

Magneto-optical spectroscopy with polarization-modulated lightE. Breschi,^{*} Z. D. Gruić, P. Knowles, and A. Weis*Department of Physics, University of Fribourg, 1700 Fribourg, Switzerland*

(Received 5 July 2013; published 16 August 2013)

We present a combined theoretical and experimental study of magnetic resonance transitions induced by polarization-modulated light in cesium vapor exposed to a transverse magnetic field. Signals are obtained by phase-sensitive analysis of the light power traversing the vapor cell at six harmonics of the polarization modulation frequency. Resonances appear whenever the Larmor frequency matches an integer multiple of the modulation frequency. We have further investigated the modifications of the spectra when varying the modulation duty cycle. The resonance amplitudes of both in-phase and quadrature components are well described in terms of the Fourier coefficients of the modulation function. The background-free signals generated by the polarization modulation scheme have a high application potential in atomic magnetometry.

DOI: [10.1103/PhysRevA.88.022506](https://doi.org/10.1103/PhysRevA.88.022506)

PACS number(s): 32.30.Dx, 07.55.Ge, 07.55.Jg, 78.20.Ls

I. INTRODUCTION

Magneto-optical spectroscopy of atomic media is a powerful tool that uses resonant optical interactions for detecting magnetic resonance in atoms (a comprehensive review is given by Alexandrov *et al.* [1]). Magnetic resonance transitions are conventionally excited by a magnetic field oscillating at the atoms' Larmor frequency, but it has been known since the seminal work by Bell and Bloom in 1961 [2] that an amplitude-modulated (AM) resonant light beam can also induce magnetic resonance transitions. More recently, the Bell-Bloom method, in combination with phase-sensitive detection, has found renewed interest [3,4], in particular for its application in high-sensitivity atomic magnetometry [5]. Alternatively, modulation of either the light's frequency or polarization at the Larmor frequency induces magnetic resonance. While frequency-modulation spectroscopy has become a well-established method for high-sensitivity magnetometry [6,7], little work has been devoted to polarization modulation [8–11]. Recently, resonant polarization modulation has been used to solve specific technical issues. The authors of Ref. [12] have demonstrated that polarization modulation between circular and linear polarization states eliminates both dead zones and heading errors in an alkali-metal-atom magnetometer. Polarization modulation between circular polarization states at the ground-state hyperfine transition frequency was shown to increase the contrast of the clock resonance as described in Ref. [13], and references therein.

A theoretical model deriving algebraic expressions for the rich resonance structure of magnetic resonance transitions driven by amplitude, frequency, and polarization modulation has recently been presented [14]. So far the model predictions were shown to give an excellent description of experimental results obtained with amplitude modulation using phase-sensitive detection [14] and for polarization modulation with low-pass-filtered detection [11]. Here, we address polarization modulation in combination with phase-sensitive detection. We characterize, both experimentally and theoretically, the magneto-optical resonance spectra observed in a transverse magnetic field, when the helicity of the exciting laser beam

is periodically reversed with different duty cycles. We first introduce our experimental and theoretical methods and then discuss the full magneto-optical resonance spectra as a function of the modulation duty cycle. We show that resonances appear when the Larmor frequency is a harmonic of the modulation frequency and that the relative resonance amplitudes are well described in terms of the Fourier coefficients of the polarization modulation function.

II. METHOD

Experiments are carried out in a conventional magneto-optical spectroscopy apparatus [11,15]. Light from a 894-nm distributed feedback diode laser passes an electro-optical modulator, which flips the light polarization between σ^+ and σ^- polarizations at a constant frequency $\omega_{\text{mod}} = 2\pi \times 267$ Hz. The degrees of circular polarization were determined to be 99.7% and 99.8%, for σ^+ and σ^- , respectively. The laser frequency can be tuned to any hyperfine component of the cesium D₁ line ($6S_{1/2} \rightarrow 6P_{1/2}$); however, the experiments reported here were carried out on the $F_g = 4 \rightarrow F_e = 3$ transition, which provides the highest magneto-optical resonance contrast [16]. The modulated beam passes an evacuated paraffin-wall-coated spherical cell containing cesium vapor at room temperature (20 °C) saturated vapor pressure. The laser beam diameter is about 10 times smaller than the cell diameter ($d = 30$ mm) so that lensing or birefringence effects from the cell itself do not critically affect the polarization quality. This is testable in that linear polarization components will create resonances at half of the Larmor frequency due to the symmetry and evolution of any created spin alignment. We did not see resonances in the experimental spectra arising from spin alignment. The final excellent agreement between model and experimental results (cf. Sec. III) strongly implies that the cell has negligible effect on the light polarization.

