8)$$

with time-dependent decoherence rates  $\gamma_k(t)$ . One asks the following question: What are the conditions for  $\gamma_k(t)$  which guarantee that the solution  $\Lambda_t = \exp(\int_0^t L_\tau d\tau)$  provides a legitimate dynamical map? Note that the solution defines a random unitary evolution with  $p_\alpha(t)$  given by

$$p_\alpha(t) = \frac{1}{4} \sum_{\beta=0}^3 H_{\alpha\beta} \lambda_\beta(t), \quad (9)$$

where  $H_{\alpha\beta}$  is the Hadamard matrix

$$H = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & 1 & -1 \\ 1 & -1 & -1 & 1 \end{pmatrix}, \quad (10)$$

and  $\lambda_\beta(t)$  are time-dependent eigenvalues of  $\Lambda_t$ ,

$$\Lambda_t[\sigma_\alpha] = \lambda_\alpha(t) \sigma_\alpha, \quad (11)$$

which read as follows:  $\Lambda_0(t) = 1$  and

$$\begin{aligned} \lambda_1(t) &= \exp(-2[\Gamma_2(t) + \Gamma_3(t)]), \\ \lambda_2(t) &= \exp(-2[\Gamma_1(t) + \Gamma_3(t)]), \\ \lambda_3(t) &= \exp(-2[\Gamma_1(t) + \Gamma_2(t)]), \end{aligned} \quad (12)$$

with  $\Gamma_k(t) = \int_0^t \gamma_k(\tau) d\tau$ . Now, the map (6) is completely positive (CP) iff  $p_\alpha(t) \geq 0$ , which is equivalent to the following set of conditions for  $\lambda_s$  [33,34]:

$$1 + \lambda_1(t) + \lambda_2(t) + \lambda_3(t) \geq 0, \quad (13)$$

and

$$\begin{aligned} \lambda_1(t) + \lambda_2(t) &\leq 1 + \lambda_3(t), \\ \lambda_3(t) + \lambda_1(t) &\leq 1 + \lambda_2(t), \\ \lambda_2(t) + \lambda_3(t) &\leq 1 + \lambda_1(t). \end{aligned} \quad (14)$$

### III. CONSTRUCTION OF LEGITIMATE MEMORY KERNELS

In this article we analyze the nonlocal description based on the following memory kernel equation:

$$\dot{\Lambda}_t = \int_0^t K_{t-\tau} \Lambda_\tau d\tau, \quad (15)$$

with

$$K_t[\rho] = \sum_{i=1}^3 k_i(t) (\sigma_i \rho \sigma_i - \rho), \quad (16)$$

where  $k_i(t)$  ( $i = 1, 2, 3$ ) represent nontrivial memory effects. Note that Eq. (15) considerably simplifies after performing the Laplace transform

$$\tilde{\Lambda}_s = \frac{1}{s - \tilde{K}_s}, \quad (17)$$

where  $\tilde{\Lambda}_s := \int_0^\infty e^{-st} \Lambda_t dt$  and similarly for  $\tilde{K}_s$ . The question we address is this: What are the conditions for  $k_i(t)$  which guarantee that the solution  $\Lambda_t$  provides a legitimate dynamical map?

Denoting by  $\kappa_\alpha(t)$  the eigenvalues of  $K_t$ ,

$$K_t[\sigma_\alpha] = \kappa_\alpha(t) \sigma_\alpha, \quad (18)$$

Eq. (15) gives rise to the following set of equations:

$$\dot{\lambda}_i(t) = \int_0^t \kappa_i(t - \tau) \lambda_i(\tau) d\tau, \quad i = 1, 2, 3. \quad (19)$$

Note that  $\kappa_0(t) = 0$  and hence  $\lambda_0(t) = 1 = \text{const}$ . In terms of the Laplace transforms  $\tilde{\lambda}_i(s)$  and  $\tilde{\kappa}_i(s)$  one finds

$$\tilde{\lambda}_i(s) = \frac{1}{s - \tilde{\kappa}_i(s)}. \quad (20)$$

In terms of  $\tilde{\lambda}_i(s)$  conditions (13) and (14) may be equivalently reformulated as follows:

$$\frac{1}{s} + \tilde{\lambda}_1(s) + \tilde{\lambda}_1(s) + \tilde{\lambda}_2(s) \text{ is CM}, \quad (21)$$

and

$$\begin{aligned} \frac{1}{s} + \tilde{\lambda}_3(s) - \tilde{\lambda}_1(s) - \tilde{\lambda}_2(s) &\text{ is CM}, \\ \frac{1}{s} + \tilde{\lambda}_2(s) - \tilde{\lambda}_1(s) - \tilde{\lambda}_3(s) &\text{ is CM}, \\ \frac{1}{s} + \tilde{\lambda}_1(s) - \tilde{\lambda}_3(s) - \tilde{\lambda}_2(s) &\text{ is CM}, \end{aligned} \quad (22)$$

where CM stands for a completely monotone function [35], i.e., a smooth function  $f : [0, \infty) \rightarrow \mathbb{R}$  satisfying the condition

$$(-1)^n \frac{d^n}{ds^n} f(s) \geq 0, \quad s \geq 0, \quad n = 0, 1, 2, \dots \quad (23)$$

The equivalence of conditions (14) and (22) results from the following.

