

Multiparty quantum secret sharing

Zhan-jun Zhang,^{1,2} Yong Li,³ and Zhong-xiao Man²

¹*School of Physics and Material Science, Anhui University, Hefei 230039, China*

²*Wuhan Institute of Physics and Mathematics, Chinese Academy of Sciences, Wuhan 430071, China*

³*Department of Physics, Huazhong Normal University, Wuhan 430079, China*

(Received 1 November 2004; published 11 April 2005)

Based on a quantum secure direct communication (QSDC) protocol [Phys. Rev. A **69** 052319 (2004)], we propose a (n, n) -threshold scheme of multiparty quantum secret sharing of classical messages (QSSCM) using only single photons. We take advantage of this multiparty QSSCM scheme to establish a scheme of multiparty secret sharing of quantum information (SSQI), in which only all quantum information receivers collaborate can the original qubit be reconstructed. A general idea is also proposed for constructing multiparty SSQI schemes from any QSSCM scheme.

DOI: 10.1103/PhysRevA.71.044301

PACS number(s): 03.67.Dd, 03.65.Ta, 89.70.+c

Suppose Alice wants to send a secret message to two distant parties Bob and Charlie. One of them, Bob or Charlie, is not entirely trusted by Alice. And she knows that if the two guys coexist, the honest one will keep the dishonest one from doing any damage. Instead of giving the total secret messages to any one of them, it may be desirable for Alice to split the secret messages into two encrypted parts and send each one a part so that no one alone is sufficient to obtain the whole original information unless they collaborate. To gain this end classical cryptography can use a technique called secret sharing [1,2], where secret messages are distributed among n users in such a way that only by combining their pieces of information can the n users recover the secret messages. Usually these kinds of protocols are divided into classes, where from the n receivers, m can collaborate to produce the desired result. In this paper, we will focus on a (n, n) scheme, where all the receivers need to collaborate to obtain the desired message.

Recently this concept has been generalized to a quantum scenario [3]. The quantum secret sharing (QSS) is likely to play a key role in protecting secret quantum information, e.g., in secure operations of distributed quantum computation, sharing difficult-to-construct ancilla states and joint sharing of quantum money [6], and so on. Hence, after the pioneering QSS work proposed by using three-particle and four-particle GHZ states [3], these kinds of works on QSS attracted a great deal of attention both theoretically and experimentally [4–16]. All these works [3–16] can be divided into two kinds, one only deals with the QSS of classical messages (i.e., bits) [5,6,8–11,13,14] or only deals with the QSS of quantum information [4,7,12,15,16], where the secret is an arbitrary unknown state in a qubit and the other [3] studies both, that is, deals with QSS of classical messages and QSS of quantum information simultaneously. In all those schemes [3,5,6,8–11,13,14] dealing with the QSS of classical messages (bits), entangled states are used with only the exception [13] of when multi-particle product states are employed. On the other hand, in all those schemes [3,4,7,12,15,16] dealing with the QSS of quantum information, multiparticulate entangled states are used.

Recently, a particular quantum secure direct communica-

tion (QSDC) protocol has been proposed by Deng and Long [17], in which only a single photon state is used. In this paper, based on Deng-Long's QSDC protocol, we propose a scheme of multiparty quantum secret sharing of classical messages (QSSCM) by using only single photons. Then we take advantage of this multiparty QSSCM scheme to establish a scheme of multiparty secret sharing of quantum information (SSQI), where the secret is an arbitrary unknown quantum state in a qubit. We will show that multiparticulate entangled states are unnecessary in our multiparty SSQI scheme. Finally, we will propose a general idea for constructing multiparty SSQI schemes from any QSSCM scheme.

Now let us turn to our multiparty QSSCM scheme. For convenience, let us first describe a three-party QSSCM scheme. Suppose Alice wants to send a secret message to two distant parties Bob and Charlie. One of them, Bob or Charlie, is not entirely trusted by Alice, and she knows that if the two guys coexist, the honest one will keep the dishonest one from doing any damage. The two receivers, Bob and Charlie, can infer the secret message only by their mutual assistance. Our following three-party QSSCM scheme can achieve this goal with five steps.