The cell is isolated in a magnetically controlled environment [15], with nominal residual field components below a few nT. The amplitude of a static magnetic field applied orthogonally to the laser propagation direction, and hence the Larmor frequency ω_L , is scanned in the range $\pm 4 \omega_{\text{mod}}$. The experimental signal is formed by detecting the light power $P(t)$ transmitted through the cell: The photocurrent signal is converted into voltage, amplified, and then demodulated via a

^{*}evelina.breschi@unifr.ch

lock-in amplifier, which allows the simultaneous recording of the in-phase and quadrature signals at six harmonics q ω_{mod} of the modulation frequency.

A. Model

For signal interpretation we apply the model developed in Ref. [14] for a spin-oriented atomic sample under periodically modulated excitation and detection. The key results of that model are summarized here. The power transmitted through the cell is given by the Lambert-Beer law, which, for an optically thin atomic vapor of length L , is well approximated by $P = P_0 - P_0 \kappa L$, where κ is the optical absorption coefficient and L is the sample length. For a laser beam resonant with a $F_g \rightarrow F_e$ hyperfine transition, incident on atoms with a longitudinal spin orientation S_z , the absorption coefficient becomes $\kappa = \kappa_0(1 - \alpha \xi S_z)$. Here, κ_0 is the peak absorption coefficient of unpolarized atoms, ξ is the light helicity, and $\alpha \equiv \alpha_{F_g, F_e}$ is the orientation analyzing power. Dropping time-independent terms, the transmitted laser power is given by

$$P(t) = (\alpha \kappa_0 L P_0) S_z(t) \xi(t), \quad (1)$$

where $S_z(t)$ is the time-dependent longitudinal spin orientation resulting from the time-dependent solutions of the Bloch equations under polarization modulation. A square-wave modulation function $\xi(t)$ is chosen to make a fast switch between the two circular polarization states without spending significant time in linearly polarized intermediate states. We allow for modulation with arbitrary duty cycle ($0 < \eta < 1$) chosen to be symmetric with respect to $t = 0$ and constrained to have $|\xi(t)| = 1$, whose Fourier expansion is given by

$$\xi(t) = \sum_{j=-\infty}^{+\infty} g_j(\eta) \cos(j\omega t), \quad (2)$$

$$g_0(\eta) = 2\eta - 1 \quad \text{and} \quad g_{j \neq 0}(\eta) = \frac{2 \sin(\pi j \eta)}{\pi j}. \quad (3)$$

In-phase I_q and quadrature Q_q parts of $P(t)$ extracted by demodulation at $q \omega_{\text{mod}}$ can be written as

$$I_q(\eta) = \sum_{m=-\infty}^{\infty} a_{q,m}(\eta) \mathcal{A}_m, \quad Q_q(\eta) = \sum_{m=-\infty}^{\infty} d_{q,m}(\eta) \mathcal{D}_m, \quad (4)$$

where \mathcal{A}_m and \mathcal{D}_m are, respectively, absorptive and dispersive Lorentzian resonances given by

$$\mathcal{A}_m = \frac{\gamma^2}{(m \omega_{\text{mod}} - \omega_L)^2 + \gamma^2}, \quad (5)$$

$$\mathcal{D}_m = \frac{\gamma(m \omega_{\text{mod}} - \omega_L)}{(m \omega_{\text{mod}} - \omega_L)^2 + \gamma^2}, \quad (6)$$

with amplitudes

$$a_{q,m}(\eta) = \sqrt{2} \alpha \kappa_0 L \frac{P_0^2}{P_s} g_m(\eta) [g_{q-m}(\eta) + g_{q+m}(\eta)], \quad (7)$$

$$d_{q,m}(\eta) = \sqrt{2} \alpha \kappa_0 L \frac{P_0^2}{P_s} g_m(\eta) [g_{q-m}(\eta) - g_{q+m}(\eta)], \quad (8)$$

where P_s is the optical pumping saturation power. We note that the model was developed for the low power limit ($a_{q,m}, d_{q,m} \propto P_0^2$), and that the experiments were carried out in that range.