*Theorem 1: Bernstein's theorem.* A function  $f : [0, \infty) \rightarrow \mathbb{R}$  is completely monotone on  $[0, \infty)$  if and only if it is a Laplace transform of a finite non-negative Borel measure  $\mu$  on  $[0, \infty)$ ; i.e.,  $f$  is of the form

$$f(s) = \int_0^\infty e^{-st} d\mu(t). \quad (24)$$

Note that the initial condition  $p_0(0) = 1$  and  $p_k(0) = 0$  for  $k = 1, 2, 3$  is equivalent to  $\lambda_k(0) = 1$  due to the following theorem.

*Theorem 2: Initial value theorem.* Let  $\tilde{f}(s)$  be the Laplace transform of  $f(t)$ . Then the following relation is true:

$$\lim_{t \rightarrow 0} f(t) = \lim_{s \rightarrow \infty} s \tilde{f}(s). \quad (25)$$

It is equivalent to

$$\lim_{s \rightarrow \infty} s \tilde{\lambda}_k(s) = 1, \quad (26)$$

for  $k = 1, 2, 3$ . This way we have proved the following theorem.

*Theorem 3.* The map  $\tilde{\Lambda}_s$  represented by the spectral decomposition

$$\tilde{\Lambda}_s[\rho] = \frac{1}{2} \sum_{\alpha=0}^3 \tilde{\lambda}_\alpha(s) \sigma_\alpha \text{tr}[\sigma_\alpha \rho], \quad (27)$$

with  $\tilde{\lambda}_0(s) = 1/s$ , defines the Laplace transform of a legitimate map  $\Lambda_t$  if and only if conditions (21), (22), and (26) are satisfied.

It is worth emphasizing that there are few analytical tools for dealing with CM functions, which is due to the fact that an infinite set of conditions (23) must be verified. Nevertheless, we found an important class of CM functions giving rise to CPTP dynamics with a straightforward interpretation. To present them, let us first observe that CM functions have the following two properties, which will not be proved.

*Property 1.* Let  $f$  and  $g$  be arbitrary completely monotone functions. Then

- (1)  $f \cdot g$  is CM,
- (2)  $\alpha f + \beta g$  is CM for any  $\alpha, \beta > 0$ .

*Property 2.* If  $s_0 \geq 0$  then  $\frac{1}{s+s_0}$  is CM.

We are now ready to prove our main result.

*Theorem 4.* Let  $W(s)$  be a function such that  $\frac{1}{s} \frac{1}{W(s)}$  is CM. Then the functions

$$\tilde{\kappa}_k(s) = -\frac{s}{a_k W(s) - 1}, \quad k = 1, 2, 3, \quad (28)$$

with  $a_1, a_2, a_3 > 0$  such that

$$\frac{1}{s} \left( 4 - \frac{1}{W(s)} \left[ \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \right] \right) \text{ is CM,} \quad (29)$$

and

$$\begin{aligned} \frac{1}{a_1} + \frac{1}{a_2} &\geq \frac{1}{a_3}, \\ \frac{1}{a_2} + \frac{1}{a_3} &\geq \frac{1}{a_1}, \\ \frac{1}{a_3} + \frac{1}{a_1} &\geq \frac{1}{a_2} \end{aligned} \quad (30)$$

define a legitimate memory kernel

$$\tilde{K}_s[\rho] = \frac{1}{2} \sum_{k=1}^3 \tilde{\kappa}_k(s) \sigma_k \text{tr}[\sigma_k \rho], \quad (31)$$

i.e., the corresponding  $\tilde{\lambda}_k(s)$  satisfy conditions (21), (22), and (26).

*Proof.* Note that formula (28) implies

$$\tilde{\lambda}_k(s) = \frac{1}{s} \left( 1 - \frac{1}{a_k W(s)} \right), \quad (32)$$

and hence

$$\frac{1}{s} + \tilde{\lambda}_3(s) - \tilde{\lambda}_1(s) - \tilde{\lambda}_2(s) = \frac{1}{s} \frac{1}{W(s)} \left( \frac{1}{a_1} + \frac{1}{a_2} - \frac{1}{a_3} \right), \quad (33)$$

which proves that  $\frac{1}{s} + \tilde{\lambda}_3(s) - \tilde{\lambda}_1(s) - \tilde{\lambda}_2(s)$  is CM due to the fact that  $\frac{1}{s} \frac{1}{W(s)}$  is CM. Similarly one proves the remaining conditions (14). ■

Note that, since  $\frac{1}{s} \frac{1}{W(s)}$  is CM, due to the Bernstein theorem, it is the Laplace transform of a positive function. Hence,

$$W(s) = \frac{1}{\tilde{f}(s)}, \quad (34)$$

where  $\tilde{f}(s)$  is the Laplace transform of  $f(t)$  satisfying  $\int_0^t f(\tau) d\tau \geq 0$  for all  $t \geq 0$ . One finds

$$\tilde{\kappa}_k(s) = \frac{-s \tilde{f}(s)}{a_k - \tilde{f}(s)}. \quad (35)$$

Note that condition (29) implies

$$\left( \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \right) \int_0^t f(\tau) d\tau \leq 4. \quad (36)$$