(a) Bob prepares a batch of N single photons randomly in one of four polarization states $|H\rangle=|0\rangle$, $|V\rangle=|1\rangle$, $|u\rangle=1/\sqrt{2}(|0\rangle+|1\rangle)$, $|d\rangle=1/\sqrt{2}(|0\rangle-|1\rangle)$. For convenience, $\{|H\rangle, |V\rangle\}$ is referred to as the rectilinear basis and $\{|u\rangle, |d\rangle\}$ the diagonal basis hereafter. Then he sends this batch of photons to Charlie.

(b) After receiving these photons, for each photon Charlie randomly chooses a unitary operation from I , U , and U_H and performs this unitary operation on it. Here $I=|0\rangle\langle 0|+|1\rangle\langle 1|$ is an identity operator, $U=|0\rangle\langle 1|-|1\rangle\langle 0|$, and $U_H=(|0\rangle\langle 0|+|1\rangle\langle 0|+|0\rangle\langle 1|-|1\rangle\langle 1|)/\sqrt{2}$ is a Hadamard gate operator. The nice feature of the U operation is that it flips the state in both measuring bases, i.e., $U|0\rangle=-|1\rangle$, $U|1\rangle=|0\rangle$, $U|u\rangle=|d\rangle$, $U|d\rangle=-|u\rangle$. The nice feature of H is that it can realize the transformation between the rectilinear basis and the diagonal basis, i.e., $U_H|H\rangle=|u\rangle$, $U_H|V\rangle=|d\rangle$, $U_H|u\rangle=|H\rangle$, $U_H|d\rangle=|V\rangle$. After his encryptions, he sends the photons to Alice. The purpose of choosing a set of three unitary operations is to

protect the channel between Alice and Charlie from Bob's interception. For example, if Charlie chooses randomly a unitary operation from only I and U , Bob could intercept the channel between Alice and Charlie. He already has all the information about the state of the photon, and can readily check the transformation Alice did, so then he can retrieve the complete message without the help of Charlie.

(c) Alice stores most of the single photons and selects randomly a subset of single photons. Alice publicly announces the position of the selected photons. For each selected photon Alice randomly selects one action from the following two choices. One is that Alice lets Bob first tell her the initial state of the photon and then lets Charlie tell her which unitary operation he has performed on it; the other is that Charlie first tells Alice which unitary operation he has performed on the photon and then Bob tells Alice the initial state of the photon. Alice's strategy of choosing two actions is to prevent either Bob's or Charlie's intercept-resend attacks. Then Alice first performs the same unitary operation as Charlie has performed on the photon and then measures the photon by using the basis the initial state belongs to. After her measurements, Alice can determine the error rate. If the error rate exceeds the threshold, the process is aborted. Otherwise, the process continues and Alice performs unitary operations (either I or U) on the stored photons to encode her secret messages. That is, if Alice wants to encode a bit "0," she performs the identity unitary operation I ; if Alice wants to encode a bit "1," she performs the unitary operation $U = |0\rangle\langle 1| - |1\rangle\langle 0|$. Alice sends these encoded photons to Charlie.

(d) After Charlie receives these encoded photons, if Bob and Charlie collaborate, both Bob and Charlie can obtain Alice's secret message by using the correct measuring basis for each encoded photon. On the other hand, if Bob and Charlie do not collaborate, then both Bob and Charlie cannot get access to Alice's secret message with 100% certainty.

(e) Alice publicly announces a small part of her secret messages for Bob and Charlie to check whether the photons traveling from Alice's site to Charlie's site have been attacked, which is called message authentication. If the photons are attacked, the eavesdropper Eve cannot get access to any useful information but interrupts the transmissions.