III. ANALYSIS OF THE RESONANCE SPECTRUM

We have studied the demodulated signals as a function of the detection harmonic q and modulation duty cycle η . Since in the Fourier expansion of a symmetric $\eta = 0.5$ square wave the even coefficients vanish, one sees from Eqs. (7) and (8) that lock-in resonances appear only at odd m and even q . To compare experimental and theoretical results we normalize the in-phase and quadrature signals to the highest amplitude absorptive and dispersive signals, respectively, found when $m = 1$, $q = 2$, and $\eta = 0.5$. No other scaling factors are needed. The normalized model signals thus read

$$\frac{I_q(\eta)}{a_{2,1}} = \sum_{m=-\infty}^{\infty} \frac{a_{q,m}}{a_{2,1}} \mathcal{A}_m, \quad (9)$$

$$\text{and} \quad \frac{Q_q(\eta)}{d_{2,1}} = \sum_{m=-\infty}^{\infty} \frac{d_{q,m}}{d_{2,1}} \mathcal{D}_m. \quad (10)$$

Theory also predicts $d_{2,1}/a_{2,1} = 1/2$.

Experimental amplitudes are obtained by fitting the recorded data with the absorptive and dispersive Lorentzians of Eqs. (5) and (6), yielding fit amplitudes that represent the model amplitudes (7) and (8), respectively. We then normalize all fitted in-phase (quadrature) amplitudes to the in-phase (quadrature) amplitude of the ($m = 1$, $q = 2$, $\eta = 0.5$) resonance. The experimental $d_{2,1}/a_{2,1}$ ratio was found to be 0.492(1), in accordance with the theory prediction. In this way, we eliminate all secondary experimental parameters affecting the signal, e.g., detector quantum efficiency, current-voltage converter gains, etc. The normalization of the model predictions, of course, removes unknown theory parameters (e.g., P_s , α_{F_g, F_e} , κ_0 , etc).

In Figs. 1 and 2 we compare experimental and theoretical spectra for different values of q when $\eta = 0.50$ and $\eta = 0.10$. No scaling has been applied beyond the normalization procedure above.

As predicted, resonances are observed when the Larmor frequency is a multiple of the modulation frequency, i.e., $\omega_L = m \omega_{\text{mod}}$. An important point to note is that in contrast to experiments with frequency- or amplitude-modulated light, polarization modulation yields background-free in-phase and quadrature signals. The amplitude of the zero-field level-crossing (Hanle) resonance is proportional to g_0 , Eq. (3), and thus vanishes for 50% duty cycle.

Figure 1 presents the signals obtained for $\eta = 0.50$, where one can see the largest signals at $m = 1$ and $q = 2$ (the absorptive and dispersive reference resonances for normalization). This resonance is the most interesting for atomic magnetometry, and work is ongoing to fully characterize the signal and study the final sensitivity limit of the method. The $\eta = 0.5$ duty cycle pumping, i.e., excitation with a symmetric square wave, essentially corresponds to the so-called push-pull optical pumping [17], whose effect is the synchronization of the pumping light's polarization modulation with the harmonic evolution of the atomic quantum state in the external field. Notable in the $\eta = 0.5$ case is that all $q = \text{odd}$ resonances are predicted to be zero by the model (as discussed at the start of the section), and the data (only one trace for $q = 1$ is plotted in the figure) show no signal at the detection limit in those cases.