Hence, to summarize, our class is characterized by a single function  $f(t)$  and three numbers  $a_1, a_2, a_3 > 0$  such that  $F(t) = \int_0^t f(\tau) d\tau \geq 0$  and conditions (30) and (36) hold. One finds for  $p_\alpha(t)$ :

$$\begin{aligned} p_1(t) &= \frac{1}{4} \left( \frac{1}{a_2} + \frac{1}{a_3} - \frac{1}{a_1} \right) F(t), \\ p_2(t) &= \frac{1}{4} \left( \frac{1}{a_3} + \frac{1}{a_1} - \frac{1}{a_2} \right) F(t), \\ p_3(t) &= \frac{1}{4} \left( \frac{1}{a_1} + \frac{1}{a_2} - \frac{1}{a_3} \right) F(t), \end{aligned} \quad (37)$$

and  $p_0(t) = 1 - p_1(t) - p_2(t) - p_3(t)$ . In particular, taking  $a_1 = a_2 = a$  and  $a_3 = \infty$  one finds

$$\tilde{\kappa}_1(s) = \tilde{\kappa}_2(s) = \frac{-s \tilde{f}(s)}{a - \tilde{f}(s)}, \quad \tilde{\kappa}_3(s) = 0, \quad (38)$$

and hence

$$\tilde{\kappa}_1(s) = \tilde{\kappa}_2(s) = 0, \quad \tilde{\kappa}_3(s) = \frac{1}{2} \frac{s \tilde{f}(s)}{a - \tilde{f}(s)} \quad (39)$$

gives rise to the legitimate memory kernel

$$K_t[\rho] = k_3(t)(\sigma_3 \rho \sigma_3 - \rho), \quad (40)$$

with arbitrary  $f(t)$  and  $a > 0$  satisfying the additional condition

$$0 \leq F(t) := \int_0^t f(\tau) d\tau \leq 2a, \quad (41)$$

for all  $t \geq 0$ . The corresponding solution reads

$$\begin{aligned} p_0(t) &= 1 - \frac{1}{2a} F(t), \\ p_1(t) &= p_2(t) = 0, \\ p_3(t) &= \frac{1}{2a} F(t). \end{aligned} \quad (42)$$

This approach resembles very much the semi-Markov construction [23,28]: For any  $f(t) \geq 0$  satisfying  $\int_0^\infty f(t) dt \leq 1$  the memory kernel (40) with

$$\tilde{\kappa}_3(s) = \frac{s \tilde{f}(s)}{1 - \tilde{f}(s)} \quad (43)$$

gives rise to CPTP evolution. In this case one finds

$$\begin{aligned} p_0(t) &= \frac{1}{2}[1 + \lambda_1(t)], \\ p_1(t) &= p_2(t) = 0, \\ p_3(t) &= \frac{1}{2}[1 - \lambda_1(t)], \end{aligned} \quad (44)$$

where

$$\tilde{\lambda}_1(s) = \tilde{\lambda}_2(s) = \frac{\tilde{f}(s) + 1}{\tilde{f}(s) - 1}. \quad (45)$$

It is therefore clear that our approach goes beyond the semi-Markov construction.

Let us recall that Markovian semigroup generated by

$$L[\rho] = \frac{1}{2} \sum_{k=1}^3 \gamma_k [\sigma_k \rho \sigma_k - \rho]. \quad (46)$$

The corresponding Bloch equation reads

$$\dot{x}_k(t) = -\frac{2}{T_k} x_k(t), \quad (47)$$

where  $x_k := \text{tr}[\rho \sigma_k]$  and the relaxation times are defined via

$$T_1 = \frac{1}{\gamma_2 + \gamma_3}, \quad T_2 = \frac{1}{\gamma_3 + \gamma_1}, \quad T_3 = \frac{1}{\gamma_1 + \gamma_2}. \quad (48)$$

It is well known [5] that complete positivity is equivalent to the following set of conditions upon  $T_k$ :

$$\begin{aligned} \frac{1}{T_1} + \frac{1}{T_2} &\geq \frac{1}{T_3}, \\ \frac{1}{T_2} + \frac{1}{T_3} &\geq \frac{1}{T_1}, \\ \frac{1}{T_3} + \frac{1}{T_1} &\geq \frac{1}{T_2}. \end{aligned} \quad (49)$$

It is therefore clear that condition (30) is an analog of condition (49). Note that condition (30) means that there exist  $b_1, b_2, b_3 > 0$  such that

$$\begin{aligned} \frac{1}{2} \frac{1}{a_1} &= \frac{1}{b_2} + \frac{1}{b_3}, \\ \frac{1}{2} \frac{1}{a_2} &= \frac{1}{b_3} + \frac{1}{b_1}, \\ \frac{1}{2} \frac{1}{a_3} &= \frac{1}{b_1} + \frac{1}{b_2}. \end{aligned} \quad (50)$$

Now, in terms of  $b_1, b_2, b_3$  our result may be reformulated as follows.

*Corollary 1.* For any  $b_1, b_2, b_3 > 0$  and the function  $f(t)$  satisfying

$$0 \leq F(t) := \int_0^t f(\tau) d\tau \leq \left( \frac{1}{b_1} + \frac{1}{b_2} + \frac{1}{b_3} \right)^{-1}, \quad (51)$$

and

$$\lim_{s \rightarrow \infty} \tilde{f}(s) = 0, \quad (52)$$

the memory kernel defined by

$$\tilde{\kappa}_k(s) = -\frac{s \tilde{f}(s)}{a_k - \tilde{f}(s)} \quad (53)$$

defines legitimate quantum evolution. Moreover, one has

$$p_k(t) = \frac{1}{b_k} F(t), \quad (54)$$

and  $p_0(1) = 1 - p_1(t) - p_2(t) - p_3(t)$ .