So far we have proposed the three-party QSSCM scheme based on Deng and Long's QSDC protocol [17] by using single photons. The security of the present three-party QSSCM scheme is the same as the security of Deng and Long's QSDC protocol [17], that is, it depends completely on the step when Charlie sends the photon batch to Alice. As proven in Ref. [17], the scheme is also unconditionally secure. Incidentally, one can easily find that if Alice sends the encoded photons to Bob instead of Charlie, then the resultant scheme also works securely.

Now let us generalize the three-party QSSCM scheme to a n -party ($n \geq 4$) QSSCM scheme. Suppose that Alice is the message sender who would like to send a message to Bob, Charlie, Dick, ..., and Zach (there is a total of n receivers). The first step of the n -party ($n \geq 4$) QSSCM scheme is the same as that in the three-party QSSCM scheme. For completeness, this step is also included as follows.

(I) Bob prepares a batch of N single photons randomly in

one of four polarization states $|H\rangle = |0\rangle$, $|V\rangle = |1\rangle$, $|u\rangle = 1/\sqrt{2}(|0\rangle + |1\rangle)$, $|d\rangle = 1/\sqrt{2}(|0\rangle - |1\rangle)$. Then he sends this batch of photons to Charlie.

(II) After receiving these photons, for each photon Charlie randomly chooses a unitary operation from I , U , and U_H and performs this unitary operation on it. After his encryptions, he sends the encrypted photons to the next receiver, say, Dick. Dick randomly encrypts the encoded photons in the same way as Charlie, then sends the photons to the next receiver, and so on. Similar procedure is repeated until Zach finishes his encryptions. After Zach's encryptions, he sends the encrypted photons to Alice.

(III) Alice stores most of the single photons and randomly selects a subset of single photons. Then Alice publicly announces the position of the selected photons. To prevent any receiver's intercept-resend attack, for each selected photon, Alice randomly selects a receiver one by one and lets him or her tell her this receiver's message till she obtains all receivers's messages. Here Bob's message is the initial state of the photon, while Charlie's (Dick's, ..., Zach's) message is which unitary operation he has performed on the photon. Alice, in turn, performs the same unitary operations as Zach's, the $(n-2)$ th receiver's, ..., Dick's and Charlie's unitary operations on the photon and then measures this photon by using the basis the initial state belongs to. After her measurements, Alice can determine the error rate. If the error rate exceeds the threshold, the process is aborted. Otherwise, the process continues and Alice performs unitary operations (either I or U) on the stored photons to encode her secret messages. Alice sends these encoded photons to Zach.

(IV) After Zach receives these encoded photons, if Zach and the other $n-1$ receivers [Bob, Charlie, Dick, ..., the $(n-1)$ th receiver] collaborate, they can obtain Alice's secret message by using correct measuring basis for each encoded photon. On the other hand, if all the receivers do not collaborate, then none of them can get access to Alice's secret message with 100% certainty.

(V) Alice publicly announces a small part of her secret messages for all the receivers to check whether the photons traveling from Alice site to Zach's site have been attacked, which is called message authentication. If the photons are attacked, the eavesdropper Eve cannot get access to any useful information but interrupt the transmissions.

So far we have established a n -party QSSCM scheme by using single photons. The security of the present n -party QSSCM scheme is the same as the security of three-party QSSCM scheme, which is also unconditionally secure. Incidentally, as mentioned previously, one can easily find that if Alice sends the encoded photons to any other receiver instead of Zach, then the resultant scheme works securely also.

