

Photoelectron angular distributions, cross sections, and branching ratios for atomic oxygen

Anthony F. Starace*

Behlen Laboratory of Physics, The University of Nebraska, Lincoln, Nebraska 68508

Steven T. Manson†

Department of Physics, Georgia State University, Atlanta, Georgia 30303

David J. Kennedy‡

School of Physics, Georgia Institute of Technology, Atlanta, Georgia 30332

(Received 30 January 1974)

Theoretical calculations of the angular distribution of photoelectrons ejected from the $2p$ subshell of atomic oxygen for photoelectron kinetic energies $0 \leq \epsilon \leq 30$ Ry are presented using both Hartree-Fock (HF) and Herman-Skillman (HS) wave functions. Owing to the weakness of anisotropic electron-ion interactions, as evidenced by small differences between HF photoelectron phase shifts for alternative outgoing channels, the simple HS calculation of the angular-distribution asymmetry parameter is found to agree to within ≤ 0.1 with the HF calculations. Photoionization cross sections and photoelectron branching ratios have also been computed in both the HF and HS approximation for the $2p$ subshell of atomic oxygen in the wavelength range $910 \leq \lambda \leq 100$ Å. Comparison is made with other theoretical calculations and with the experimental branching-ratio measurements of Samson and Petrosky at $\lambda = 584.3$ Å.

I. INTRODUCTION

Dill, Manson, and Starace¹ have shown that the Cooper-Zare theory² for photoelectron angular distributions, while correct for closed-shell atoms, is not appropriate for open-shell atoms except in those cases where the phase shifts for alternative photoelectron channels are the same. We present here theoretical calculations that confirm atomic oxygen as one of these special cases. We find that the use of Herman-Skillman³ (HS) wave functions and the Cooper-Zare theory provides angular distribution asymmetry parameters that are in substantial agreement with those obtained using Hartree-Fock (HF) wave functions and a more rigorous theory¹ for photoelectron angular distributions. We expect similar results for the other atoms having open $2p$ subshells. Thus, the angular distributions of photoelectrons from $2p$ subshells calculated by Manson⁴ using HS wave functions and the Cooper-Zare theory should prove reliable.

We present also new calculations, using both HS and HF wave functions, of the total and partial photoionization cross sections of atomic oxygen. Because these cross sections are of fundamental importance for upper atmosphere physics, they have been calculated by many authors.⁵⁻¹⁰ Following Henry⁷ we plot also photoelectron branching ratios since these are the quantities measured directly in photoelectron spectroscopy. At $\lambda = 584.3$ Å our branching ratios are in reasonable agreement with both the experimental values of Samson

and Petrosky¹¹ and the more accurate close-coupling calculations of Henry.⁷ The calculated angular distributions, cross sections, and branching ratios presented in this paper thus extend and complement existing theoretical data on the photoionization of atomic oxygen.

II. THEORY

Dill and Fano¹²⁻¹⁴ have expressed the differential cross section for photoionization of an unpolarized target as an incoherent sum over contributions corresponding to alternative values of the angular momentum \vec{j}_i transferred by the incident photon to the target:

$$\frac{d\sigma}{d\Omega} = \sum_{j_i} \frac{\sigma(j_i)}{4\pi} [1 + \beta(j_i)P_2(\cos\theta)]. \quad (1)$$

Detailed expressions for the partial cross sections $\sigma(j_i)$ and asymmetry parameters $\beta(j_i)$ are given in Ref. 14 in terms of scattering amplitudes $S_l(j_i)$, where l is the orbital angular momentum of the outgoing photoelectron. The experimentally measured asymmetry parameter β is then given by the following weighted average¹⁴:

$$\beta = \left(\sum_{j_i} \sigma(j_i)\beta(j_i) \right) / \sum_{j_i} \sigma(j_i). \quad (2)$$