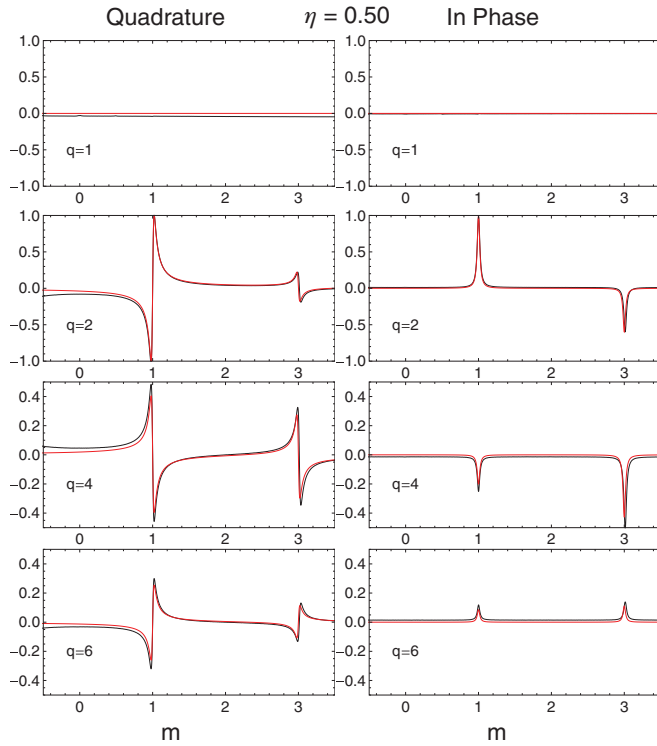


FIG. 1. (Color online) Magnetic field dependence at 50% duty cycle ($\eta = 0.50$) of normalized in-phase and quadrature lock-in signals, demodulated at even harmonics q of ω_{mod} ; the odd harmonic signal is predicted to be zero, and the data (here $q = 1$ is shown) support the prediction. The ratio between the Larmor frequency and the modulation frequency is m , with $m = \omega_L/\omega_{\text{mod}}$, where $\omega_{\text{mod}} = 2\pi \times 267$ Hz. Black lines are measurements, and red (gray) lines are the model prediction.

Figure 2 presents the signals obtained for $\eta = 0.10$, where $q =$ odd resonances are not null.

The magnetic resonance linewidths show no dependence on m and q , consistent with the fact that the linewidth is determined by relaxation mechanisms and power broadening common to all resonances. In the antirelaxation wall-coated cell, the primary relaxation mechanisms arise from spin-exchange collisions and loss of the atoms to the alkali reservoir via the entrance channel, followed in importance by collisions between alkali atoms and the antirelaxation coating, and magnetic field inhomogeneities, none of which has an expected dependence on m or q .

We have investigated the dependence of the spectra on the duty cycle η by varying η between 0.1 and 0.9 in steps of 0.05. For each value of η we measured the amplitudes of the resonances occurring at $\omega_L = \omega_{\text{mod}}, 2\omega_{\text{mod}},$ and $3\omega_{\text{mod}}$ (i.e., for $m = 1, 2, 3$) in each of the six demodulation channels $q = (1, \dots, 6)\omega_{\text{mod}}$. Figure 3 shows the experimental amplitudes (after normalization) as a function of η and q , together with the model predictions that are in excellent agreement with experiment. We note the alternating symmetric-antisymmetric η dependence with respect to $\eta = 0.5$ for odd and even q , respectively.

The magnetic resonance linewidth depends on η because the effective power driving either the σ^+ or σ^- transition—and hence the power broadening of the line—depends on the duty

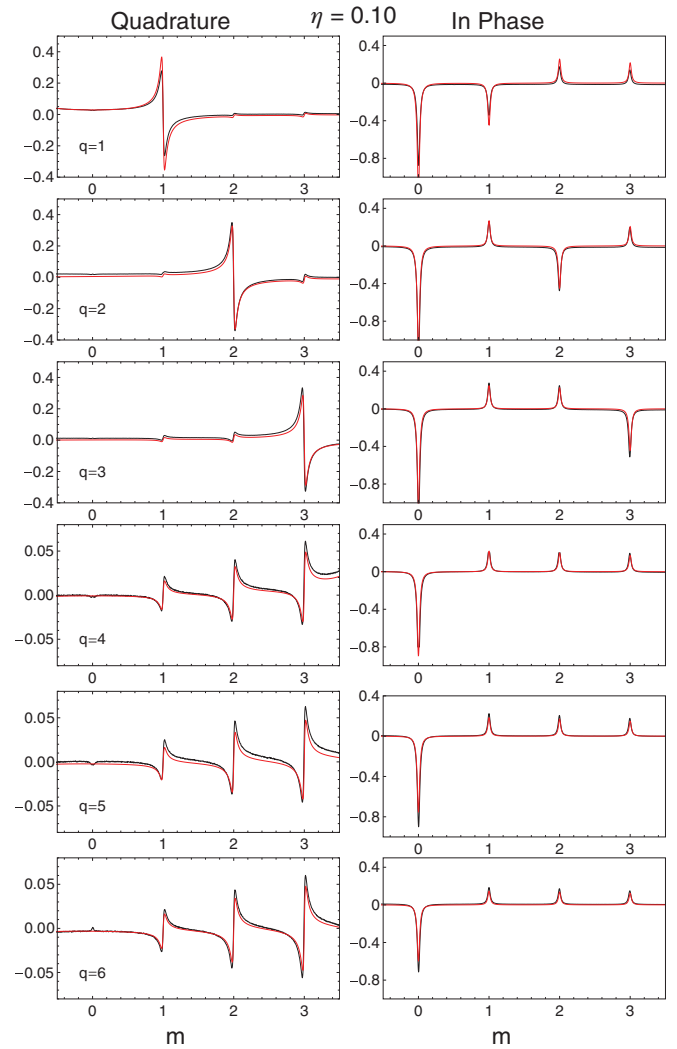


FIG. 2. (Color online) Magnetic field dependence of normalized in-phase and quadrature lock-in signals, demodulated at harmonic q of ω_{mod} , for laser light resonant with the $F_g = 4 \rightarrow F_e = 3$ transition. Black lines are measurements, and red (gray) lines are the model prediction. The duty cycle was 10% ($\eta = 0.10$).

cycle η . For this reason the balanced σ^+/σ^- excitation with $\eta = 0.5$ is expected to produce the narrowest resonances, as observed in the experiment.