Now let us move to propose a multiparty SSQI scheme. Before this, let us briefly review the secure teleportation of an unknown quantum state [18,19]. Suppose that Alice wants to send to Bob an unknown state $\alpha|H\rangle_u + \beta|V\rangle_u$ in her qubit. Bob prepares a photon pair in any Bell state, say, $|\Phi^+\rangle_{ht} = (1/\sqrt{2})(|H\rangle_h|H\rangle_t + |V\rangle_h|V\rangle_t) = (1/\sqrt{2})(|u\rangle_h|u\rangle_t + |d\rangle_h|d\rangle_t)$. Bob sends the t photon to Alice. By randomly selecting one of the two sets of measuring basis, both Alice and Bob can check whether the quantum channel for photon transmission is at-

tacked or not according to their joint actions [24]. Suppose the quantum channel is safe and Bob successfully transmits a t photon to Alice. The state of the whole system can be rewritten as

$$\begin{aligned}
 & (\alpha|H\rangle_u + \beta|V\rangle_u)|\Phi^+\rangle_{ht} \\
 &= (\alpha|H\rangle_u + \beta|V\rangle_u) \frac{1}{\sqrt{2}}(|H\rangle_h|H\rangle_t + |V\rangle_h|V\rangle_t) \\
 &= \frac{1}{2}|\Phi^+\rangle_{ut}(\alpha|H\rangle_h + \beta|V\rangle_h) + \frac{1}{2}|\Psi^+\rangle_{ut}(\alpha|V\rangle_h + \beta|H\rangle_h) \\
 &\quad + \frac{1}{2}|\Phi^-\rangle_{ut}(\alpha|H\rangle_h - \beta|V\rangle_h) + \frac{1}{2}|\Psi^-\rangle_{ut}(\alpha|V\rangle_h - \beta|H\rangle_h),
 \end{aligned} \tag{1}$$

where $|\Psi^+\rangle = (|H\rangle_h|V\rangle_t + |V\rangle_h|H\rangle_t)/\sqrt{2}$, $|\Psi^-\rangle = (|H\rangle_h|V\rangle_t - |V\rangle_h|H\rangle_t)/\sqrt{2}$, and $|\Phi^-\rangle = (|H\rangle_h|H\rangle_t - |V\rangle_h|V\rangle_t)/\sqrt{2}$. Hence, if Alice performs a Bell-state measurement on the two photons in her lab and tells Bob her measurement outcome, say, $|\Phi^+\rangle$ ($|\Psi^+\rangle$, $|\Psi^-\rangle$, $|\Phi^-\rangle$), then Bob can perform a unitary operation $I = |H\rangle\langle H| + |V\rangle\langle V|$ ($u_1 = |H\rangle\langle V| + |V\rangle\langle H|$, $u_2 = |H\rangle\langle H| - |V\rangle\langle V|$, $u_3 = |H\rangle\langle V| - |V\rangle\langle H|$) to reconstruct the unknown state in the qubit h . Since the teleportation is based on EPR pairs, so the proof of the security is the same in essence as those in Refs. [20–24]. This is the secure teleportation of an unknown state in a qubit. In such teleportation, Alice's public announcement of the Bell-state measurement outcome is a necessary step, otherwise, Bob cannot reconstruct the unknown state in his retained qubit.

Our multiparty SSQI scheme ($n \geq 3$) is almost the same as the secure teleportation of an unknown quantum state as mentioned above, except for one point. Alice would like to send an unknown quantum state to Bob, Charlie, Dick, ..., and Zach (there is totally n receivers). To do this, she sends the unknown quantum state to Bob by teleportation. But in-

stead of public announcement of the Bell-state measurement outcome, Alice distributes her Bell-state measurement outcome to $n-1$ receivers without Bob by use of the QSSCM (quantum secret sharing of classical messages) protocol we just proposed. To reconstruct an unknown state in a qubit, all n receivers must collaborate.

In our multiparty SSQI (secret sharing of quantum information) scheme, the multi-particle GHZ states in all other existing multiparty SSQI schemes [3,4,7] are not necessary. Although in Ref. [15] it is claimed that only Bell states are needed, the identification of multiparticle GHZ state is necessary. In our multiparty SSQI protocol, only during the teleportation step are the use and identification of Bell states needed. In all other steps, single photon states are enough. Hence, the present multiparty SSQI scheme is more feasible with present-day technique [25].