For the particular case of LS coupling the general photoionization process via electric dipole interaction,

$$A(l_0^n L_0 S_0 J_0) + h\nu \rightarrow A^+(l_0^{n-1} L_c S_c J_c) + e^-(l, s, j), \quad (3)$$

has been shown by Dill, Manson, and Starace¹ to have the following amplitude for transfer of j_i units of angular momentum (when the separation between fine-structure levels J_c is ignored):

$$S_i(j_i) = \frac{4\pi}{\lambda} \left(\frac{\pi\alpha h\nu}{3} \right)^{1/2} i^{-l} e^{i\sigma_{\epsilon l}} \hat{J}_0 \hat{l}_0 \begin{pmatrix} l & 1 & l_0 \\ 0 & 0 & 0 \end{pmatrix} (l_0^n L_0 S_0 \{ |l_0^{n-1} L_c S_c \rangle \times \sum_L \exp(i\delta_{\epsilon l}^{L_c S_c L}) R_{\epsilon l}^{L_c S_c L} \hat{L}^2 \begin{Bmatrix} L_0 & L_c & j_i \\ l & 1 & L \end{Bmatrix} \begin{Bmatrix} L_0 & L_c & l_0 \\ l & 1 & L \end{Bmatrix} \right). \quad (4)$$

Here $\sigma_{\epsilon l}$ is the Coulomb phase shift, dependent on the photoelectron orbital momentum l and kinetic energy ϵ , $\hat{x} \equiv (2x+1)^{1/2}$, $\nu\lambda = c$, and $\delta_{\epsilon l}^{L_c S_c L}$ is the photoelectron phase shift relative to Coulomb waves. The radial dipole matrix element is given by

$$R_{\epsilon l}^{L_c S_c L} \equiv \int_0^\infty \psi_i(n_0 l_0 L_0 S_0 | r) r \psi_f(\epsilon l L_c S_c L | r) dr \quad (5)$$

and the final-state continuum wave function ψ_f is normalized per unit energy.

Three points emphasized in Ref. 1 are important for understanding the calculations presented below. First, only in the limit that the final state wave functions ψ_f (and thus the phase shifts $\delta_{\epsilon l}^{L_c S_c L}$ and dipole matrix elements $R_{\epsilon l}^{L_c S_c L}$) are independent of L , and hence of $L_c S_c$, does the asymmetry parameter β in Eq. (2) reduce to the form obtained by Cooper and Zare.² For in this special case the sum over L in Eq. (4) may be performed analytically, yielding a δ function that restricts j_i to the single value $j_i = l_0$. In other words, the Cooper-Zare formula obtains only in the limit when there is no dynamical coupling of the orbital motion of the photoelectron to the net orbital motion of the residual ion. Second, Ref. 14 shows that the asymmetry parameter β depends on products of the scattering amplitudes in Eq. (4): specifically, on $|S_+(j_i)|^2$, $|S_-(j_i)|^2$, and $[S_+(j_i)S_-^\dagger(j_i) + \text{c.c.}]$, where the subscripts \pm indicate $l = j_i \pm 1$. Each of these products contains cross terms having a factor $\exp[i(\delta_{\epsilon l}^{L_c S_c L} - \delta_{\epsilon l'}^{L_c S_c L})]$ where $l' = l$ for the first two products and $l' = l - 2$ for the product $S_+(j_i) \times S_-^\dagger(j_i)$. The Cooper-Zare theory, however, assumes the dynamical phase shifts to be independent of L_c , S_c , and L ; hence only the product $S_+(j_i)$

$\times S_-^\dagger(j_i)$ involves a phase shift difference, which occurs in a factor $\exp[i(\delta_{\epsilon l} - \delta_{\epsilon l'})]$ for $l' = l - 2$. Therefore examination of the phase-shift differences for alternative photoelectron channels $L_c S_c L$ provides an estimate of the magnitude of anisotropic electron-ion interactions and hence of the validity of the Cooper-Zare formula. Third, for open-shell atoms the Cooper-Zare theory predicts the asymmetry parameters for alternative ion levels $L_c S_c$ to be identical when plotted as a function of photoelectron kinetic energy ϵ . Again, the validity of this result depends on the strength of anisotropic electron-ion interactions.