IV. SUMMARY AND CONCLUSIONS

We have studied magneto-optical resonances that occur at multiples of the Larmor frequency when the polarization of a resonant laser beam traversing an alkali atom vapor is switched between left and right circular polarization. The lock-in demodulated signals have a rich spectral structure that is well reproduced by algebraic model predictions. The linear zero crossings of the quadrature resonances can be used as discriminator signals for magnetic field measurements. The optimal resonance for magnetometry applications is the one with $q = 2, m = 1$, and $\eta = 0.5$, since it has the largest amplitude and simultaneously the narrowest linewidth. The application of the polarization-modulation approach reported

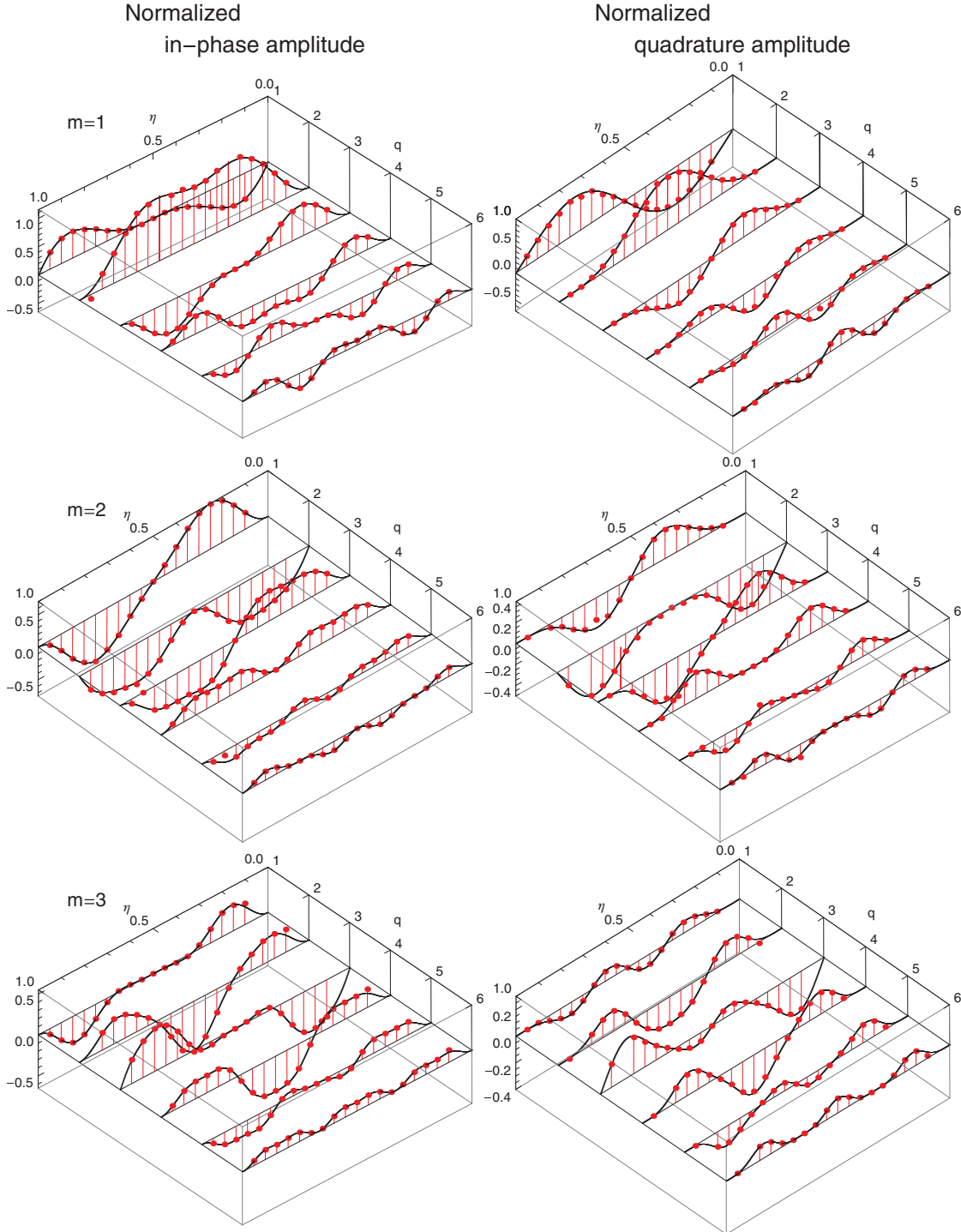


FIG. 3. (Color online) Normalized amplitude of the in-phase (left) and quadrature (right) signal for $m = 1, 2, 3$ (from top to bottom). The amplitudes are function of the detection harmonic ($q = 1, \dots, 6$) and the duty cycle ($0 < \eta < 1$).

here to atomic magnetometry has a distinctive advantage compared to similar approaches using amplitude or frequency modulation, viz., the absence of a dc background on the in-phase signal. In a feedback-locked magnetometer, an imperfect phase setting transfers a part of the dc background onto the quadrature signal, thereby introducing a shift in the magnetic field measure and increasing power noise.

Finally polarization modulation with variable duty cycle can also find applications in the preparation and manipulation

of specific atomic states, i.e., in metrology [18] and in quantum information processing [19].

ACKNOWLEDGMENTS

This work is supported by SNF-Ambizione Grant No. PZ00P2_131926. We thank the mechanical workshop and the electronics pool of the Physics Department for expert technical support.