As a matter of fact, till now there have been many existing multiparty QSSCM (quantum secret sharing of classical messages) schemes [3,5,6,8–11,13,14]. Each of them can be combined with the secure quantum teleportation of an unknown state to establish a multiparty SSQI scheme. Hence the idea of combining a secure teleportation of an unknown state with a QSSCM scheme to set up a SSQI scheme in the present paper is a general one.

To summarize, in this paper by using single photon state instead of entangled states (Bell states or multiparticle GHZ states) or of multiphoton product states we have presented a multiparty QSSCM scheme based on Deng and Long's QSDC protocol. We have also proposed a multiparty SSQI scheme by taking advantage of our multiparty QSSCM scheme. The idea of combining a multiparty QSSCM scheme with the secure quantum teleportation to establish a multiparty SSQI scheme is general.

Zhang thanks Professor Baiwen Li for his encouragement. This work is supported by the National Natural Science Foundation of China under Grant No. 10304022.

-
- [1] B. Schneier, *Applied Cryptography* (Wiley, New York, 1996), p. 70.
- [2] J. Gruska, *Foundations of Computing* (Thomson Computer Press, London, 1997), p. 504.
- [3] M. Hillery, V. Buzk, and A. Berthiaume, *Phys. Rev. A* **59**, 1829 (1999).
- [4] R. Cleve, D. Gottesman, and H. K. Lo, *Phys. Rev. Lett.* **83**, 648 (1999).
- [5] A. Karlsson, M. Koashi, and N. Imoto, *Phys. Rev. A* **59**, 162 (1999).
- [6] D. Gottesman, *Phys. Rev. A* **61**, 042311 (2000).
- [7] S. Bandyopadhyay, *Phys. Rev. A* **62**, 012308 (2000).
- [8] W. Tittel, H. Zbinden, and N. Gisin, *Phys. Rev. A* **63**, 042301 (2001).
- [9] V. Karimipour and A. Bahraminasab, *Phys. Rev. A* **65**, 042320 (2002).
- [10] H. F. Chau, *Phys. Rev. A* **66**, 060302 (2002).
- [11] S. Bagherinezhad and V. Karimipour, *Phys. Rev. A* **67**, 044302 (2003).
- [12] L. Y. Hsu, *Phys. Rev. A* **68**, 022306 (2003).
- [13] G. P. Guo and G. C. Guo, *Phys. Lett. A* **310**, 247 (2003).
- [14] L. Xiao, G. L. Long, F. G. Deng and J. W. Pan, *Phys. Rev. A* **69**, 052307 (2004).
- [15] Y. M. Li, K. S. Zhang and K. C. Peng, *Phys. Lett. A* **324**, 420 (2004).
- [16] A. M. Lance, T. Symul, W. P. Bowen, B. C. Sanders and P. K. Lam, *Phys. Rev. Lett.* **92**, 177903 (2004).
- [17] F. G. Deng and G. L. Long, *Phys. Rev. A* **69**, 052319 (2004).
- [18] Z. X. Man, Z. J. Zhang, and Y. Li, *Chin. Phys. Lett.* **22**, 18 (2005).
- [19] C. H. Bennett, G. Brassard C. Crepeau, R. Jozsa, A. Peres, and W. K. Wothers, *Phys. Rev. Lett.* **70**, 1895 (1993).
- [20] C. H. Bennett, G. Brassard, and N. D. Mermin, *Phys. Rev. Lett.* **68**, 557 (1992).

- [21] H. Inamori, L. Rallan, and V. Verdral, *J. Phys. A* **34**, 6913 (2001).
- [22] G. L. Long and X. S. Liu, *Phys. Rev. A* **65**, 032302 (2002).
- [23] E. Waks, A. Zeevi, and Y. Yamamoto, *Phys. Rev. A* **65**, 052310 (2002).
- [24] F. G. Deng, G. L. Long, and X. S. Liu, *Phys. Rev. A* **68**, 042317 (2003).
- [25] Y. H. Kim, S. P. Kulik, and Y. Shih, *Phys. Rev. Lett.* **86**, 1370 (2001).