We are concerned here with the following photoionization processes in atomic oxygen:

$$O(2p^4 {}^3P) + h\nu \rightarrow O^+(2p^3 {}^4S, {}^2D, {}^2P) + e^-. \quad (6)$$

Binding energies for the 4S , 2D , and 2P ion terms are respectively 0.50079 a.u. (909.83 Å), 0.62300 a.u. (731.35 Å), and 0.68519 a.u. (664.97 Å).¹⁵ These experimental energies were used in computing the theoretical cross sections presented below.

Our discrete HF single-particle orbitals for the neutral atom and for the ion were obtained from the tabulation of Clementi.¹⁶ Continuum HF orbitals were obtained by the procedures described fully in Refs. 6 and 17. These continuum wave functions depend on both the ionic term level and on the total orbital angular momentum and thus the asymmetry parameter in Eq. (2) must be computed using the scattering amplitude in Eq. (4).

The HS continuum wave functions, on the other hand, do not depend on the ionic term levels or on the total orbital angular momentum and thus the Cooper-Zare formula for the asymmetry parameter is appropriate. These wave functions are

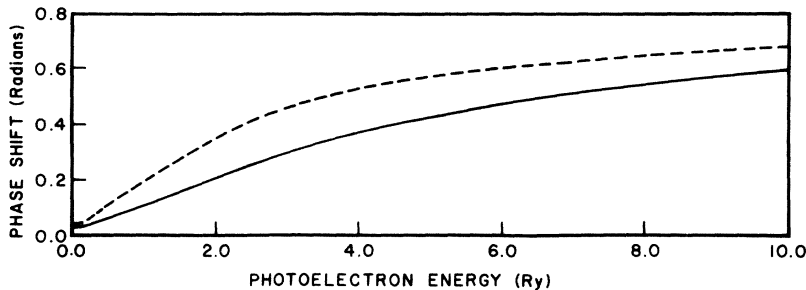


FIG. 1. Hartree-Fock d -wave phase shifts $\delta_{\epsilon d}^{L_c S_c L}$ for the 2D ion term vs photoelectron kinetic energy ϵ for two allowed values of L . Solid line corresponds to $L = 0$ [i.e., the state $2p^3({}^2D)\epsilon d^3S$] and the dashed line to $L = 2$ [i.e., the state $2p^3({}^2D)\epsilon d^3D$].

TABLE I. HF asymmetry parameters for the reactions $O(2p^4^3P) + h\nu \rightarrow O^+(2p^3^4S, ^2D, ^2P) + e^-$ as a function of photoelectron kinetic energy ϵ using dipole length (velocity) formula and comparison with HS asymmetry parameter.

ϵ (Ry)	$\beta(^3P \rightarrow ^4S)$	$\beta(^3P \rightarrow ^2D)$	$\beta(^3P \rightarrow ^2P)$	HS β
0.00	-0.243(-0.140)	-0.205(-0.114)	-0.199(-0.89)	-0.110
0.05	-0.247(-0.138)	-0.248(-0.149)	-0.233(-0.114)	-0.105
0.10	-0.091(0.010)	-0.104(-0.011)	-0.086(0.025)	+0.033
0.20	0.183(0.267)	0.159(0.237)	0.181(0.274)	0.277
0.40	0.548(0.603)	0.520(0.572)	0.544(0.606)	0.605
0.60	0.784(0.815)	0.756(0.787)	0.782(0.817)	0.799
0.80	0.944(0.958)	0.918(0.932)	0.946(0.961)	0.946
1.00	1.061(1.062)	1.038(1.038)	1.067(1.065)	1.052
1.50	1.250(1.230)	1.230(1.211)	1.261(1.234)	1.218
2.00	1.360(1.329)	1.342(1.313)	1.376(1.335)	1.317
3.00	1.473(1.437)	1.458(1.427)	1.495(1.444)	1.424
4.00	1.523(1.489)	1.508(1.482)	1.545(1.496)	1.476
6.00	1.541(1.521)	1.525(1.518)	1.560(1.525)	1.505
8.00	1.519(1.509)	1.502(1.507)	1.532(1.509)	1.495
10.00	1.479(1.479)	1.462(1.477)	1.485(1.475)	1.461
20.00	1.264(1.267)	1.258(1.265)	1.259(1.258)	1.250
30.00	1.069(1.069)	1.067(1.068)	1.059(1.060)	1.075

computed in the tabulated potential field of Herman and Skillman,³ which does not account for core relaxation. Discrete states for both the ion and the atom were taken to be the tabulated HS neutral-atom discrete orbitals. In computing the partial cross sections for photoionization to alternative ionic term levels the only difference is the experimental binding energy used: For a given photoelectron kinetic energy the wave functions employed are identical. Further details on the HS independent particle model are provided by Manson and Cooper.¹⁸