Let us observe that it is very hard, in general, to invert formula (35) to the time domain. Now, we provide a family of  $W(s)$  which enables one to easily compute  $\kappa_i(t)$  and have the memory kernel in time domain.

*Theorem 5.* Let  $W(s)$  be a polynomial

$$W(s) = (s + z_1) \cdots (s + z_n), \quad (55)$$

with  $z_i > 0$ . If  $a_1, a_2, a_3$  satisfy (30) and

$$\prod_{i=1}^n z_i \geq \frac{1}{4} \left( \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \right), \quad (56)$$

then  $\kappa_i(t)$  defined via Eq. (28) define a legitimate memory kernel.

*Proof.* It is clear that it is enough to prove condition (21).

*Lemma 1.* One has the following decomposition:

$$\frac{1}{s \prod_{i=1}^n (s + z_i)} = A \left( \frac{1}{s} - \sum_{i=1}^n \frac{\prod_{j=1}^{i-1} z_j}{\prod_{j=1}^i (s + z_j)} \right), \quad (57)$$

where

$$A = \frac{1}{\prod_{i=1}^n z_i}. \quad (58)$$

For the proof see the Appendix. Now we show that condition (21) holds. According to Eq. (57) one has

$$\begin{aligned} \frac{1}{s} + \tilde{\lambda}_1(s) + \tilde{\lambda}_2(s) + \tilde{\lambda}_3(s) &= \frac{1}{s} \left( 4 - \left[ \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \right] \frac{1}{W(s)} \right) \\ &= \frac{1}{s} \left( 4 - \left[ \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \right] \frac{1}{\prod_{i=1}^n z_i} \right) \\ &\quad + \left[ \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \right] \frac{1}{\prod_{i=1}^n z_i} \sum_{j=1}^n \frac{\prod_{i=1}^{j-1} z_i}{\prod_{i=1}^j (s + z_i)}. \end{aligned} \quad (59)$$

A second term in Eq. (59) is CM due to the fact that it is a sum of CM functions. Hence, if condition (56) is satisfied then condition (21) holds. ■

Note that

$$\tilde{\kappa}_i(s) = -\frac{s}{a_k W(s) - 1} = -\frac{1}{a_k} \frac{s}{(s - s_1) \cdots (s - s_m)}, \quad (60)$$

where  $\{s_1, \dots, s_m\}$  are the roots of the polynomial  $a_k W(s) - 1$ . It is therefore clear that formula (60) may be easily inverted to the time domain.

*Remark 1.* Note that  $W(s)$  defined in Eq. (55) implies that  $\frac{1}{W(s)}$  is CM and hence  $\frac{1}{s} \frac{1}{W(s)}$  is CM as well.

#### IV. CHECKING FOR NON-MARKOVIANITY

Let us recall that according to [30] the evolution represented by  $\Lambda_t$  is non-Markovian if the condition

$$\frac{d}{dt} \|\Lambda_t[\rho_1 - \rho_2]\|_{\text{tr}} \leq 0 \quad (61)$$

is violated for some initial states  $\rho_1$  and  $\rho_2$ . One defines [30] a well-known non-Markovianity measure

$$\mathcal{N}_{\text{BLP}}[\Lambda_t] = \sup_{\rho_1, \rho_2} \int \frac{d}{dt} \|\Lambda_t[\rho_1 - \rho_2]\|_{\text{tr}} dt, \quad (62)$$

where the integral is evaluated over the region where  $\frac{d}{dt} \|\Lambda_t[\rho_1 - \rho_2]\|_{\text{tr}} > 0$  (and BLP stands for Breuer-Laine-Piilo). Now, it has been proved [33] that for random unitary qubit evolution if all eigenvalues  $\lambda_k(t) \geq 0$ , then condition (61) is equivalent to

$$\frac{d}{dt} \lambda_k(t) \leq 0, \quad k = 1, 2, 3. \quad (63)$$

*Proposition 1.* For  $a_1, a_2, a_3$  satisfying conditions (30) and  $W(s) = \frac{1}{\tilde{f}(s)}$ , where  $\tilde{f}(s)$  is CM and

$$\int_0^t f(\tau) d\tau \leq a_{\min}, \quad a_{\min} = \min\{a_1, a_2, a_3\}, \quad (64)$$

the corresponding memory kernel gives rise to the dynamical map  $\Lambda_t$  such that  $\mathcal{N}_{\text{BLP}}[\Lambda_t] = 0$ .

*Proof.* Let us observe that condition (64) implies condition (29). Indeed, from condition (64) one has

$$\left( \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \right) \int_0^t f(\tau) d\tau \leq 3, \quad (65)$$

and hence condition (29) follows. Now, observe that

$$\lambda_k(t) = 1 - \frac{1}{a_k} \int_0^t f(\tau) d\tau \geq 0,$$

due to condition (64). Hence, it is sufficient to show that  $\frac{d}{dt} \lambda_k(t) \leq 0$ . It is clear  $\frac{d}{dt} \lambda_k(t) \leq 0$  if and only if  $1 - s\tilde{\lambda}_k(s)$  is CM and hence taking into account Eq. (20) it is equivalent to the requirement that  $-\tilde{\kappa}_k(s)\tilde{\lambda}_k(s)$  is CM. One has, therefore,

$$-\tilde{\kappa}_k(s)\tilde{\lambda}_k(s) = \frac{\tilde{f}(s)}{a_k}, \quad (66)$$

which ends the proof since  $\tilde{f}(s)$  is CM and  $a_k > 0$ . ■

*Remark 2.* If  $W(s) = (s + z_1) \cdots (s + z_n)$  with  $z_k > 0$  and  $a_1, a_2, a_3$  satisfying conditions (30) together with

$$\prod_{i=1}^n z_i \geq \frac{1}{a_k}, \quad k = 1, 2, 3, \quad (67)$$

then the corresponding dynamical map  $\Lambda_t$  satisfies  $\mathcal{N}_{\text{BLP}}[\Lambda_t] = 0$ .