- [1] E. B. Alexandrov, M. Auzinsh, D. Budker, D. F. Kimball, S. M. Rochester, and V. V. Yashchuk, *J. Opt. Soc. Am. B* **22**, 7 (2005).
- [2] W. E. Bell and A. L. Bloom, *Phys. Rev. Lett.* **6**, 280 (1961).
- [3] W. Gawlik, L. Krzemień, S. Pustelny, D. Sangla, J. Zachorowski, M. Graf, A. O. Sushkov, and D. Budker, *Appl. Phys. Lett.* **88**, 131108 (2006).
- [4] V. Schultze, R. IJsselsteijn, T. Scholtes, S. Woetzel, and H.-G. Meyer, *Opt. Exp.* **20**, 14201 (2012).
- [5] D. Budker and M. Romalis, *Nat. Phys.* **3**, 227 (2007).
- [6] D. Budker, D. F. Kimball, V. V. Yashchuk, and M. Zolotarev, *Phys. Rev. A* **65**, 055403 (2002).
- [7] C. Andreeva, S. Cartaleva, Y. Dancheva, V. Biancalana, A. Burchianti, C. Marinelli, E. Mariotti, L. Moi, and K. Nasyrov, *Phys. Rev. A* **66**, 012502 (2002).
- [8] E. B. Aleksandrov, *Sov. Phys. Usp.* **15**, 436 (1973).
- [9] H. Gilles, B. Cheron, and J. Hamel, *Opt. Commun.* **81**, 369 (1991).
- [10] H. Klepel and D. Suter, *Opt. Commun.* **90**, 46 (1992).
- [11] I. Fescenko, P. Knowles, A. Weis, and E. Breschi, *Opt. Exp.* **21**, 15121 (2013).
- [12] A. Ben-Kish and M. V. Romalis, *Phys. Rev. Lett.* **105**, 193601 (2010).
- [13] Y.-Y. Jau and W. Happer, *Phys. Rev. Lett.* **99**, 223001 (2007).
- [14] Z. D. Grujić and A. Weis, *Phys. Rev. A* **88**, 012508 (2013).
- [15] E. Breschi, Z. D. Grujić, and A. Weis, *Appl. Phys. B* 1 (2013).
- [16] M. Auzinsh, R. Ferber, F. Gahbauer, A. Jarmola, and L. Kalvans, *Phys. Rev. A* **78**, 013417 (2008).
- [17] Y.-Y. Jau, E. Miron, A. B. Post, N. N. Kuzma, and W. Happer, *Phys. Rev. Lett.* **93**, 160802 (2004).
- [18] G. Di Domenico, L. Devenoges, C. Dumas, and P. Thomann, *Phys. Rev. A* **82**, 053417 (2010).
- [19] B. Wang, Y. Han, J. Xiao, X. Yang, C. Zhang, H. Wang, M. Xiao, and K. Peng, *Phys. Rev. A* **75**, 051801 (2007).