TABLE II. HF asymmetry parameters for the reactions $O(2p^4^1D) + h\nu \rightarrow O^+(2p^3^2D, ^2P) + e^-$ and $O(2p^4^1S) + h\nu \rightarrow O^+(2p^3^2P) + e^-$ as a function of photoelectron kinetic energy ϵ using dipole length (velocity) formula.

ϵ (Ry)	$\beta(^1D \rightarrow ^2D)$	$\beta(^1D \rightarrow ^2P)$	$\beta(^1S \rightarrow ^2P)$
0.00	-0.202(-0.127)	-0.155(-0.090)	-0.184(-0.099)
0.05	-0.230(-0.147)	-0.221(-0.147)	-0.218(-0.125)
0.10	-0.088(-0.007)	-0.092(-0.018)	-0.081(0.009)
0.20	0.169(0.240)	0.154(0.220)	0.169(0.247)
0.40	0.519(0.569)	0.498(0.547)	0.509(0.564)
0.60	0.749(0.779)	0.728(0.758)	0.734(0.768)
0.80	0.907(0.921)	0.888(0.903)	0.888(0.906)
1.00	1.023(1.025)	1.005(1.010)	1.002(1.009)
1.50	1.209(1.193)	1.197(1.184)	1.188(1.177)
2.00	1.318(1.294)	1.310(1.289)	1.297(1.279)
3.00	1.430(1.406)	1.427(1.408)	1.412(1.396)
4.00	1.478(1.463)	1.479(1.468)	1.464(1.457)
6.00	1.498(1.504)	1.502(1.511)	1.489(1.503)
8.00	1.479(1.500)	1.481(1.507)	1.472(1.501)
10.00	1.448(1.475)	1.450(1.481)	1.443(1.478)
20.00	1.259(1.275)	1.265(1.278)	1.263(1.280)
30.00	1.077(1.078)	1.080(1.081)	1.085(1.084)

Finally, the theory of atomic photoionization and appropriate formulas have been presented by Bates,¹⁹ Dalgarno, Henry, and Stewart,⁶ and Manson and Cooper.¹⁸

III. RESULTS

In Fig. 1 we have plotted HF phase shifts $\delta_{ed}^{L, S_c, L}$ as a function of photoelectron kinetic energy ϵ for the 2D ion term and for two values of L . These two phases differ from one another by about 0.16 rad in the energy region $2.0 \leq \epsilon \leq 4.0$. This difference is the largest found among any two oxygen phase shifts. It compares with phase-shift differences of up to 0.6 rad found between different photoelectron channels in atomic sulfur.¹ Because of these relatively small phase-shift differences in atomic oxygen, we expect the asymmetry parameter calculated with HS wave functions and the Cooper-Zare formula² to be in reasonably good agreement with asymmetry parameters calculated with HF wave functions and the more rigorous angular distribution formulas of Ref. 1.

In Table I we compare our HF asymmetry parameters with the HS one as a function of photoelectron kinetic energy. Note that if the oxygen photoelectron phase shifts were all identical then the HF asymmetry parameters for each ionic term would be the same. As it is, we see that the differences between the HF asymmetry parameters for different ionic term levels are generally smaller than the length and velocity values for a given term level. If the HF length values are regarded as the correct ones,²⁰ then the HS values are about 0.1 too high for $0.0 \leq \epsilon \leq 0.4$ Ry, but are