*Remark 3.* It was shown [14,36] that BLP condition (61) is equivalent to so-called  $P$  divisibility. This means that

$$\Lambda_t = V_{t,s} \Lambda_s, \quad (68)$$

and for any  $t > s$  the propagator  $V_{t,s}$  is positive (but not necessarily completely positive).

Interestingly, our construction provides a class of legitimate random unitary qubit evolutions generated by the nontrivial memory kernel but still satisfying BLP condition (61) (cf. also [37]). It is clear that to violate condition (61) one needs a more refined construction such that  $\frac{1}{W(s)}$  is not CM but  $\frac{1}{s} \frac{1}{W(s)}$  is already CM. It deserves further analysis.

Consider now the question of CP divisibility which is fully controlled by the local decoherence rates in Eq. (8). One may

easily compute them in terms of  $f(t)$ :

$$\begin{aligned} \gamma_1(t) &= \frac{f(t)}{4} \left( \frac{-1}{a_1 - F(t)} + \frac{1}{a_2 - F(t)} + \frac{1}{a_3 - F(t)} \right), \\ \gamma_2(t) &= \frac{f(t)}{4} \left( \frac{1}{a_1 - F(t)} - \frac{1}{a_2 - F(t)} + \frac{1}{a_3 - F(t)} \right), \\ \gamma_3(t) &= \frac{f(t)}{4} \left( \frac{1}{a_1 - F(t)} + \frac{1}{a_2 - F(t)} - \frac{1}{a_3 - F(t)} \right). \end{aligned}$$

The dynamical map  $\Lambda_t$  is CP divisible iff  $\gamma_k(t) \geq 0$  for  $k = 1, 2, 3$ . Let us assume that

$$a_1 \leq a_2 \leq a_3. \quad (69)$$

*Proposition 2.* If  $a_1, a_2, a_3$  and  $f(t) \geq 0$  satisfy conditions (29) and (30) the corresponding memory kernel,

$$\tilde{\kappa}_k(s) = -\frac{s\tilde{f}(s)}{a_k - \tilde{f}(s)}, \quad (70)$$

leads to a CP-divisible dynamical map iff

$$F(t) \leq a_1 - \sqrt{(a_2 - a_1)(a_3 - a_1)}. \quad (71)$$

*Proof.* Due to condition (69) it is sufficient to show that  $\gamma_1(t) \geq 0$  which, for  $f(t) \geq 0$ , is equivalent to

$$\frac{-1}{a_1 - F(t)} + \frac{1}{a_2 - F(t)} + \frac{1}{a_3 - F(t)} \geq 0. \quad (72)$$

Let us assume that  $F(t) < a_1$ , which means that  $\gamma_1(t)$  is not singular. Inequality (72) is satisfied iff

$$F(t) \in (-\infty, F_-] \cup [F_+, +\infty)$$

with

$$F_{\pm} = a_1 \pm \sqrt{(a_2 - a_1)(a_3 - a_1)}.$$

Now, taking into account that  $F(t) < a_1$ , one finally proves condition (71). ■

Proposition 2 shows that positivity of the function  $f(t)$  is not sufficient for CP divisibility. One needs an extra condition (71) which involves not only  $f(t)$  but  $\{a_1, a_2, a_3\}$  as well.

## V. EXAMPLES

*Example 1.* Consider the simplest case with a polynomial of degree 1,

$$W(s) = s + z, \quad (73)$$

with  $z > 0$ . One finds

$$\tilde{\kappa}_k(s) = -\frac{s}{a_k(s + z) - 1}, \quad (74)$$

and the inverse Laplace transform gives

$$\kappa_k(t) = -\frac{1}{z} \left( \delta(t) - \left[ z - \frac{1}{a_k} \right] e^{-[z - \frac{1}{a_k}]t} \right). \quad (75)$$

Note that if  $a_k = 1/z$ , then the dynamics is purely local. One easily finds

$$\lambda_k(t) = 1 - \frac{1}{za_k} (1 - e^{-zt}), \quad (76)$$

and finally the solution for  $p_k(t)$  is defined by Eqs. (37) with

$$F(t) = \frac{1}{z}(1 - e^{-zt}). \quad (77)$$

Note that condition (56) implies the following relation between  $z$  and  $a_1, a_2, a_3$ :

$$4z \geq \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3}, \quad (78)$$

which guarantees that  $p_0(t) \geq 0$ . In the symmetric case  $a_1 = a_2 = a_3 = a$  one finds  $p_1(t) = p_2(t) = p_3(t) =: p(t)$  with

$$p(t) = \frac{1}{4za}[1 - e^{-zt}], \quad (79)$$

and  $p_0(t) = 1 - 3p(t)$  with  $4za \geq 3$ . One finds that asymptotically

$$p_0(t) \rightarrow 1 - \frac{3}{4za}. \quad (80)$$

Note that for  $za > 1$  one has asymptotically  $p_0(\infty) < 1/4$ . This property cannot be reproduced within the local approach with regular generators  $L_t$ . Indeed, it follows from Eq. (9) (see also [33] for more details) that

$$p_0(t) = \frac{1}{4}[1 + \lambda_1(t) + \lambda_2(t) + \lambda_3(t)], \quad (81)$$

and, hence, using conditions (14), one finds

$$p_0(t) \geq \frac{1}{4}. \quad (82)$$

This example shows that local and memory kernel approaches may lead to essentially different evolutions.