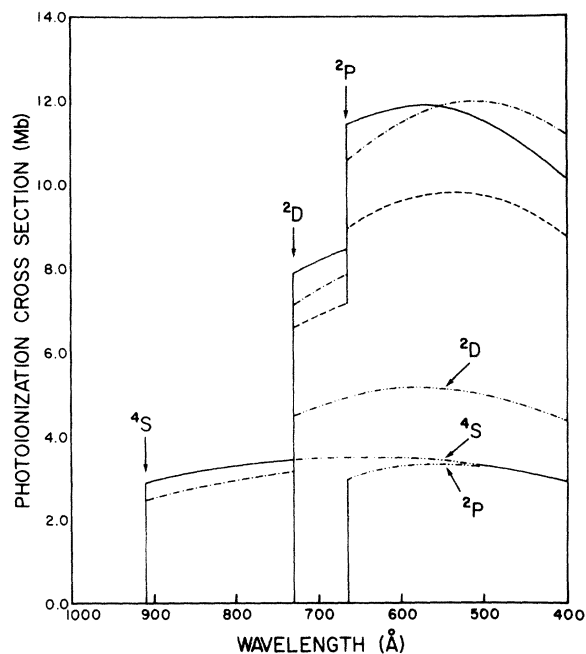


FIG. 2. Total and partial photoionization cross sections for atomic oxygen. Solid line: HS total cross section. Dash-dot line: HF (length) total cross section. Dashed line: HF (velocity) total cross section. Dash-double-dot line: HS partial cross sections. Note that for $910 \leq \lambda \leq 731$ the HS and HF velocity cross sections coincide.

in excellent agreement with HF for higher kinetic energies. In Table II we have listed HF asymmetry parameters for photoionization from the 1D and 1S excited term levels of atomic oxygen. These are also in agreement with the HS asymmetry parameter in Table I.

In Fig. 2 we have plotted HS and HF total cross

sections for atomic oxygen. The HS partial cross sections are also shown and the HF partial cross sections are listed in Table III. We see that near each threshold the HS cross sections are higher than either HF length or HF velocity. For higher kinetic energies, the HS cross sections fall between HF length and velocity. Note that at the 4S threshold our HF length (velocity) cross section of 2.46 (2.89) agrees best with the values 2.3 (2.8) of Bates and Seaton⁵ rather than the values 2.7 (3.4) of Dalgarno, Henry, and Stewart.⁶ The differences among these three HF calculations probably arise from the different discrete orbitals used. Above the 2P threshold our HF length and velocity cross sections are slightly lower (≈ 0.2 – 0.3 Mb) than those of Ref. 6 but not as low as the more accurate close-coupling cross sections of Henry.⁷ It is unclear why the cross sections of Dalgarno, Henry, and Stewart⁶ and of this paper are in closer agreement with the absolute experimental measurements of Cairns and Samson²¹ than is the more complete calculation of Henry.⁷

Finally, in Fig. 3 we have plotted our HS photoelectron branching ratios $\sigma(^2D)/\sigma(^4S)$ and $\sigma(^2P)/\sigma(^4S)$ and compared them with those of Henry.⁷ In Table III we also list the HF branching ratios we have obtained. For $731 \leq \lambda \leq 300$ Å our HF-length branching ratios agree nearly identically with those of HS. For $\lambda < 300$ Å the HF-length branching ratios drop below the values obtained by Henry.⁷ The HF-velocity branching ratios are much lower than the HF-length ratios for $\lambda > 300$ Å and generally equal to the length ratios below 300 Å. At 584.3 Å both our HS- and HF-length branching ratios are equal. Our $\sigma(^2D)/\sigma(^4S)$ branching ratio of 1.48–1.49 compares with the experimental ratio of 1.57 ± 0.14 .¹¹ Our $\sigma(^2P)/\sigma(^4S)$ ratio of 0.95–0.96

TABLE III. HF partial cross sections and branching ratios for the reactions $O(2p^4^3P) + h\nu \rightarrow O^+(2p^4^1S, ^2D, ^2P) + e^-$ as a function of wavelength using dipole length (velocity) formula.

λ (Å)	$\sigma(^4S)$	$\sigma(^2D)$	$\sigma(^2P)$	$\sigma(^2D)/\sigma(^4S)$	$\sigma(^2P)/\sigma(^4S)$
909.8	2.463(2.892)
731.4	3.122(3.399)	3.982(3.177)	...	1.28(0.93)	...
665.0	3.285(3.450)	4.578(3.685)	2.682(1.806)	1.39(1.07)	0.82(0.52)
650.0	3.321(3.461)	4.685(3.775)	2.802(1.904)	1.41(1.09)	0.84(0.55)
600.0	3.374(3.400)	4.948(3.985)	3.122(2.173)	1.47(1.17)	0.93(0.64)
584.3	3.391(3.381)	5.013(4.034)	3.205(2.246)	1.48(1.19)	0.95(0.66)
550.0	3.400(3.308)	5.142(4.129)	3.315(2.348)	1.51(1.25)	0.97(0.71)
500.0	3.374(3.159)	5.183(4.122)	3.421(2.451)	1.54(1.31)	1.01(0.78)
450.0	3.275(2.943)	5.085(3.992)	3.400(2.449)	1.55(1.36)	1.04(0.83)
400.0	3.097(2.668)	4.809(3.721)	3.239(2.338)	1.55(1.39)	1.05(0.88)
350.0	2.827(2.332)	4.336(3.307)	2.926(2.116)	1.53(1.42)	1.04(0.91)
300.0	2.461(1.940)	3.674(2.760)	2.444(1.772)	1.49(1.42)	0.99(0.91)
250.0	1.979(1.496)	2.813(2.092)	1.829(1.335)	1.42(1.40)	0.92(0.89)
200.0	1.405(1.028)	1.850(1.380)	1.157(0.864)	1.32(1.34)	0.82(0.84)
150.0	0.815(0.588)	0.991(0.751)	0.591(0.457)	1.21(1.28)	0.72(0.78)
100.0	0.308(0.222)	0.339(0.268)	0.193(0.159)	1.10(1.21)	0.63(0.72)

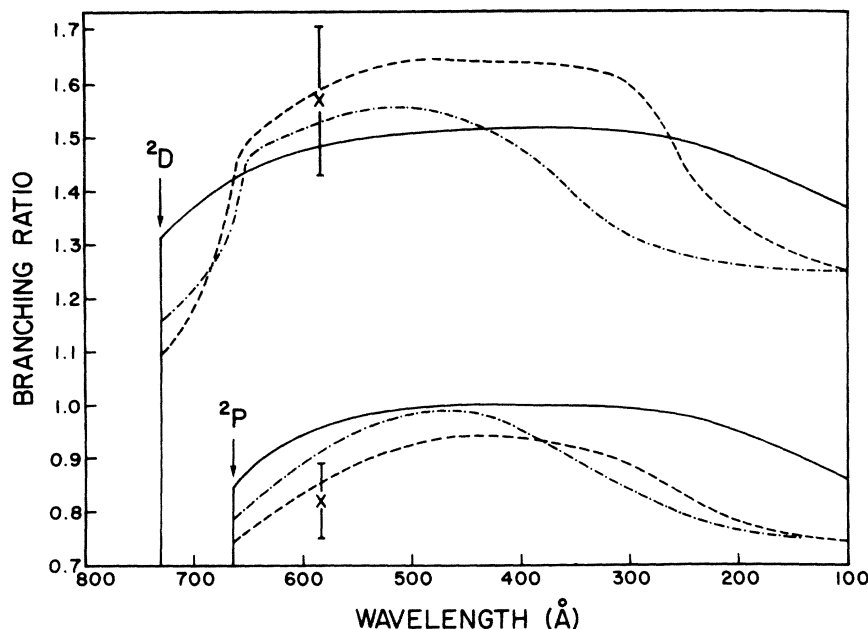


FIG. 3. Photoelectron branching ratios $\sigma(^2D)/\sigma(^4S)$ and $\sigma(^2P)/\sigma(^4S)$ for atomic oxygen. Solid line: HS. Dash-dot line: Henry (length). Dashed line: Henry (velocity). Crosses: experimental values of Samson and Petrosky.

compares with the experimental value of 0.82 ± 0.07 .¹¹ It is clear from Fig. 3 and Table III that experimental measurements of photoelectron branching ratios at a variety of wavelengths would be of great interest to atomic theorists.