*Example 2.* Consider now the same polynomial  $W(s) = s + z$  but let  $z = 2c > 0$ . Moreover,

$$a_1 = a_2 = \frac{1}{c}, \quad a_3 = \frac{1}{2c}. \quad (83)$$

One finds

$$\tilde{\kappa}_1(s) = \tilde{\kappa}_2(s) = -\frac{sc}{s+c}, \quad \tilde{\kappa}_3(s) = -2c,$$

and hence

$$\kappa_1(t) = \kappa_2(t) = -c\delta(t) + c^2e^{-ct}, \quad \kappa_3(t) = -2c\delta(t).$$

Finally, one finds the following formula for the memory kernel:

$$K_t[\rho] = \frac{c}{2}\delta(t)[\sigma_1\rho\sigma_1 + \sigma_2\rho\sigma_2 - 2\rho] - \frac{c^2}{2}e^{-ct}[\sigma_3\rho\sigma_3 - \rho]. \quad (84)$$

One has

$$\lambda_1(t) = \lambda_2(t) = \frac{1}{2}(1 + e^{-2ct}), \quad \lambda_3(t) = e^{-2ct}.$$

Interestingly, this evolution reproduces the time-local description with

$$\gamma_1(t) = \gamma_2(t) = \frac{c}{2}, \quad \gamma_3(t) = -\frac{c}{2} \tanh(ct), \quad (85)$$

as discussed in [34]. It was shown [36] that  $\Lambda_t$  is a convex combination of two Markovian semigroups  $\Lambda_t^{(1)}$  and  $\Lambda_t^{(2)}$

generated by

$$L_k[\rho] = \frac{c}{2}[\sigma_k\rho\sigma_k - \rho], \quad k = 1, 2, \quad (86)$$

that is,

$$\Lambda_t = \frac{1}{2}(e^{tL_1} + e^{tL_2}). \quad (87)$$

This simple example shows that a convex combination of Markovian semigroups leads to a quantum evolution displaying essential memory effects.

*Example 3.* Consider now a polynomial of degree 2,

$$W(s) = (s + c_1)(s + c_2), \quad (88)$$

with  $c_2 > c_1 > 0$ . Our construction gives rise to a legitimate memory kernel if conditions (30) hold and

$$4c_1c_2 \geq \frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3}. \quad (89)$$

One finds

$$\begin{aligned} \tilde{\kappa}_k(s) &= -\frac{1}{a_k} \frac{s}{(s+c_1)(s+c_2) - \frac{1}{a_k}} \\ &= -\frac{1}{a_k} \frac{s}{(s+s_1)(s+s_2)}, \end{aligned} \quad (90)$$

with

$$s_1 + s_2 = c_1 + c_2, \quad s_1s_2 = c_1c_2 - \frac{1}{a_k}.$$

Hence, the solution has the form (37) with the function  $F(t)$  given by

$$F(t) = \frac{1}{c_2 - c_1} \left( \frac{1}{c_1}[1 - e^{-c_1t}] - \frac{1}{c_2}[1 - e^{-c_2t}] \right). \quad (91)$$

*Example 4.* Let

$$W(s) = s^2 + \omega^2. \quad (92)$$

Note that  $\frac{1}{s} \frac{1}{W(s)}$  is CM since

$$\frac{1}{s} \frac{1}{W(s)} = \frac{1}{\omega} \frac{1}{s} \left( \frac{\omega}{s^2 + \omega^2} \right)$$

is the Laplace transform of  $\int_0^t \sin(\omega\tau) d\tau$  which is positive for all  $t \geq 0$ . Condition (29) implies

$$\frac{1}{a_1} + \frac{1}{a_2} + \frac{1}{a_3} \leq 2\omega^2. \quad (93)$$

The corresponding eigenvalues of the memory kernel read

$$\kappa_i(t) = -\frac{1}{a_i} \cos \left( \sqrt{\omega^2 - \frac{1}{a_i}} t \right), \quad (94)$$

for  $\omega^2 \geq 1/a_i$ , and

$$\kappa_i(t) = -\frac{1}{a_i} \cosh \left( \sqrt{\frac{1}{a_i} - \omega^2} t \right), \quad (95)$$

for  $\omega^2 < 1/a_i$ . Moreover, one finds

$$\lambda_k(t) = 1 + \frac{1}{a_k \omega^2} [\cos(\omega t) - 1], \quad (96)$$

and hence

$$\begin{aligned} p_1(t) &= \frac{1}{4\omega^2} \left( \frac{1}{a_2} + \frac{1}{a_3} - \frac{1}{a_1} \right) [1 - \cos(\omega t)], \\ p_2(t) &= \frac{1}{4\omega^2} \left( \frac{1}{a_3} + \frac{1}{a_1} - \frac{1}{a_2} \right) [1 - \cos(\omega t)], \\ p_3(t) &= \frac{1}{4\omega^2} \left( \frac{1}{a_1} + \frac{1}{a_2} - \frac{1}{a_3} \right) [1 - \cos(\omega t)], \end{aligned} \quad (97)$$

together with  $p_0(t) = 1 - p_1(t) - p_2(t) - p_3(t)$ . In particular taking

$$a_1 = a_2 = \frac{1}{\omega^2}, \quad a_3 = \infty, \quad (98)$$

one finds

$$\kappa_1(t) = \kappa_2(t) = -\omega^2, \quad \kappa_3(t) = 0, \quad (99)$$

and hence

$$k_1(t) = k_2(t) = 0, \quad \kappa_3(t) = \frac{\omega^2}{2}, \quad (100)$$

which proves that the constant (time-independent)

$$K_t[\rho] = \frac{k}{2} (\sigma_3 \rho \sigma_3 - \rho) \quad (101)$$

provides a legitimate memory kernel for arbitrary  $k = \omega^2 > 0$ . Moreover, one finds for the local decoherence rates

$$\begin{aligned} \gamma_1(t) &= \frac{\omega \sin(\omega t)}{4} \left( \frac{-1}{a_1 \omega^2 - 1 + \cos(\omega t)} \right. \\ &\quad \left. + \frac{1}{a_2 \omega^2 - 1 + \cos(\omega t)} + \frac{1}{a_3 \omega^2 - 1 + \cos(\omega t)} \right), \end{aligned}$$

and similarly for  $\gamma_2(t)$  and  $\gamma_3(t)$ . Note that if for some  $k$  one has  $a_k \omega^2 < 1$  then local decoherence rates are singular and hence in this case the nonlocal approach is more suitable.

## VI. CONCLUSIONS

We analyzed random unitary evolution of a qubit within a memory kernel approach. Our main result formulated in Theorem 4 allows one to construct legitimate memory kernels leading to CPTP dynamical maps. The power of this method is based on the fact that (1) it allows one to reconstruct well-known examples of legitimate qubit evolution and (2) the structure of polynomials  $W_k(s)$  enables one to perform the inverse Laplace transform and to find a formula for the kernel in the time domain. The mathematical analysis heavily uses the notion of completely monotone functions. These functions are not commonly used in theoretical physics and knowledge of their properties is rather limited. There are no known effective methods allowing one to check whether a given function is CM. We stress that Theorem 4 provides only a sufficient condition, and further analysis is needed to cover physically interesting cases which do not fit the assumptions of the theorem. Interestingly, it turns out that the quantum evolution with a memory kernel generated by our approach gives rise to a vanishing non-Markovianity measure based on the distinguishability of quantum states [30]. We also have shown when the corresponding dynamical map is CP divisible. It shows that the evolution satisfying the nonlocal master equation does not necessarily lead to a non-Markovian evolution. It would be also interesting to analyze the relation between semi-Markov evolution and the one governed by our approach in more detail.

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## APPENDIX: PROOF OF LEMMA 1

Let us observe that Eq. (57) may be represented in the following form:

$$\frac{1}{s \prod_{i=1}^n (s + z_i)} = A \frac{\prod_{i=1}^n (s + z_i) - s \left( \prod_{i=2}^n (s + z_i) + z_1 \prod_{i=3}^n (s + z_i) + \cdots + \prod_{j=1}^{n-2} z_j (s + z_n) + \prod_{j=1}^{n-1} z_j \right)}{s \prod_{i=1}^n (s + z_i)}, \quad (A1)$$

therefore, to prove Lemma 1 it suffices to show that

$$\prod_{i=1}^n z_i = \prod_{i=1}^n (s + z_i) - s \left( \prod_{i=2}^n (s + z_i) + z_1 \prod_{i=3}^n (s + z_i) + \cdots + \prod_{j=1}^{n-2} z_j (s + z_n) + \prod_{j=1}^{n-1} z_j \right). \quad (A2)$$

We prove this by induction. For  $n = 1$  it is clear that LHS = RHS =  $z_1$ . We assume that Eq. (A2) is true for  $n$  and prove it is also true for  $(n + 1)$ . The left-hand side (LHS) may be written as

$$\text{LHS} = \prod_{i=1}^n z_i z_{n+1}, \quad (A3)$$

while the right-hand side (RHS) reads

$$\begin{aligned}
 \text{RHS} &= \prod_{i=1}^n (s + z_i)(s + z_{n+1}) - s \left( \prod_{i=2}^n (s + z_i)(s + z_{n+1}) + z_1 \prod_{i=3}^n (s + z_i)(s + z_{n+1}) + \dots \right. \\
 &\quad \left. + \prod_{j=1}^{n-2} z_j (s + z_n)(s + z_{n+1}) + \prod_{j=1}^{n-1} z_j (s + z_{n+1}) + \prod_{j=1}^{n-1} z_j \cdot z_n \right) \\
 &= (s + z_{n+1}) \left( \prod_{i=1}^n (s + z_i) - s \left( \prod_{i=2}^n (s + z_i) + z_1 \prod_{i=3}^n (s + z_i) + \dots + \prod_{j=1}^{n-2} z_j (s + z_n) + \prod_{j=1}^{n-1} z_j \right) \right) \\
 &\quad - s \prod_{j=1}^{n-1} z_j z_n = s \prod_{i=1}^n z_i + z_{n+1} \prod_{i=1}^n z_i - s \prod_{i=1}^n n z_i, \tag{A4}
 \end{aligned}$$

which proves that  $\text{RHS} = \text{LHS}$ . ■

[1] H.-P. Breuer and F. Petruccione, *The Theory of Open Quantum Systems* (Oxford University Press, Oxford, 2007).