IV. CONCLUSIONS

We have confirmed the prediction of Dill, Manson, and Starace¹ that examination of photoelectron phase-shift differences is the appropriate criterion for determining how elaborate a calculation must be made for obtaining accurate angular distribution asymmetry parameters for open-shell atoms. We have found that HF photoelectron phase shifts in atomic oxygen are sufficiently close to one another to permit a simple calculation of the asymmetry parameters for $2p$ photoionization. We expect our results to apply generally to atoms having open $2p$

subshells.

We have also calculated total and partial photoionization cross sections for atomic oxygen using both Herman-Skillman and Hartree-Fock wave functions. At 584.3 \AA our calculated branching ratios are in good agreement with the experimental measurements of Samson and Petrosky.¹¹ They compare well with other theoretical cross-section calculations except the more detailed close-coupling calculations of Henry.⁷ It would be of great theoretical interest therefore to employ close-coupling wave functions to calculate asymmetry parameters within the framework of Ref. 1.

ACKNOWLEDGMENTS

One of us (A.F.S.) wishes to thank Professor J. A. R. Samson for suggesting the study of atomic oxygen.

*Supported in part under National Aeronautics and Space Administration Contract No. NGR 28-004-021.

†Supported in part under National Science Foundation Contract No. GP-38905.

‡Present address: School of Physical Sciences, New University of Ulster, Coleraine, Northern Ireland.

¹D. Dill, S. T. Manson, and A. F. Starace, *Bull. Am. Phys. Soc.* **18**, 1514 (1973); *Phys. Rev. Lett.* (to be published).

²J. Cooper and R. N. Zare, *Lectures in Theoretical Physics*, edited by S. Geltman, K. T. Mahanthappa, and W. E. Britton (Gordon and Breach, New York, 1969), Vol. XI-C, pp. 317-37. The limitations of the Cooper-Zare theory, presumably well-known to the authors, arise essentially from the authors' initial as-

sumption that the radial part of the photoelectron wave function is independent of photoelectron channel quantum numbers.

³F. Herman and S. Skillman, *Atomic Structure Calculations* (Prentice-Hall, Englewood Cliffs, N. J., 1963).

⁴S. T. Manson, *J. Electron Spectroscopy and Related Phenomena* **1**, 413 (1973).

⁵D. R. Bates and M. J. Seaton, *Mon. Not. Roy. Astron. Soc.* **109**, 698 (1949).

⁶A. Dalgarno, R. J. W. Henry, and A. L. Stewart, *Planet. Space Sci.* **12**, 235 (1964).

⁷R. J. W. Henry, *Planet. Space Sci.* **15**, 1747 (1967).

⁸G. M. Thomas and T. M. Helliwell, *J. Quant. Spectrosc. Radiat. Trans.* **10**, 423 (1970).

⁹H. Kähler, *J. Quant. Spectrosc. Radiat. Transfer* **11**,

- 1521 (1971).
- ¹⁰P. S. Ganas, Phys. Rev. A 7, 928 (1973).
- ¹¹J. A. R. Samson and V. E. Petrosky, preceding paper, Phys. Rev. A 9, 2449 (1974).
- ¹²U. Fano and D. Dill, Phys. Rev. A 6, 185 (1972).
- ¹³D. Dill and U. Fano, Phys. Rev. Lett. 29, 1203 (1972).
- ¹⁴D. Dill, Phys. Rev. A 7, 1976 (1973).
- ¹⁵C. E. Moore, *Ionization Potentials and Ionization Limits Derived from the Analysis of Optical Spectra* (U. S. GPO, Washington, D.C., 1970).
- ¹⁶E. Clementi, IBM J. Res. Develop. 9, 2 (1965), and supplement entitled "Tables of Atomic Functions."
- ¹⁷D. J. Kennedy and S. T. Manson, Phys. Rev. A 5, 227 (1972).
- ¹⁸S. T. Manson and J. W. Cooper, Phys. Rev. 165, 126 (1968).
- ¹⁹D. R. Bates, Mon. Not. Roy. Astron. Soc. 106, 432 (1946).
- ²⁰A. F. Starace, Phys. Rev. A 3, 1242 (1971); 8, 1141 (1973).
- ²¹R. B. Cairns and J. A. R. Samson, Phys. Rev. 139, A1403 (1965).