[2] U. Weiss, *Quantum Dissipative Systems* (World Scientific, Singapore, 2000).

[3] V. Gorini, A. Kossakowski, and E. C. G. Sudarshan, *J. Math. Phys.* **17**, 821 (1976).

[4] G. Lindblad, *Commun. Math. Phys.* **48**, 119 (1976).

[5] R. Alicki and K. Lendi, *Quantum Dynamical Semigroups and Applications* (Springer, Berlin, 1987).

[6] C. W. Gardiner and P. Zoller, *Quantum Noise* (Springer-Verlag, Berlin, 1999).

[7] M. B. Plenio and P. L. Knight, *Rev. Mod. Phys.* **70**, 101 (1998).

[8] H. J. Carmichael, *Statistical Methods in Quantum Optics I: Master Equations and Fokker-Planck Equations* (Springer, Berlin, 1999).

[9] Á. Rivas, S. F. Huelga, and M. B. Plenio, *Rep. Prog. Phys.* **77**, 094001 (2014).

[10] F. Benatti, R. Floreanini, and G. Scholes, *J. Phys. B: At. Mol. Opt. Phys.* **45**, 150201 (2012).

[11] C. Addis, B. Bylicka, D. Chruściński, and S. Maniscalco, *Phys. Rev. A* **90**, 052103 (2014).

[12] M. M. Wolf, J. Eisert, T. S. Cubitt, and J. I. Cirac, *Phys. Rev. Lett.* **101**, 150402 (2008).

[13] Á. Rivas, S. F. Huelga, and M. B. Plenio, *Phys. Rev. Lett.* **105**, 050403 (2010).

[14] D. Chruściński and S. Maniscalco, *Phys. Rev. Lett.* **112**, 120404 (2014).

[15] S. Nakajima, *Prog. Theor. Phys.* **20**, 948 (1958); R. Zwanzig, *J. Chem. Phys.* **33**, 1338 (1960).

[16] S. Chaturvedi and J. Shibata, *Z. Phys. B* **35**, 297 (1979); N. H. F. Shibata and Y. Takahashi, *J. Stat. Phys.* **17**, 171 (1977).

[17] F. Benatti and R. Floreanini, *Mod. Phys. Lett. A* **12**, 1465 (1997).

[18] P. Hänggi and H. Thomas, *Z. Phys. B: Condens. Matter* **26**, 85 (1977); H. Grabert, P. Talkner, and P. Hänggi, *ibid.* **26**, 389 (1977); A. Fuliński and W. J. Kramarczyk, *Physica* **39**, 575 (1968).

[19] D. Chruściński and A. Kossakowski, *Phys. Rev. Lett.* **104**, 070406 (2010).

[20] S. M. Barnett and S. Stenholm, *Phys. Rev. A* **64**, 033808 (2001).

[21] A. A. Budini, *Phys. Rev. A* **69**, 042107 (2004).

[22] J. Wilkie, *Phys. Rev. E* **62**, 8808 (2000); J. Wilkie and Yin Mei Wong, *J. Phys. A: Math. Theor.* **42**, 015006 (2009).

[23] H.-P. Breuer and B. Vacchini, *Phys. Rev. Lett.* **101**, 140402 (2008); *Phys. Rev. E* **79**, 041147 (2009).

[24] S. Daffer, K. Wódkiewicz, J. D. Cresser, and J. K. McIver, *Phys. Rev. A* **70**, 010304 (2004).

[25] A. Shabani and D. A. Lidar, *Phys. Rev. A* **71**, 020101(R) (2005).

[26] S. Maniscalco, *Phys. Rev. A* **72**, 024103 (2005); S. Maniscalco and F. Petruccione, *ibid.* **73**, 012111 (2006).

[27] A. Kossakowski and R. Rebolledo, *Open Syst. Inf. Dyn.* **14**, 265 (2007); **16**, 259 (2009).

[28] D. Chruściński and A. Kossakowski, *Europhys. Lett.* **97**, 20005 (2012).

[29] B. Vacchini, *Phys. Rev. A* **87**, 030101(R) (2013).

[30] H.-P. Breuer, E.-M. Laine, and J. Piilo, *Phys. Rev. Lett.* **103**, 210401 (2009).

[31] K. M. R. Audenaert and S. Scheel, *New J. Phys.* **10**, 023011 (2008).

[32] In [5] random unitary evolution is called “dynamics for a system in random external fields.” In this paper we use “random unitary channel/evolution,” following, e.g., [31].

[33] D. Chruściński and F. A. Wudarski, *Phys. Lett. A* **377**, 1425 (2013).

[34] M. J. W. Hall, J. D. Cresser, L. Li, and E. Andersson, *Phys. Rev. A* **89**, 042120 (2014).

[35] K. S. Miller and S. G. Samko, *Integr. Transforms Spec. Funct.* **12**, 389 (2001).

[36] D. Chruściński and F. A. Wudarski, *Phys. Rev. A* **91**, 012104 (2015).

[37] L. Mazzola, E.-M. Laine, H.-P. Breuer, S. Maniscalco, and J. Piilo, *Phys. Rev. A* **81**, 062120 (2010).